

# CosmiXs: Cosmic messenger spectra for indirect dark matter searches

**M. Di Mauro**

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arXiv:2411.04815



**Darktools, June 19th Torino 2025**

[https://theconversation.com/  
why-do-astronomers-believe-in-  
dark-matter-122864](https://theconversation.com/why-do-astronomers-believe-in-dark-matter-122864)

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Istituto Nazionale di Fisica Nucleare



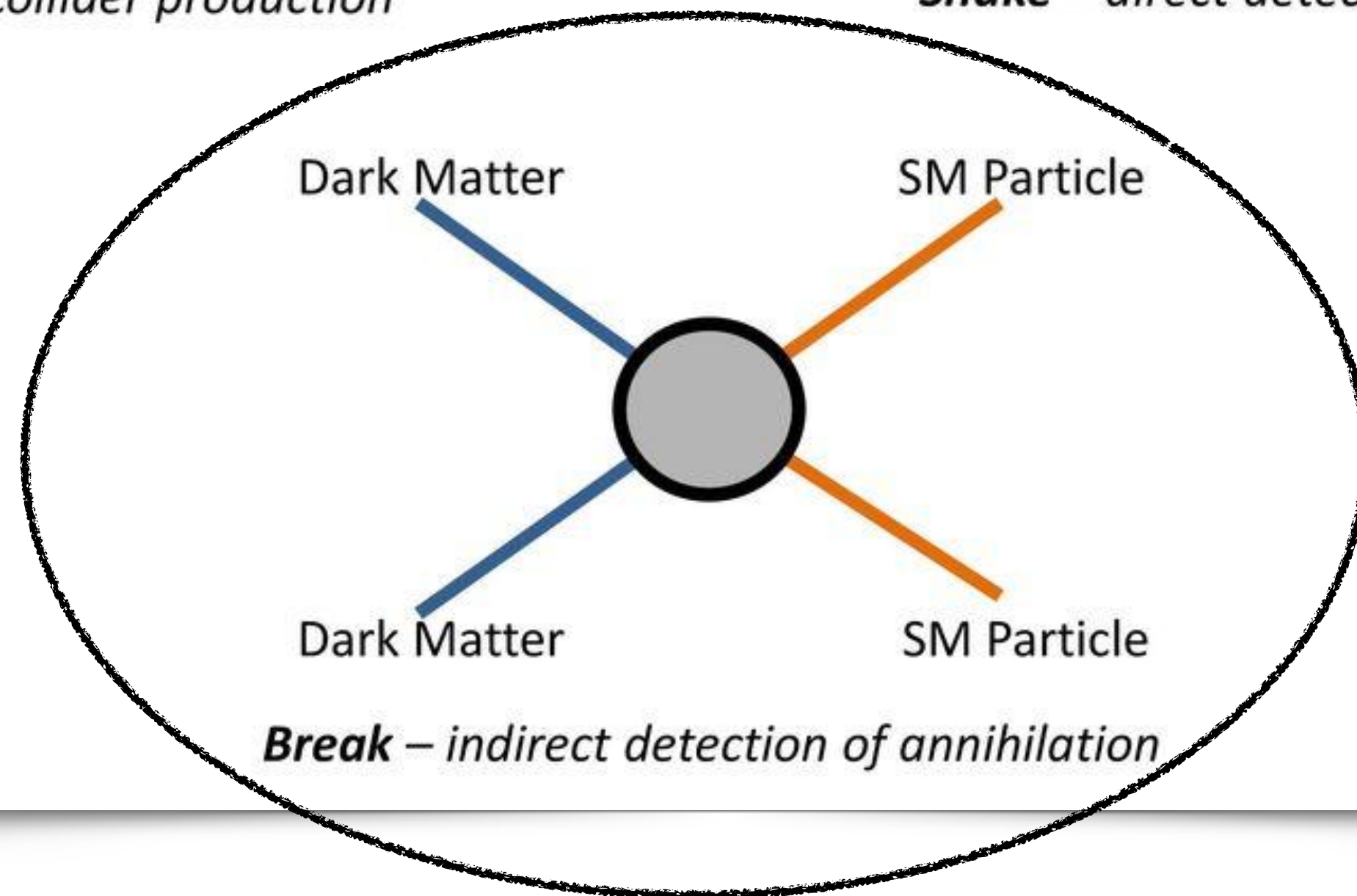
# Dark matter search strategies

## Ways to Detect Dark Matter – *Make, Shake and Break*

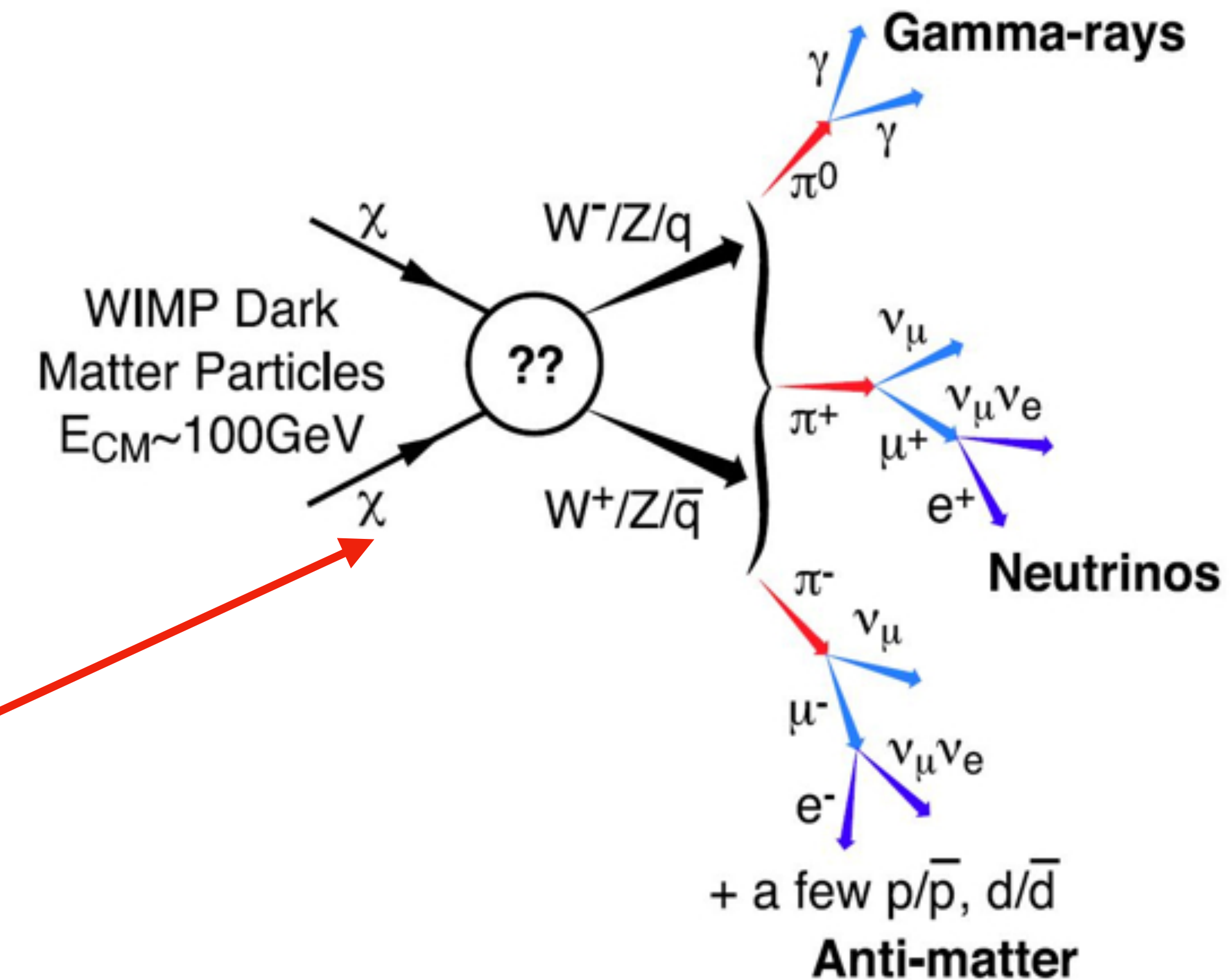


**Make** – collider production

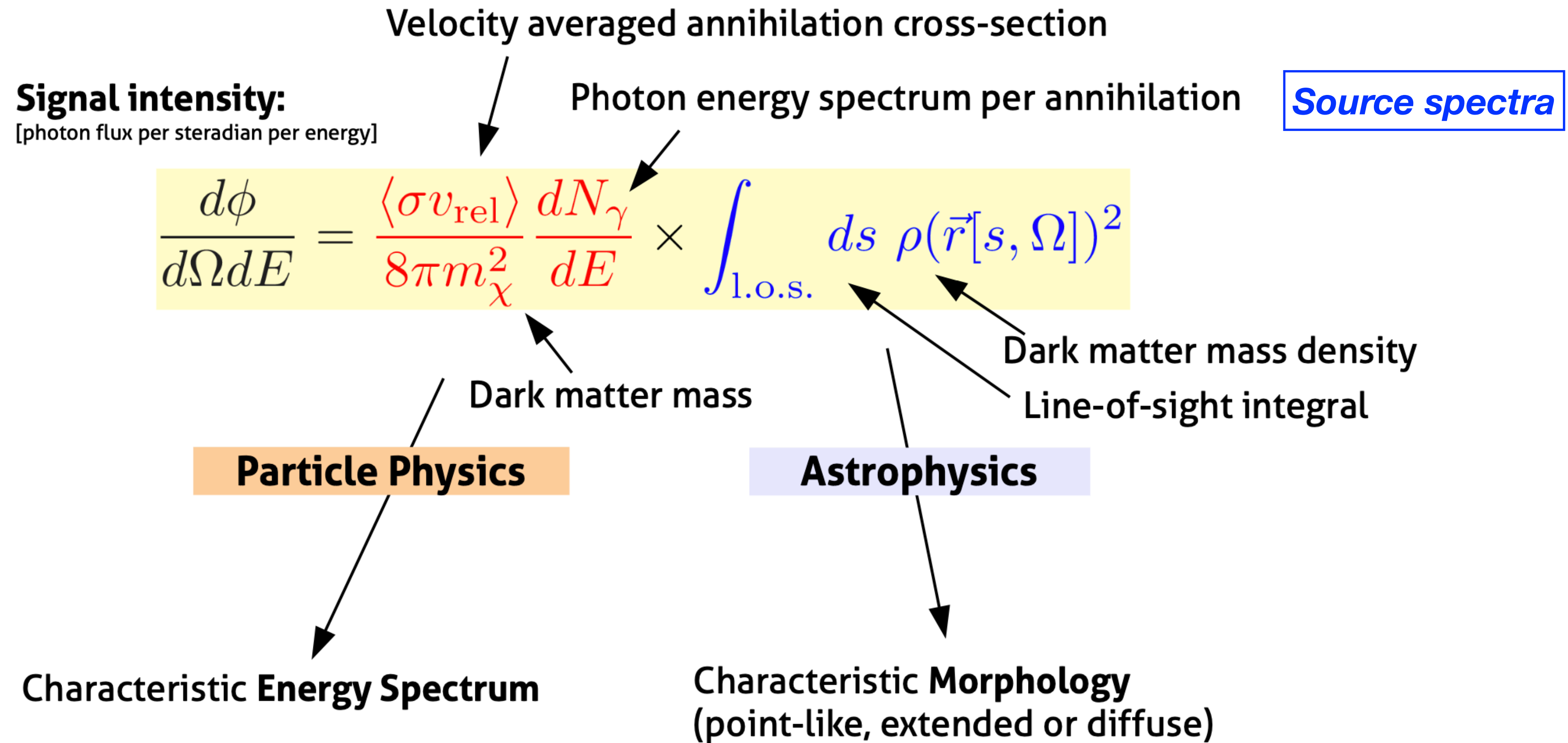
**Shake** – direct detection scattering



**Break** – indirect detection of annihilation



# Gamma-ray flux from dark matter



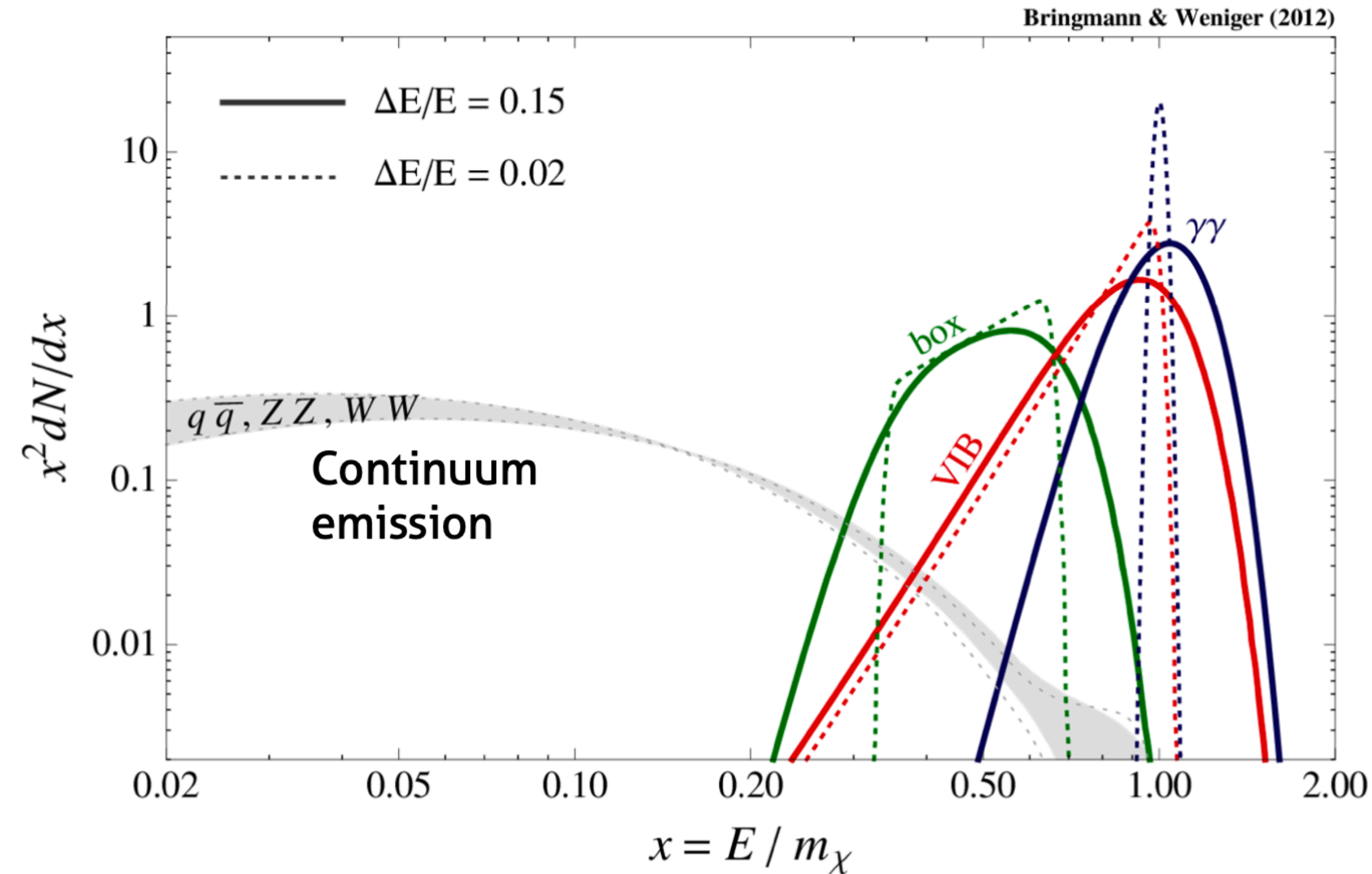
[review DM searches with gamma rays: Bringmann & Weniger (2012)]

It is convenient to define a "J-value":

$$J_{\Delta\Omega} \equiv \int_{\Delta\Omega} d\Omega \int_{\text{l.o.s.}} ds \rho(r[s, \vec{\Omega}])^2$$



# Spectral features of gamma rays from dark matter



## Box-shaped spectra

- Cascade-decay into monochromatic photons
- already at tree level

## Internal Bremsstrahlung (IB)

- radiative correction to processes with charged final states
- Generically suppressed by  $O(\alpha)$

$$\chi\chi \rightarrow \bar{f}f\gamma$$

## Gamma-ray lines

- from two-body annihilation into photons
- forbidden at tree-level, generically suppressed by  $O(\alpha^2)$

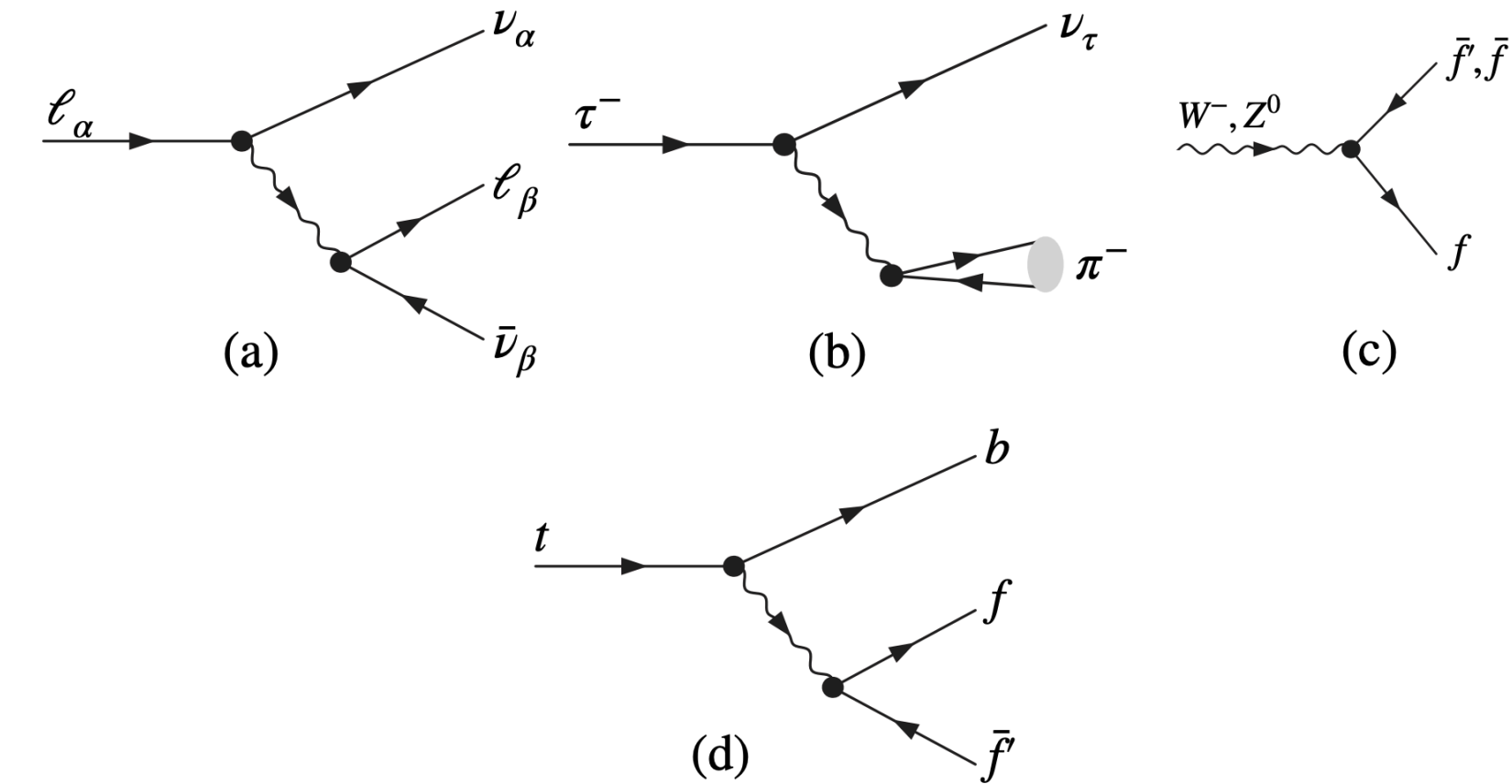
$$\chi\chi \rightarrow \gamma\gamma$$



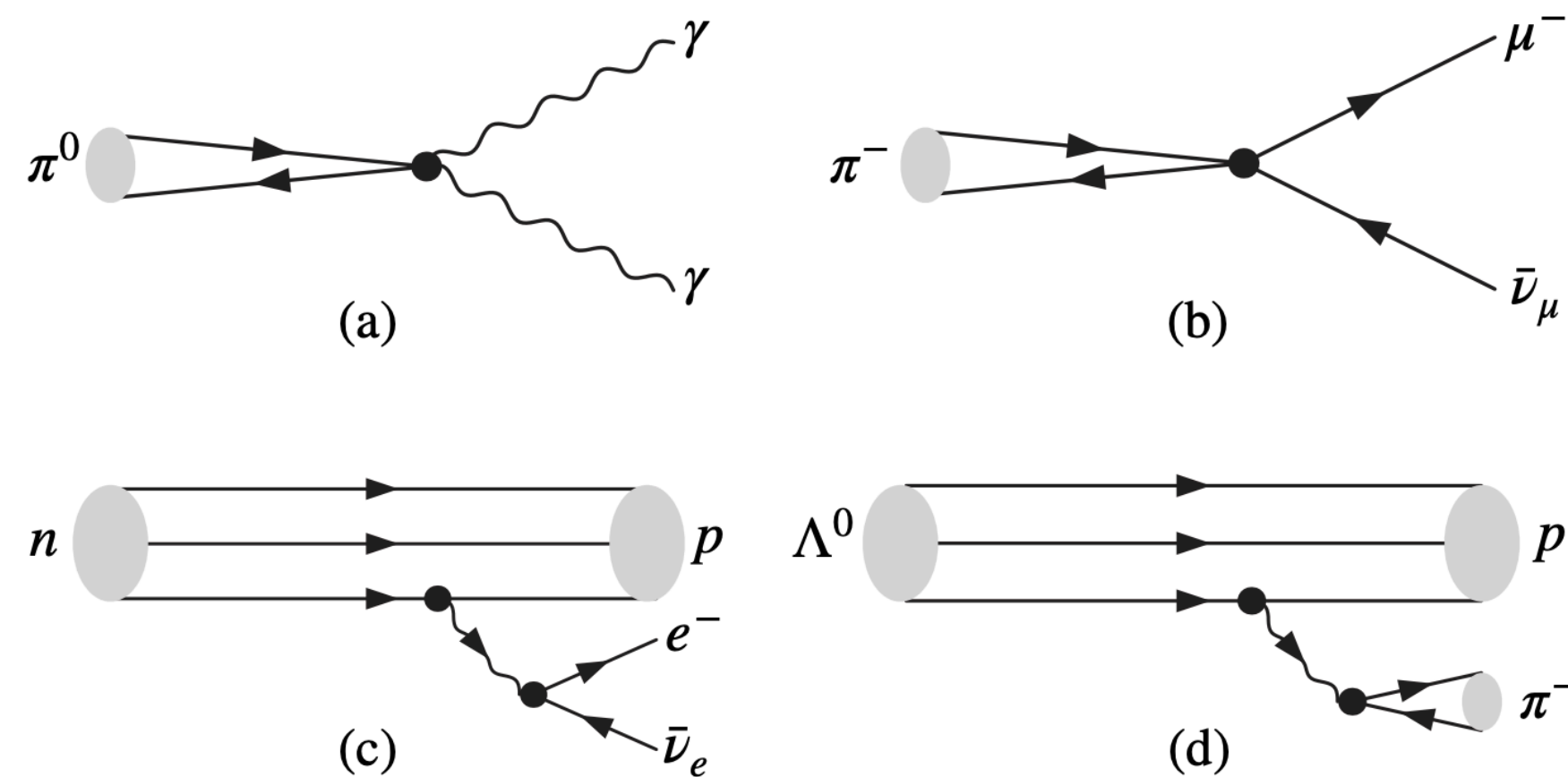
# Summary of the possible processes

$$\chi\chi \rightarrow \underbrace{\left[ X_1 X_2 \dots X_N \right]}_{\text{Intermediate states}} \rightarrow \overbrace{\left( Y_{11} \dots Y_{1a_1} \right) \dots \left( Y_{N1} \dots Y_{Na_N} \right)}^{\text{Stable particles}}.$$

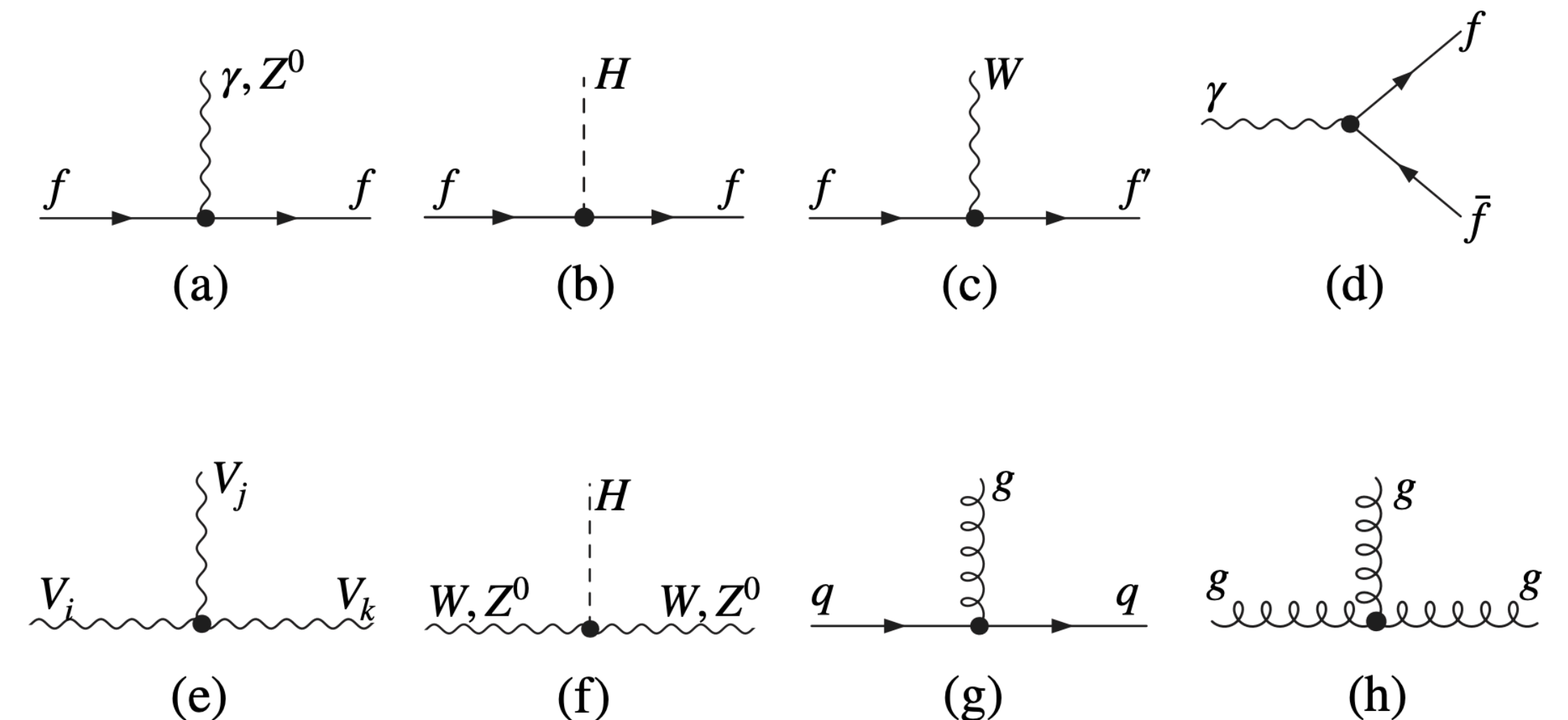
## Leading order EW interactions



## Baryon and meson decays



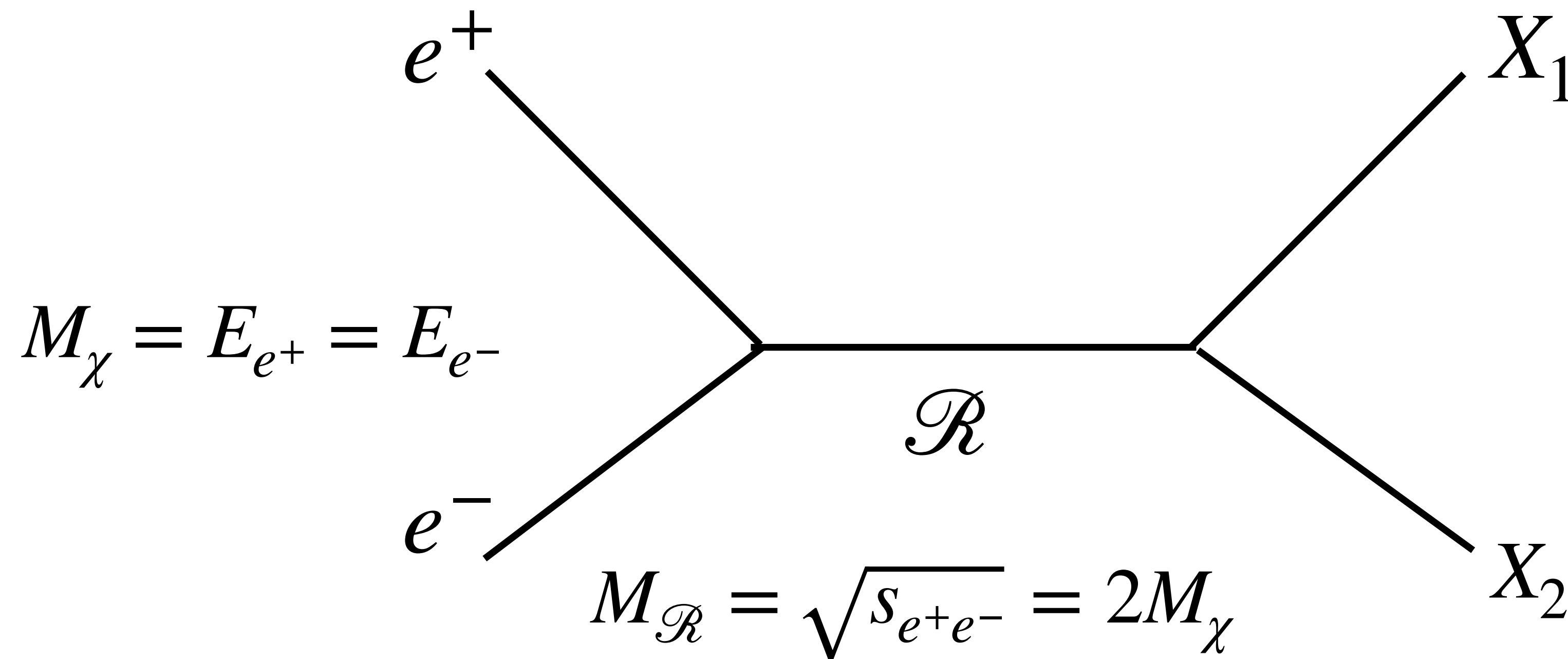
## Bremsstrahlung and EW corrections





# Spectra calculated with the *resonance approach*

- Standard tool to calculate the spectra using PYTHIA (PPPC4DMID).
- This case resemble the annihilation of fermionic DM (e.g. neutralinos).
- Spin information is lost as the outgoing particles do not have assigned helicities and polarisation.
- Electroweak corrections will not be taken properly into account.





# State of the art: PPPC 4 DM ID

CERN-PH-TH/2010-057

SACLAY-T10/025

IFUP-TH/2010-44

## PPPC 4 DM ID: A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection

Marco Cirelli<sup>a,b</sup>, Gennaro Corcella<sup>c,d,e</sup>, Andi Hektor<sup>f</sup>,  
Gert Hütsi<sup>g</sup>, Mario Kadastik<sup>f</sup>, Paolo Panci<sup>a,h,i,j</sup>,  
Martti Raidal<sup>f</sup>, Filippo Sala<sup>d,e</sup>, Alessandro Strumia<sup>a,e,f,k</sup>

- They used Pythia 8.135 (about 13 years old) to calculate, with the resonance approach, DM spectra for different annihilation channels and masses from 5 GeV to 100 TeV.
- EW corrections are added, without resummation, by hand on top of Pythia results (matching issue?) (Ciafaloni et al. 2010).
- *Large cutoff on the minimum transverse momentum for photons emitted off lepton lines in the shower.*
- Polarization and helicity information is absent during the showering.
- Off-shell effects for the EW Gauge boson channels were not taken into account.



# State of the art: HDMS

## Dark Matter Spectra from the Electroweak to the Planck Scale

Will be denoted by HDMS in what follows

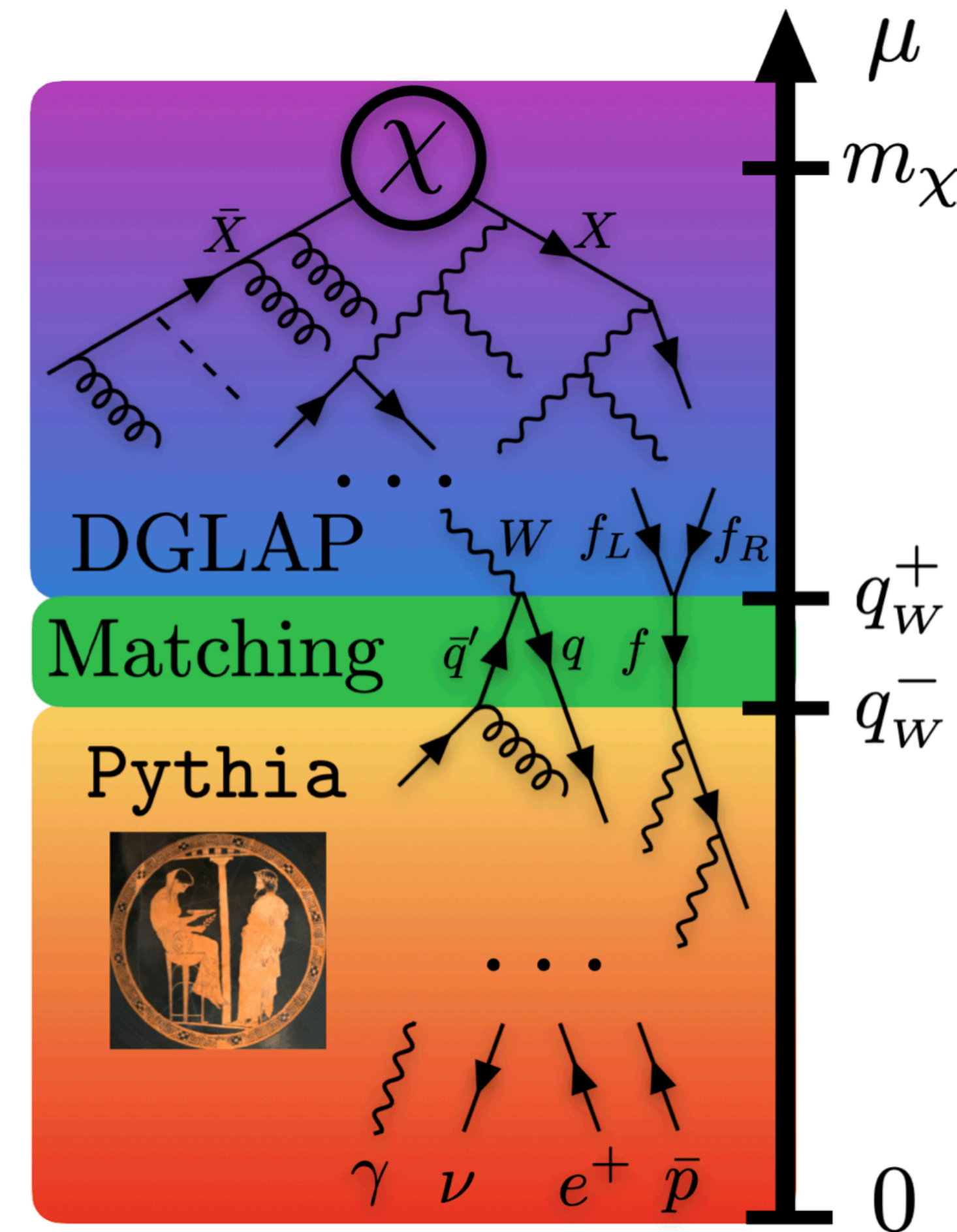
Christian W. Bauer,<sup>1,2</sup> Nicholas L. Rodd,<sup>1,2</sup> Bryan R. Webber<sup>3</sup>

<sup>1</sup>*Berkeley Center for Theoretical Physics, University of California, Berkeley, CA 94720, USA*

<sup>2</sup>*Theoretical Physics Group, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA*

<sup>3</sup>*University of Cambridge, Cavendish Laboratory, J.J. Thomson Avenue, Cambridge, UK*

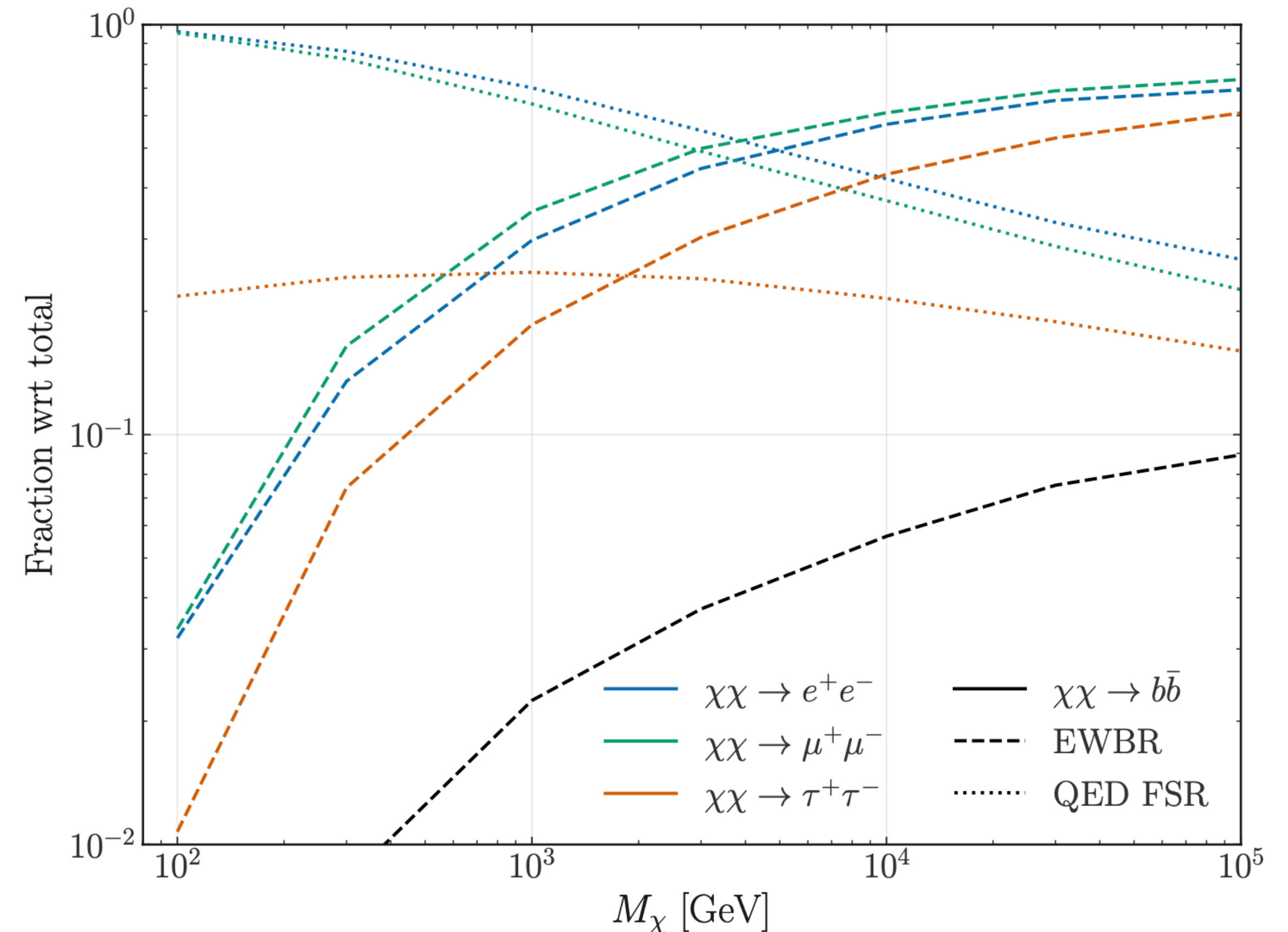
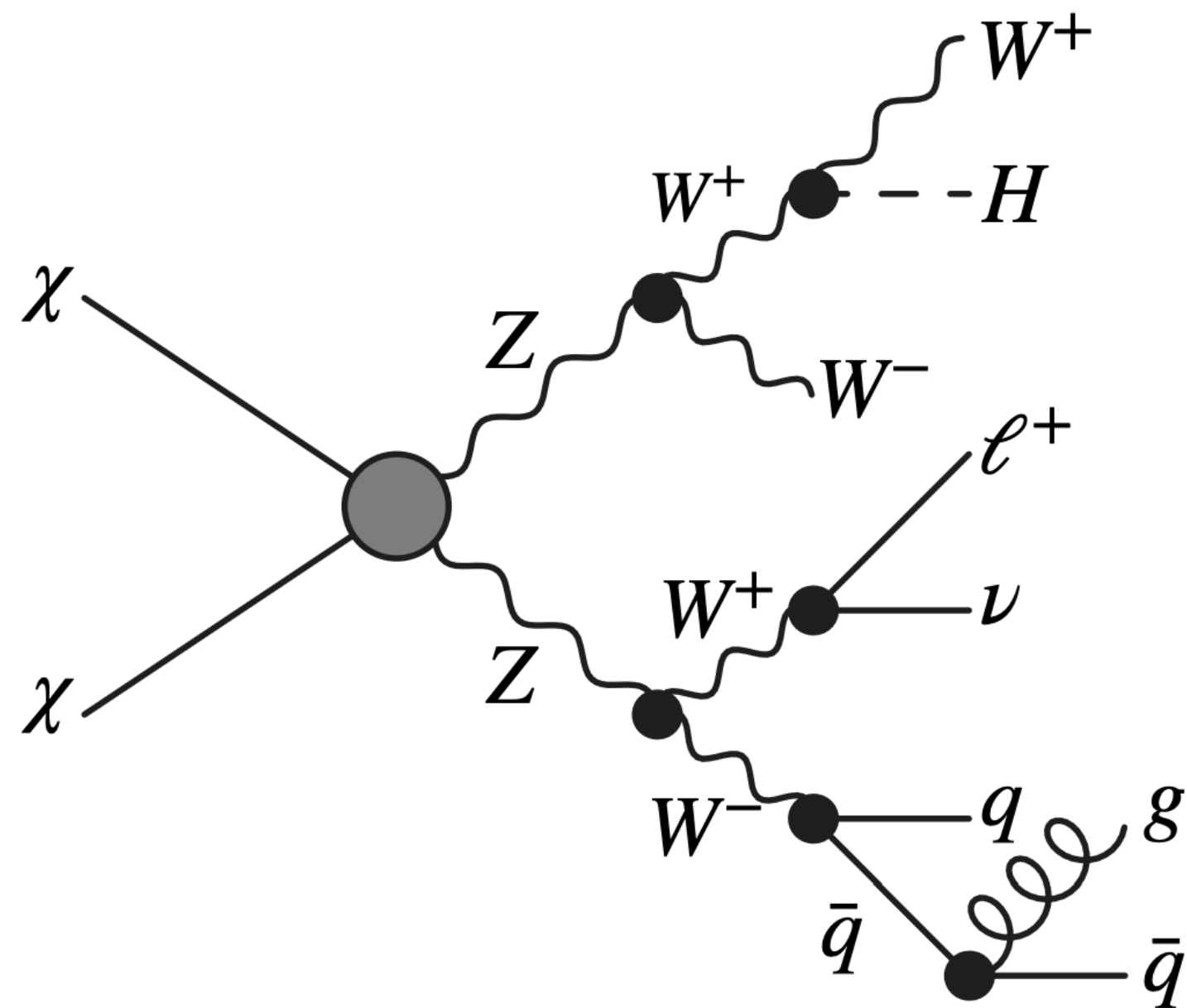
- They have provided decay spectra for DM with masses  $>500$  GeV.
- Spectra of dark matter annihilation/decay were calculated using analytical methods (DGLAP evolution equations) and matched to PYTHIA at the electroweak scale.
- The results of HDMS can lead to theoretical problems in the matching between the physics generated by the DGLAP formalism in the unbroken phase and the physics described by Pythia





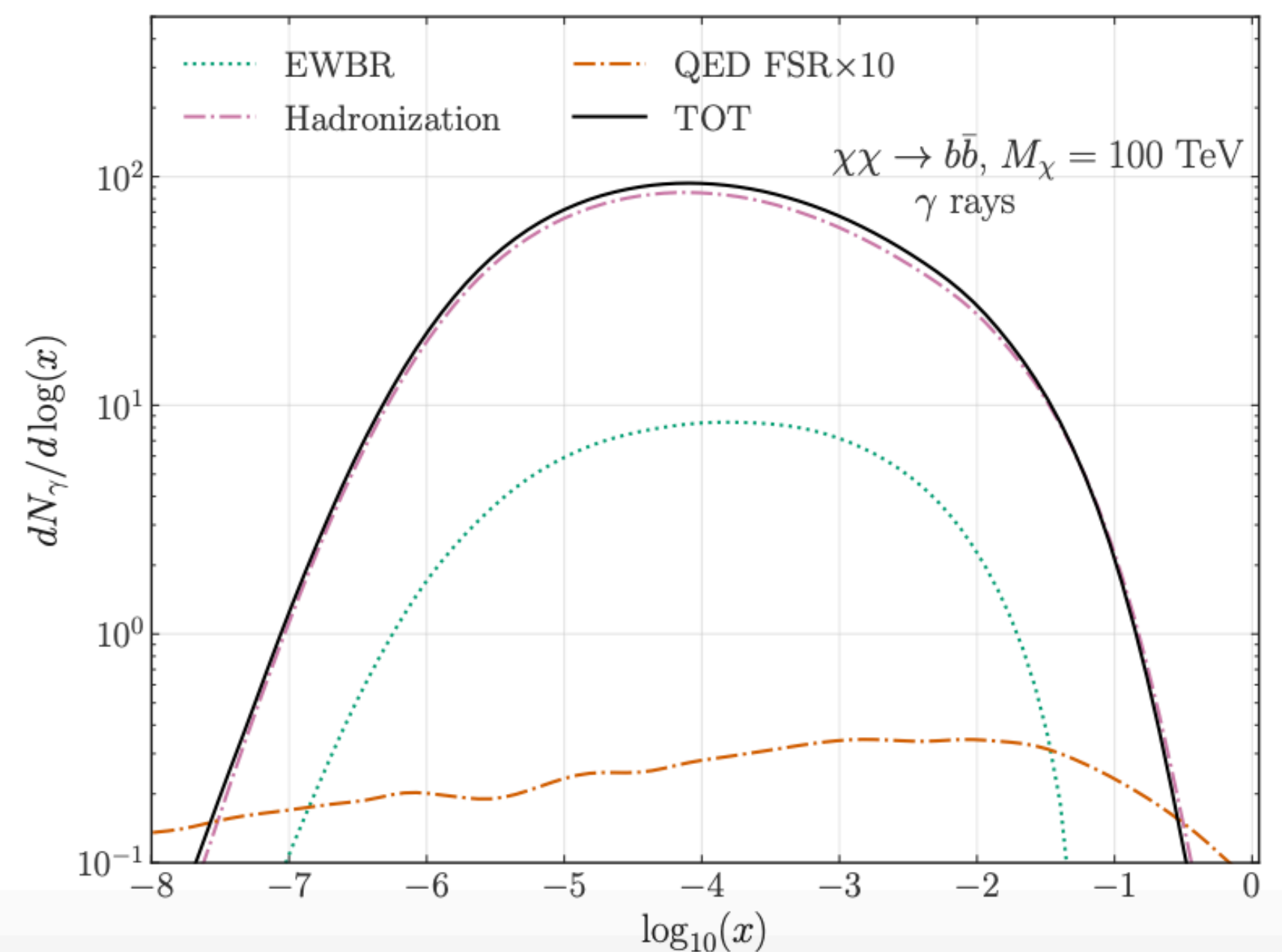
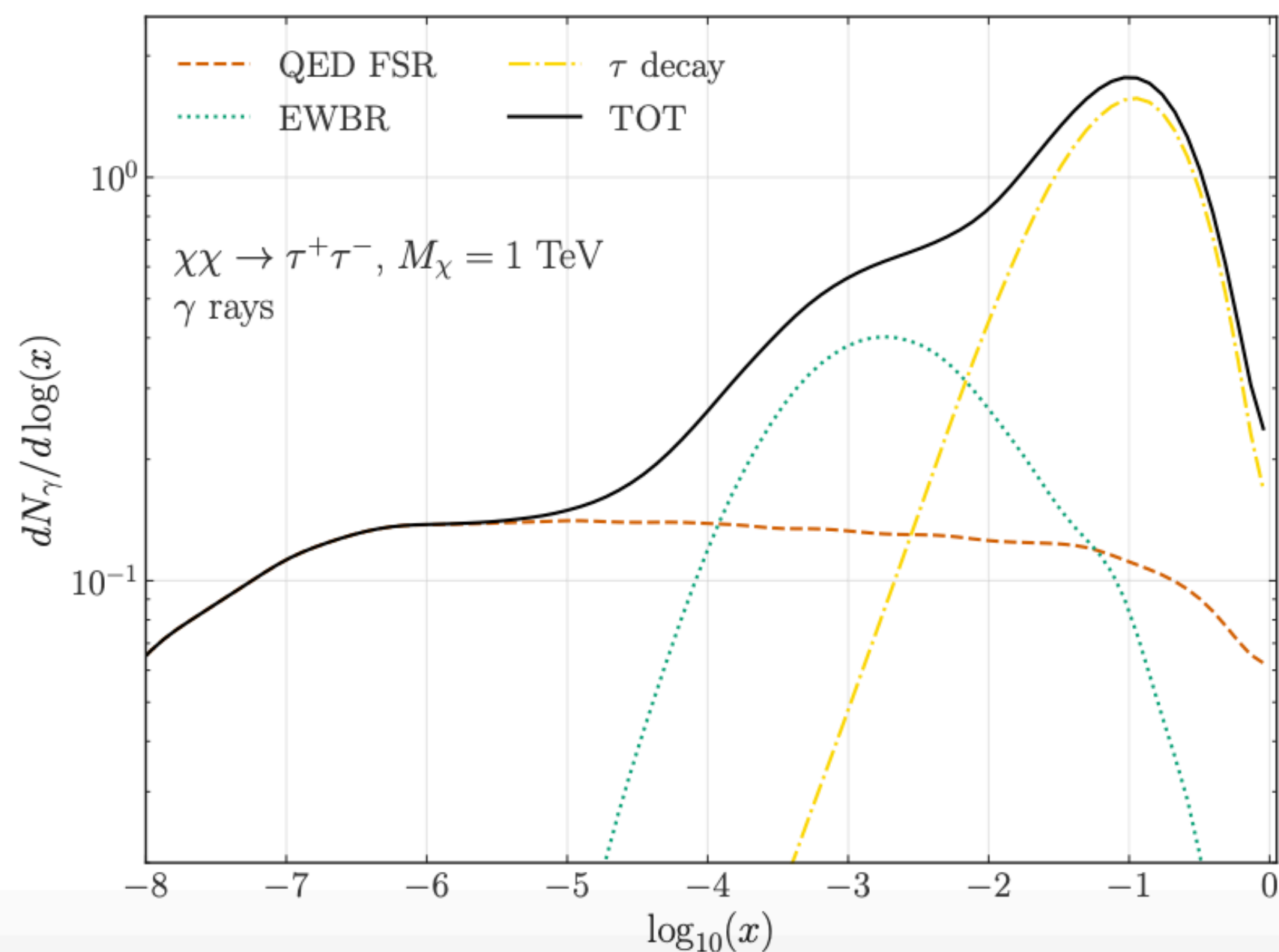
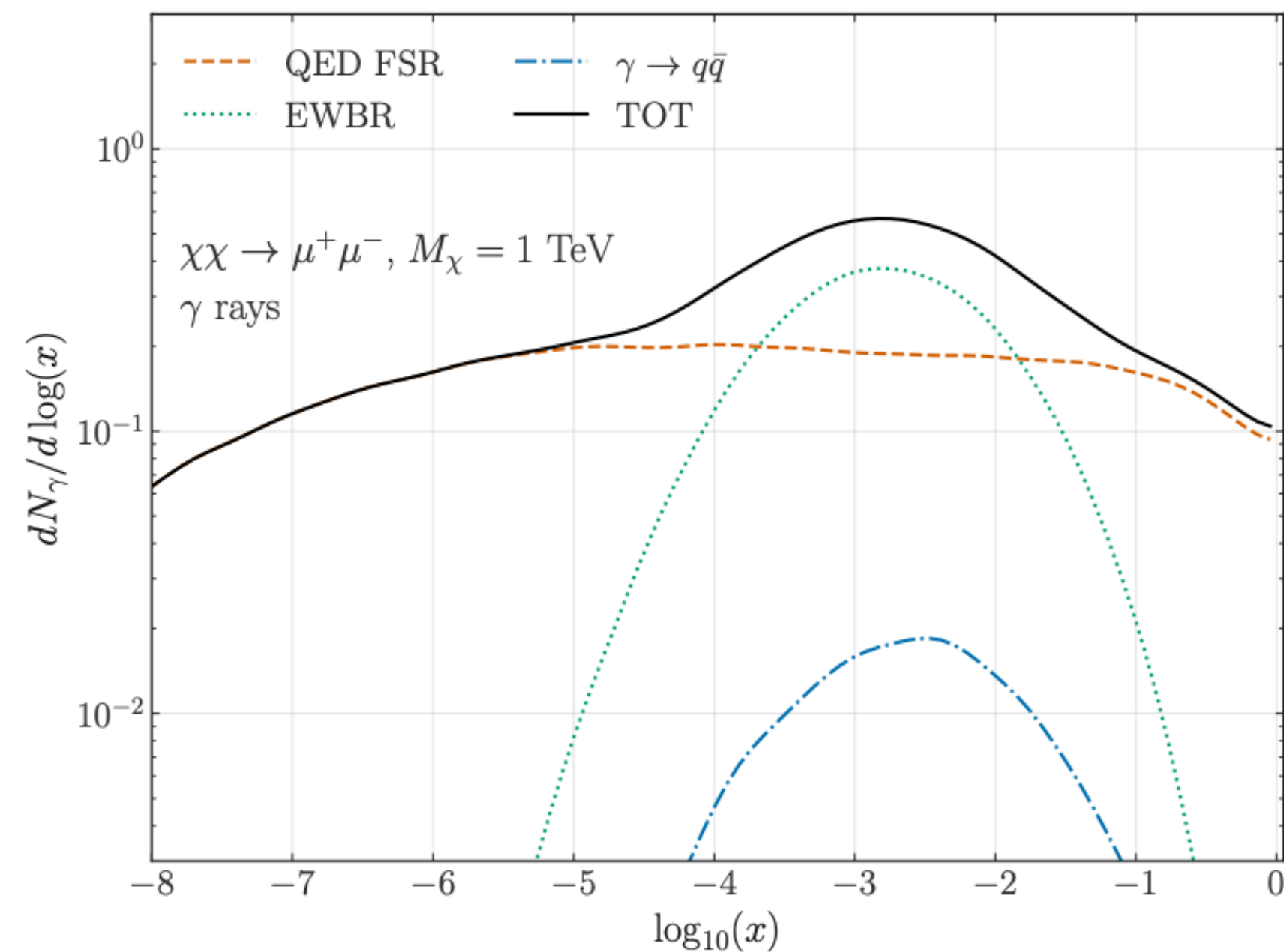
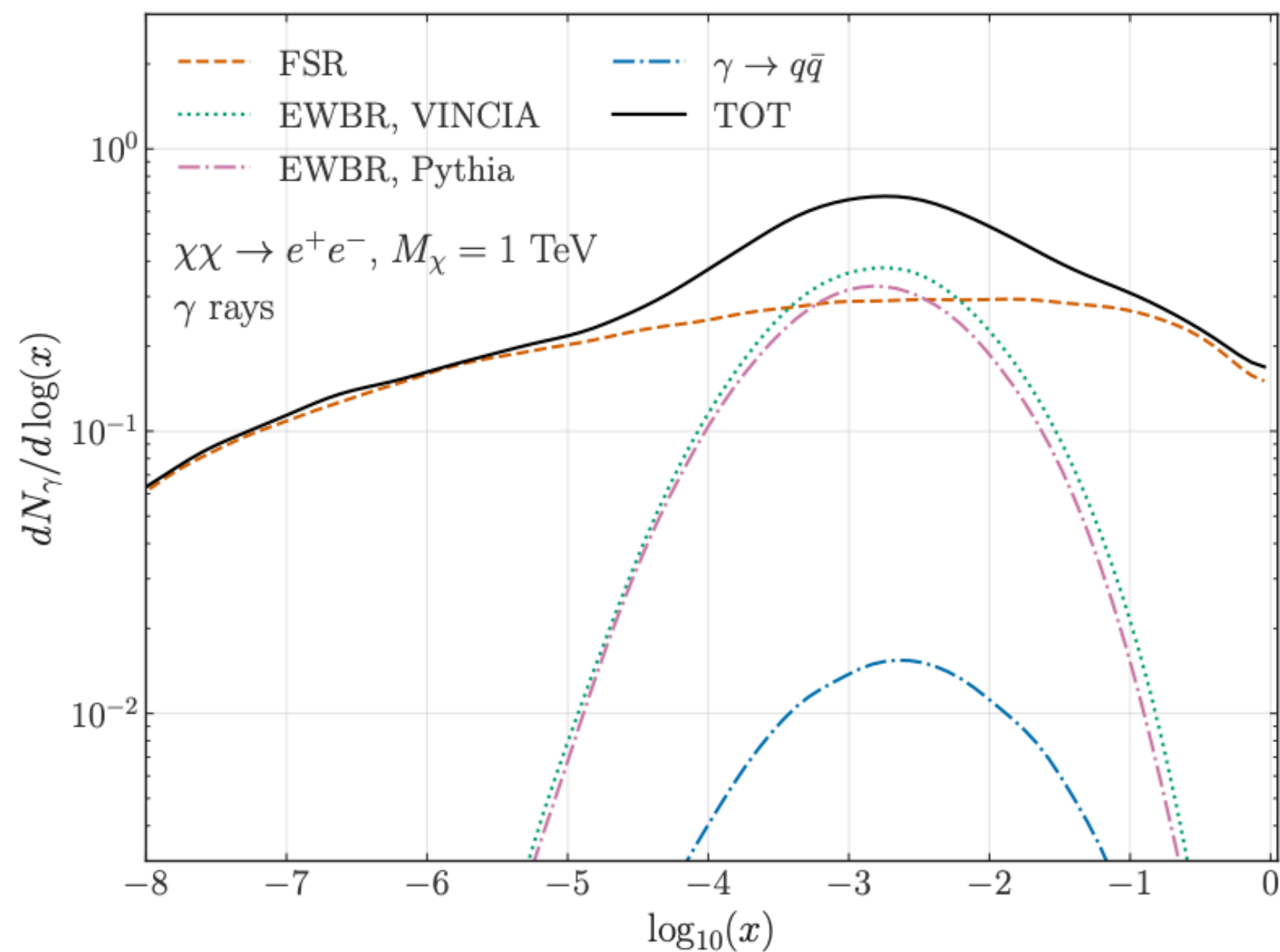
# Electroweak boson radiation

- We use the Vincia algorithm (Pythia v. 8.309) which we interface with MadDM.
- We include contributions from triple gauge boson interaction.
  - We include subsequent radiation of Gauge bosons
  - Helicity and polarization info are considered across the entire showering.





# Components in the spectra

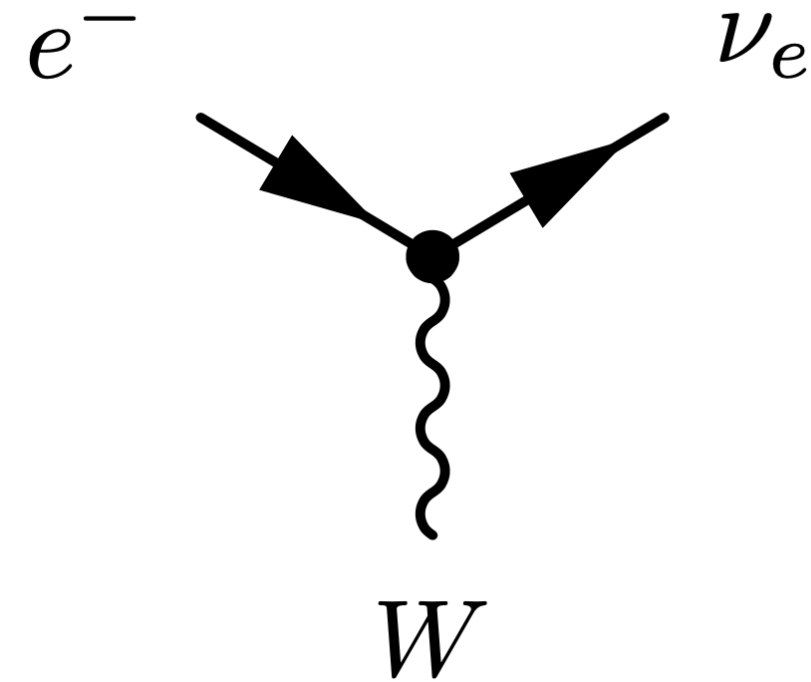




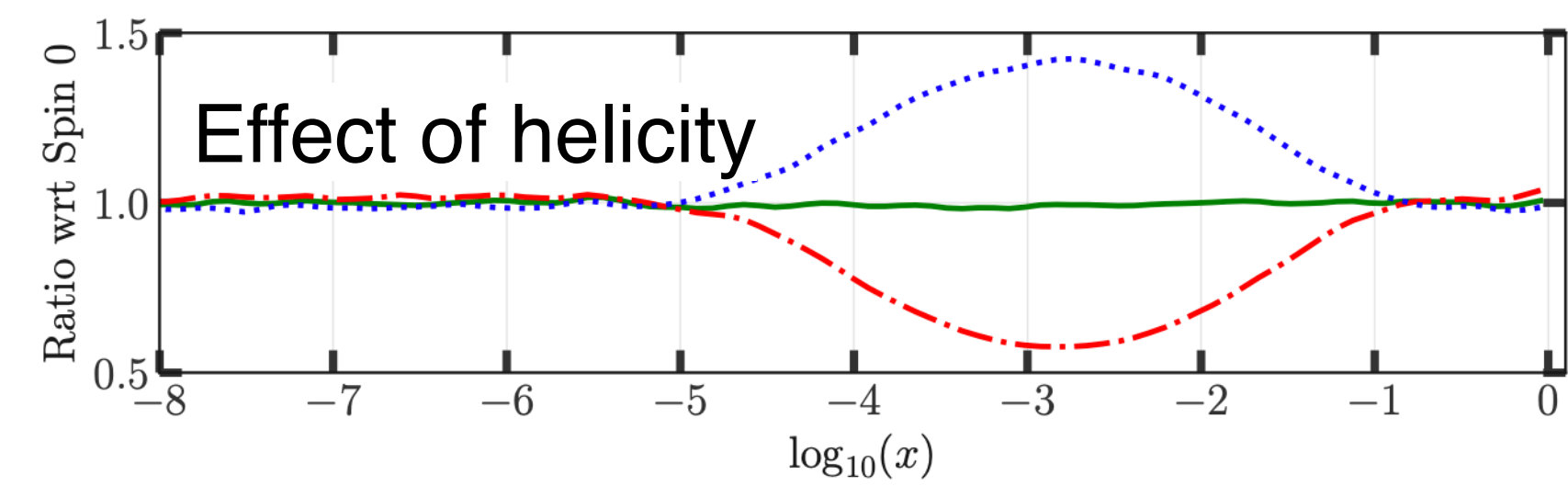
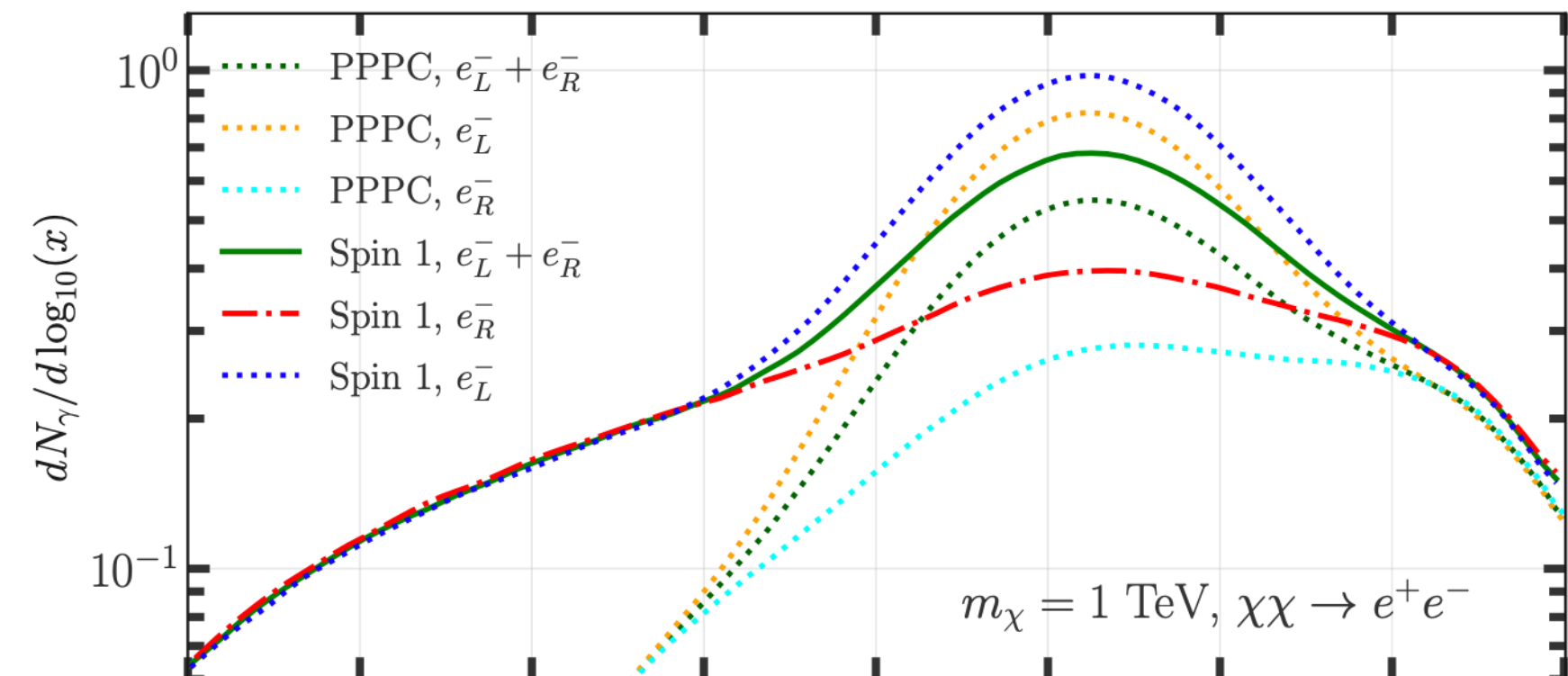
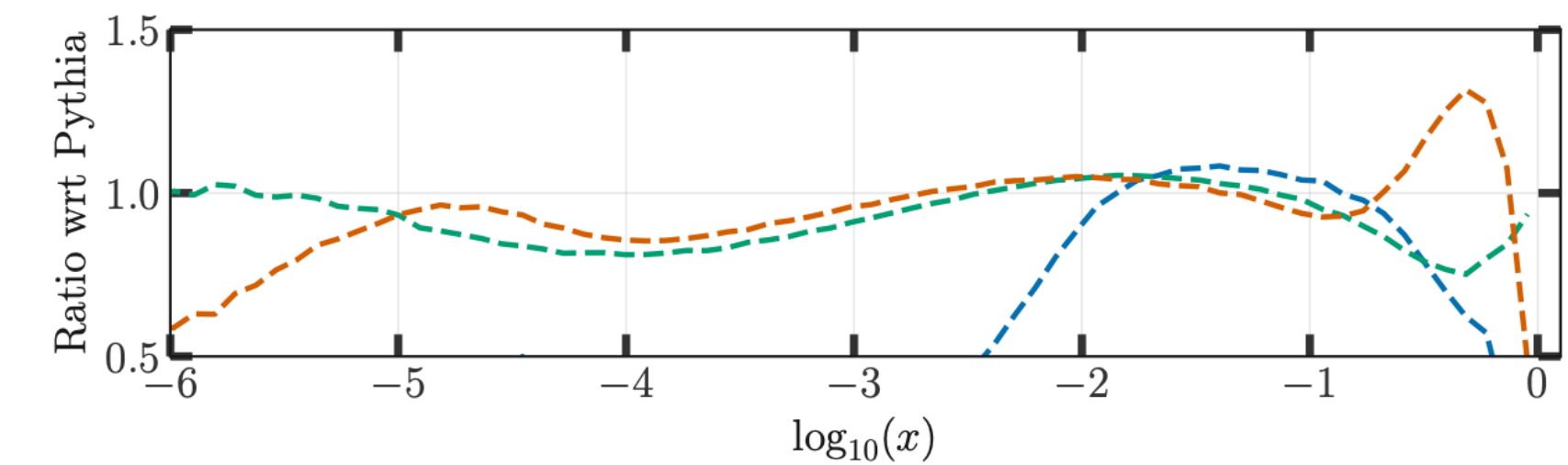
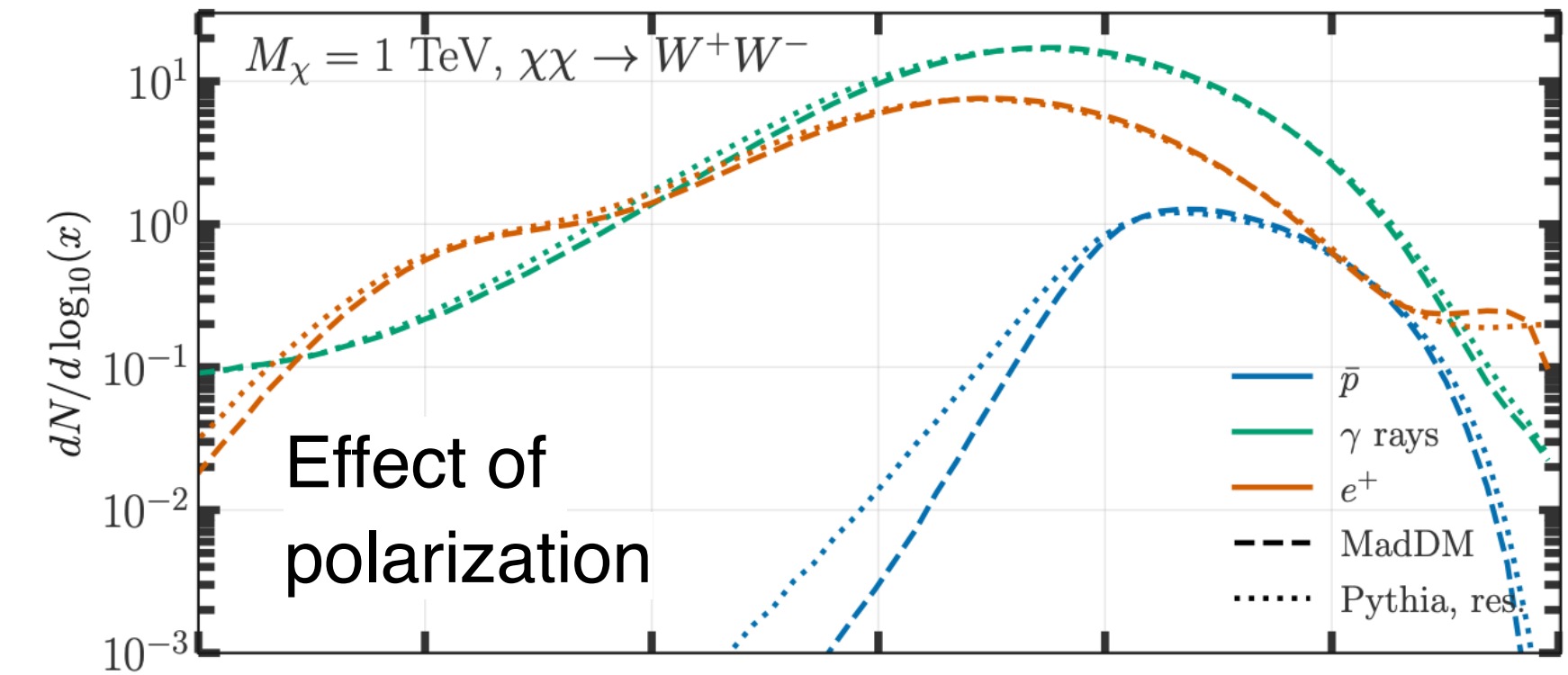
# Polarisation and helicity information

- We include polarisation and helicity of particles.
- We include annihilation channels with off-shell gauge bosons.

$$\begin{aligned}\mathcal{L}_{\text{I,L}}^{(\text{CC})} &= -\frac{g}{\sqrt{2}} \{ \bar{\nu}_{eL} W e_L + \bar{e}_L W^\dagger \nu_{eL} \} \\ &= -\frac{g}{2\sqrt{2}} \bar{\nu}_e \gamma^\mu (1 - \gamma^5) e W_\mu + \text{H.c.}\end{aligned}$$



$$\bar{\chi} \gamma_\mu (g_\chi^V + \gamma_5 g_\chi^A) \chi Y_1^\mu + \sum_i^{N_f} \bar{F}_i \gamma_\mu (g_{ij}^V + \gamma_5 g_{ij}^A) F_j Y_1^\mu,$$

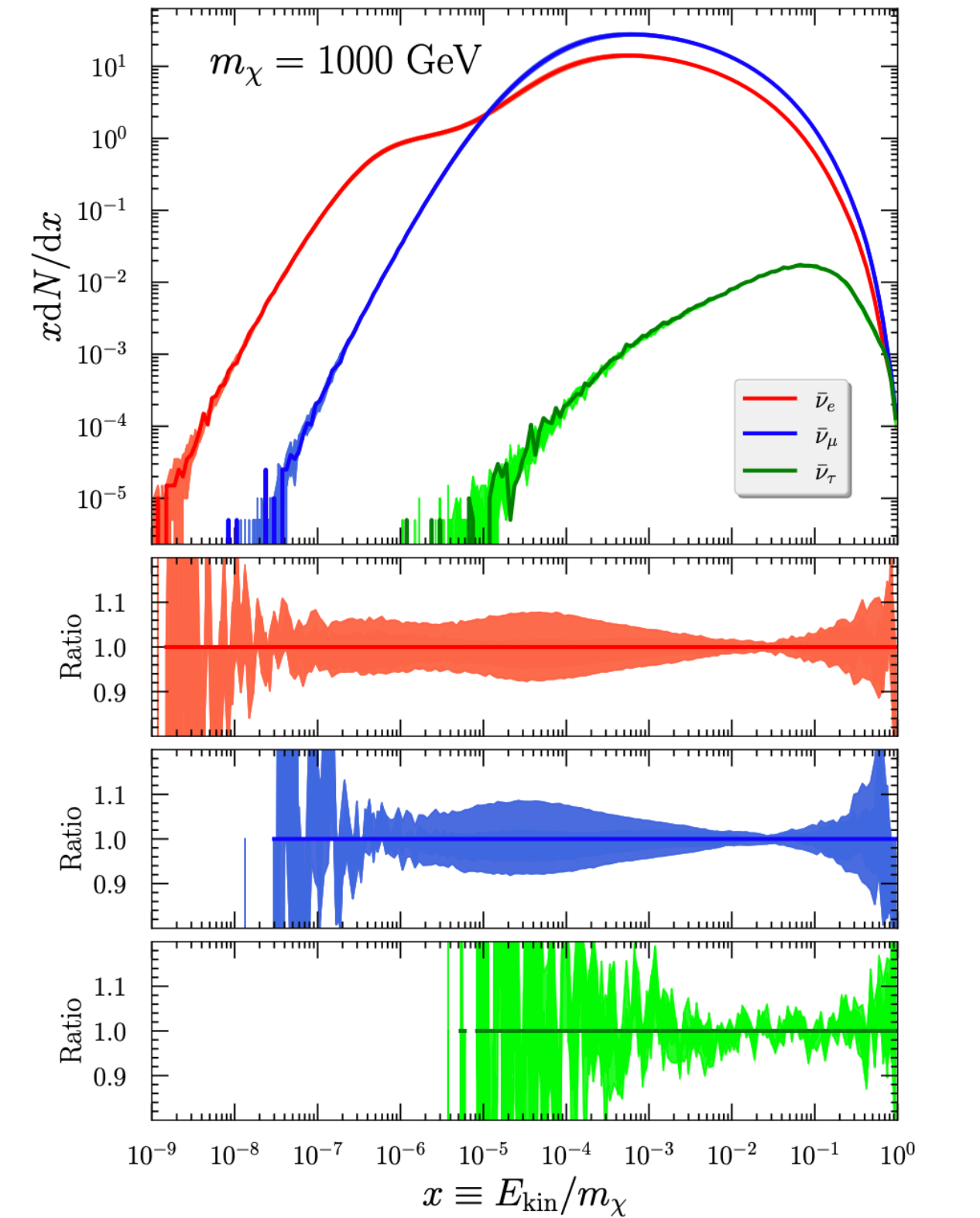
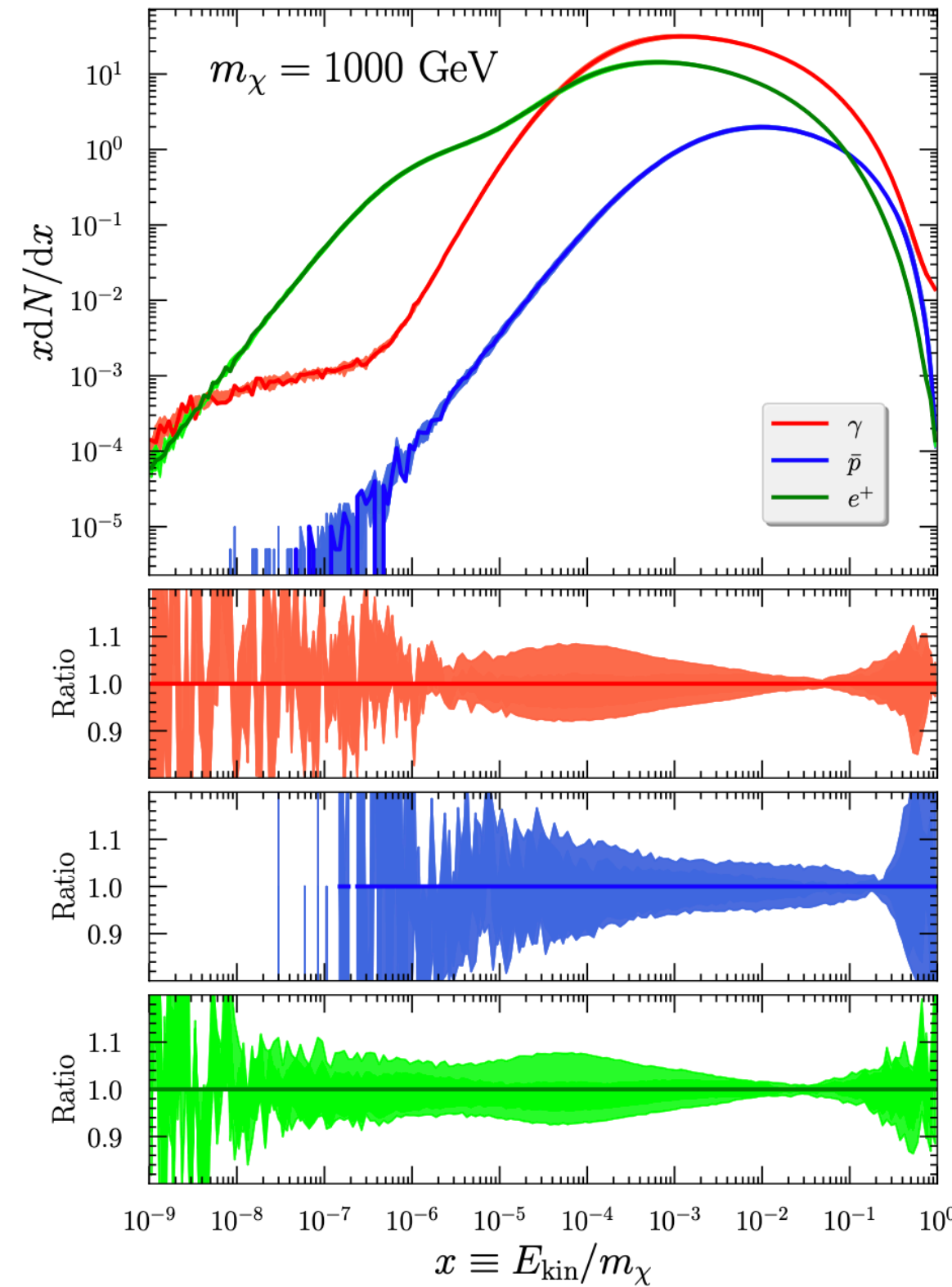




# VINCIA Tuning

Measurement	Experiment	$\chi^2/N_{\text{bins}}$	Measurement	Experiment	$\chi^2/N_{\text{bins}}$
$1 - T$	ALEPH [54]	0.13	$C$ -parameter	ALEPH [54]	0.39
$\log(1/x_p)$	ALEPH [54]	0.19	$\langle N_{\text{ch}} \rangle$	ALEPH [54]	0.028
$\langle N_{\text{ch}} \rangle$ ( $ Y  < 0.5$ )	ALEPH [54]	0.012	$\langle N_{\text{ch}} \rangle$ ( $ Y  < 1.0$ )	ALEPH [54]	0.028
$\langle N_{\text{ch}} \rangle$ ( $ Y  < 1.5$ )	ALEPH [54]	0.030	$\langle N_{\text{ch}} \rangle$ ( $ Y  < 2.0$ )	ALEPH [54]	0.040
$\pi^\pm$ spectrum	ALEPH [54]	0.67	$\pi^0$ spectrum	ALEPH [54]	0.24
$\Lambda^0$ spectrum	ALEPH [55]	1.24	$\Lambda^0$ spectrum (2-jet events)	ALEPH [55]	1.31
Thrust	ALEPH [56]	0.097	$C$ -parameter	ALEPH [56]	0.35
$N_{\text{ch}}$ ( $y_{\text{cut}} = 0.01$ )	DELPHI [57]	5.99	$N_{\text{ch}}$ ( $y_{\text{cut}} = 0.02$ )	DELPHI [57]	4.88
$\Lambda^0$ spectrum	DELPHI [58]	1.34	$\langle N_{\Lambda^0} \rangle$	DELPHI [58]	0.53
$\pi^0$ momentum	DELPHI [59]	0.41	$\log(1/x_p)$	DELPHI [59]	0.33
$1 - T$	DELPHI [59]	0.18	$C$ -parameter	DELPHI [59]	0.34
$\langle N_{\text{ch}} \rangle$	DELPHI [59]	0.031	$\langle N_{\pi^\pm} \rangle$	DELPHI [59]	0.063
$\langle N_{\pi^0} \rangle$	DELPHI [59]	0.39	$\langle N_\rho \rangle$	DELPHI [59]	3.40
$\langle N_p \rangle$	DELPHI [59]	2.30	$\langle N_{\Lambda^0} \rangle$	DELPHI [59]	1.54
$\langle N_{\text{ch}} \rangle$	DELPHI [60]	0.005	$\langle N_{\pi^\pm} \rangle$	DELPHI [60]	0.10
$\langle N_p \rangle$	DELPHI [60]	0.05	$N_{p/\bar{p}}/N_{\text{ch}}$	DELPHI [60]	0.27
$\pi^\pm$ momentum	DELPHI [60]	0.46	$p/\bar{p}$ momentum	DELPHI [60]	0.43
Thrust (udsc events)	L3 [61]	0.34	$C$ -parameter (udsc events)	L3 [61]	0.22
Charged multiplicity	L3 [61]	3.39	$\log(1/x_p)$	L3 [61]	0.96
$x_p$ (udsc events)	L3 [61]	0.78			
$\langle N_{\text{ch}} \rangle$	OPAL [62]	0.37	$\pi^\pm$ spectrum	OPAL [63]	0.25
$\Lambda^0$ scaled energy	OPAL [64]	1.49	$\pi^0$ scaled momentum	OPAL [65]	0.12
All events $\log(1/x_p)$	OPAL [66]	0.38	$\langle N_{\text{ch}} \rangle$	OPAL [66]	0.16
$1 - T$	OPAL [67]	0.10	$C$ -parameter	OPAL [67]	0.35

Parameter	MONASH	VINCIA (default)	PYTHIA [17, 18]	This work
$a_L$	0.68	0.45	0.601	$0.337 \pm 0.015$
$b_L$	0.98	0.80	0.897	$0.784 \pm 0.020$
$\sigma_\perp$ (GeV)	0.335	0.305	0.307	$0.296 \pm 0.003$
$a_{QQ}$	0.97	0.90	1.671	$1.246 \pm 0.082$
$\chi^2/N_{\text{df}}$	1034.52/852	786.11/852	676.69/852	660.21/852



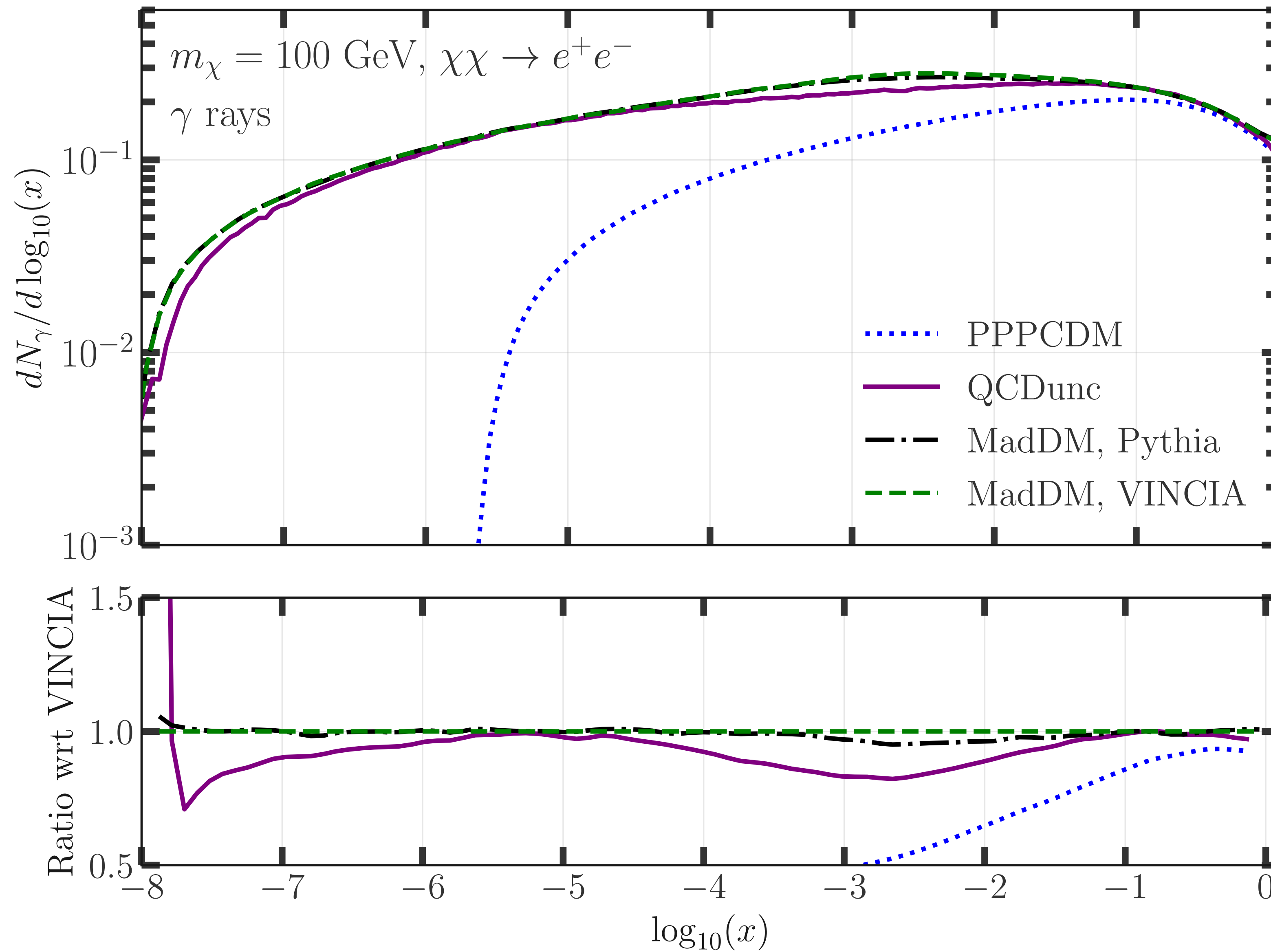


# Summary of the novelties

- **Polarization and helicity effects:** We use MadDM which we interface with PYTHIA 8 and VINCIA shower plugin being the default option.
- **Resummed electroweak corrections and interleaved resonance decays:** The electroweak corrections are modeled with helicity-dependent Antenna showers and Sudakov form factors.
- **Running quark masses and full mass effects:** We use running quark masses instead of pole masses.
- **New annihilation channels:** We also calculate the spectra for two new annihilation channels ( $\chi\chi \rightarrow \gamma Z, HZ$ ).
- **Off-shell effects:** We take into account off-shell effects. For the case of  $WW, ZZ, HZ$  we generate the spectra of the four-body decays and DM masses down to 5 GeV.
- **Full one-loop effects:** For one-loop induced annihilation channels ( $\gamma\gamma, \gamma Z, gg$ ), we take into account the full one-loop effects instead of effective couplings.
- **Improved hadronization model:** We carry out a new tuning of the hadronization model parameters using a set of measurements performed at the Z-boson pole.

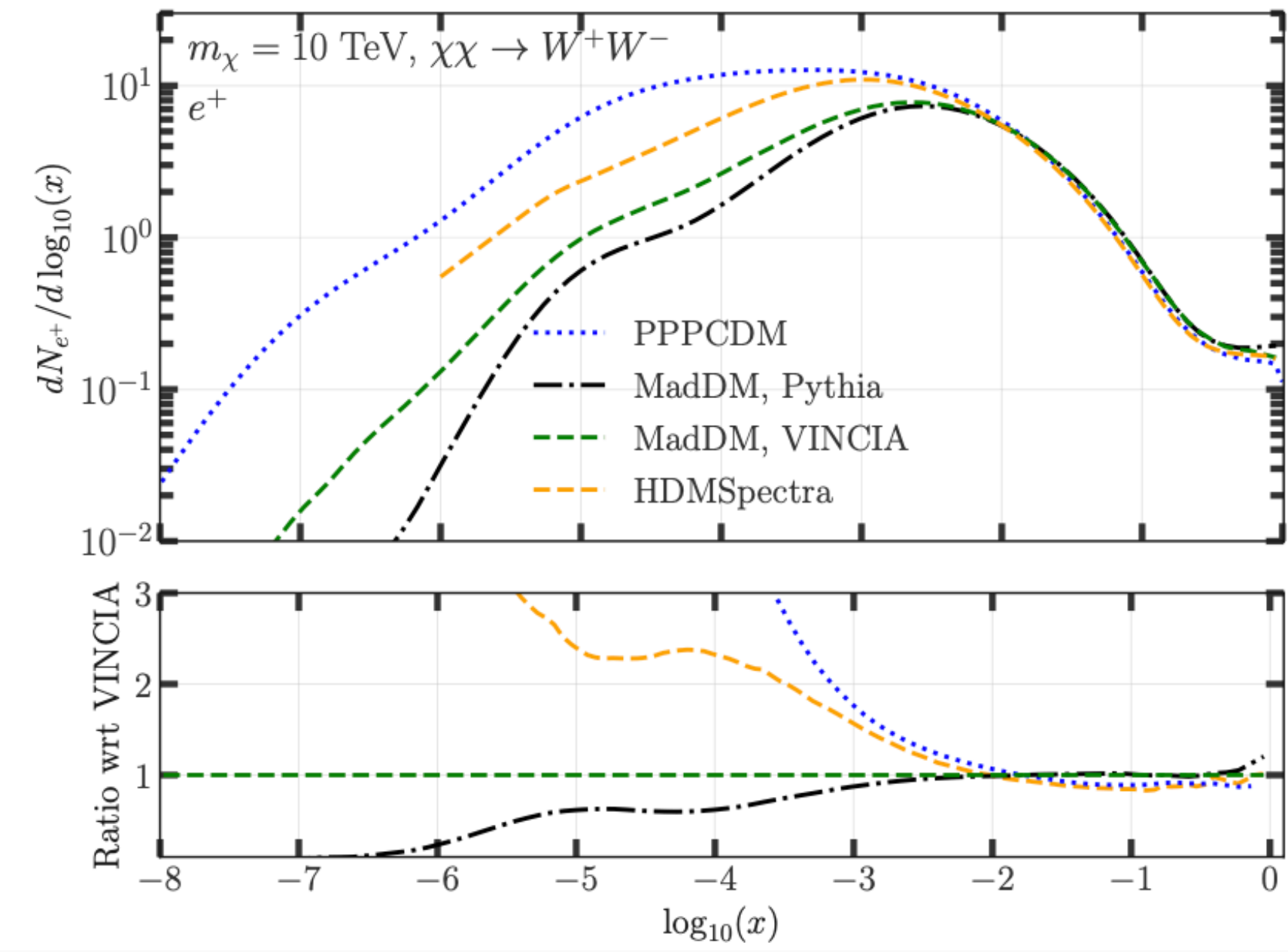
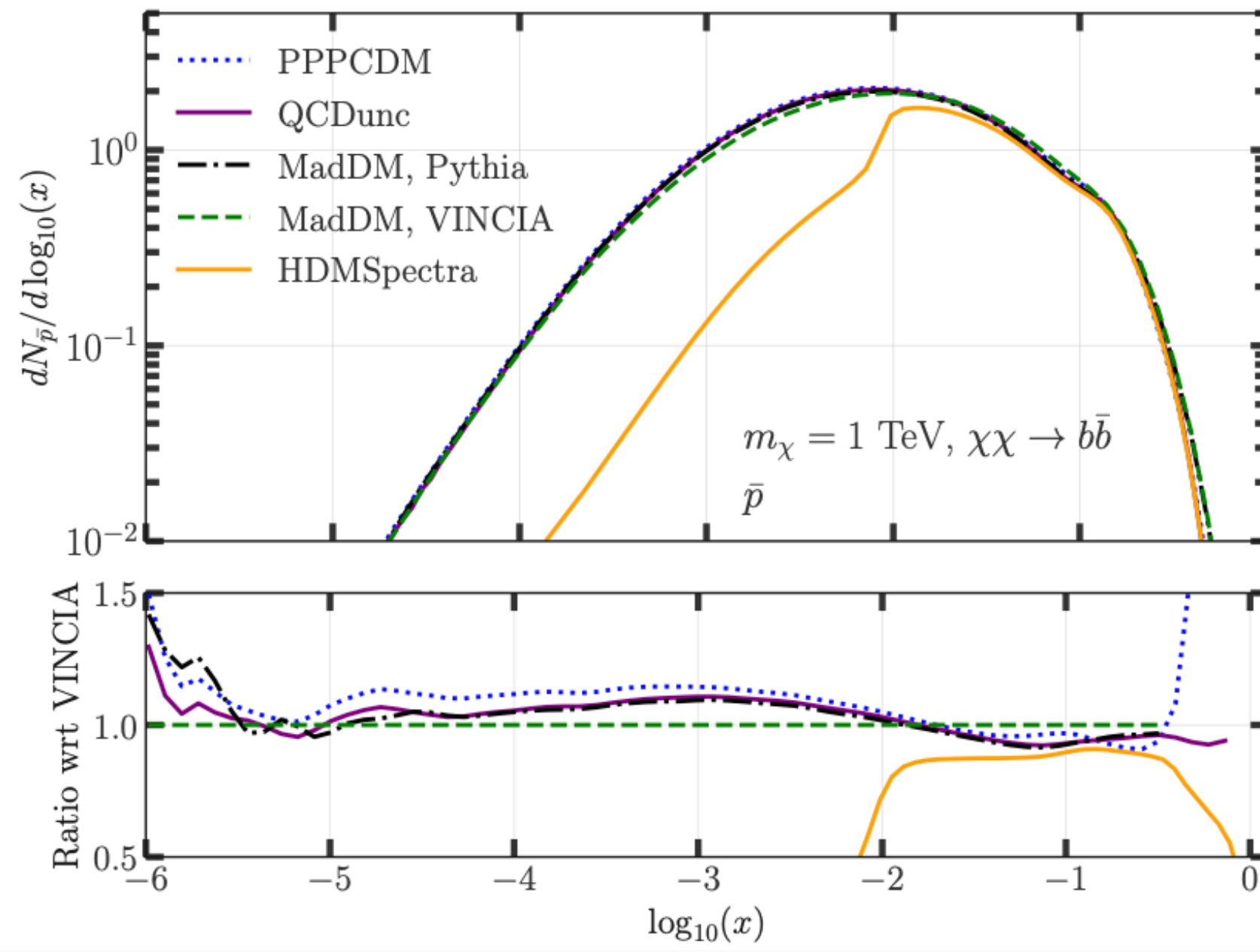
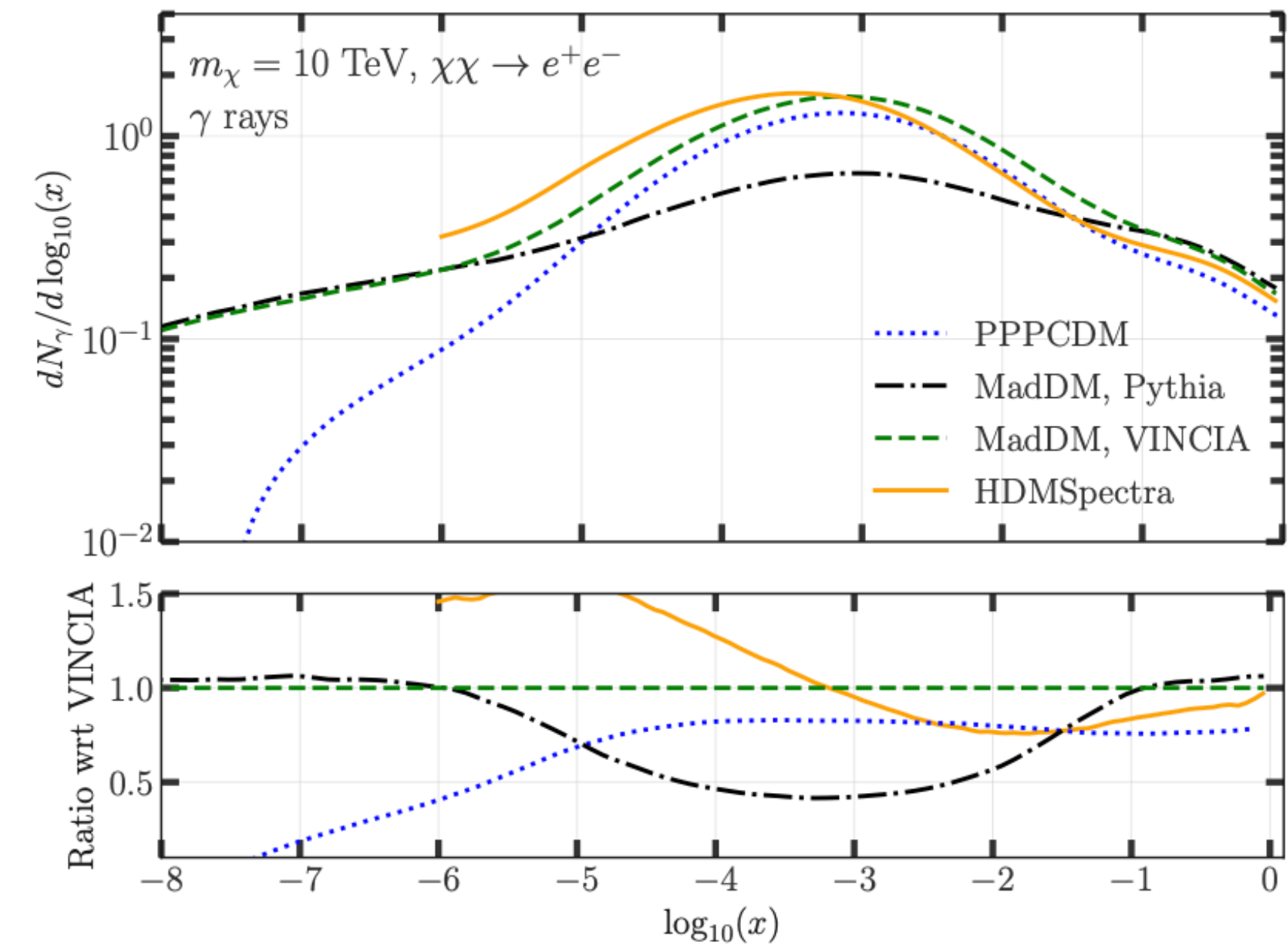
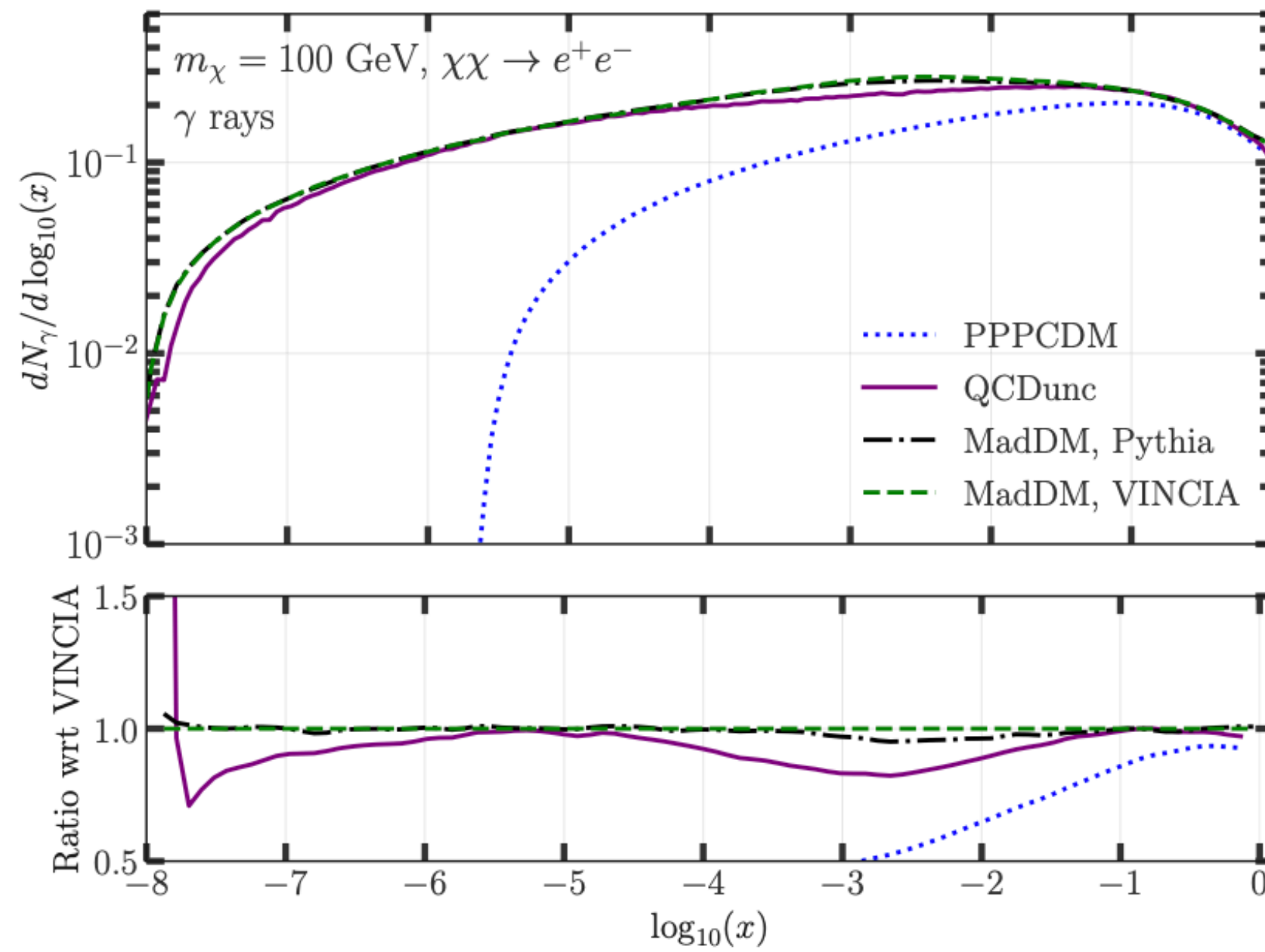


# Comparison with PPPC and HDMS





# Comparison with PPPC and HDMS

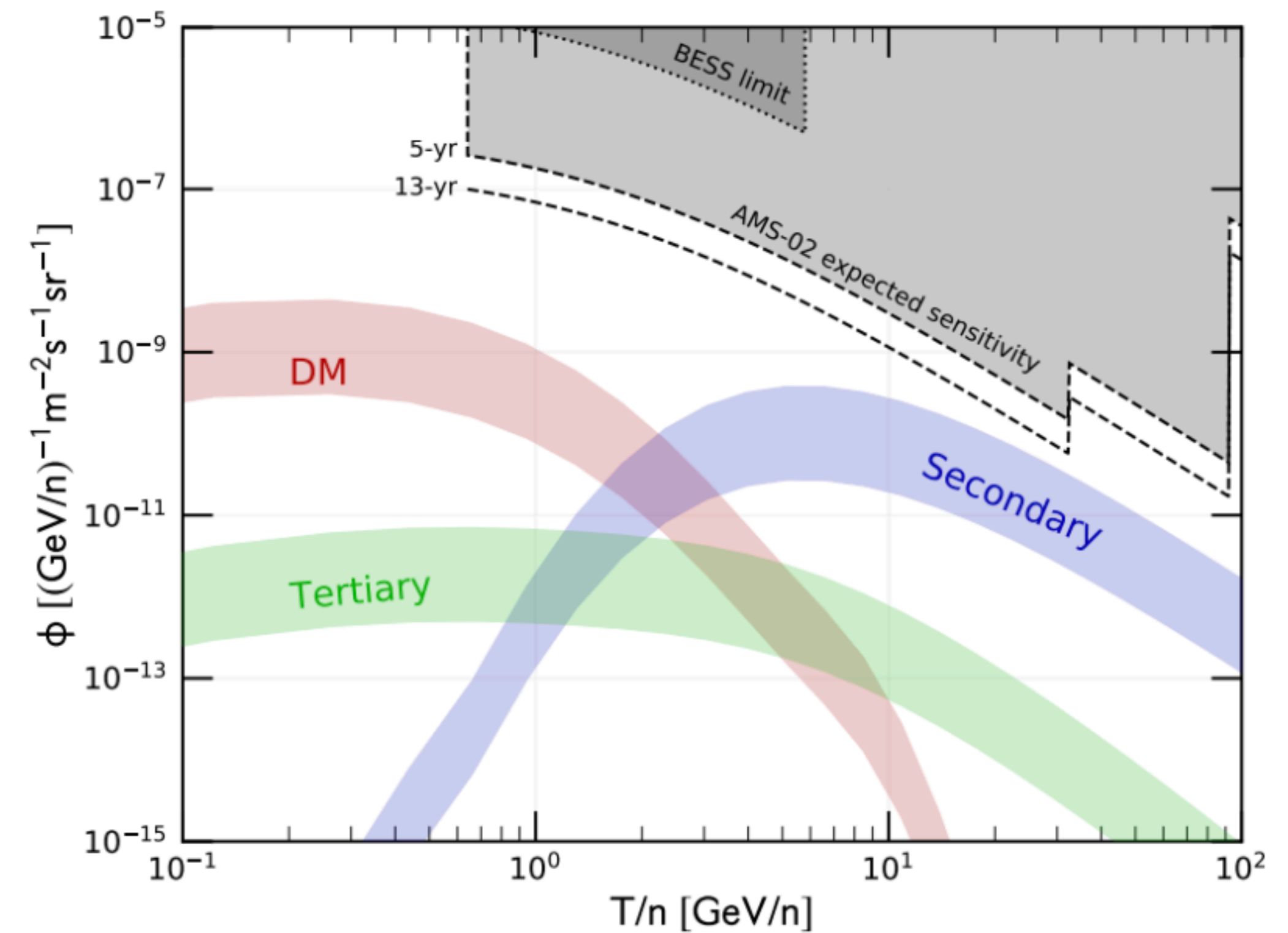
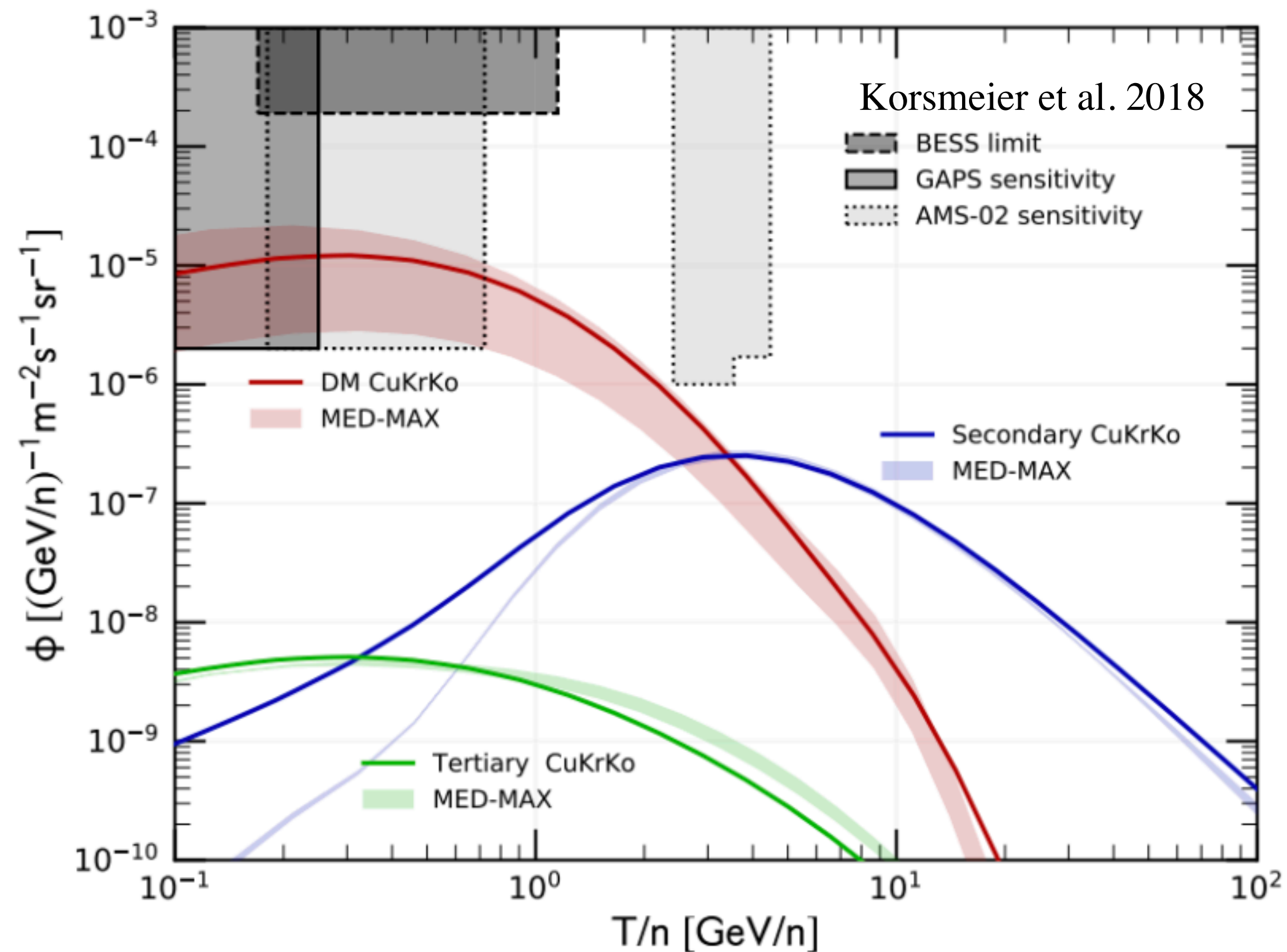




# Dark matter search with cosmic antinuclei

$$pp \rightarrow \bar{D}pppn$$

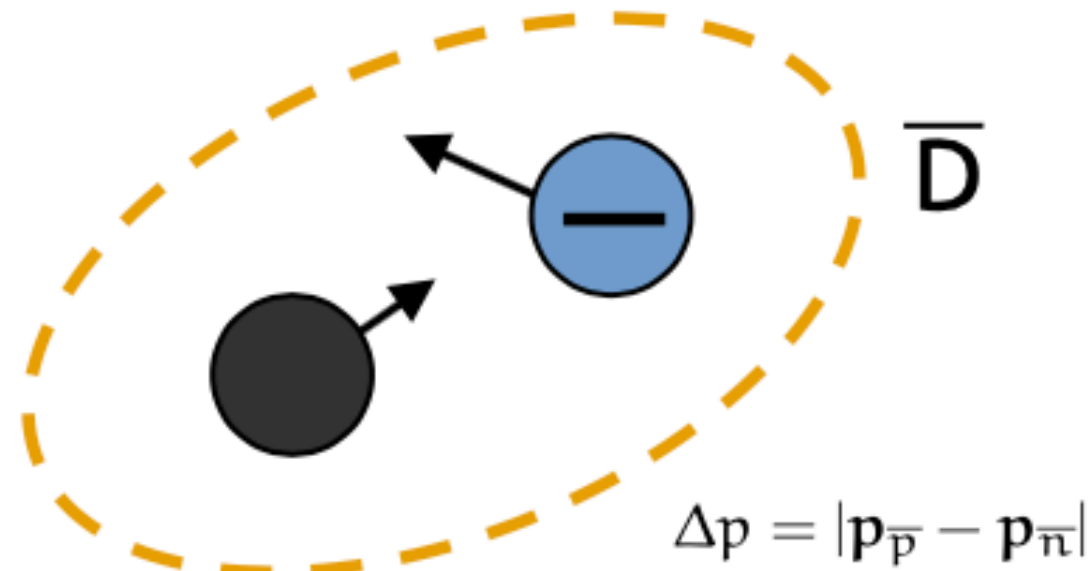
$$\sqrt{s} = 2E_{\text{CM}} = \sqrt{2m_p \cdot (E_{\text{LAB}} + m_p)} \geq 6m_p \rightarrow E_{\text{LAB}} \geq 17m_p$$



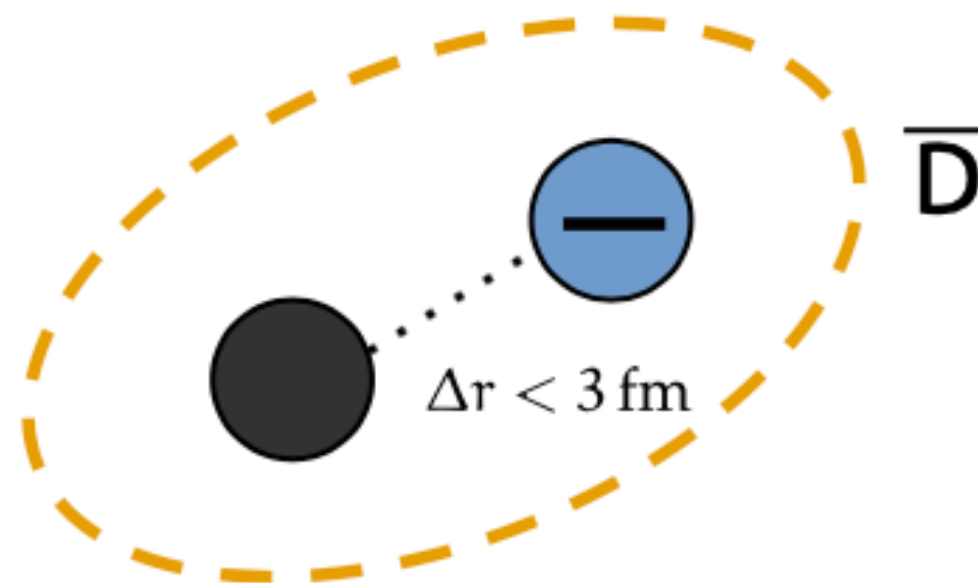


# Coalescence models

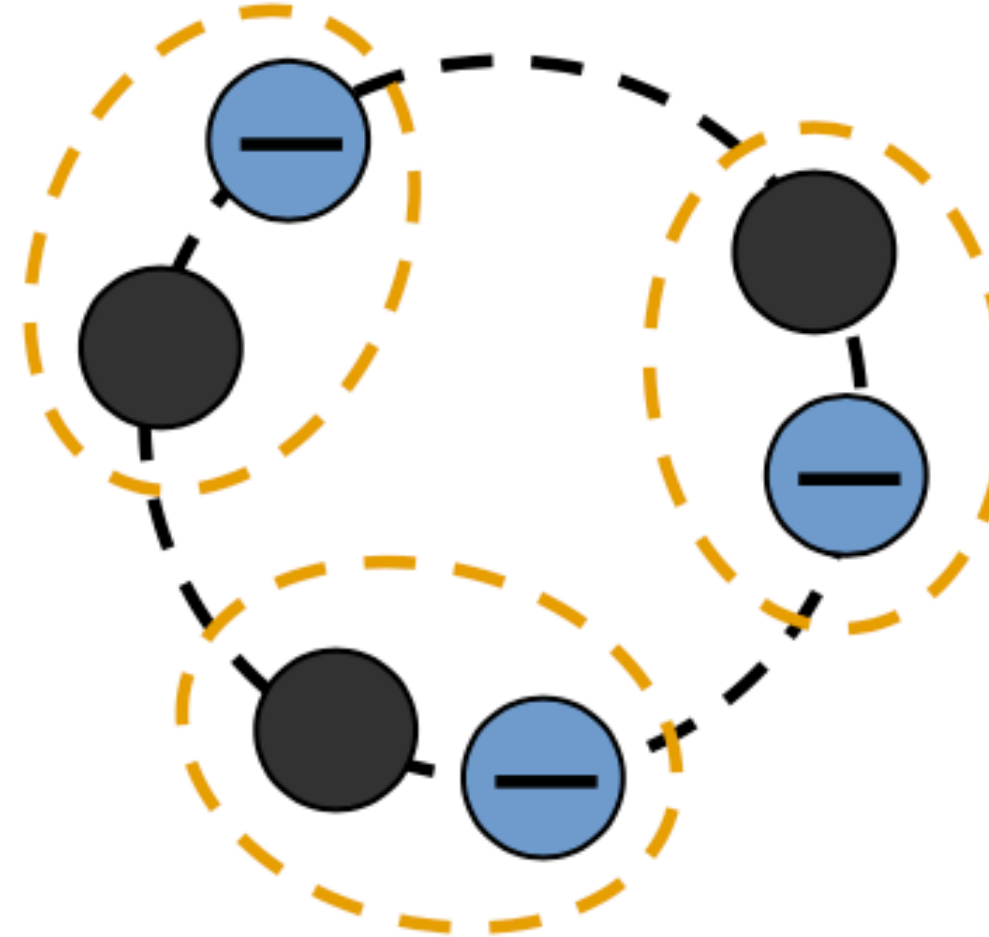
## Simple coalescence model



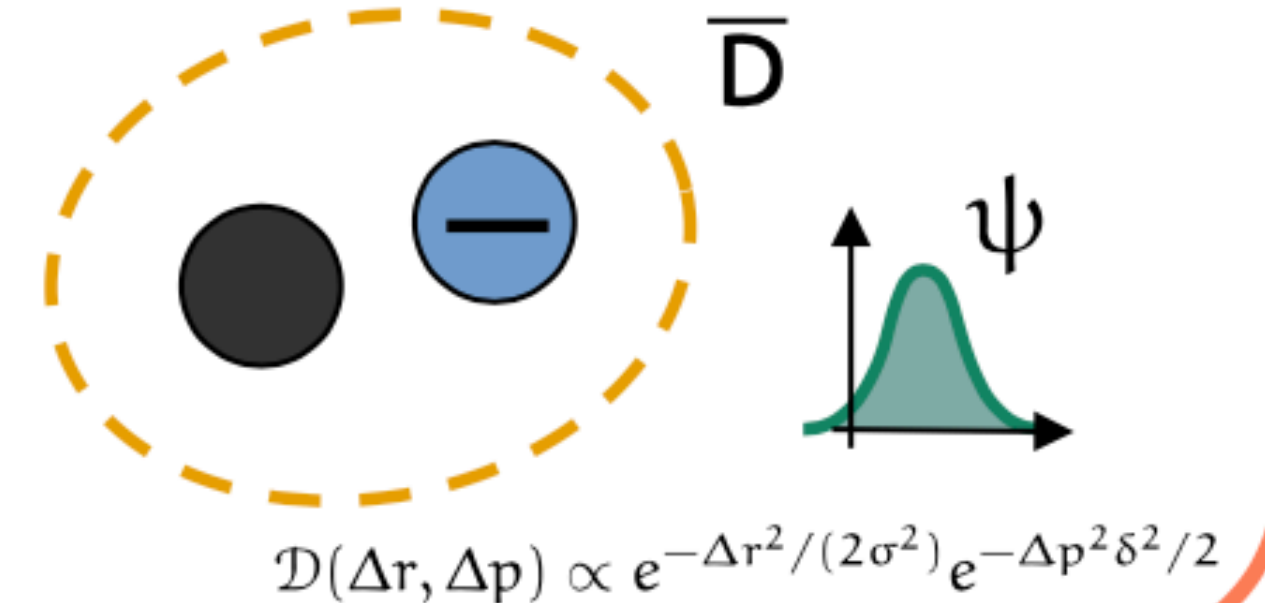
## Simple coalescence model + sharp cutoff in distance



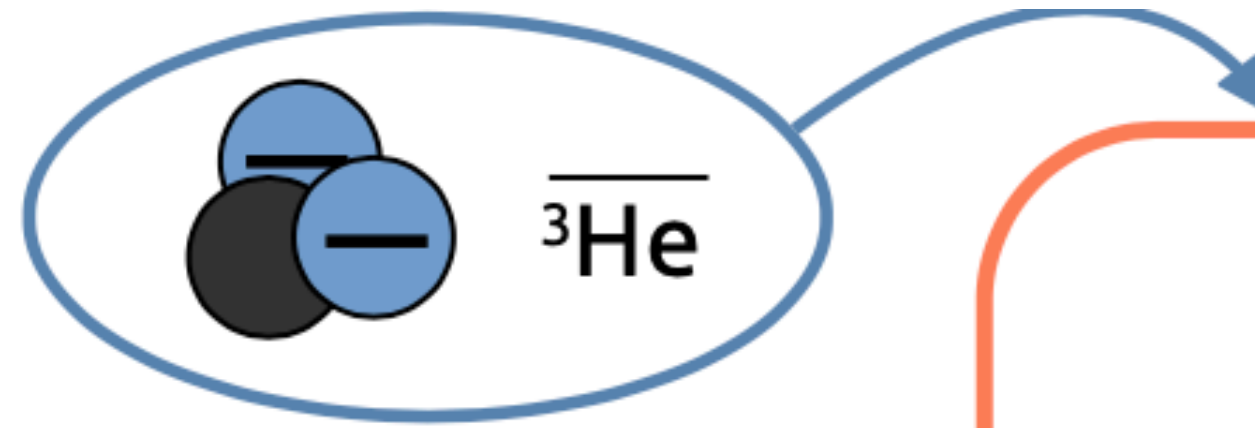
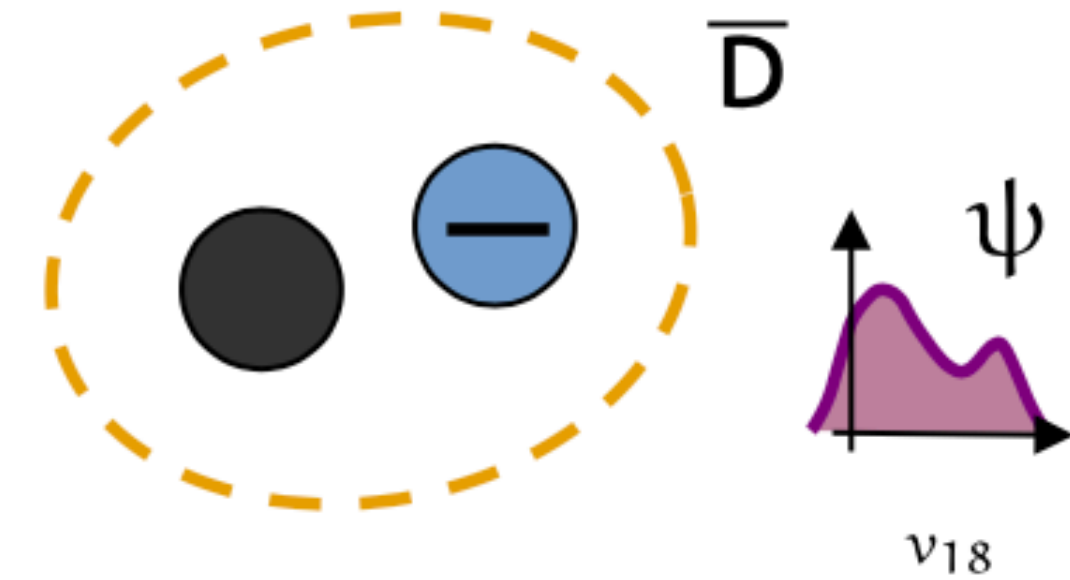
## Spherical model



## Wigner + Gaussian wavefunction

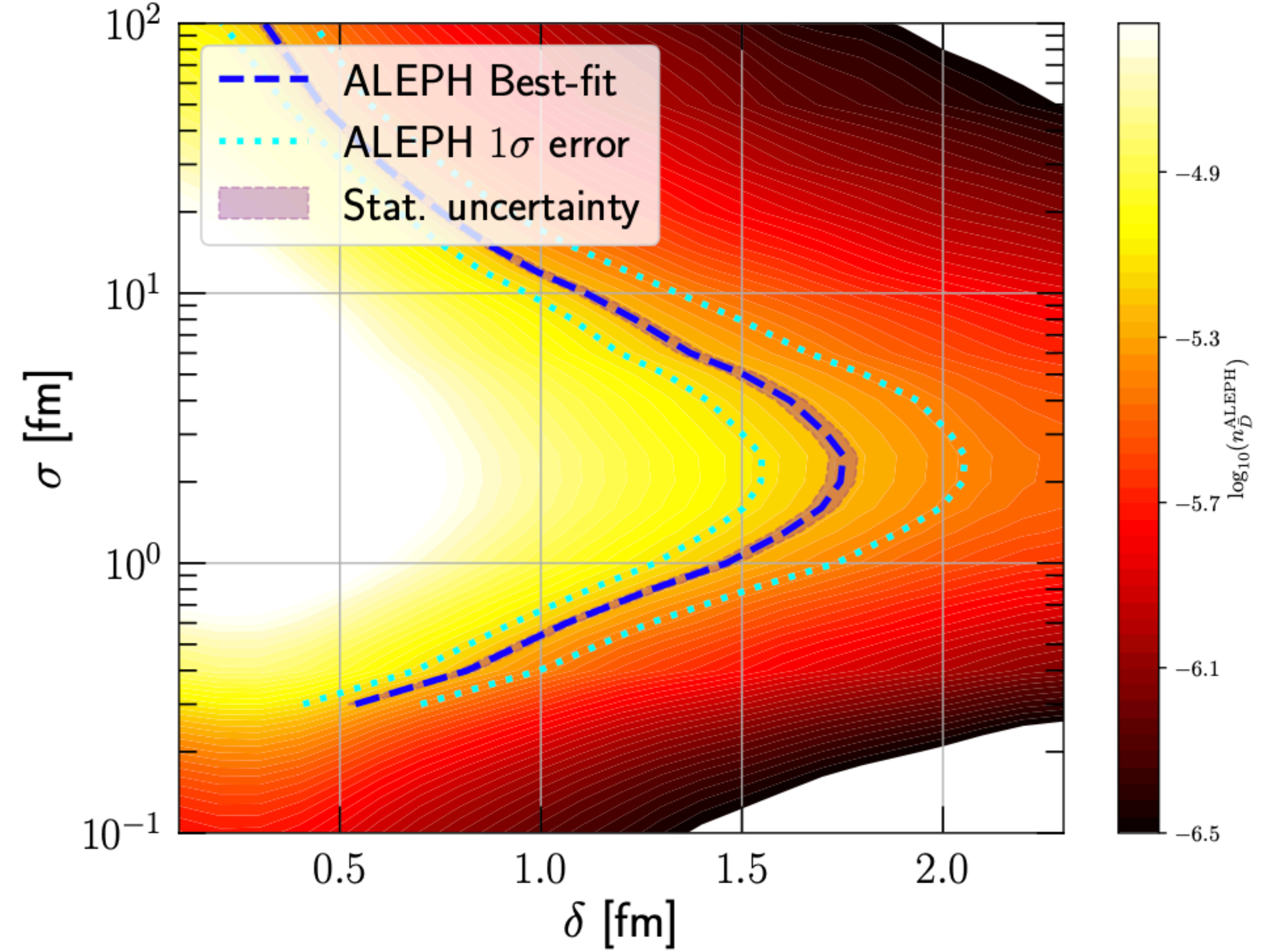
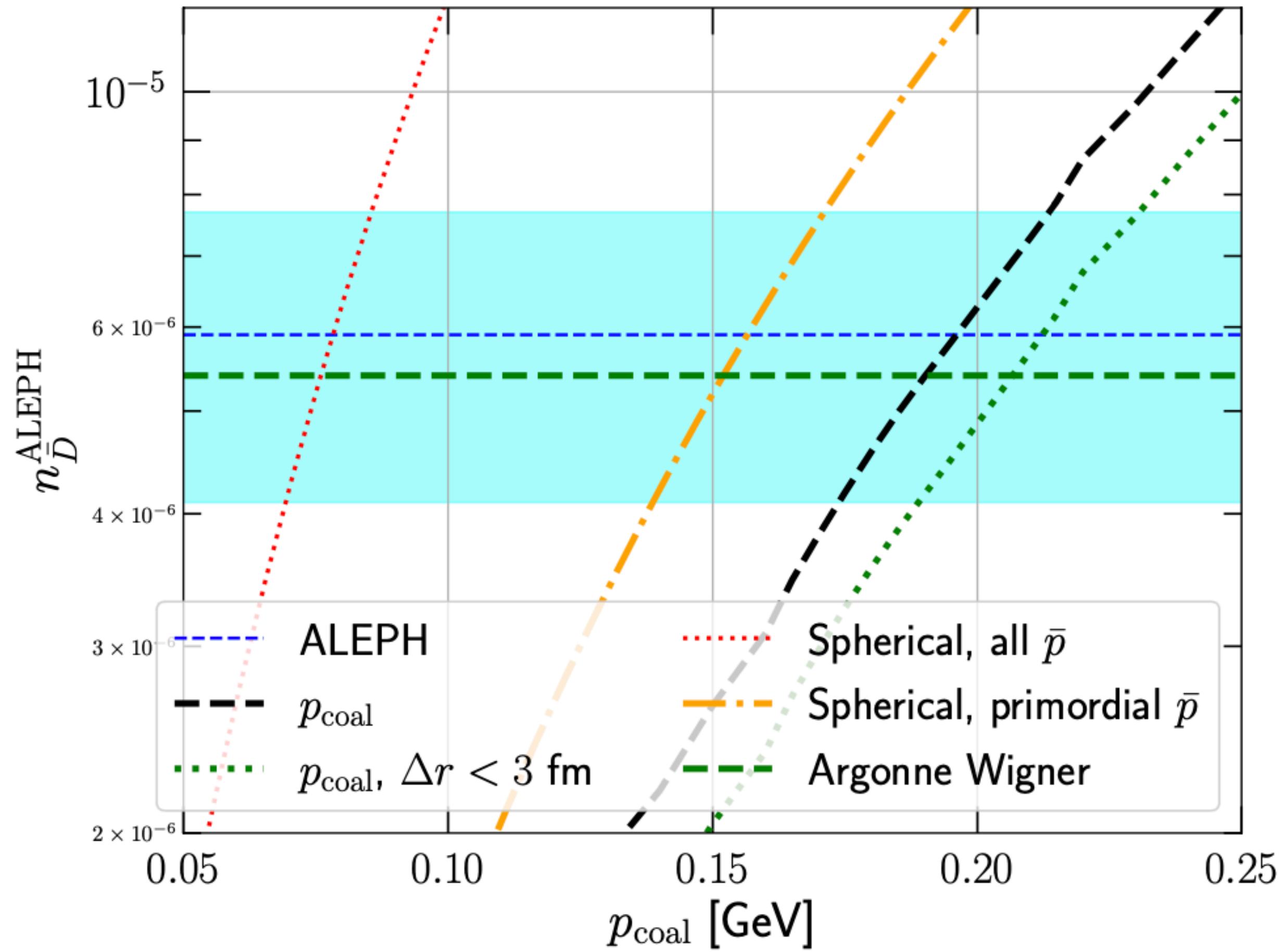


## Wigner + Argonne wavefunction

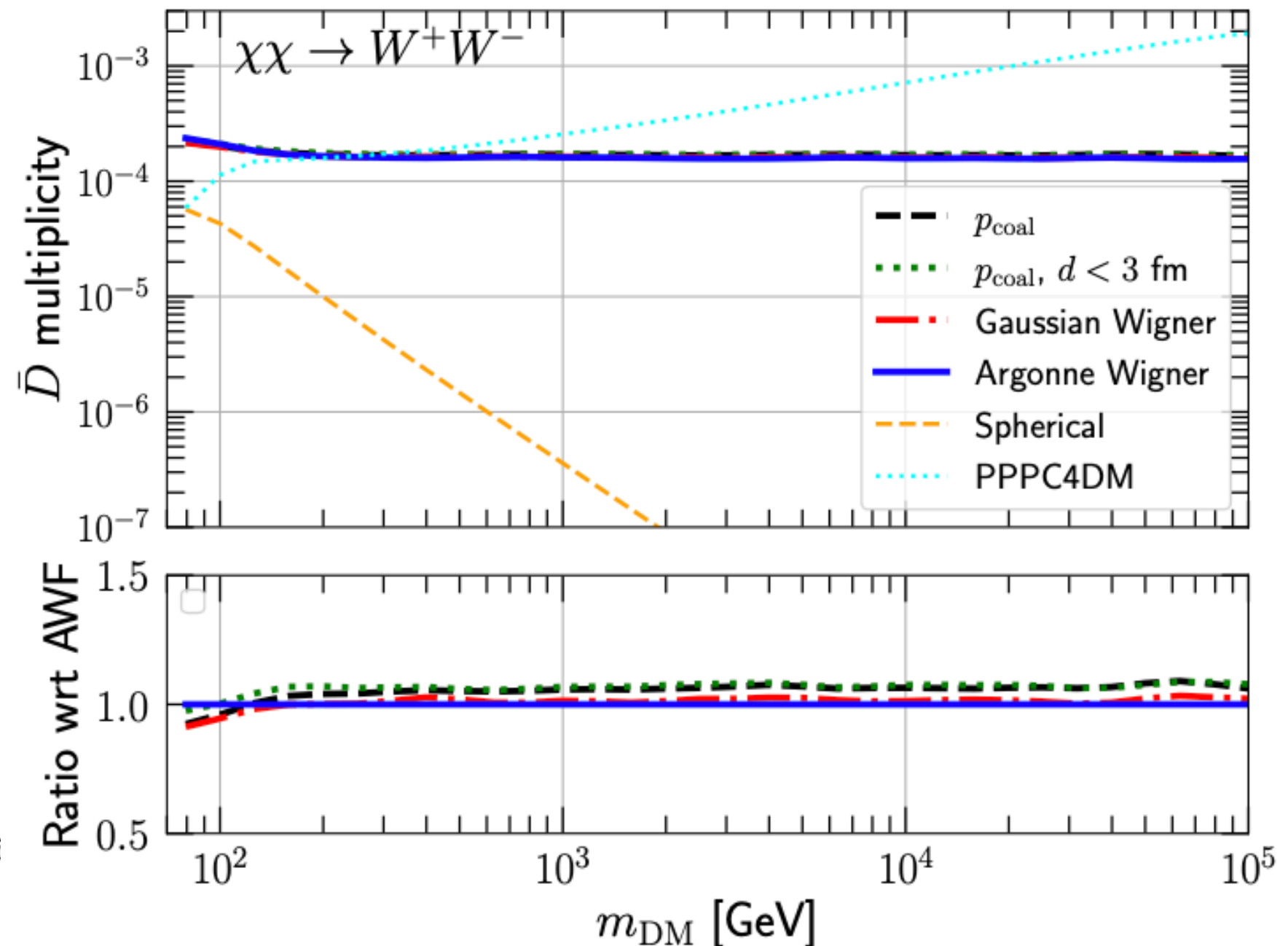
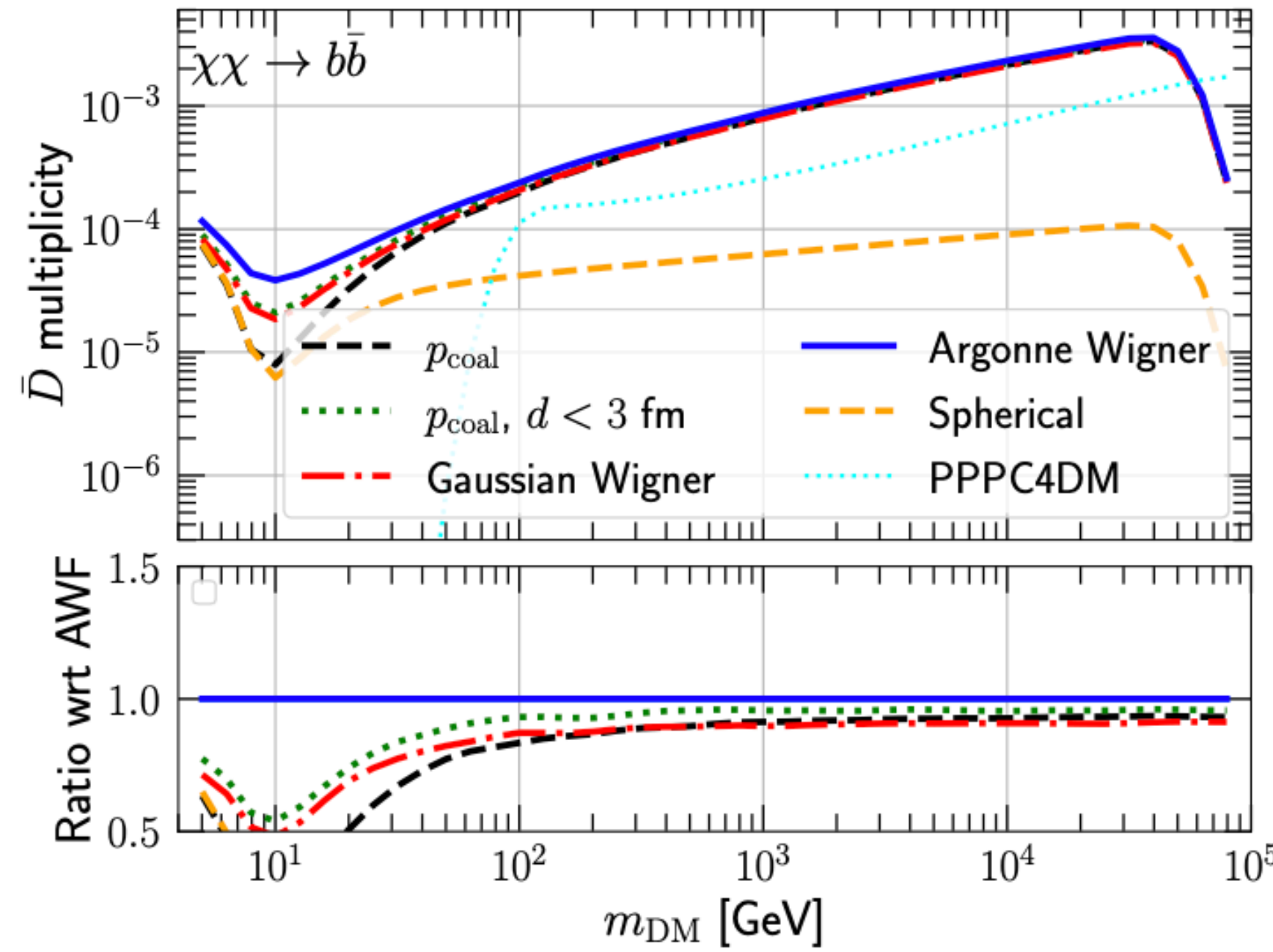
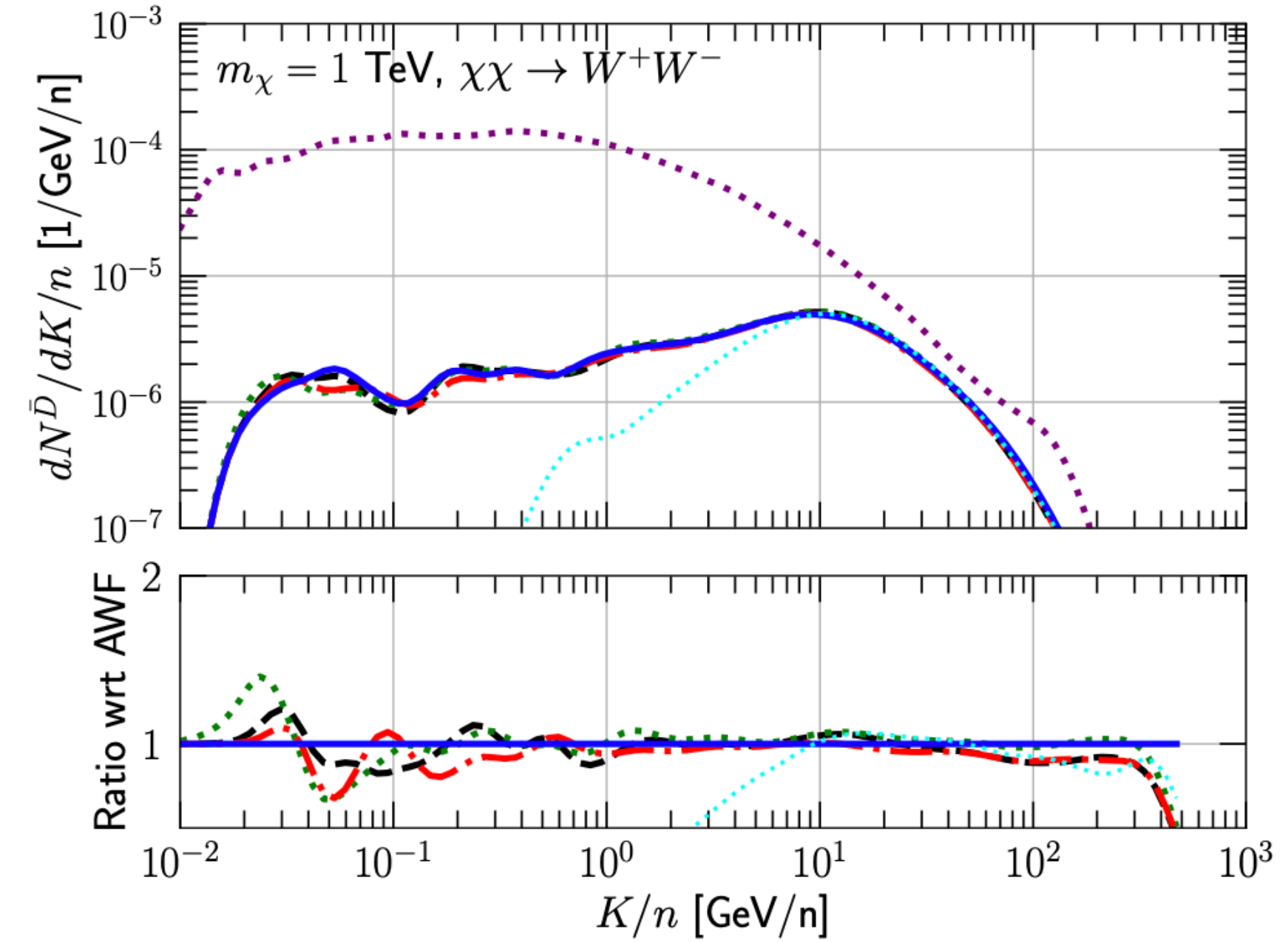
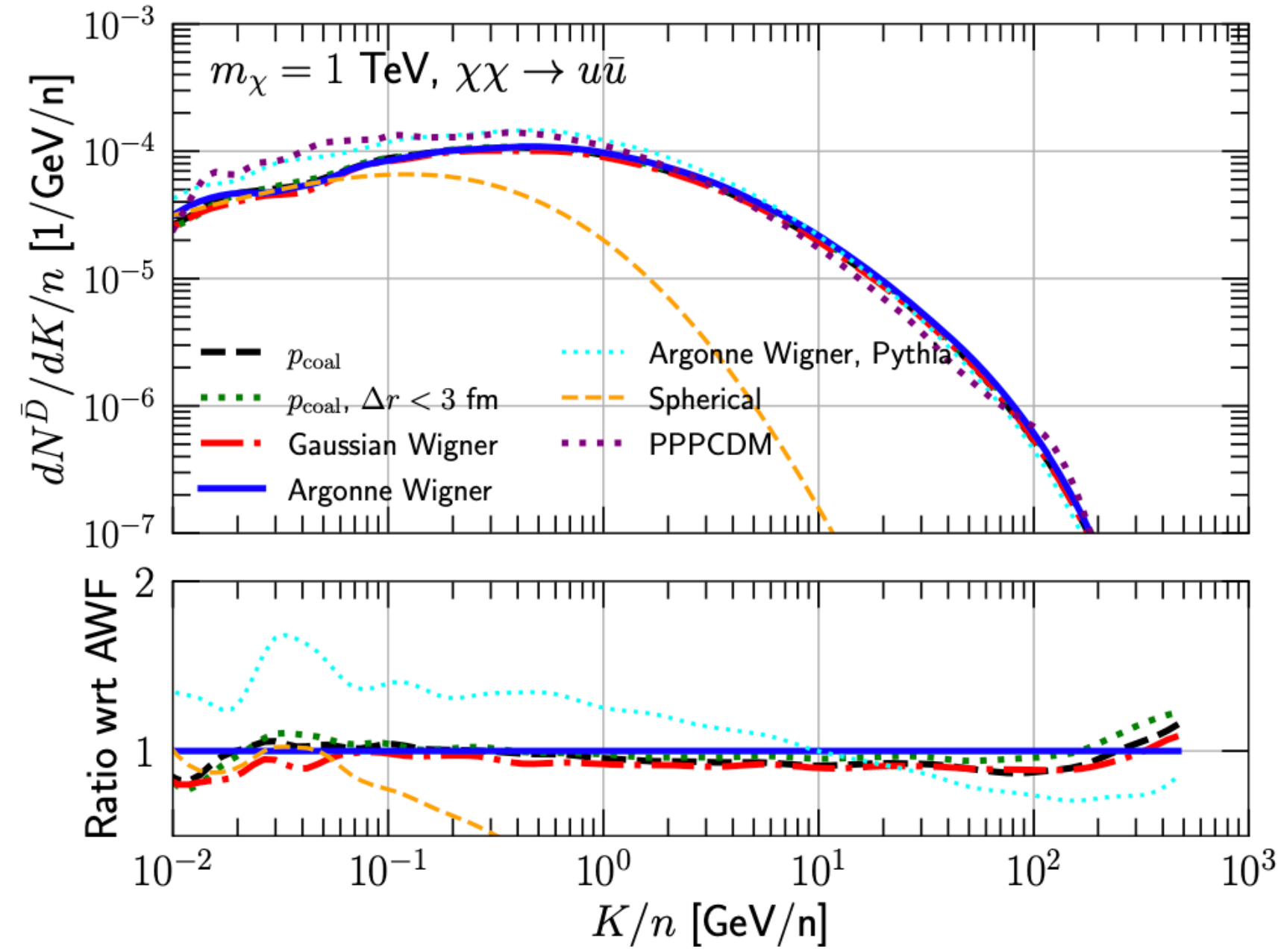




# Tuning on ALEPH data point

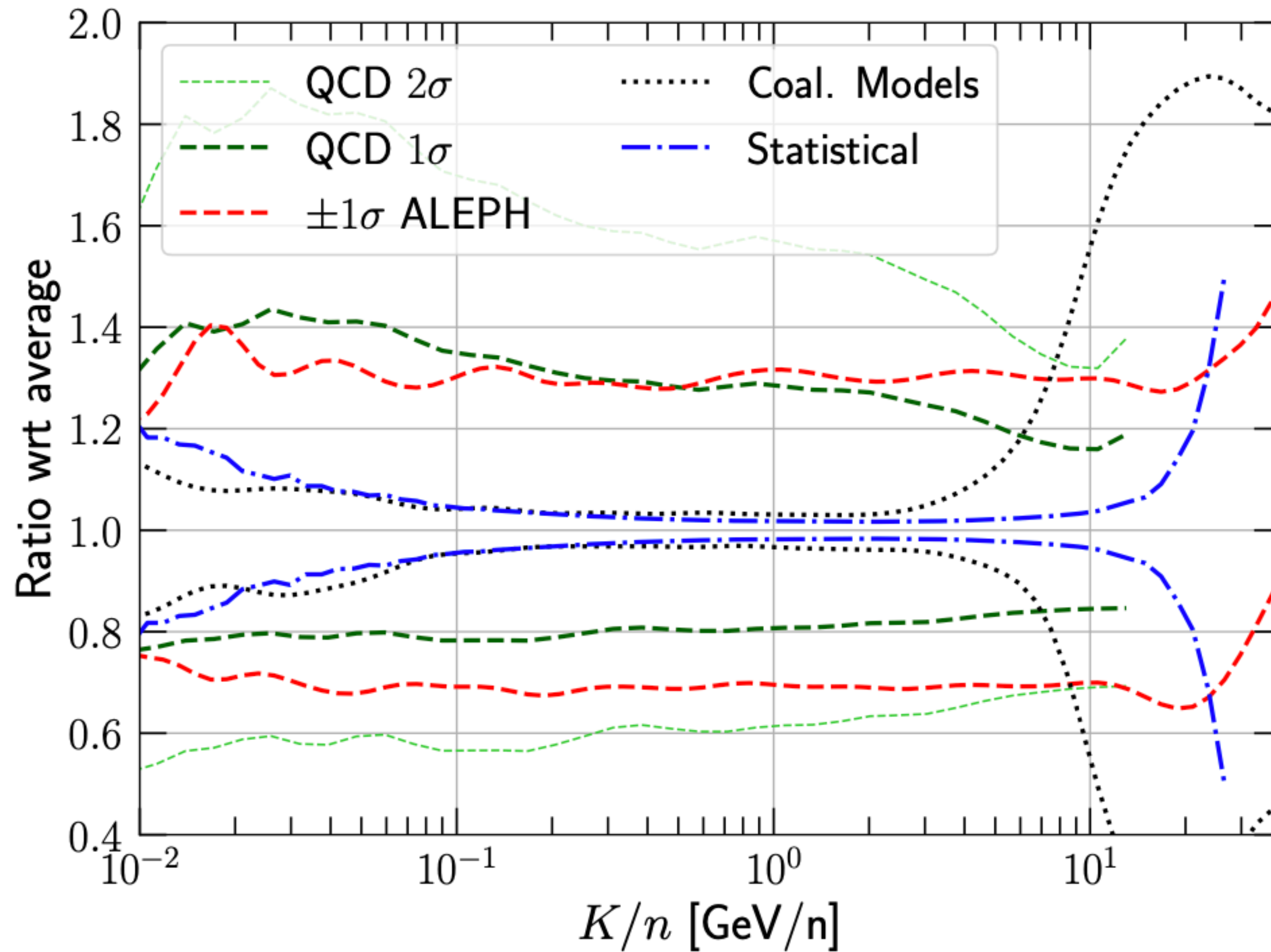


# Cosmixs predictions for Dbar





# Current Uncertainties on Dbar source spectrum

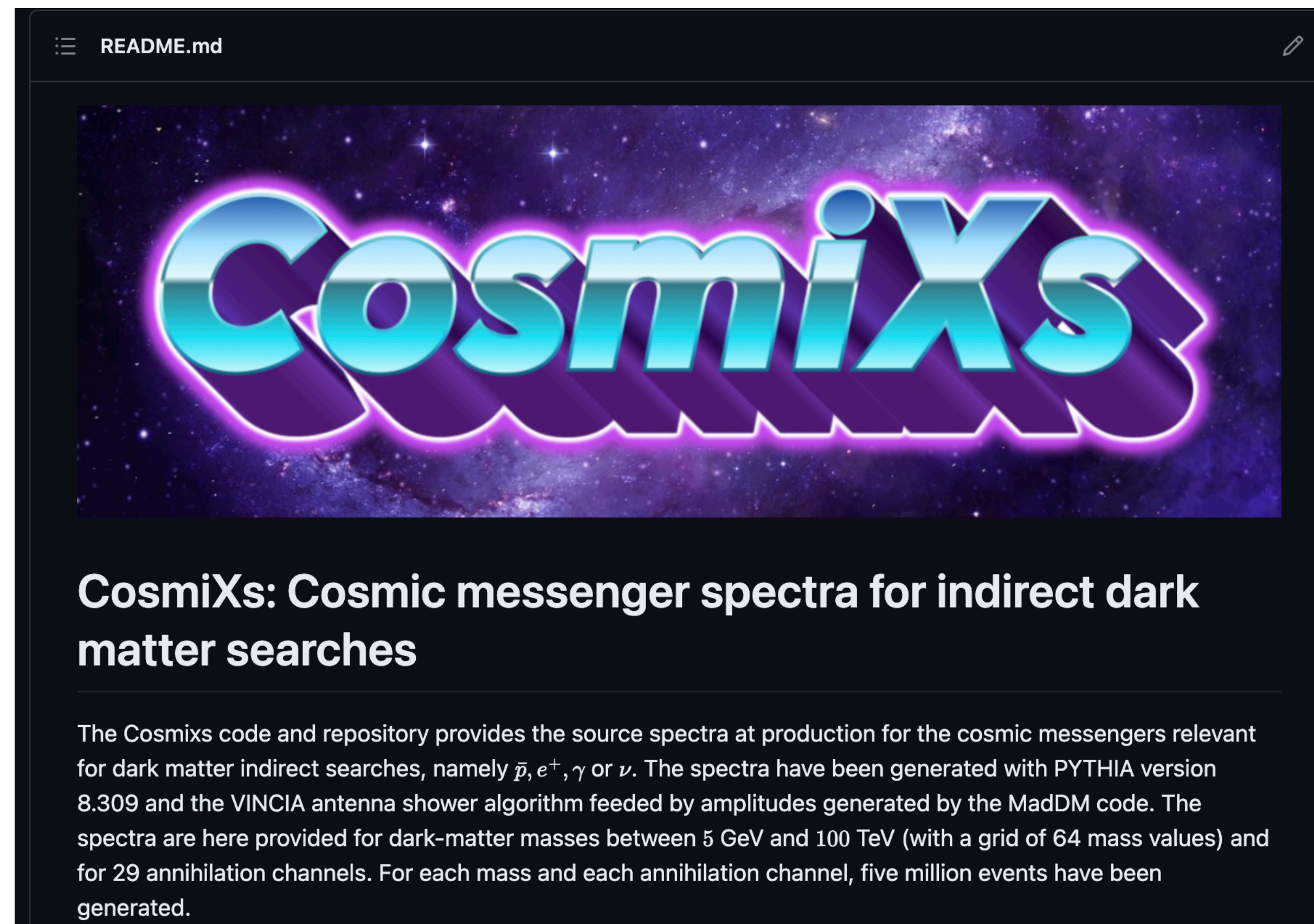


# Spectra are publicly available

$e^+e^-$ ,  $\mu^+\mu^-$ ,  $\tau^+\tau^-$ ,  $\nu_e\nu_e$ ,  $\nu_\mu\nu_\mu$ ,  $\nu_\tau\nu_\tau$ ,  $u\bar{u}$ ,  $d\bar{d}$ ,  $c\bar{c}$ ,  $s\bar{s}$ ,  $t\bar{t}$ ,  $b\bar{b}$ ,  $\gamma\gamma$ ,  $gg$ ,  $W^+W^-$ ,  $ZZ$ ,  $HH$ ,  $ZH$ ,  $\gamma Z$ .

**Dark Matter mass between 5 GeV and 100 TeV**

**<https://github.com/ajueid/CosmiXs>**





# Conclusions

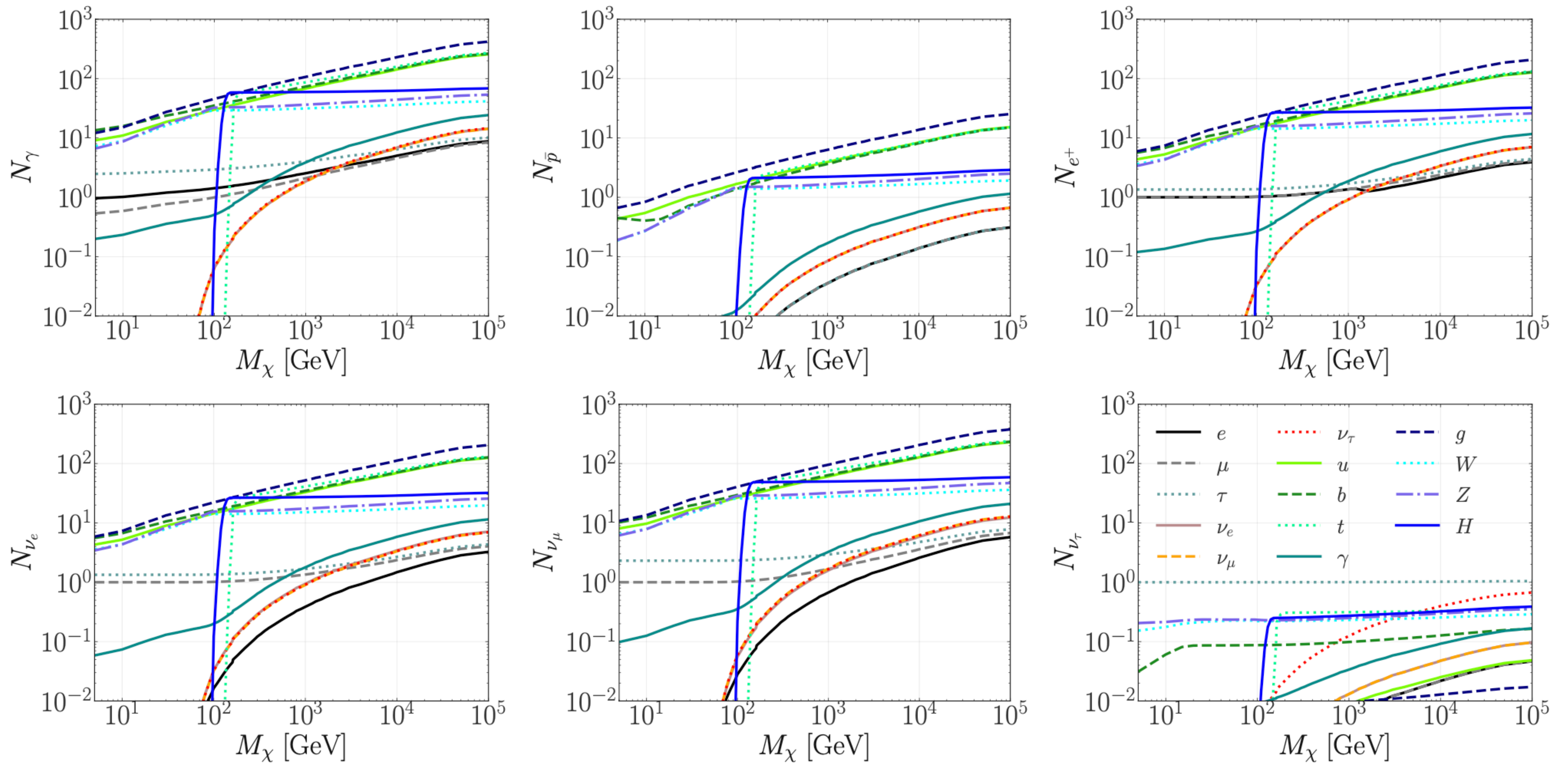
- We have significantly improved the predictions for particle spectra from dark matter annihilation.
- Our results include several effects not previously accounted for in PPPC4DMID.
- These improvements are relevant across all indirect detection strategies.
- CosmiXs tables are already implemented in *MadDM* and *MicrOMEGAs*, and will soon be adopted as the default in the *FermiTools*.
- We encourage developers of other tools to adopt our results.
- We are extending CosmiXs up to 1 PeV in dark matter mass.
- The tables will be periodically updated, especially following major *Pythia* releases.

# Backup slides

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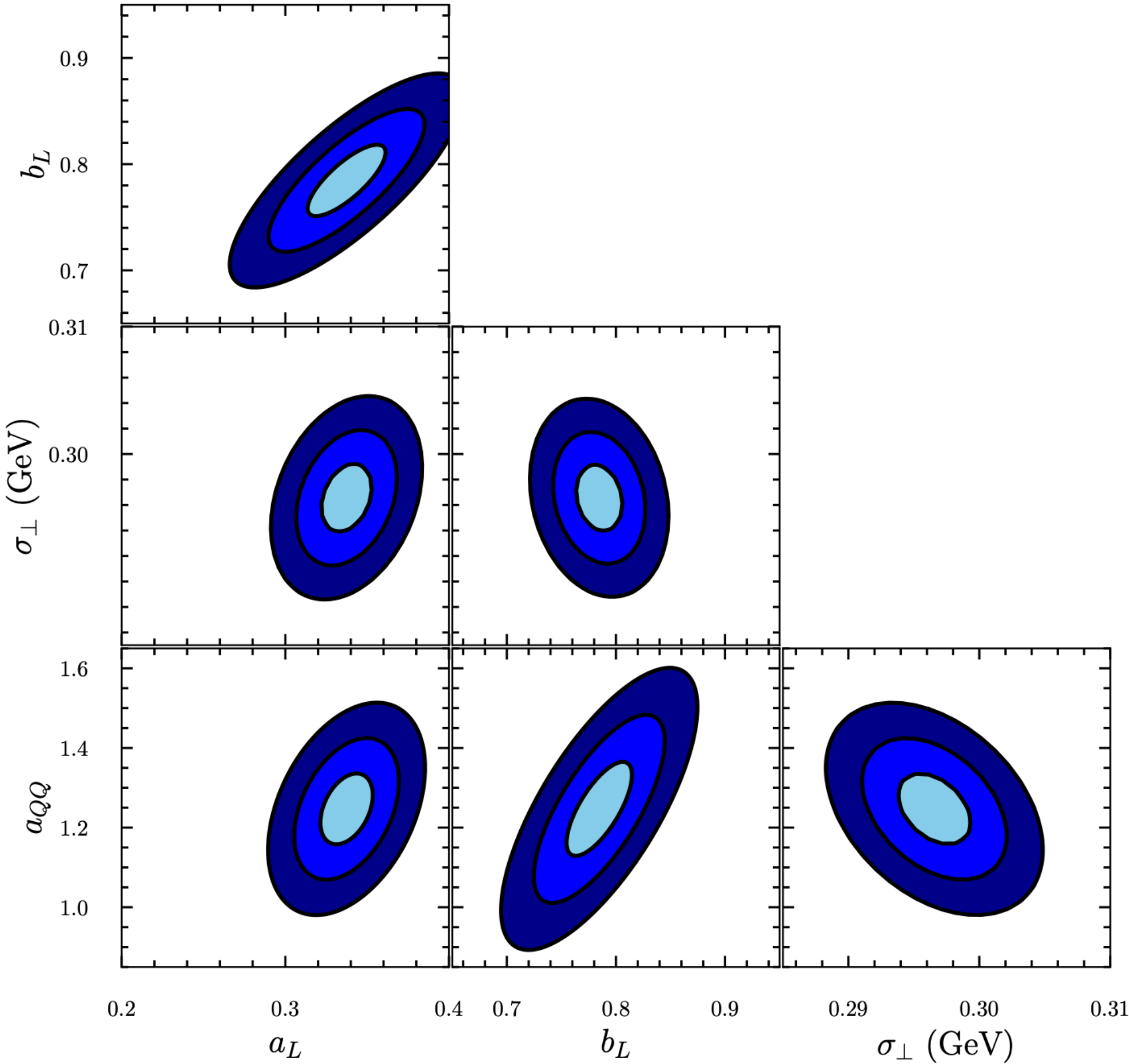


# Multiplicities



# Backup slides

parameter	PYTHIA 8 setting	Variation range	VINCIA
$\sigma_{\perp}$ (GeV)	StringPT:Sigma	0.0 – 1.0	0.305
$a_L$	StringZ:aLund	0.0 – 2.0	0.45
$b_L$	StringZ:bLund	0.2 – 2.0	0.80
$a_{QQ}$	StringZ:aExtraDiquark	0.0 – 2.0	0.90
$r_c$	StringZ:rFactC	0.0 – 2.0	0.85
$r_b$	StringZ:rFactB	0.0 – 2.0	1.15





# State of the art: QCDUnc

PREPARED FOR SUBMISSION TO JHEP

Estimating QCD uncertainties in Monte Carlo event generators for gamma-ray dark matter searches

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Simone Amoroso,<sup>a</sup> Sascha Caron,<sup>b,c</sup> Adil Jueid,<sup>d</sup> Roberto Ruiz de Austri<sup>e</sup> and Peter Skands<sup>f</sup>

(arXiv: 1812.07424)

PREPARED FOR SUBMISSION TO JHEP

CTPU-PTC-23-08

The Strong Force meets the Dark Sector: a robust estimate of QCD uncertainties for anti-matter dark matter searches

---

Adil Jueid,<sup>a</sup> Jochem Kip,<sup>b</sup> Roberto Ruiz de Austri<sup>c</sup> and Peter Skands<sup>d</sup>

(arXiv: 2303.11363)

Will be denoted by QCDUnc in what follows

PREPARED FOR SUBMISSION TO JCAP

Impact of QCD uncertainties on antiproton spectra from dark-matter annihilation

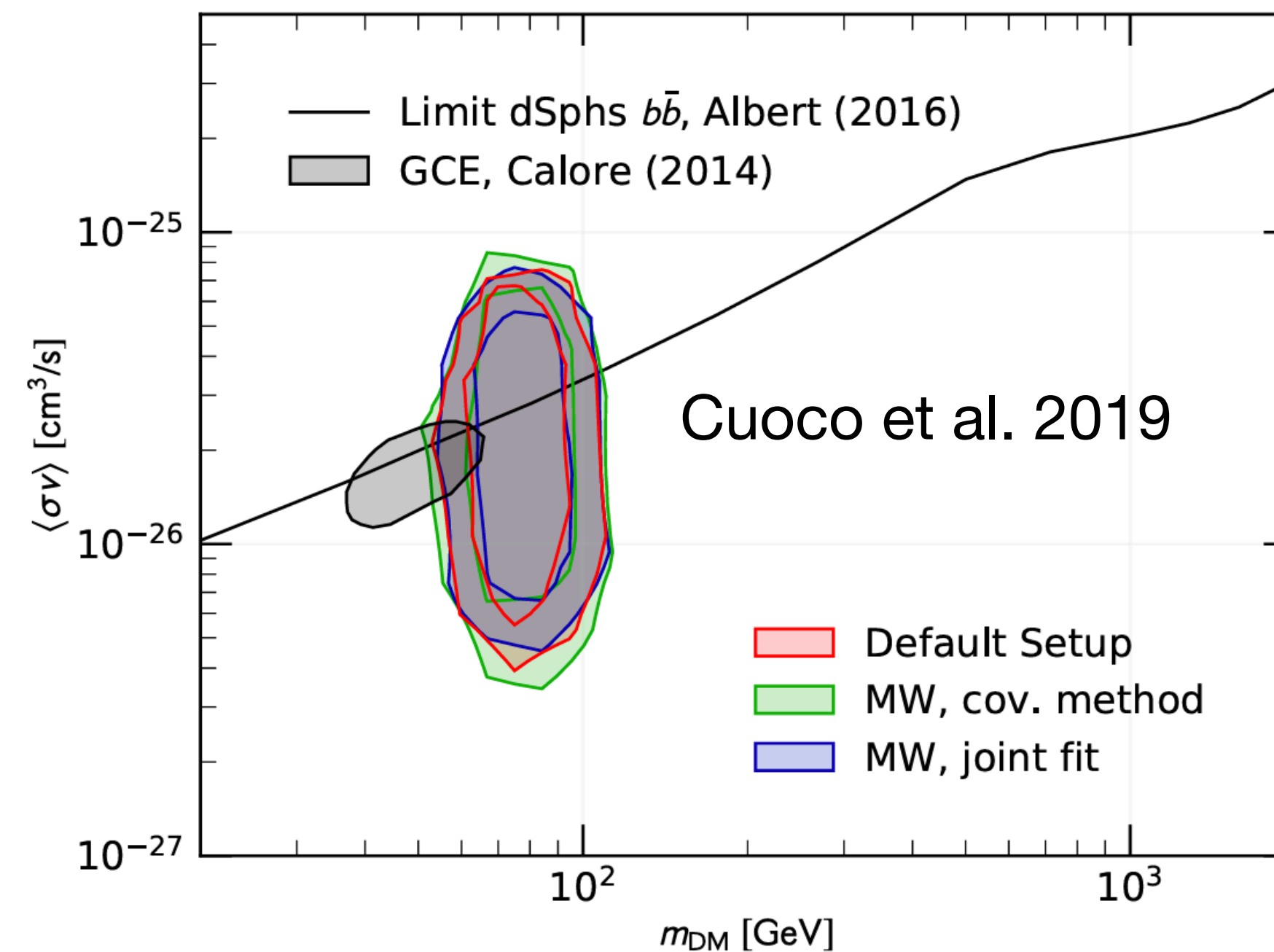
Adil Jueid,<sup>a</sup> Jochem Kip,<sup>b</sup> Roberto Ruiz de Austri<sup>c</sup> and Peter Skands<sup>d</sup>

(arXiv: 2202.11546)

- They New spectra of DM cosmic messengers using new tunes of PYTHIA 8 (version 2.19 and 3.07)
- *Estimated QCD uncertainties using parametric variations of the hadronization parameters (about 10%).*
- Estimated the impact on the best-fit point of the fitted DM mass and thermally-averaged annihilation cross section (in a two-parameter model).

# Importance of precise predictions

- There are hints for possible excess over the astrophysical backgrounds in various experiments, especially Fermi-LAT and AMS.
- These excesses triggered a plethora of phenomenological analyses aiming to explain it with dark matter.
- An important finding is that the precision in the determination of the particle spectra from DM annihilation is important in the fitting procedure.

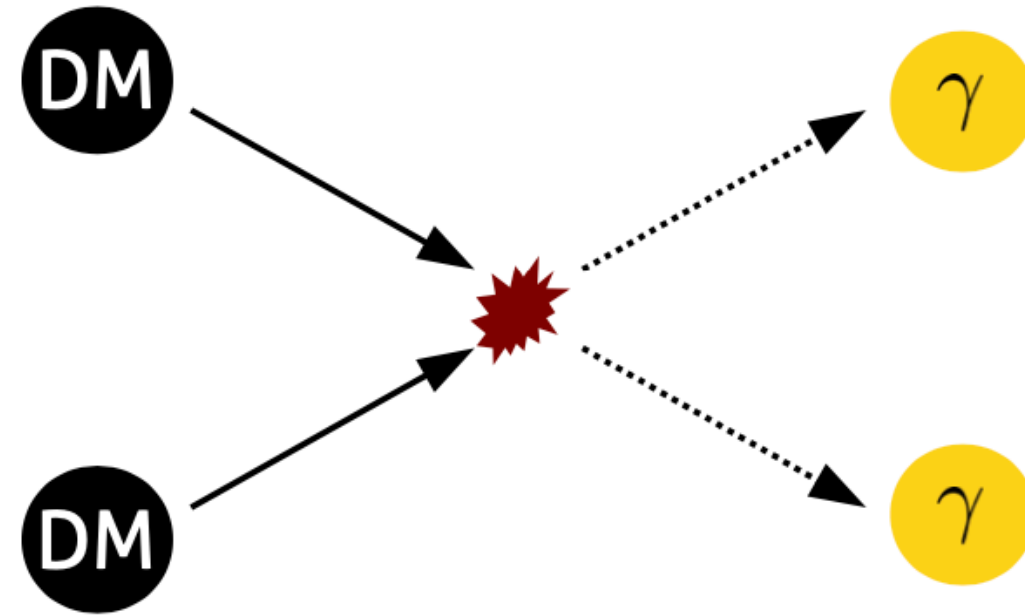




# Gamma rays from dark matter annihilation

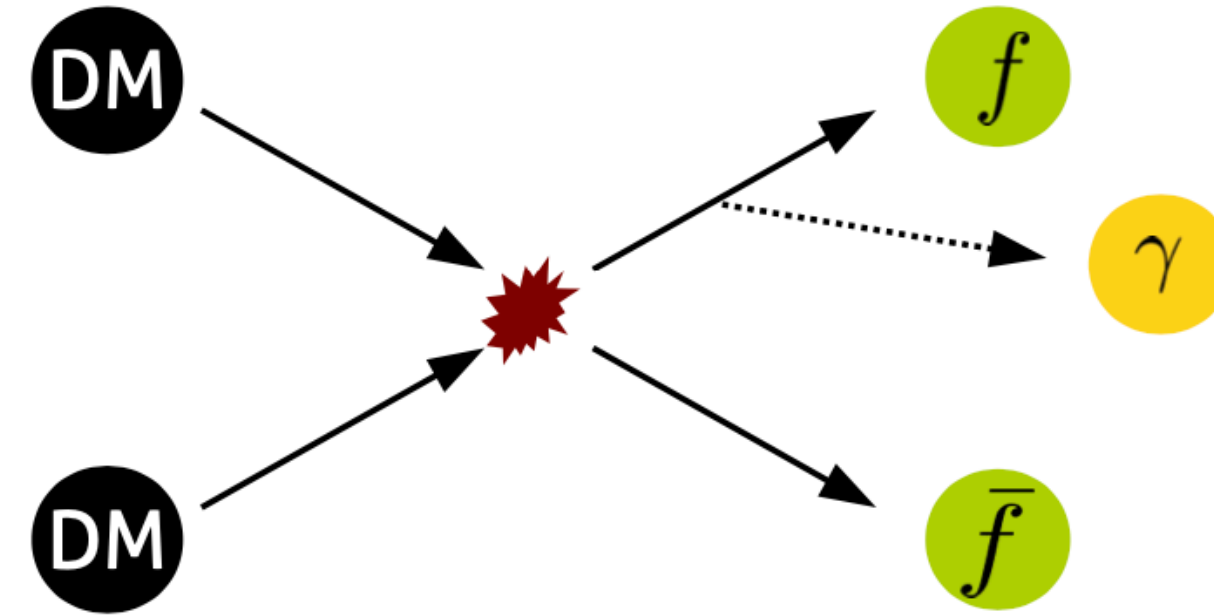
## Gamma-ray lines:

Two-body annihilation into photons



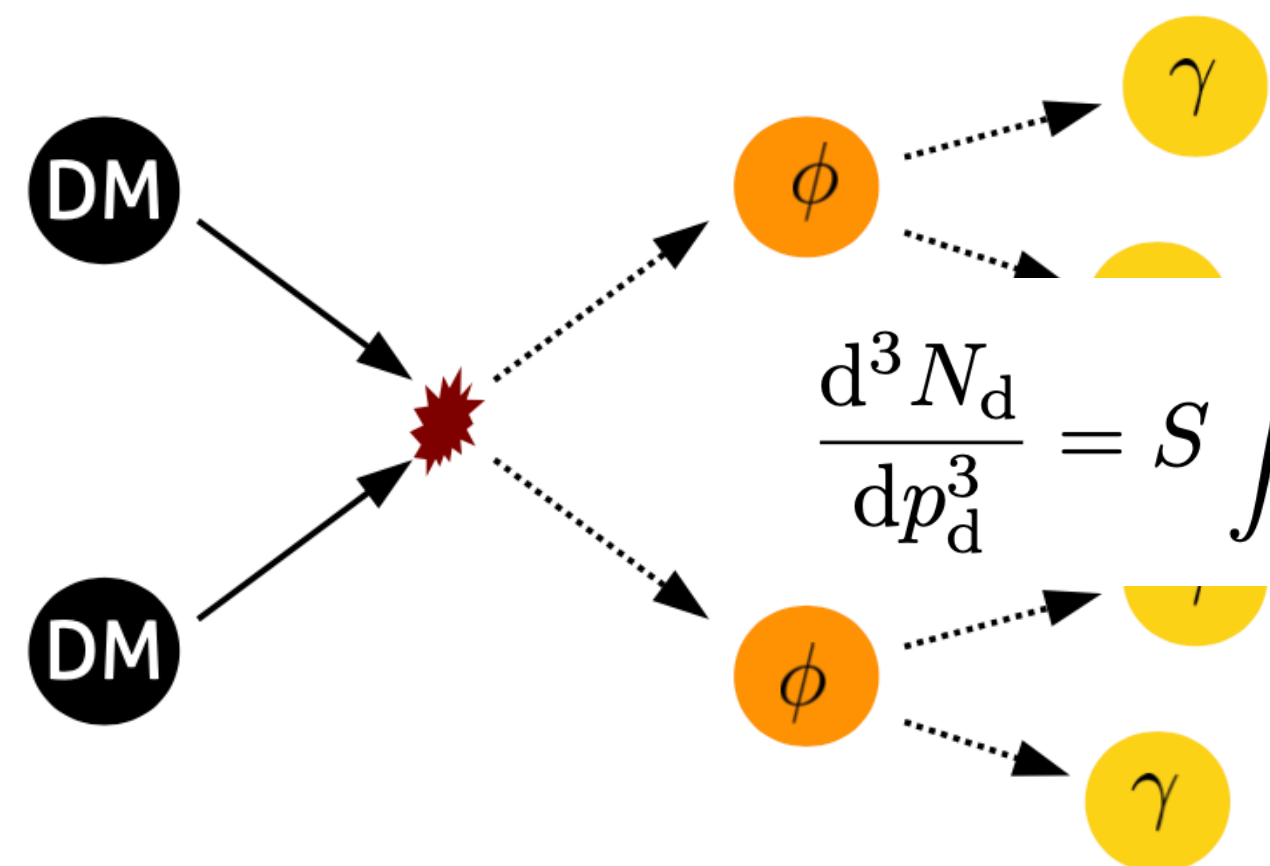
## Bremsstrahlung:

Photon production in "hard process"



## Box-shaped spectra:

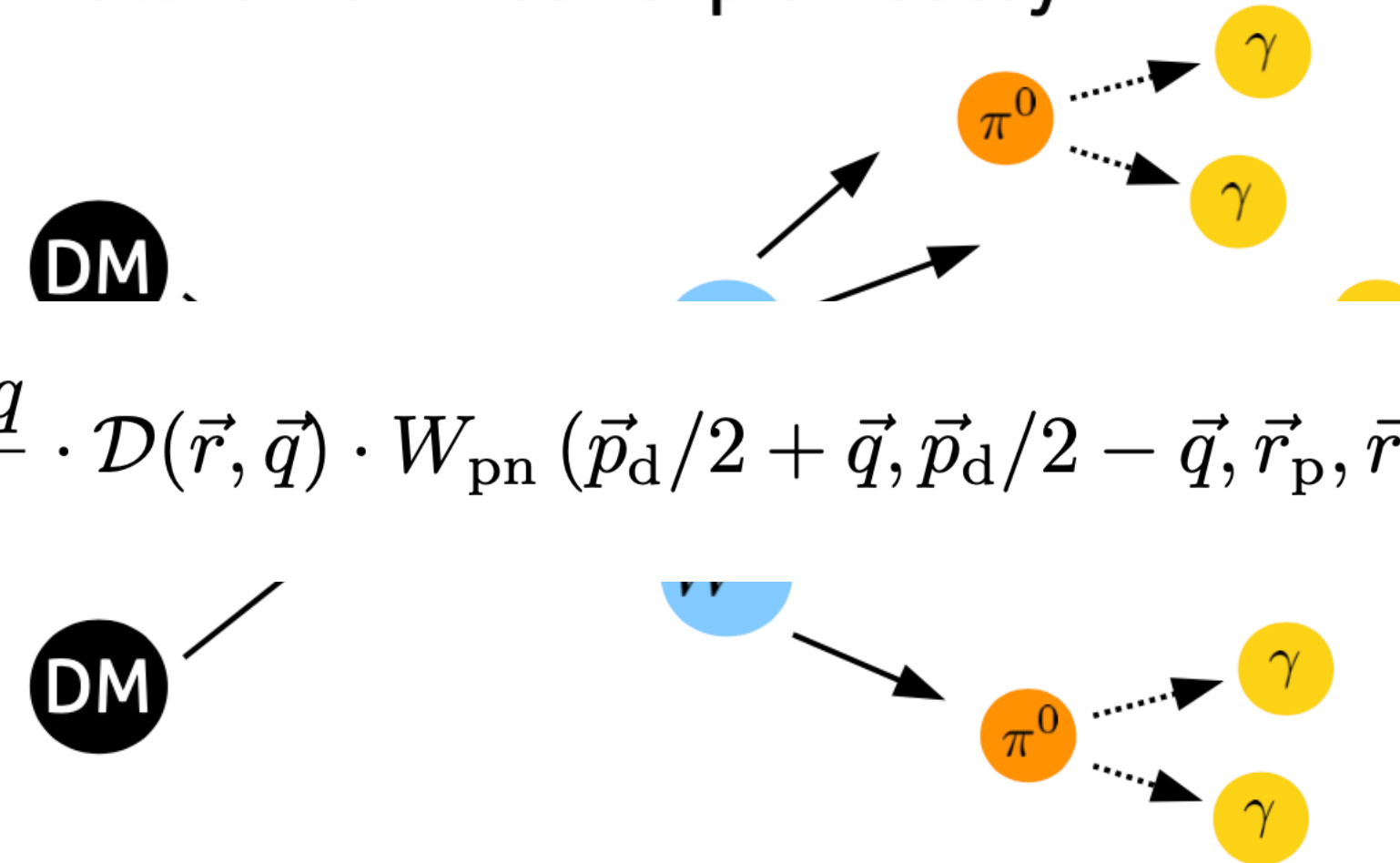
Photons from cascade decay



$$\frac{d^3 N_d}{dp_d^3} = S \int \frac{d^3 r_d d^3 r d^3 q}{(2\pi)^6} \cdot \mathcal{D}(\vec{r}, \vec{q}) \cdot W_{pn}(\vec{p}_d/2 + \vec{q}, \vec{p}_d/2 - \vec{q}, \vec{r}_p, \vec{r}_n)$$

## Continuum emission: (Prompt)

Photons from neutral pion decay



# Wigner approach

$$\frac{d^3 N_d}{dp_d^3} = S \int \frac{d^3 r_d d^3 r d^3 q}{(2\pi)^6} \cdot \mathcal{D}(\vec{r}, \vec{q}) \cdot W_{pn}(\vec{p}_d/2 + \vec{q}, \vec{p}_d/2 - \vec{q}, \vec{r}_p, \vec{r}_n)$$

$$\mathcal{D}(\vec{r}, \vec{q}) = \int d^3 \xi e^{-i\vec{q} \cdot \vec{\xi}} \varphi_d(\vec{r} + \vec{\xi}/2) \varphi_d^*(\vec{r} - \vec{\xi}/2)$$

$$W_{pn} = H_{pn}(\vec{r}_p, \vec{r}_n) G_{pn}(\vec{p}_d/2 + \vec{q}, \vec{p}_d/2 - \vec{q})$$

$$\varphi_d(r) = (\pi d^2)^{-3/4} e^{-r^2/(2d^2)}$$

$$\mathcal{D}(r, q) = 8e^{-r^2/d^2} e^{-q^2 d^2} \quad \frac{d^3 N_d}{dp_d^3} = \frac{3}{(2\pi)^6} \left( \frac{d^2}{d^2 + 4\sigma^2} \right)^{3/2} \int d^3 q e^{-q^2 d^2} G_{np}(\vec{p}_d, \vec{q})$$

$$h(r_{p/n}) = (2\pi\sigma^2)^{-3/2} e^{-r_{p/n}^2/(2\sigma^2)}$$