

Multi-Component Dark Matter from Minimal Flavor Violation

Based on Arxiv[2408.16812]

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Motivation: Why Flavor and Dark Matter

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Summary

Motivation: Why Flavor and Dark Matter

Flavor Violation

- The flavor symmetry of SM:

$$SU(3)^5 = SU(3)_{q_L} \times SU(3)_{u_R} \times SU(3)_{d_R} \times SU(3)_{l_L} \times SU(3)_{e_R},$$

broken only by Yukawa couplings.

$$\mathcal{L}_{\text{yuk}} = -\bar{q}_L Y_u \tilde{H} u_R - \bar{q}_L Y_d H d_R - \bar{\ell}_L Y_e H e_R + h.c.$$

- Minimal Flavor Violation (MFV)**: new physics respects SM flavor structure. In MFV, quark Yukawa matrices are treated as **spurions** transforming as

$$Y_u \sim (3, \bar{3}, 1), Y_d \sim (3, 1, \bar{3}) \text{ under } \mathcal{G}_F = SU(3)_{q_L} \times SU(3)_{u_R} \times SU(3)_{d_R}$$

For a new flavor-violating interaction like

$$\mathcal{L} = C_{ij} (\bar{u}_{Ri} \gamma^\mu u_{Rj}) O_{\alpha\beta}$$

MFV requires that the coupling C_{ij} take the form,

$$C_{ij} = c_0 \delta_{ij} + \epsilon c_1 (Y_u^\dagger Y_u)_{ij} + \epsilon^2 \left[c_2 (Y_u^\dagger Y_u Y_u^\dagger Y_u)_{ij} + c'_2 (Y_u^\dagger Y_d Y_d^\dagger Y_u)_{ij} \right] + \dots$$

Stability of flavored DM

Flavored DM candidate:

- Consider a colorless but flavored field:

$$\chi \sim (n_{q_L}, m_{q_L}) \times (n_{u_R}, m_{u_R}) \times (n_{d_R}, m_{d_R})$$



(n_i, m_i) count the number of $(\mathbf{3}_i, \bar{\mathbf{3}}_i)$ representations under $SU(3)_i$

- General decay vertices of χ into SM:

$$O_{\text{decay}} = \chi \times \underbrace{q_L \dots \bar{q}_L}_{\bar{A}} \dots \underbrace{u_R \dots \bar{u}_R}_{\bar{B}} \dots \underbrace{d_R \dots \bar{d}_R}_{\bar{C}} \dots \underbrace{Y_u \dots Y_u^\dagger}_{\bar{D}} \dots \underbrace{Y_d \dots Y_d^\dagger}_{\bar{E}} \dots + O_{\text{weak}}$$



A and $\bar{A}$ denote the number of q and $\bar{q}$ fields

Stability of flavored DM

Stability of flavored DM:

- Allowing O_{decay} requires:

$$SU(3)_c : (A - \bar{A} + B - \bar{B} + C - \bar{C}) \bmod 3 = 0,$$

$$SU(3)_{q_L} : (n_{q_L} - m_{q_L} + A - \bar{A} + D - \bar{D} + E - \bar{E}) \bmod 3 = 0,$$

$$SU(3)_{u_R} : (n_{u_R} - m_{u_R} + B - \bar{B} - D + \bar{D}) \bmod 3 = 0,$$

$$SU(3)_{d_R} : (n_{d_R} - m_{d_R} + C - \bar{C} - E + \bar{E}) \bmod 3 = 0,$$

So for χ to be unstable is $(n - m) \bmod 3 = 0$, where $n = n_{q_L} + n_{u_R} + n_{d_R}$ and $m = m_{q_L} + m_{u_R} + m_{d_R}$. Also it follows that the O_{decay} is forbidden and χ is stable if

$$(n - m) \bmod 3 \neq 0$$

Flavored Representations and Stability

(n, m)	$SU(3)_Q \times SU(3)_{u_R} \times SU(3)_{d_R}$	Stable?
(0, 0)	(1, 1, 1)	
(1, 0)	(3, 1, 1), (1, 3, 1), (1, 1, 3)	Yes
(0, 1)	($\bar{3}$, 1, 1), (1, $\bar{3}$, 1), (1, 1, $\bar{3}$)	Yes
(2, 0)	(6, 1, 1), (1, 6, 1), (1, 1, 6) (3, 3, 1), (3, 1, 3), (1, 3, 3)	Yes
(0, 2)	($\bar{6}$, 1, 1), (1, $\bar{6}$, 1), (1, 1, $\bar{6}$) ($\bar{3}$, $\bar{3}$, 1), ($\bar{3}$, 1, $\bar{3}$), (1, $\bar{3}$, $\bar{3}$)	Yes
(1, 1)	(8, 1, 1), (1, 8, 1), (1, 1, 8) (3, $\bar{3}$, 1), (3, 1, $\bar{3}$), (1, 3, $\bar{3}$) ($\bar{3}$, 3, 1), ($\bar{3}$, 1, 3), (1, $\bar{3}$, 3)	

Key Points:

- Representations are assigned under flavor group: $SU(3)_Q \times SU(3)_{u_R} \times SU(3)_{d_R}$
- Stability depends on triality:
 $(n - m) \bmod 3 \neq 0$
- Our model: red-highlighted entry is a flavor triplet under $SU(3)_{u_R}$

Source: Adapted from [arXiv:1105.1781v2]

Model Setup: Flavoured Dark Matter

A Flavored Scalar as Dark Matter

- Introduce scalar field $S_i \sim (1,3,1)$ under $SU(3)_{qL} \times SU(3)_{uR} \times SU(3)_{dR}$.
- scalar potential allowed by the MFV:

$$\begin{aligned} V(H, S) = & m_S^2 S_i^* \left(a_0 \delta_{ij} + \epsilon a_1 (Y_u^\dagger Y_u)_{ij} + \dots \right) S_j \\ & + \lambda S_i^* \left(b_0 \delta_{ij} + \epsilon b_1 (Y_u^\dagger Y_u)_{ij} + \dots \right) S_j (H^\dagger H) \\ & + \left(\lambda_0 \delta_{ij} \delta_{kl} + \epsilon \lambda_1 \delta_{ij} (Y_u^\dagger Y_u)_{kl} + \dots \right) S_i^* S_j S_k^* S_l \end{aligned}$$

with $(Y_u)_{ij} = (V^\dagger \hat{Y}_u)_{ij}$ and $(\hat{Y}_u)_{ij} = y_u^i \delta_{ij}$.

$$\Rightarrow \text{mass splitting:} \quad M_j^2 - M_i^2 = \epsilon \left(a_1 m_S^2 + b_1 \frac{\lambda v^2}{2} \right) \left\{ (y_u^j)^2 - (y_u^i)^2 \right\}$$

Dimension-6 operators

- Five classes of dim-6 operators involving quarks and the scalar field S_i

$$\mathcal{L}_{d=6} = \frac{1}{\Lambda^2} \sum_I c_{ijkl}^I \mathcal{O}_{ijkl}^I$$

$$\begin{aligned} \mathcal{O}_{ijkl}^1 &= (\bar{q}_{Li} \gamma^\mu q_{Lj}) (S_k^* i \overleftrightarrow{\partial}_\mu S_l), & \mathcal{O}_{ijkl}^2 &= (\bar{u}_{Ri} \gamma^\mu u_{Rj}) (S_k^* i \overleftrightarrow{\partial}_\mu S_l), \\ \mathcal{O}_{ijkl}^3 &= (\bar{d}_{Ri} \gamma^\mu d_{Rj}) (S_k^* i \overleftrightarrow{\partial}_\mu S_l), & \mathcal{O}_{ijkl}^4 &= (\bar{q}_{Li} \tilde{H} u_{Rj}) (S_k^* S_l) \\ \mathcal{O}_{ijkl}^5 &= (\bar{q}_{Li} H q_{Lj}) (S_k^* S_l), \end{aligned}$$

- MFV dictates coefficients c_{ijkl}^4 as series expansions in Yukawa spurions

$$\begin{aligned} c_{ijkl}^4 &= c_1 (Y_u)_{il} \delta_{kj} + c_2 (Y_u)_{ij} \delta_{kl} \\ &+ \epsilon \left[c_3 (Y_u Y_u^\dagger Y_u)_{ij} \delta_{kl} + c_4 (Y_u Y_u^\dagger Y_u)_{il} \delta_{kj} + c_5 (Y_u)_{ij} (Y_u^\dagger Y_u)_{kl} + c_6 (Y_u)_{il} (Y_u^\dagger Y_u)_{jl} \right] \\ &+ \mathcal{O}(\epsilon^2) \end{aligned}$$

- Focus on the $\mathcal{O}^4$ operator at the order of ϵ^0 .

$$\mathcal{L}_{d=6} = \frac{1}{\Lambda^2} \left[c_1 \bar{u}_i (m_i P_R + m_j P_L) u_j (S_j^* S_i) + c_2 (m_i \bar{u}_i u_i) (S_j^* S_j) \right]$$

Decay of Heavy Components

Decay processes

The decay interaction of the heavy component takes the form:

$$\mathcal{L}_{d=6} = \frac{1}{\Lambda^2} \left[c_1 \bar{u}_i (m_i P_R + m_j P_L) u_j (S_j^* S_i) + c_2 (m_i \bar{u}_i u_i) (S_j^* S_j) \right]$$

Mass splittings arise from MFV:

$$M_j^2 - M_i^2 = \epsilon \left(a_1 m_S^2 + b_1 \frac{\lambda v^2}{2} \right) \left\{ (y_u^j)^2 (y_u^i)^2 \right\}, \quad \frac{M_3^2 - M_1^2}{M_2^2 - M_1^2} = \frac{y_t^2 - y_u^2}{y_c^2 - y_u^2} \simeq \frac{y_t^2}{y_c^2}$$

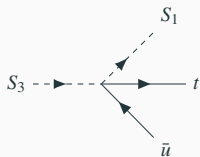
$$\Delta M = M_3 - M_1, \quad \delta M = M_2 - M_1$$

Mass spectrum :

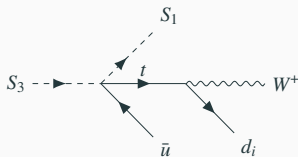
$$\text{normal} : M_1 < M_2 < M_3, \quad \text{inverted} : M_3 < M_2 < M_1$$

Decay processes

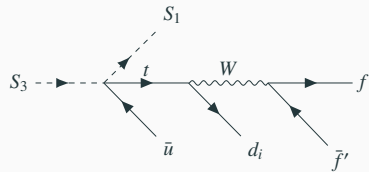
- For a normal mass hierarchy ($M_1 < M_2 < M_3$), the heaviest state S_3 can decay through the following channels depending on ΔM :



(a) $\Delta M \gtrsim m_t$



(b) $m_W + m_d^i \lesssim \Delta M \lesssim m_t$



(c) $m_u + m_d^i + m_f + m_{f'}^i \lesssim \Delta M \lesssim m_W + m_d^i$

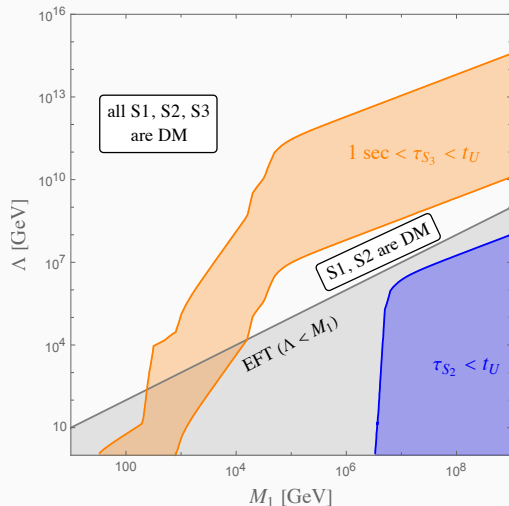
Decay widths: (a) $\sim \frac{m_t^2 (\Delta M)^5}{480\pi^3 \Lambda^4 M_3^2}$

(b) $\sim \frac{(\Delta M)^{11} |V_{ti}|^2}{41472\pi^5 \Lambda^4 M_3^2 m_t^2 v^2}$

(c) $\sim \frac{(\Delta M)^{13} |V_{ti}|^2}{11612160\pi^7 \Lambda^4 M_3^2 m_t^2 v^4}$

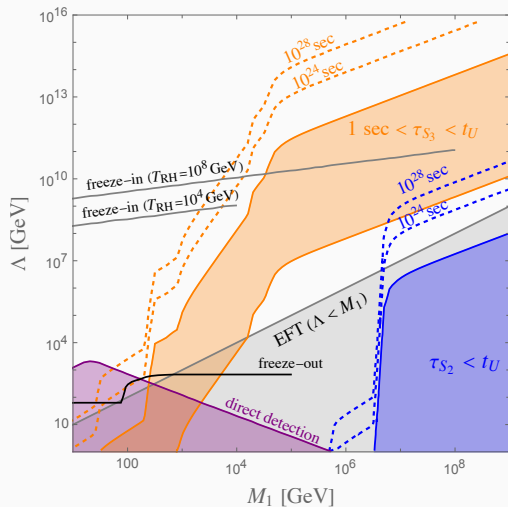
Parameter Space and Phenomenology

Parameter Space and DM Regions



- The model is mainly defined by four parameters: M_1 , Δ , λ , ϵ
- We take $\epsilon \ll 1$ to justify the MFV expansion: $\epsilon \sim \frac{1}{(4\pi)^2} \sim 10^{-2}$
- $\lambda = 0$
- Mass splittings: $\epsilon \simeq \frac{2\Delta M}{y_t^2 M_1} \simeq \frac{2\delta M}{y_c^2 M_1}$
- $\tau_{S_i} > \tau_U \rightarrow \text{DM}$
- $\tau_{S_i} < \tau_U \rightarrow \text{not DM and have to decay}$

Experiment constraints



- Direct detection (purple region)
- Relic abundance from freeze-out and freeze-in (black line)
- Indirect detection (dotted lines)

Summary

Summary:

- We investigated a possibility that, under the MFV framework, the heavy components of such a new flavored field are also stable over cosmological timescales and constitute a significant portion of DM.
- We studied parameter space where two or three components of the flavored scalar field have sufficient longevity to serve as DM.

Thank you for your attention!