

Astrophysical neutrino point sources as a probe of new physics

arXiv:2304.08533 in collaboration with <u>Stefan Vogl</u>

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Brief Overview

- Introduction
 - Neutrinos and Cosmic Rays
 - Observations: TXS 0506+056 and NGC 1068
- New Physics and the Neutrino Signal
 - Secret neutrino interaction and the cosmic neutrino background
- Mean free path and flux
 - Massfull case
 - Massless case
- Results
 - Universal Coupling
 - Tauphilic Coupling
 - Future Sources
- Conclusion



Cosmic Rays and Neutrinos



Galactic Neutrino Signals



What if...? How to "stop" Neutrinos



... there would be Sterile Neutrinos?

20 26 See e.g.: arXiv:2211.16520v2

arXiv:2212.00737v2



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2c

See e.g.: arXiv:2301.08756v2



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... there would be strong Neutrino-Neutrino interactions?



In the standard model the interaction is negligibley small but...



Cosmic Neutrino Background as a Milk Glass



How far can they come? The mean free path

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Our Neutrino Sector (Assumptions):

- -Flavor universal coupling
- -Normal mass ordering
- -Majorana fermion

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 $m(\nu_{\rm light}) \stackrel{\rm non-relativistic today}{\scriptstyle (m_\nu \gg T_\nu)}$

$$\sum_{i} m_i = 0.1 \,\mathrm{eV}$$
$$\Delta m_{12}^2, \Delta m_{23}^2$$

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Interaction Rate for the Massless Neutrinos



Mean Free Path

Mean free path: $\lambda_{\rm MFP} = 1 / \sum_i \Gamma_i(E_{\rm a})$

y=0.05 $m_{\phi} \in \{0.25, 2.5\} \text{ MeV}$

Illustrative example:

single neutrino species



Flux

$$\Phi_0(E) = \hat{\Phi}_0 \cdot \left(\frac{E}{1 \,\text{TeV}}\right)^{-\gamma} \cdot \mathbf{T}(E)$$
Transmittance
$$\mathbf{T}(E) = e^{-\frac{d}{\lambda_{\text{MFP}}(E)}}$$



Flux



Problem: We <u>don't</u> know the original amount of neutrinos emitted by the source...

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Estimate:

$$\frac{n}{n_0} = \frac{\int_{E_{\min}}^{E_{\max}} dE A_{\text{eff}}(E) \Phi(E)}{\int_{E_{\min}}^{E_{\max}} dE A_{\text{eff}}(E) \Phi_0(E)} \ge q \qquad \text{with absorption} \\ (\text{milky}) \qquad \text{measured number} \\ (\text{transparent}) \qquad \text{measured} \qquad \text{measured}$$

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Here:
$$q=0.5$$



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Color Code: NGC 1068 TXS 0506+056



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- **Brown Line:** Lab K⁻ decay (flavor dependent)



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- Pink Line: Lab constrain Z-decay



Tauphilic interactions

Example: NGC 1068 Flavor Universal Coupling Coupling only to Tau-Neutrinos



Future Sources – Outlook to PKS 1424+240

Color Code:

PKS 1424+240 Combined estimated limits: NGC 1068 and TXS 0506+056

Facts:Distance: $1.8 \, \text{Gpc}$ Flux: $\hat{\Phi}_0 \approx \hat{\Phi}_{0,\text{NGC 1068}}$ Spectral index: $\gamma = 3.5$ Energy: $[E_{\min}, E_{\max}]_{\text{NGC 1068}}$

Relativistic today



2

Conclusion

- Neutrinos from astrophysical point sources have been measured by IceCube and are great messengers for astro- and particlephysics
- New physics (e.g. a scalar) can lead to interactions with the CNuB and thus turn the Universe opaque for Neutrinos
- Using the two observed sources TXS 0506+056 and NGC 1068 we put new estimated constraints on light scalar masses and neutrino coupling
- Two cases: lightest neutrino relativistic vs non-relativistic today
- Only estimate: the original neutrino emission at the source is not known
- More sources and higher energetic neutrinos could improve the constraints as well as a better understanding of the original neutrino luminosity of these sources

THANK YOU!



Crosssection

Flavor universal neutrino scattering cross section

$$\sigma_{\nu\nu}(s) = \frac{y^4}{32\pi((m_{\phi}^2 - s)^2 + m_{\phi}^2\Gamma_{\phi}^2)s^2} \left(\frac{s(5m_{\phi}^6 - 9m_{\phi}^4s + 6s^3)}{m_{\phi}^2 + s} + \frac{2(5m_{\phi}^8 - 9m_{\phi}^6s + 4m_{\phi}^2s^3)\log(\frac{m_{\phi}^2}{m_{\phi}^2 + s})}{2m_{\phi}^2 + s}\right)$$

 ϕ -pair production $E_{CM} \ge m_{\phi}$

$$\sigma_{\phi\phi}(s) = \frac{y^4}{64\pi s^2} \left(\frac{s^2 - 4m_{\phi}^2 s + 6m_{\phi}^4}{s - 2m_{\phi}^2} \log \left[\left(\frac{(s(s - 4m_{\phi}^2))^{1/2} + s - 2m_{\phi}^2}{(s(s - 4m_{\phi}^2))^{1/2} - s + 2m_{\phi}^2} \right)^2 \right] - 6(s(s - 4m_{\phi}^2))^{1/2} \right)$$



See also arXiv:2107.13568

Massless Neutrino

Rate approximations in different limit cases:

Heavy mediator mass:

Small mediator mass:

Resonance:



$$\Gamma_{\text{light}} \approx \frac{\pi y^4}{192\,\zeta(3)} \frac{1}{E_{\text{a}} T_{\nu}} n_{\nu_1}$$

$$\Gamma_{\rm NWA} \approx \frac{y^4}{384\zeta(3)} \frac{m_{\phi}^3}{E_{\nu}^2 T_{\nu}^2 \Gamma_{\phi}} \log[1 + e^{-\frac{m_{\phi}^2}{4E_{\nu} T_{\nu}}}] n_{\nu_1}$$

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 $E_{a} [GeV]$

Dashed-Black: Analytical Approximation Green: Numerical Result

Redshift broadening

In expanding Universe: Flux evolves according to transport equation

$$\frac{\partial \Phi(t, E_{\rm a})}{\partial t} = \frac{\partial}{\partial E_{\rm a}} [H(t)E_{\rm a}\Phi(t, E_{\rm a})] - \Phi(t, E_{\rm a})\Gamma(E_{\rm a}, t)$$

Which becomes $\frac{\partial Z(z, E_{\rm a})}{\partial z} = \frac{Z(z, E_{\rm a})\Gamma(E_{\rm a}, z)}{H(z)(1+z)}$ with $Z(z, E_{\rm a}) := (1+z)\Phi(z, E_{\rm a}[1+z])$

The redshift dependent rate is: $\Gamma_i(E_{\rm a},z) = \int \frac{\mathrm{d}^3 p}{(2\pi)^3} (1+z)^3 f_i(\vec{p}(1+z)) \, v_{M \not o l} \sigma_{\nu\nu}(s(E_{\rm a}(1+z),\vec{p}(1+z)))$

Transmittance:





Labconstraints



See: arXiv:1802.00009 arXiv:2003.05339

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Scalar mediated Nu-Nu Interactions Details on: Tauphilic Coupling vs Flavor-Universal Coupling

Relation between oberserved flux and
$$\phi_{obs} = \sum_{i} \exp(-\tau_{i}) \frac{\phi_{source}}{3}$$

Case I: Flavor-Universal
Coupling:
Case II: Flavor-Spezific
Coupling:
 $\phi_{obs} = \exp(-d\sum_{j} \sigma_{u}n_{j})\phi_{source}$
Attention: depend also
on neutrino-mass m_j !!!
Case II: Flavor-Spezific
Coupling:
 $\phi_{obs} = \sum_{i} \exp\left(-d\sum_{j} \sigma_{ij}n_{j}\right) \frac{\phi_{source}}{3}$
Comparison:
 $\exp\left(d\sum_{j} \sigma_{u}n_{j}\right) = \frac{1}{3}\sum_{i} \exp\left(-d|U_{\tau i}|^{2}\sum_{j} |U_{j,\tau}|^{2}\sigma_{\tau}n_{j}\right)$.
Approximation: small
Couplings:
Tails: $\sigma \propto y^{4}$
 $y_{\frac{h}{y_{\tau}}} \approx \left(\frac{|U_{\tau j}|^{2}}{3}\right)^{1/4} \approx 0.58$
 $\sum_{j} \sigma_{u}n_{j} = \frac{1}{3}\sum_{j} |U_{\tau j}|^{2}\sigma_{\tau}n_{j}$
Relation of
Reson $\sigma \propto y^{2}$
 $y_{\frac{h}{y_{\tau}}} \approx \sqrt{\frac{|U_{\tau j}|^{2}}{3}} \approx 0.28$
 $\sum_{j} \sigma_{u}n_{j} = \frac{1}{3}\sum_{j} |U_{\tau j}|^{2}\sigma_{\tau}n_{j}$
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