

A closer look at dark photon production modes

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25/06/24

Work in progress with

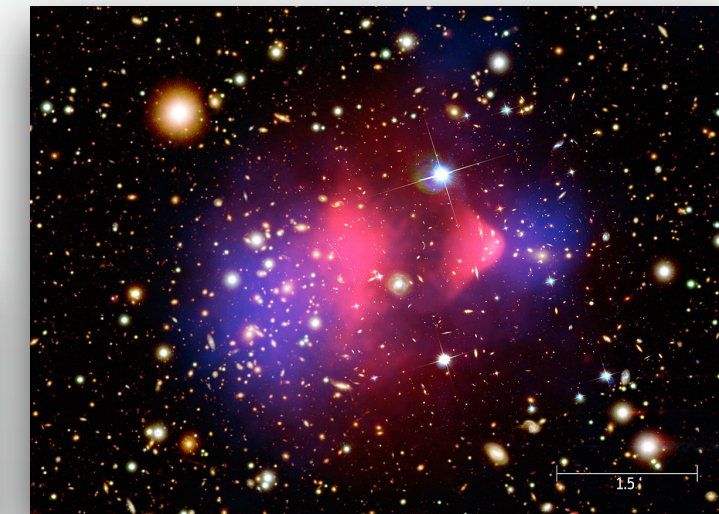
Adam Ritz (Victoria), Saeid Foroughi-Abari (Carleton)

Aidin Masouminia (Durham), Simon Plätzer (Vienna)
(Herwig Collaboration)

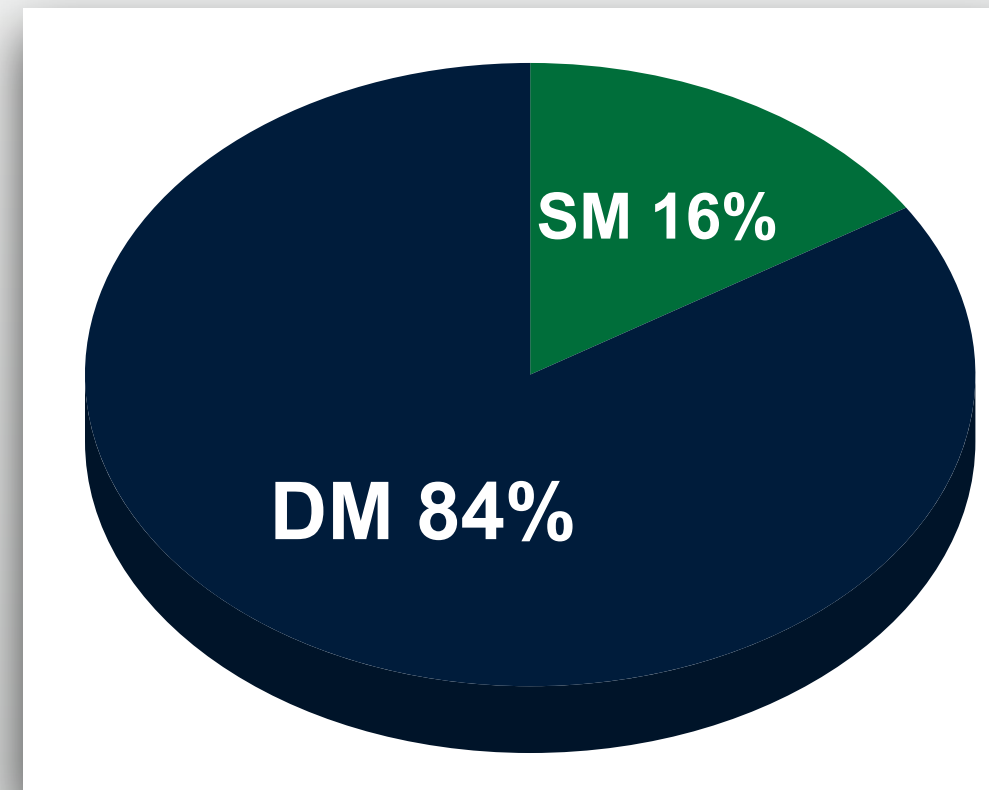


@  UCLouvain

Hot Topics and Problems in Particle Physics



Dark Matter



New Physics

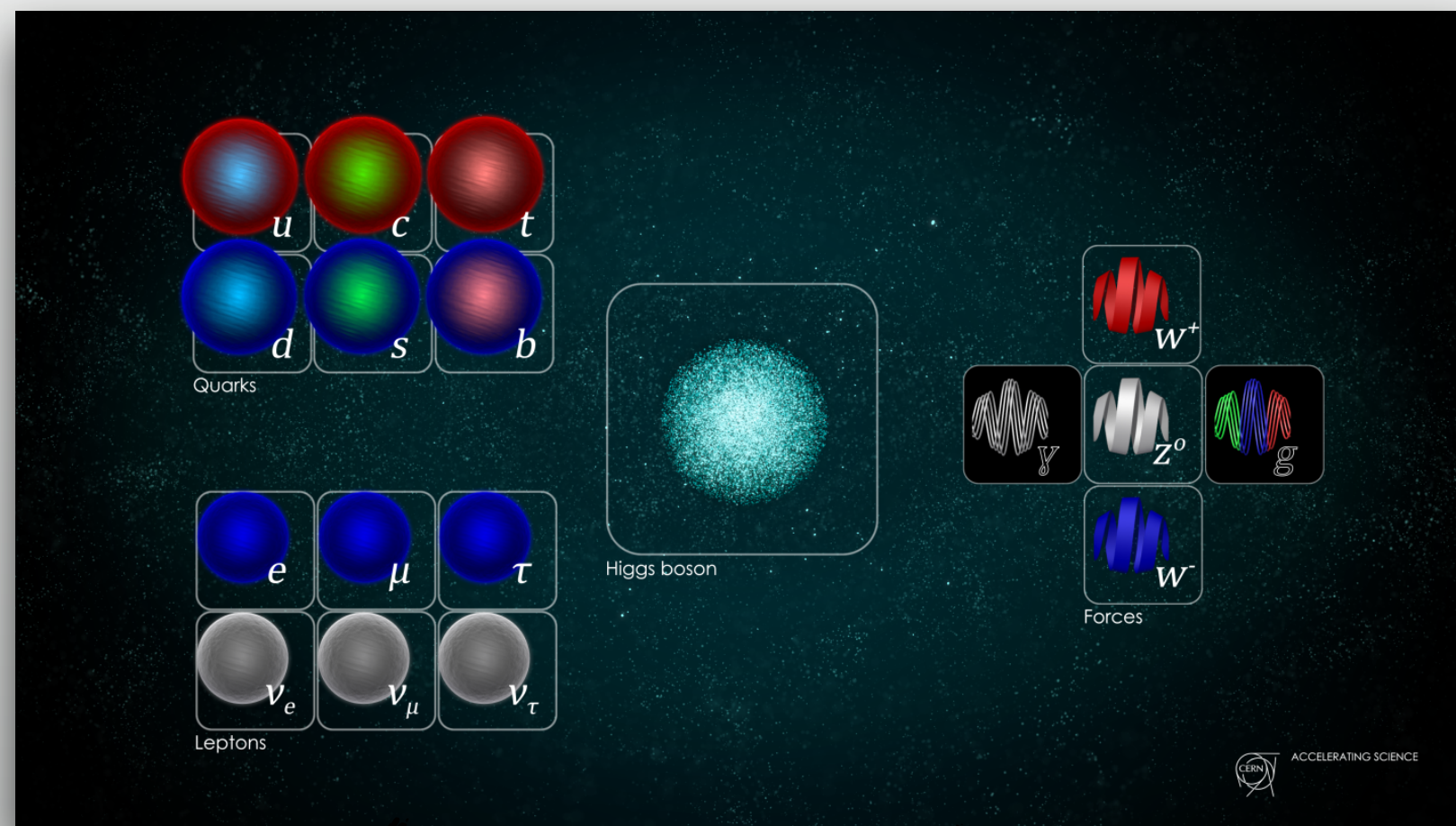


Baryon asymmetry



Neutrino masses

Standard Model



Tensions



Flavour anomalies



$g-2$

Hot Topics and Problems in Particle Physics

SM

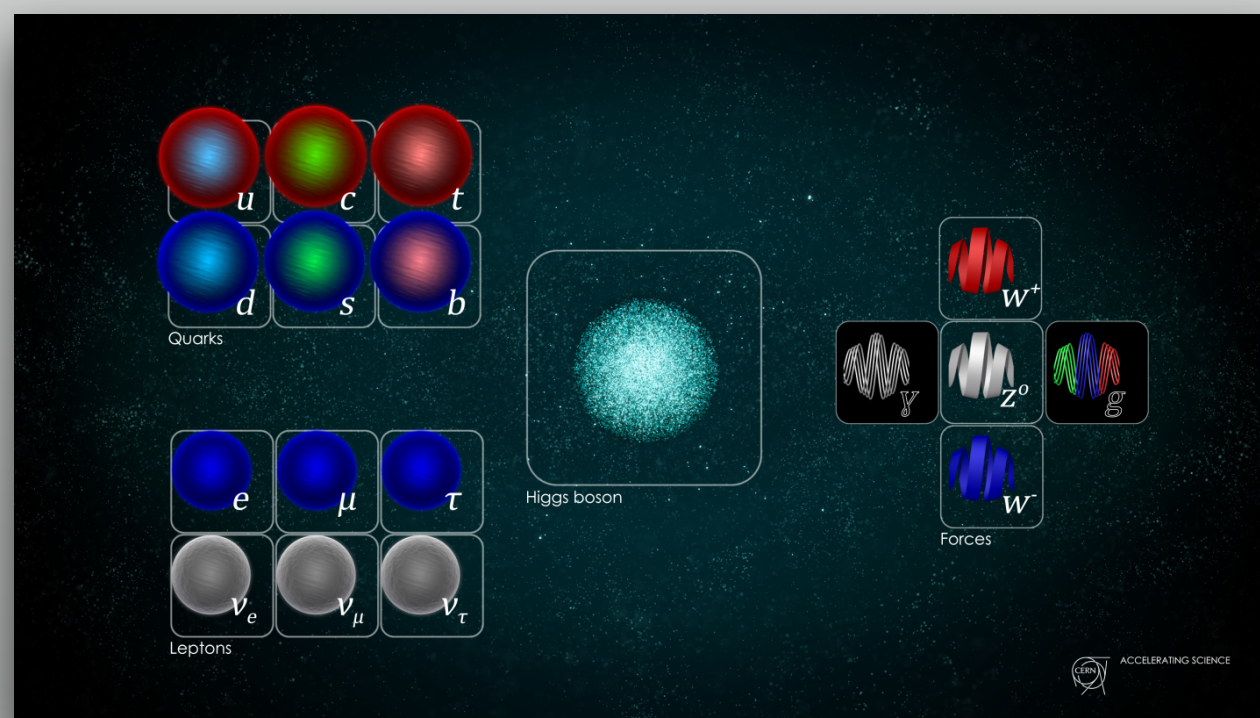
?

Beyond the
Standard Model

(BSM):

Neutrino masses
Baryon asymmetry

DM



Hot Topics and Problems in Particle Physics

SM

Renormalizable
+ gauge invariant

1. *Neutrino Portal*
2. *Higgs Portal*
3. *Vector Portal*

BSM:

DM

Neutrino masses

Baryon asymmetry

Dark Photon - Model Basics

Kinetic mixing: $\mathcal{L} = -\frac{\epsilon}{2} F_{\mu\nu} X^{\mu\nu}$
photon field strength \curvearrowright $F_{\mu\nu}$ \curvearrowleft $X^{\mu\nu}$ new vector field strength

$\rightarrow \mathcal{L} \supset -\frac{1}{4} A'^{\mu\nu} A'_{\mu\nu} + \frac{1}{2} m_{A'} A'^2 + \mathcal{L}_{\text{int}}^{\text{D}} + \mathcal{L}_{\text{int}}^{\text{SM}}$

mass range: $10 \text{ MeV} \lesssim m_{A'} \lesssim 2 \text{ GeV}$

...but could be anything!

Dark Photons + DM

SM

$$\mathcal{L} = -\frac{\epsilon}{2} F_{\mu\nu} X^{\mu\nu}$$

$$\mathcal{L}_{\text{int}}^{\text{SM}} = \epsilon e A'_\mu J_{\text{em}}^\mu$$

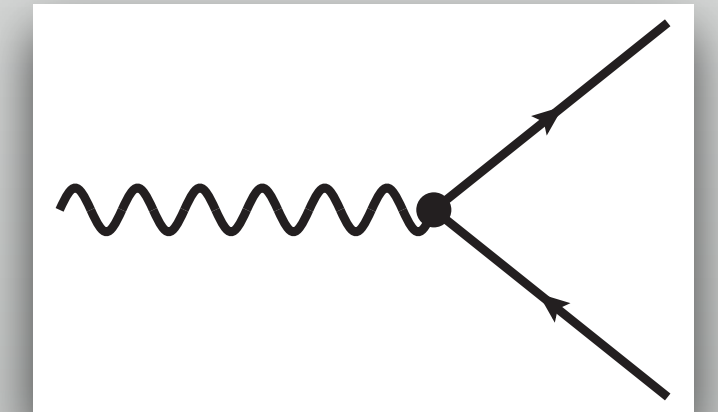
DM

Majorana fermion DM

$$\mathcal{L}_{\text{int}}^{\text{D}} = \frac{g_D}{2} A'_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi$$

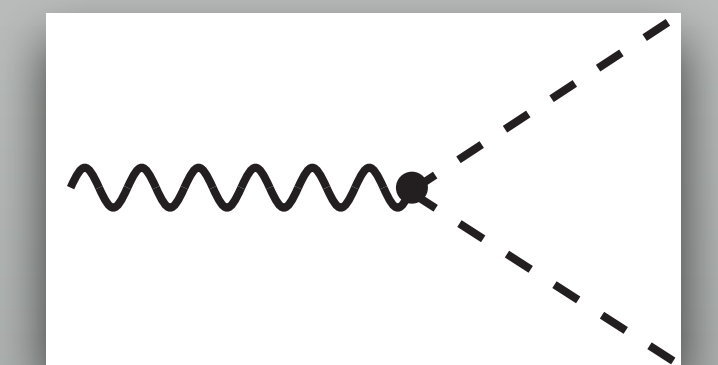
inelastic DM

$$\mathcal{L}_{\text{int}}^{\text{D}} = \frac{i}{2} g_D A'_\mu \bar{\chi}_2 \gamma^\mu \chi_1 + \text{h.c.}$$



complex-scalar DM

$$\mathcal{L}_{\text{int}}^{\text{D}} = i g_D A'_\mu \chi^* \overleftrightarrow{\partial}^\mu \chi$$



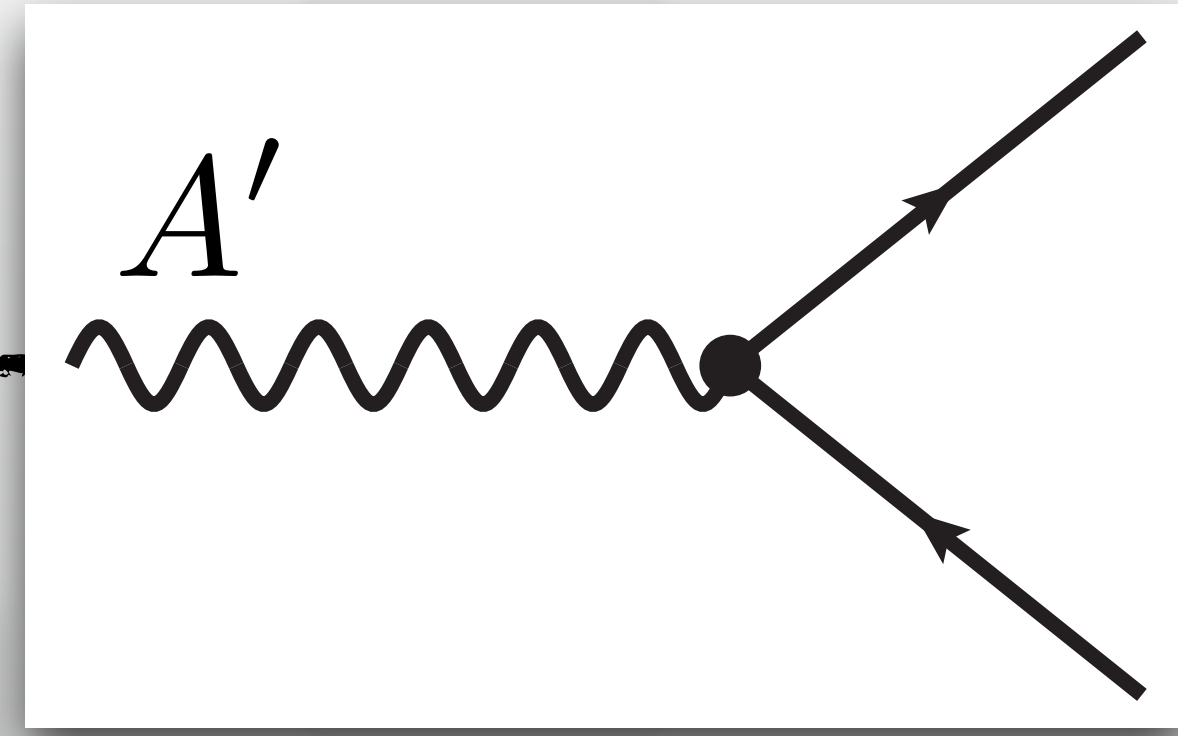
$$\mathcal{L} \supset -\frac{1}{4} A'^{\mu\nu} A'_{\mu\nu} + \frac{1}{2} m_{A'} A'^2 + \mathcal{L}_{\text{int}}^{\text{D}} + \mathcal{L}_{\text{int}}^{\text{SM}}$$

Dark Photons + SM

SM

$$\mathcal{L} = -\frac{\epsilon}{2} F_{\mu\nu} X^{\mu\nu}$$

$$\mathcal{L}_{\text{int}}^{\text{SM}} = \epsilon e A'_\mu J_{\text{em}}^\mu$$



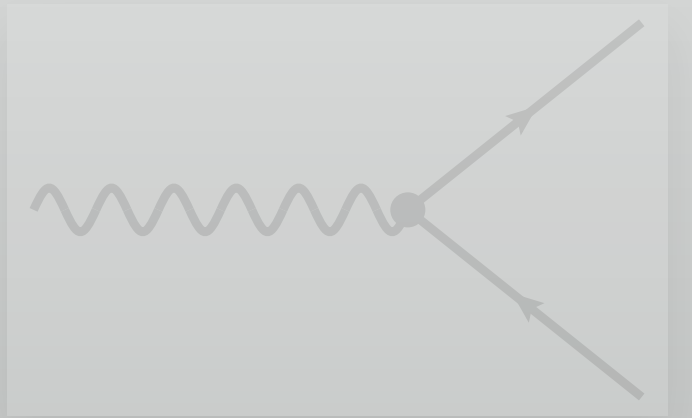
SM

SM

$$\mathcal{L} \supset -\frac{1}{4} A'^{\mu\nu} A'_{\mu\nu} + \frac{1}{2} m_{A'} A'^2 + \mathcal{L}_{\text{int}}^{\text{D}} + \mathcal{L}_{\text{int}}^{\text{SM}}$$

Majorana fermion DM

$$\mathcal{L}_{\text{int}}^{\text{D}} = \frac{g_{\text{D}}}{2} A'_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi$$



inelastic DM

$$\mathcal{L}_{\text{int}}^{\text{D}} = \frac{i}{2} g_{\text{D}} A'_\mu \bar{\chi}_2 \gamma^\mu \chi_1 + \text{h.c.}$$

complex-scalar DM

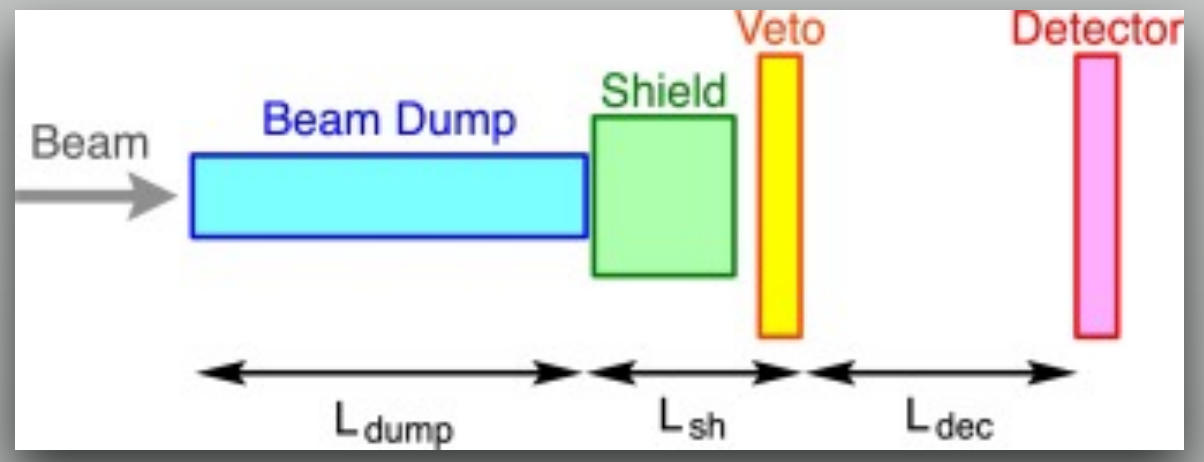
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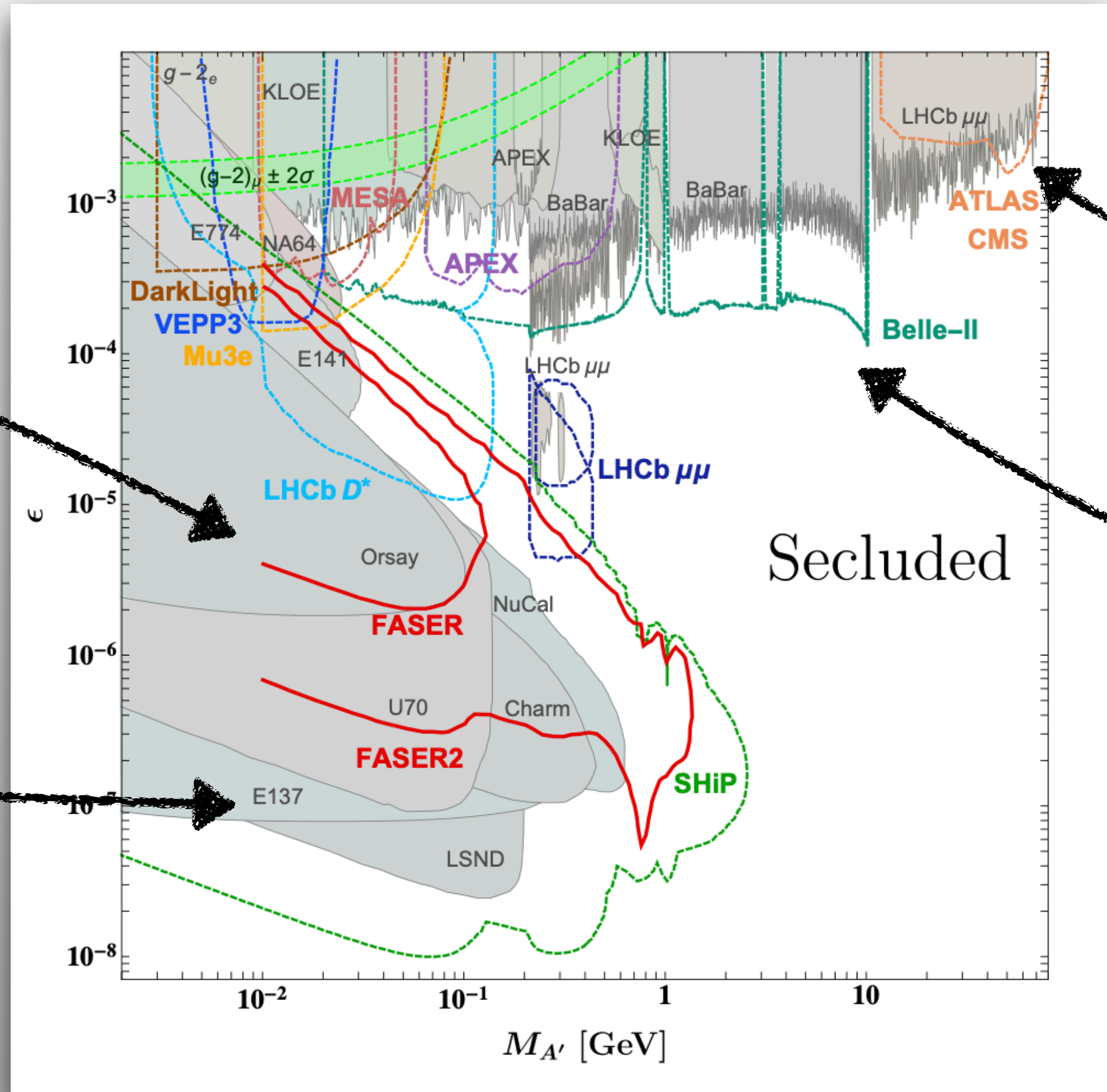
Search Strategy for Dark Photons

Proton Beam dumps:
 CHARM, NuCal, NA62
 U70, LSND

Electron beam dump:
 E137
 Orsay Linac

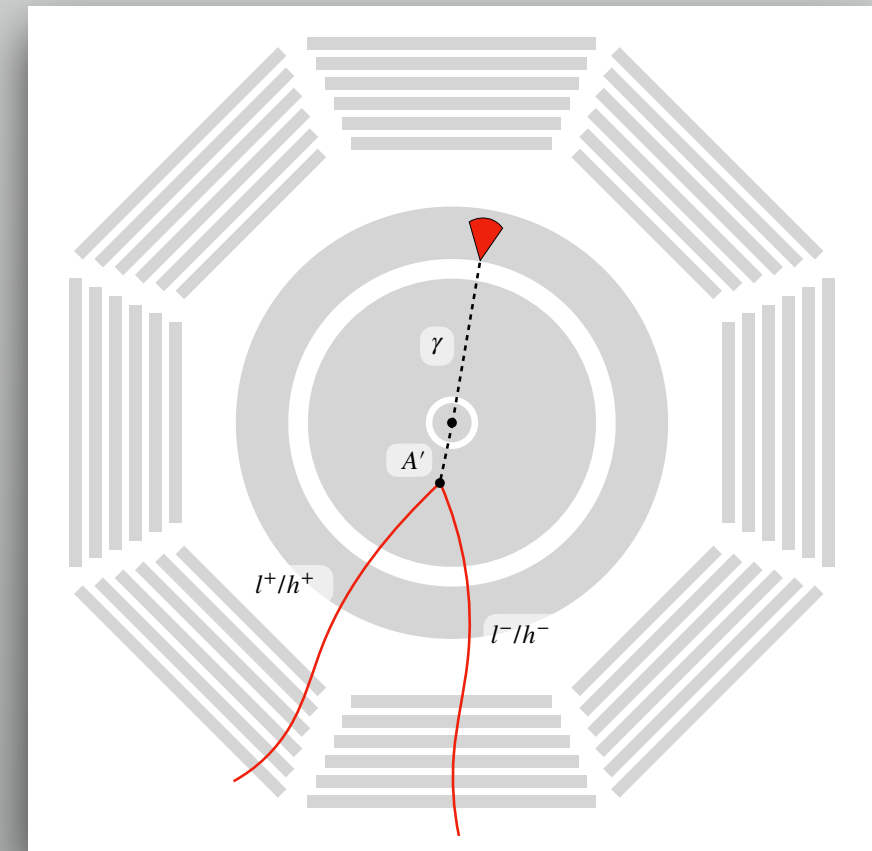


1507.02809



Colliders:
 pp LHCb, ATLAS, CMS

e^+e^- BaBar, Belle II, KLOE



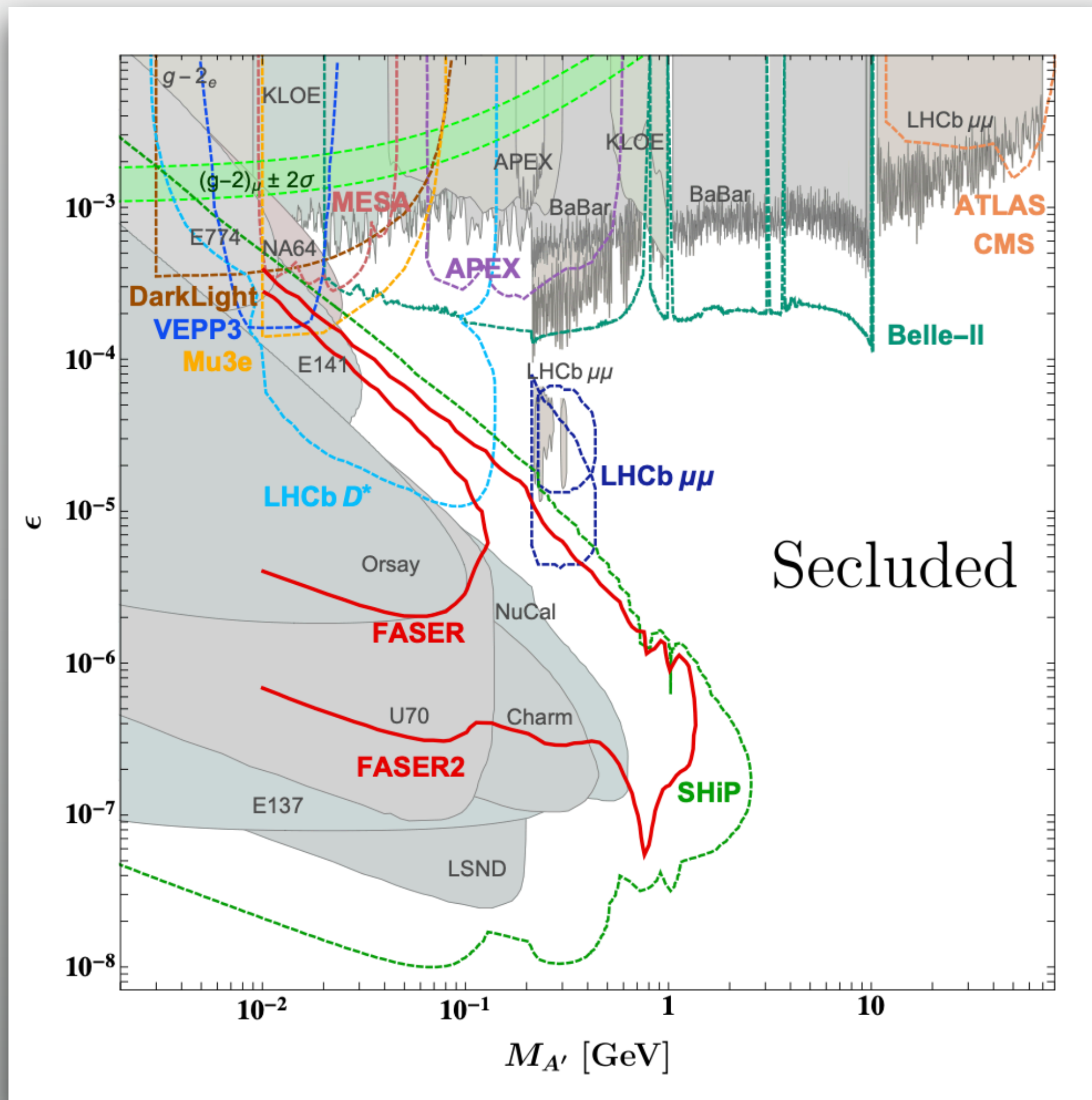
2202.03452

2203.05090

Search Strategy for Dark Photons

$$\mathcal{L} = -\frac{\epsilon}{2} F_{\mu\nu} X^{\mu\nu}$$

$$\mathcal{L}_{\text{int}}^{\text{SM}} = \epsilon e A'_\mu J_{\text{em}}^\mu$$

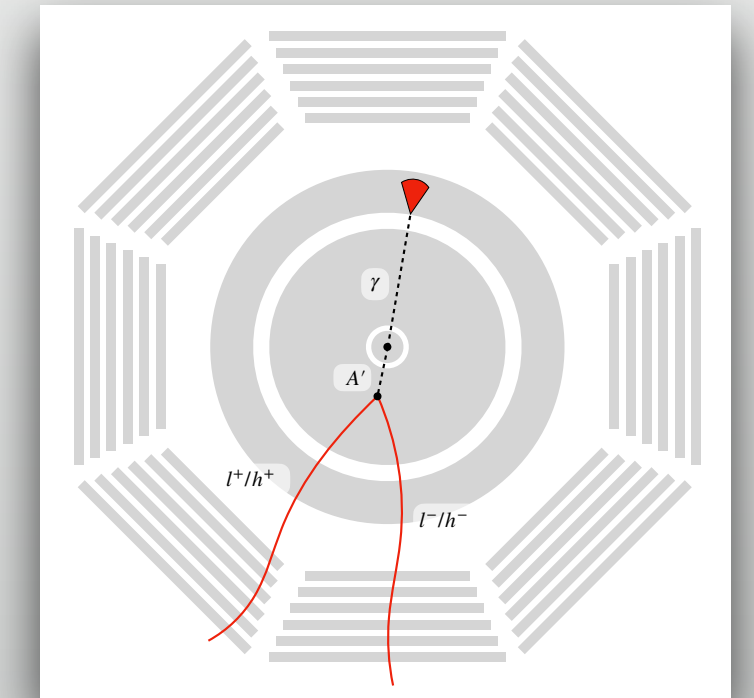


Decay:

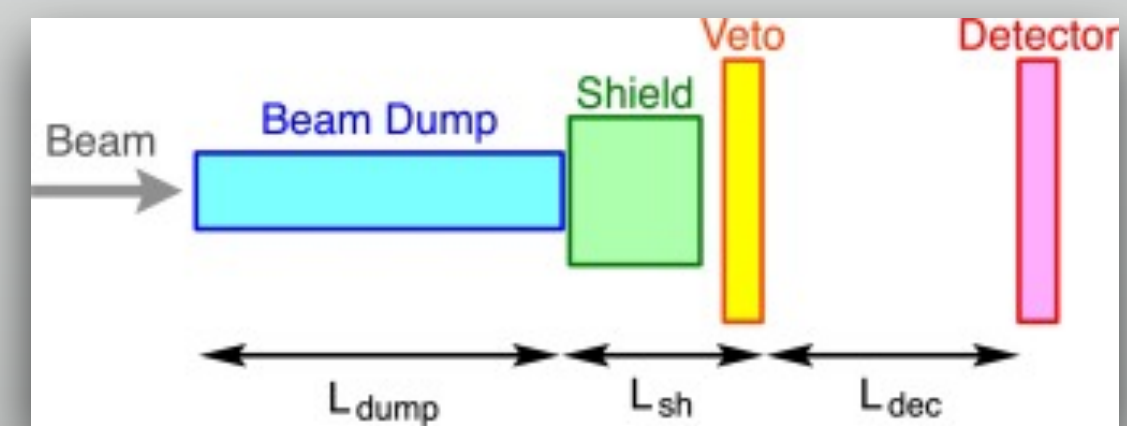
Prompt decay

Displaced decay

Long-lived particle



2202.03452



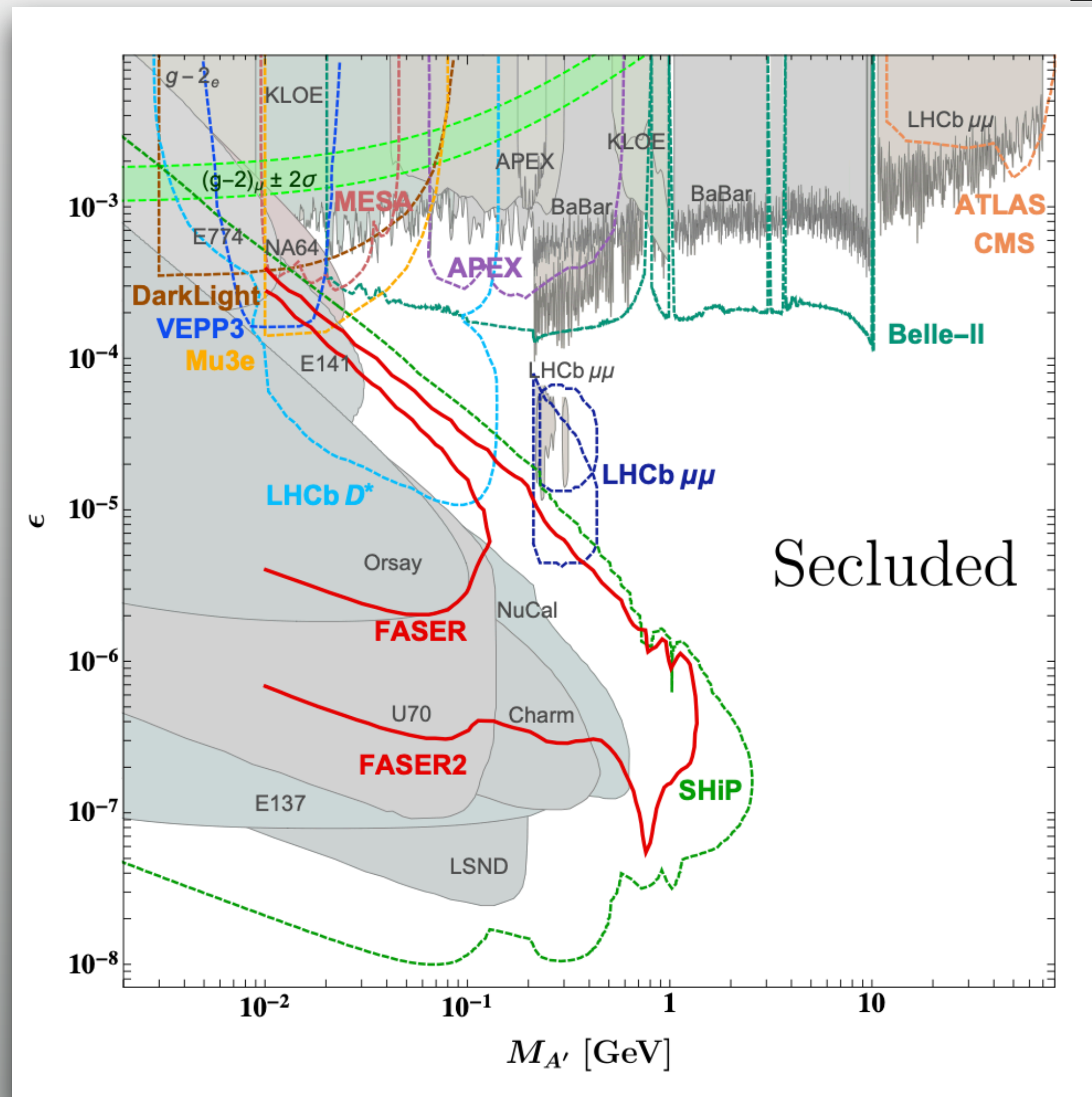
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Search Strategy for Dark Photons

$$\mathcal{L} = -\frac{\epsilon}{2} F_{\mu\nu} X^{\mu\nu}$$

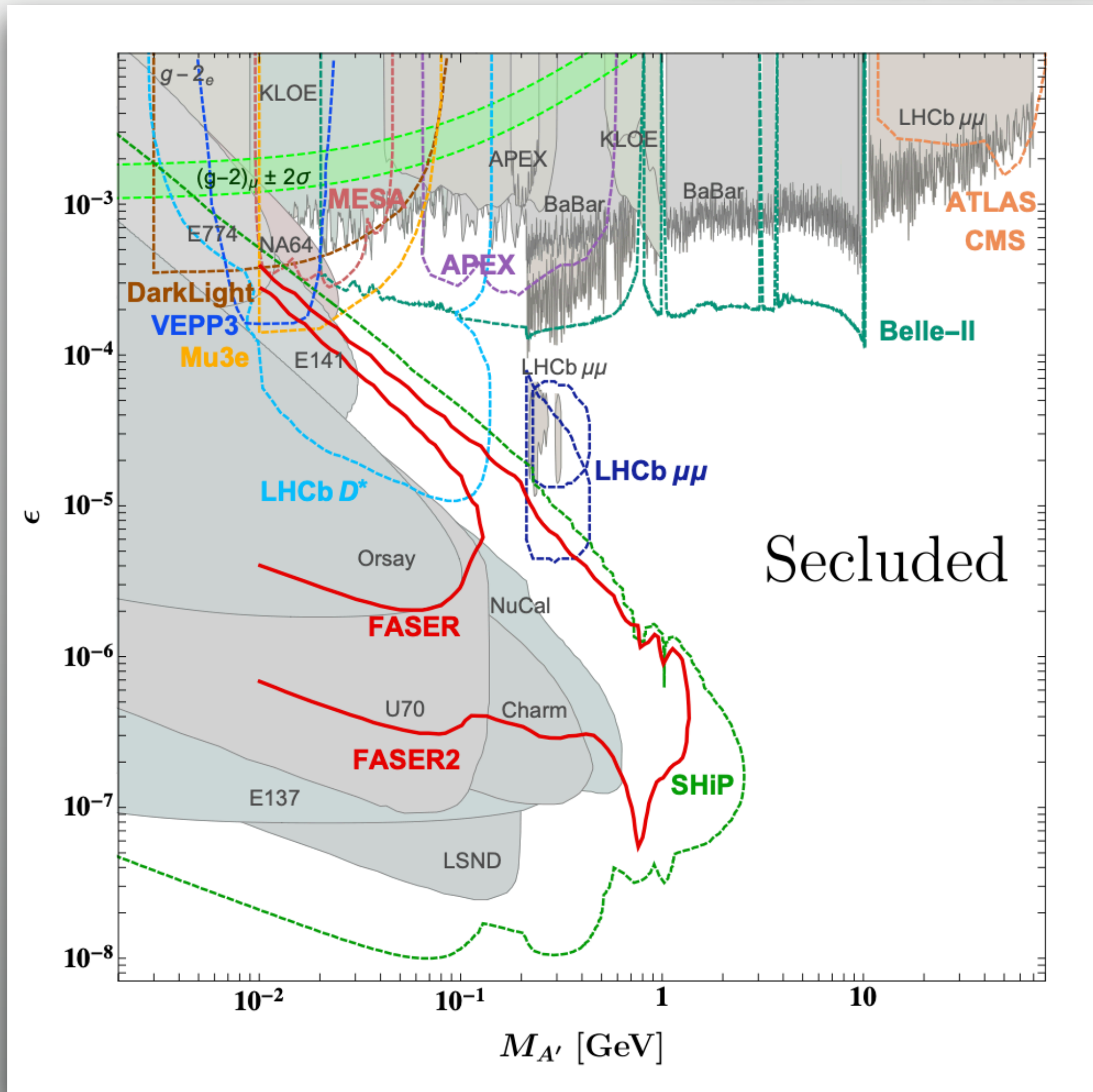
$$\mathcal{L}_{\text{int}}^{\text{SM}} = \epsilon e A'_\mu J_{\text{em}}^\mu$$

Decay



Production

Search Strategy for Dark Photons



Meson decays



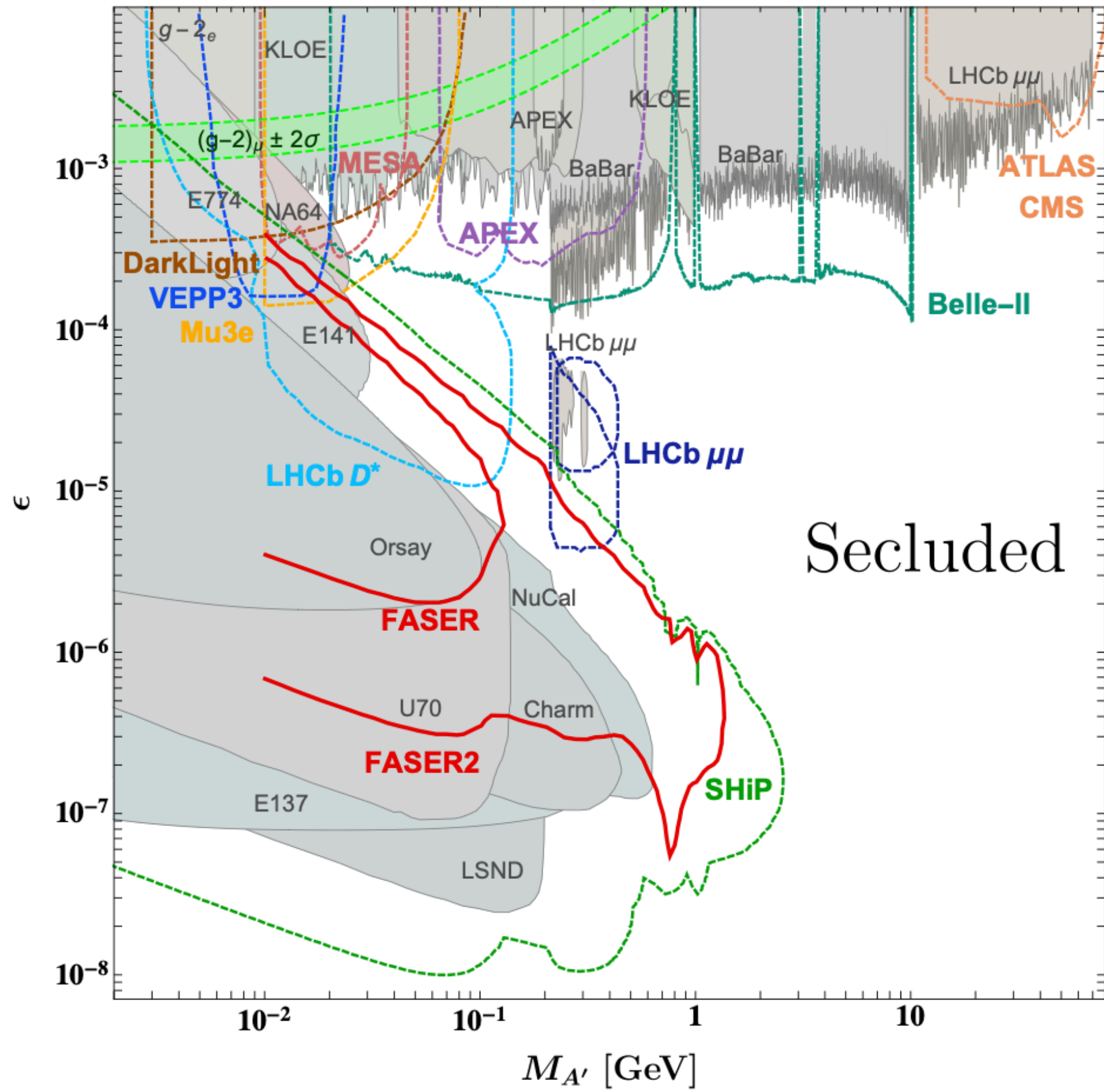
$$\text{BR}(\pi^0/\eta \rightarrow \gamma A') = 2\varepsilon^2 \left(1 - \frac{m_{A'}^2}{m_{\pi^0, \eta}^2} \right) \text{BR}(\pi^0/\eta \rightarrow \gamma\gamma)$$

Uncertainties in
meson production

2309.08604

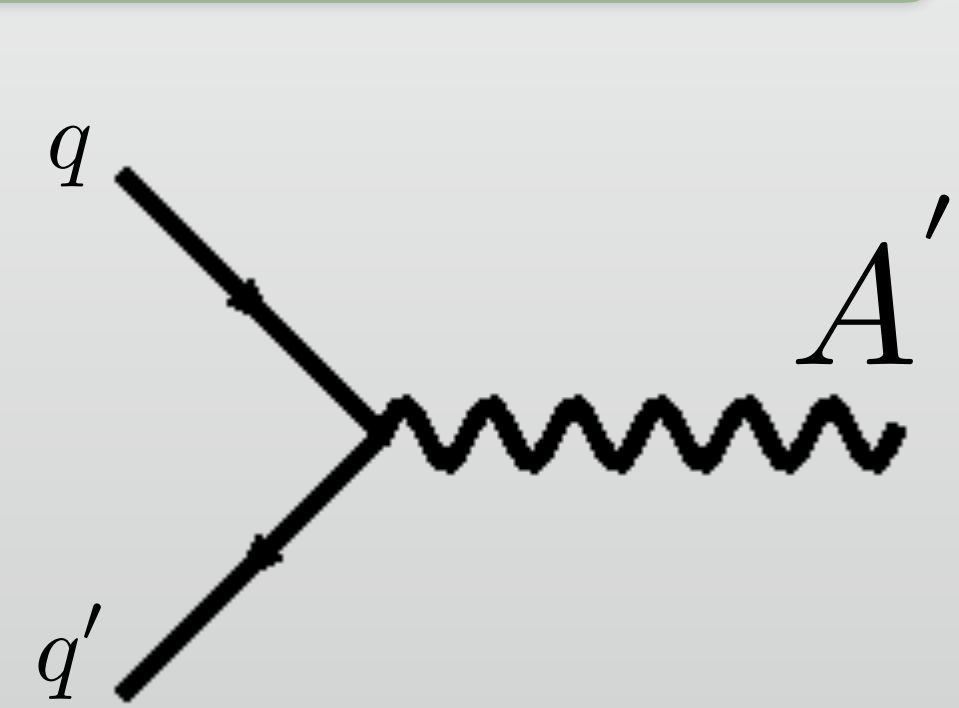
Production

Search Strategy for Dark Photons

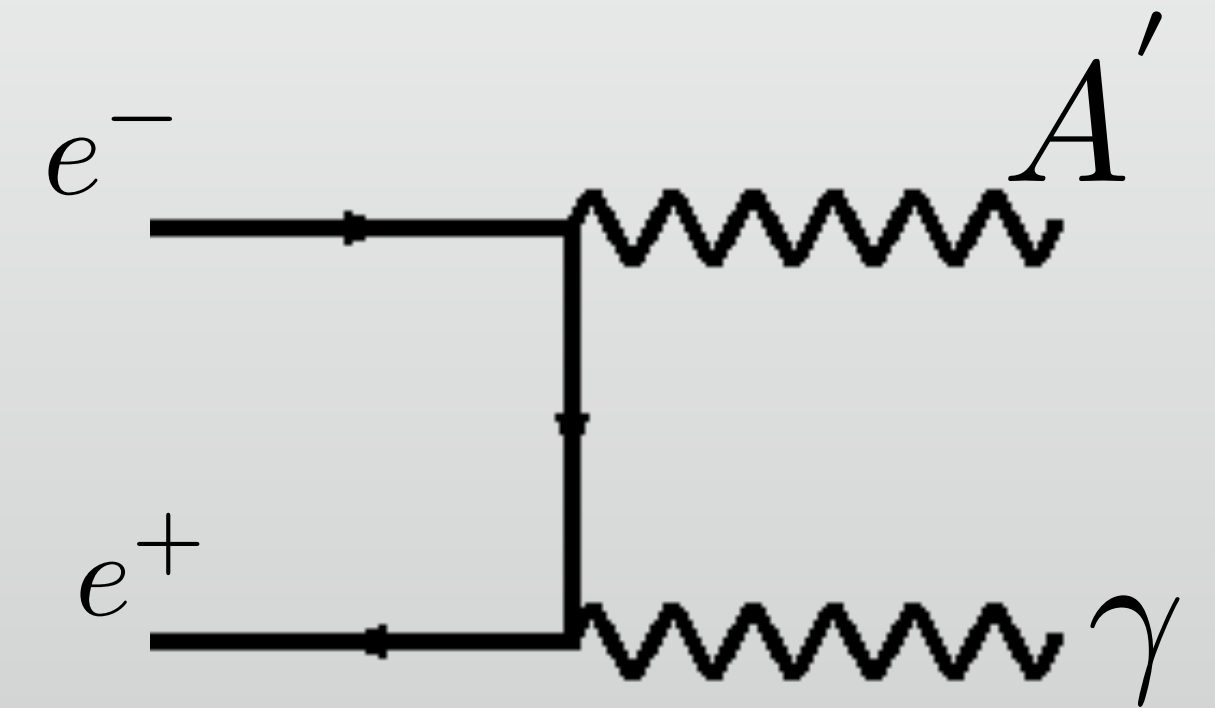


Secluded

Direct Production

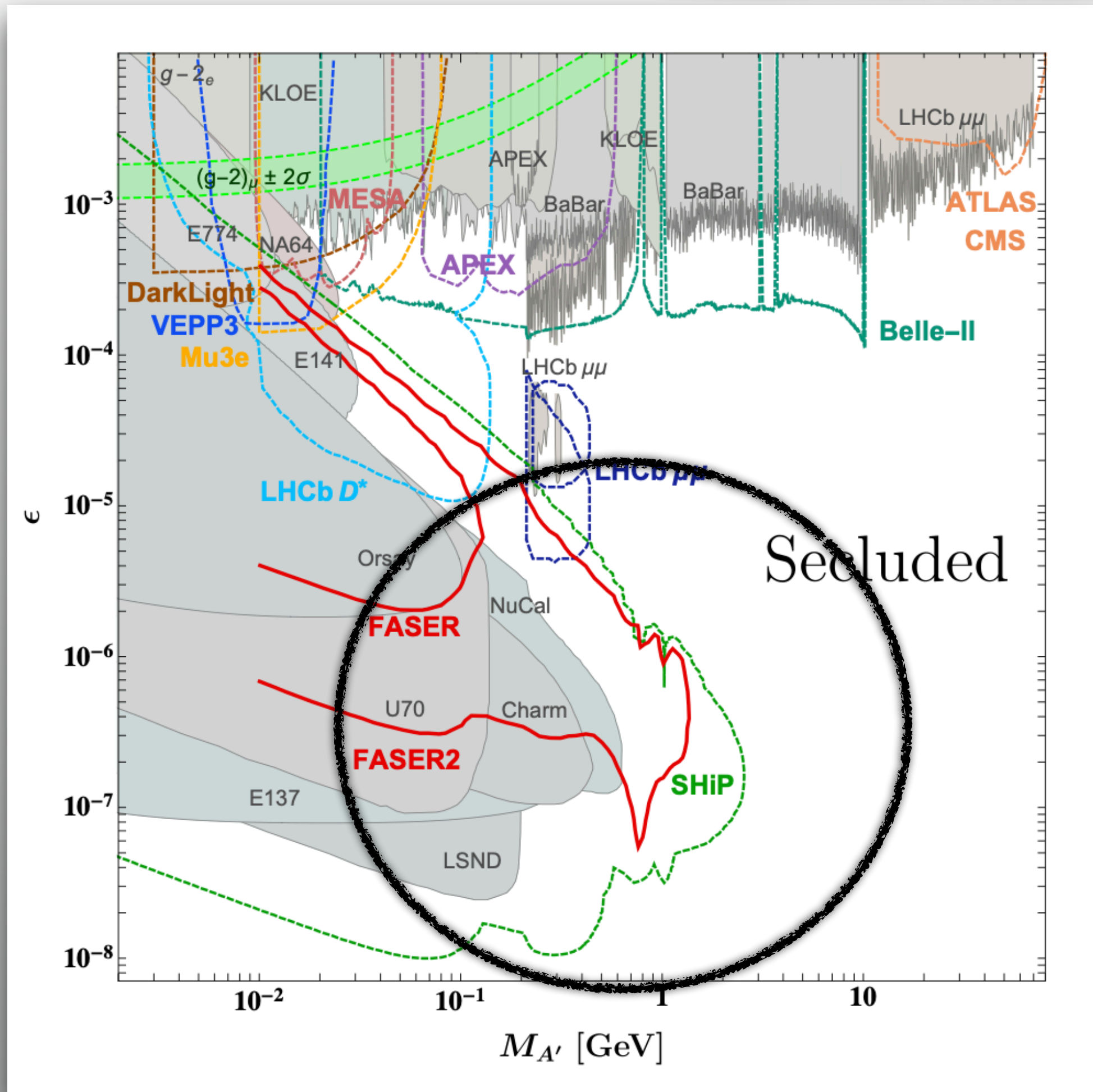


$$= \frac{12\pi}{m_{A'}^2} \text{BR}(A' \rightarrow q\bar{q})$$

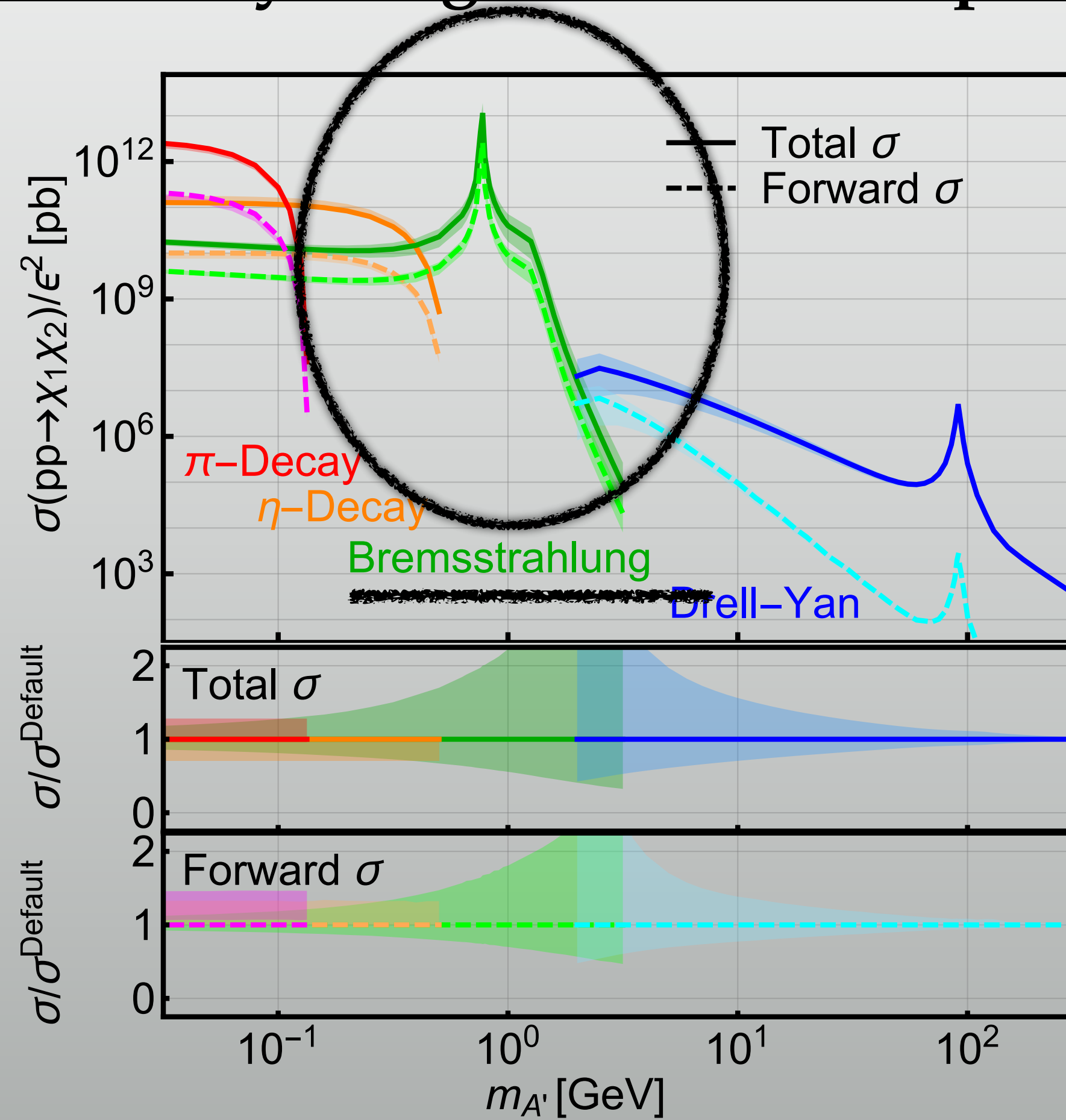


Production

Search Strategy for Dark Photons



Sensitivity Range of (future) experiments



Production

arXiv 1810.01879

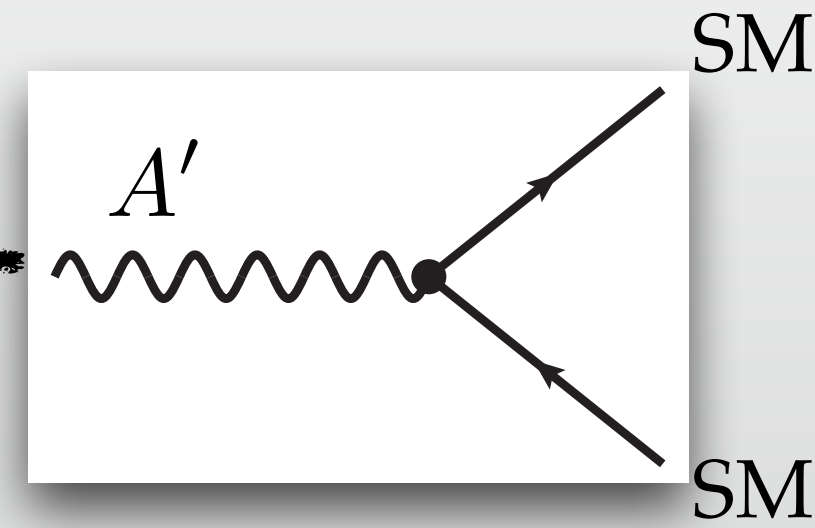
A. Berlin, F. Kling

Dark Photons from proton beam experiments

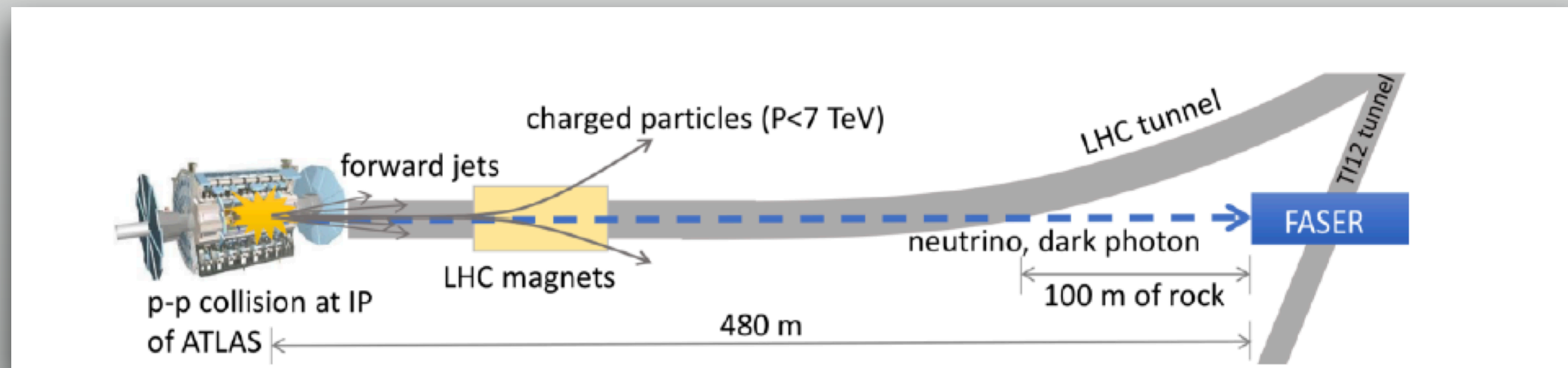
SM

$$\mathcal{L} = -\frac{\epsilon}{2} F_{\mu\nu} X^{\mu\nu}$$

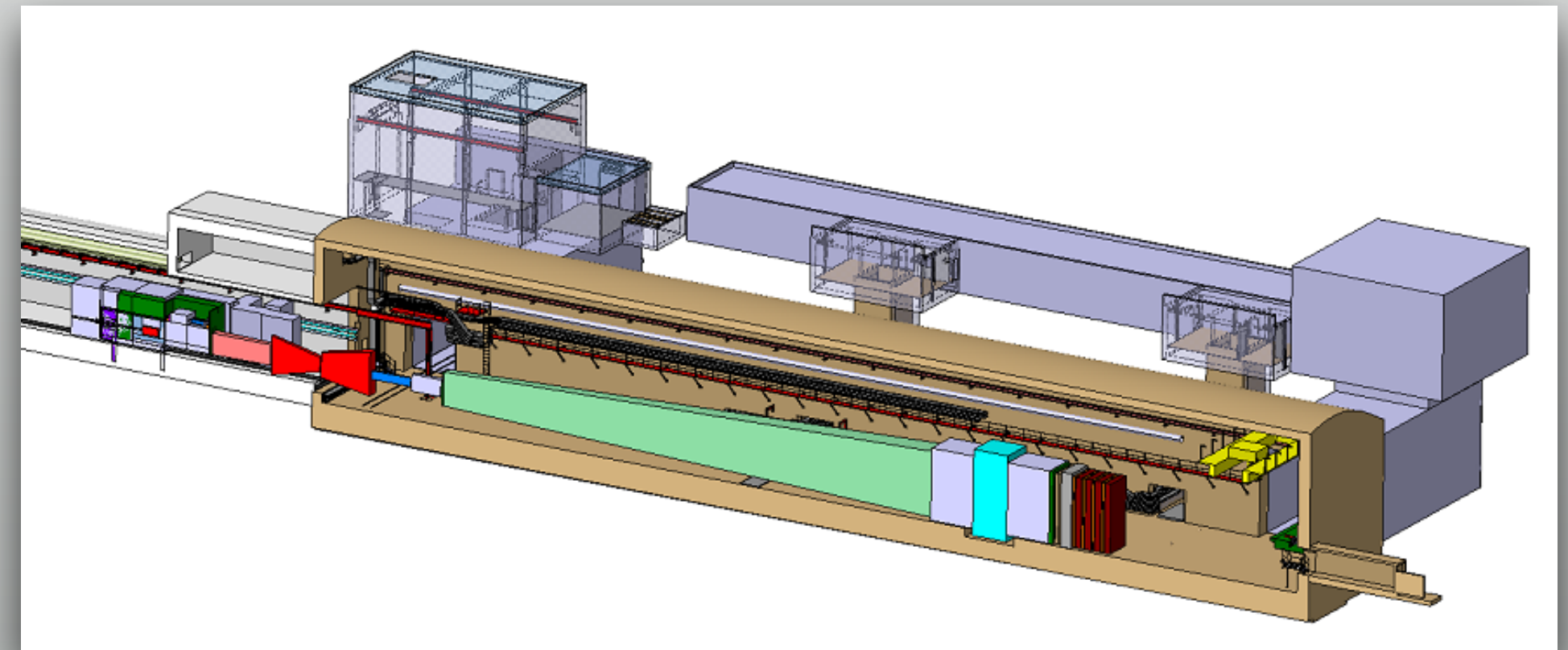
$$\mathcal{L}_{\text{int}}^{\text{SM}} = \epsilon e A'_\mu J_{\text{em}}^\mu$$



Discovery potential in (future) experiments



[FASER Collaboration]



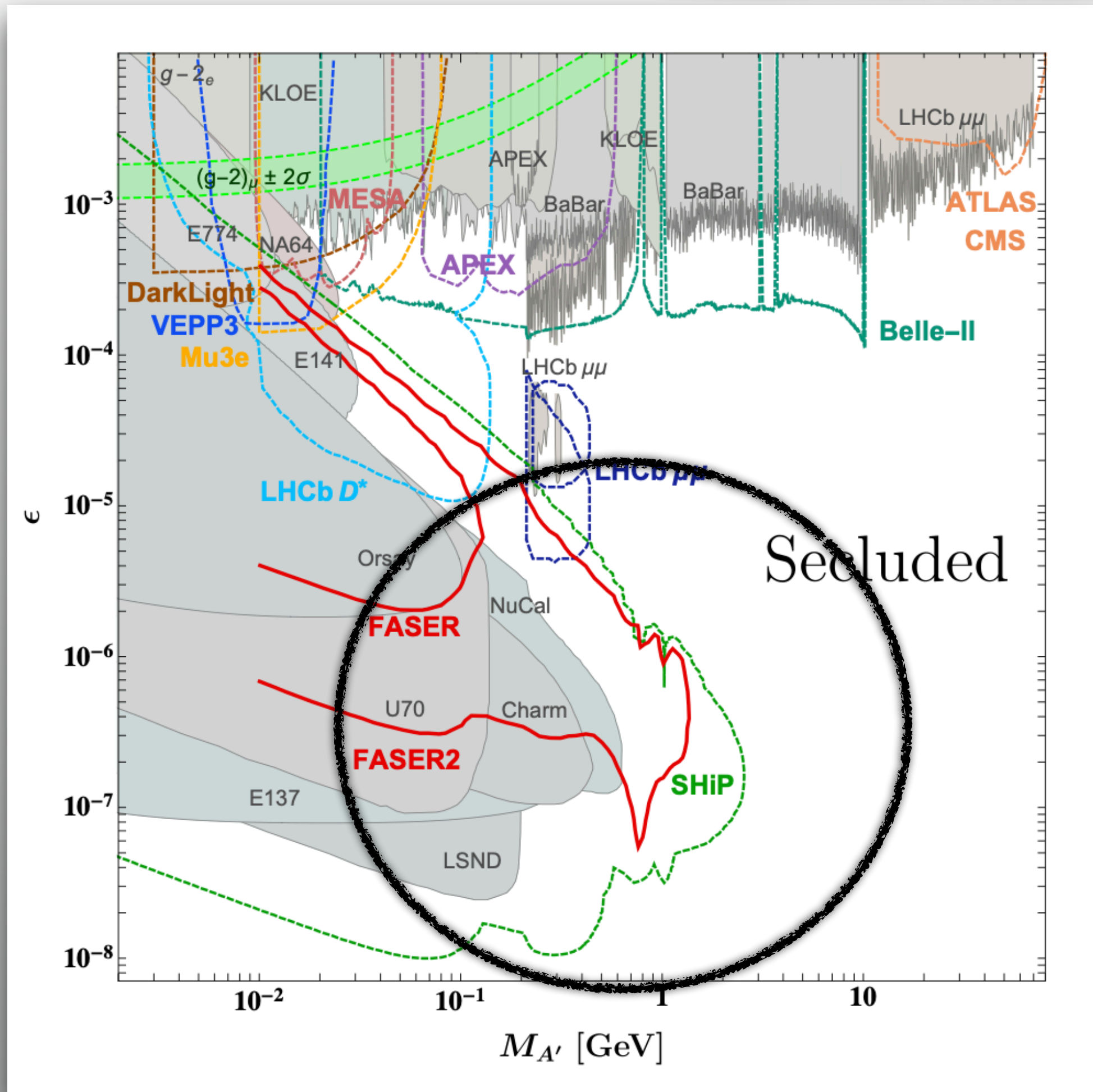
[SHiP Collaboration]

+

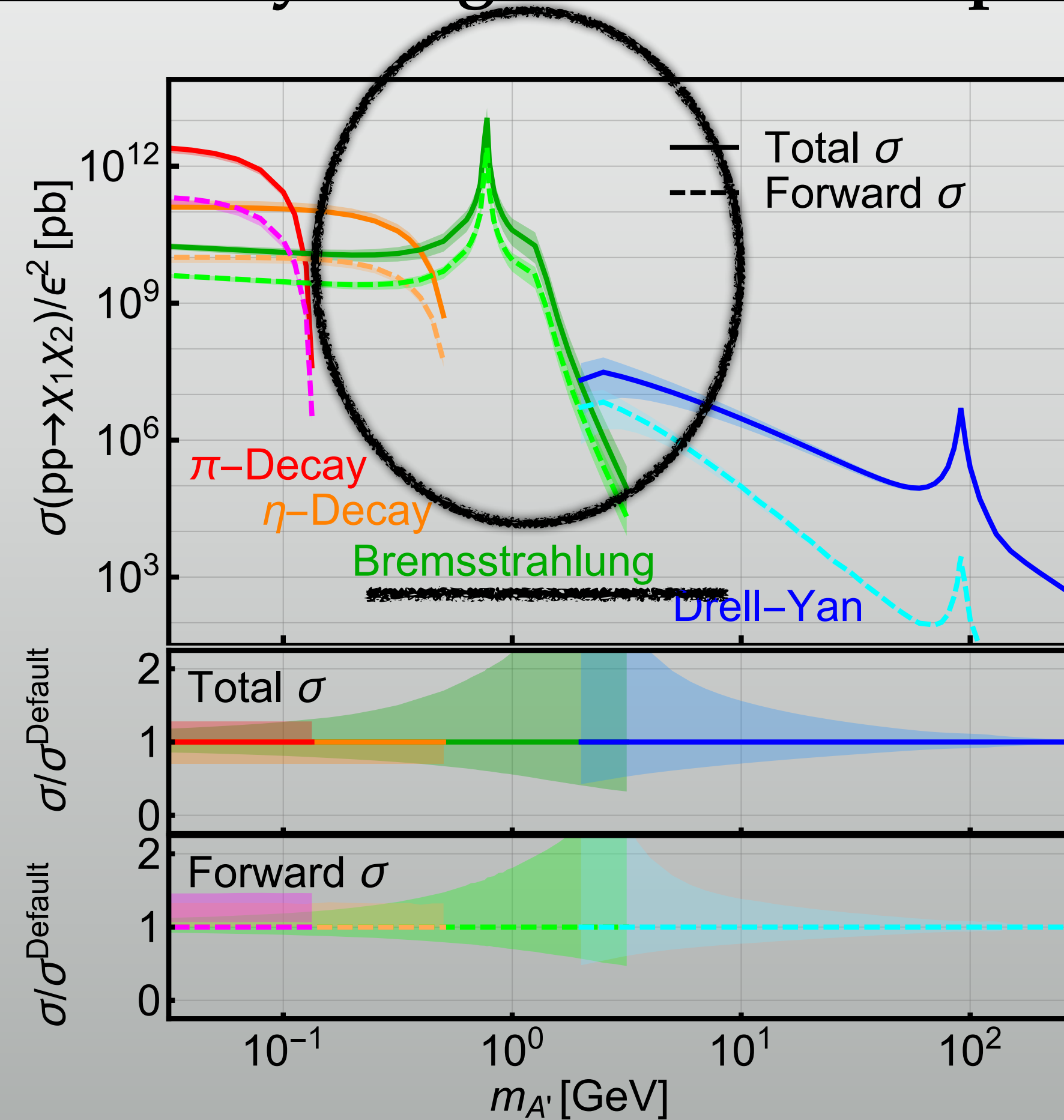


Approved!

Search Strategy for Dark Photons



Sensitivity Range of (future) experiments



**Proton
bremsstrahlung!**

Production

arXiv 1810.01879

A. Berlin, F. Kling

Production of BSM particles

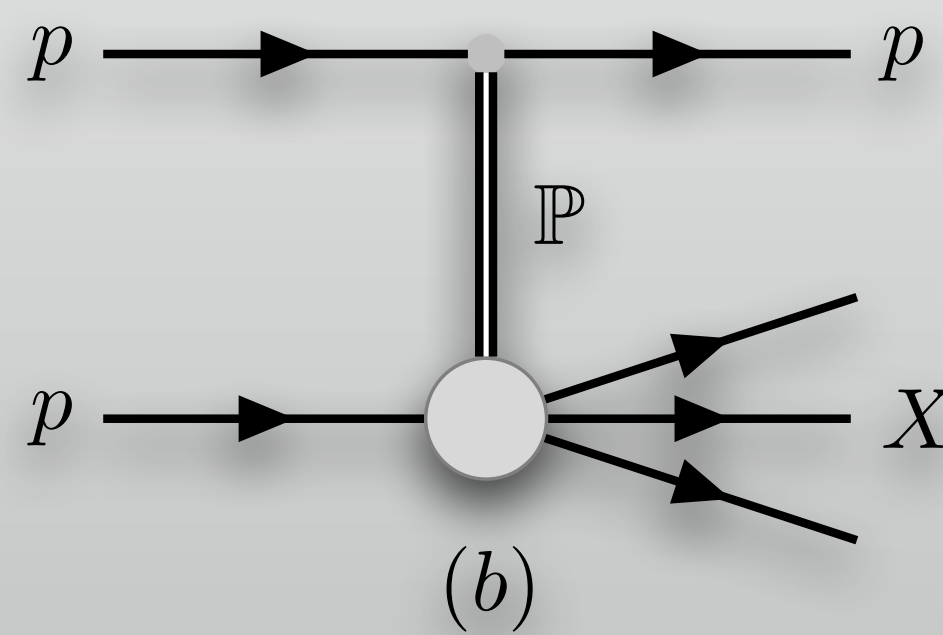
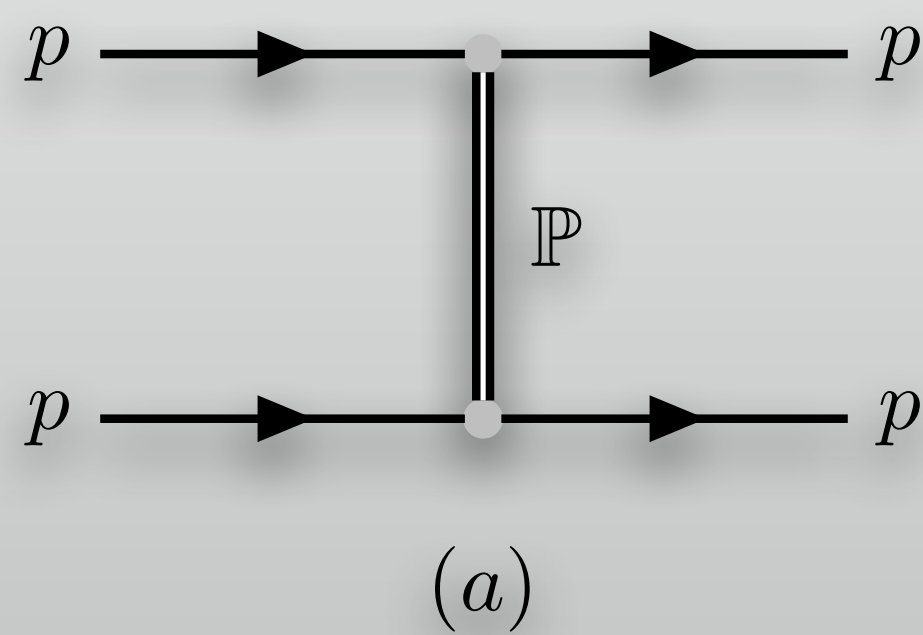
Proton
Bremsstrahlung

$$\mathcal{L}_{\text{eff}} \supset \varepsilon e A'_\mu \bar{p} \gamma^\mu p$$

Any process including protons

elastic

~ 25%

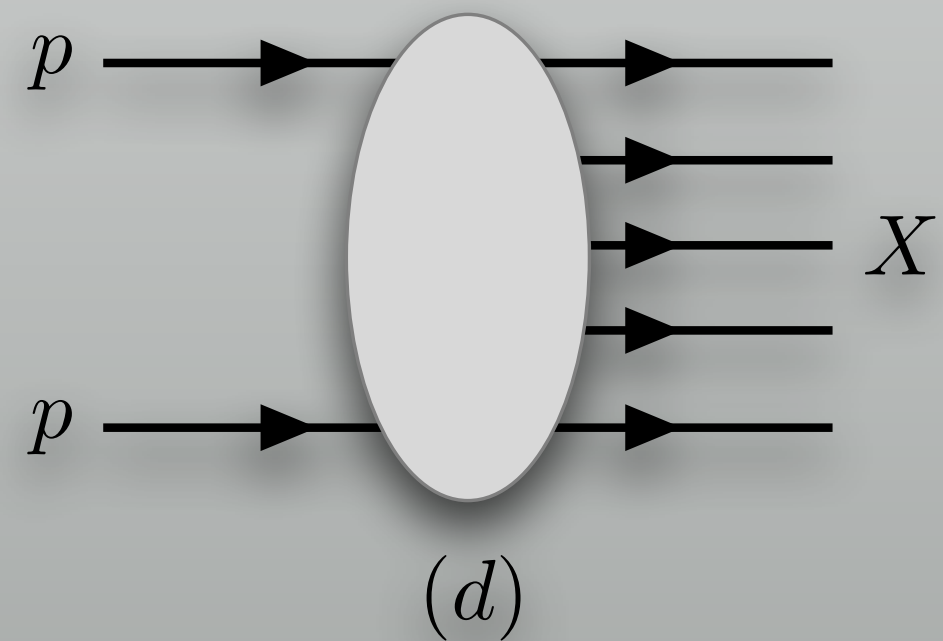
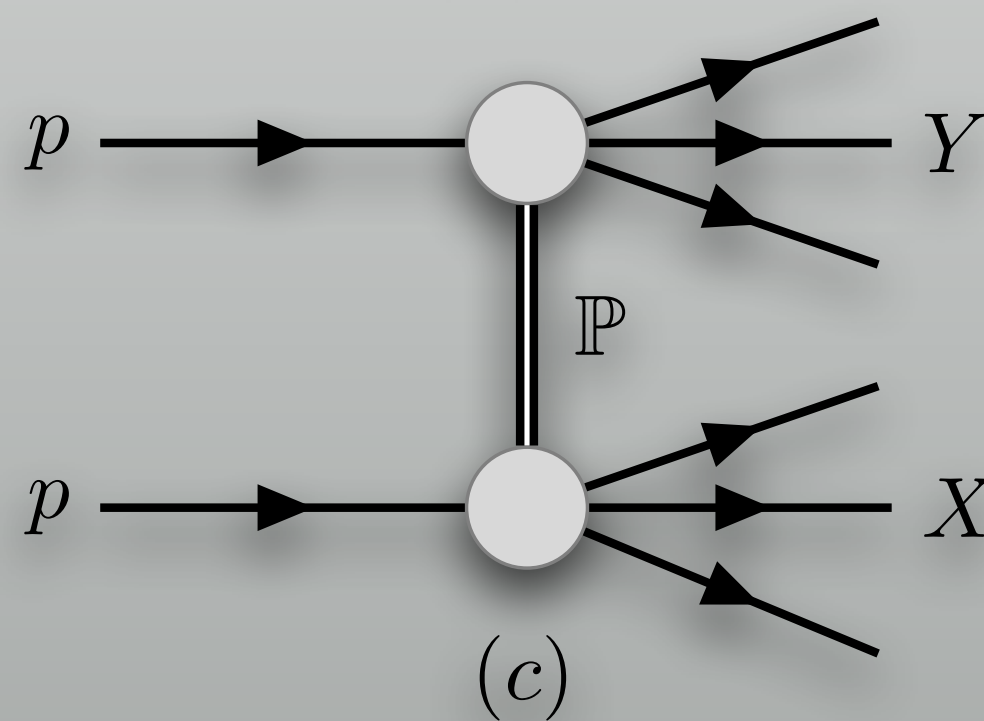


~ 11%

semi
diffractive

double
diffractive

~ 8%



~ 56%

inelastic

Production of BSM particles

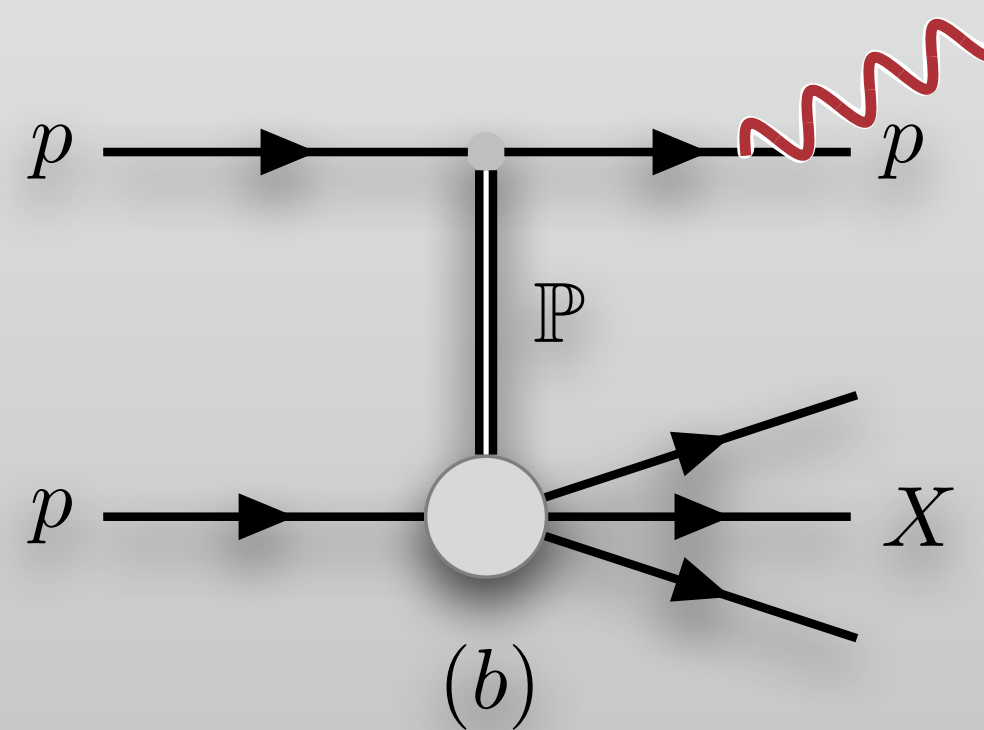
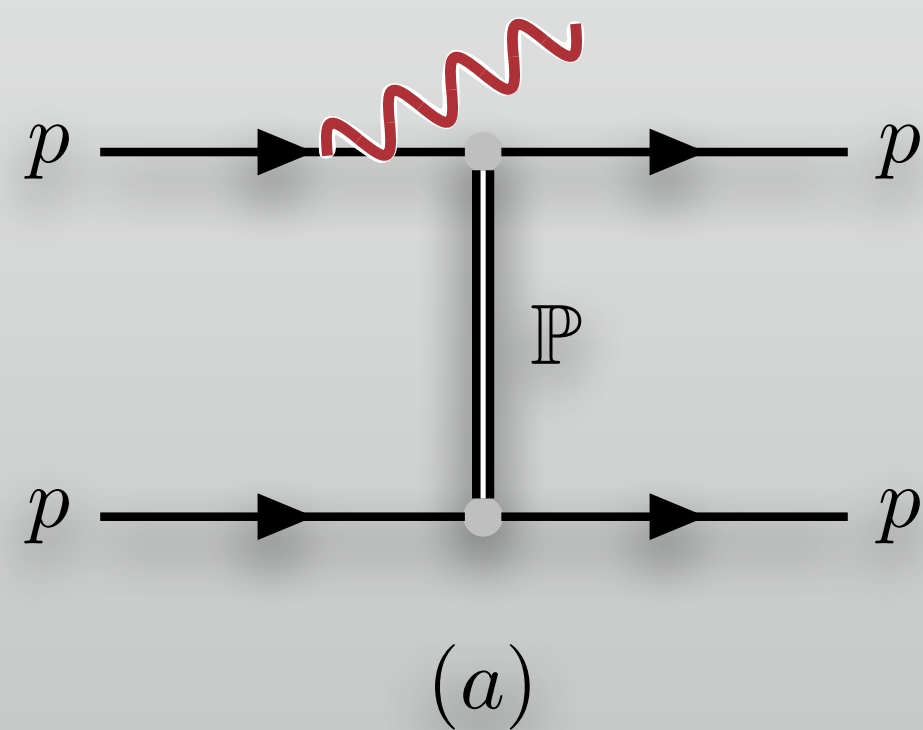
Proton
Bremsstrahlung

$$\mathcal{L}_{\text{eff}} \supset \varepsilon e A'_{\mu} \bar{p} \gamma^{\mu} p$$

Any process including protons

elastic

~ 25%

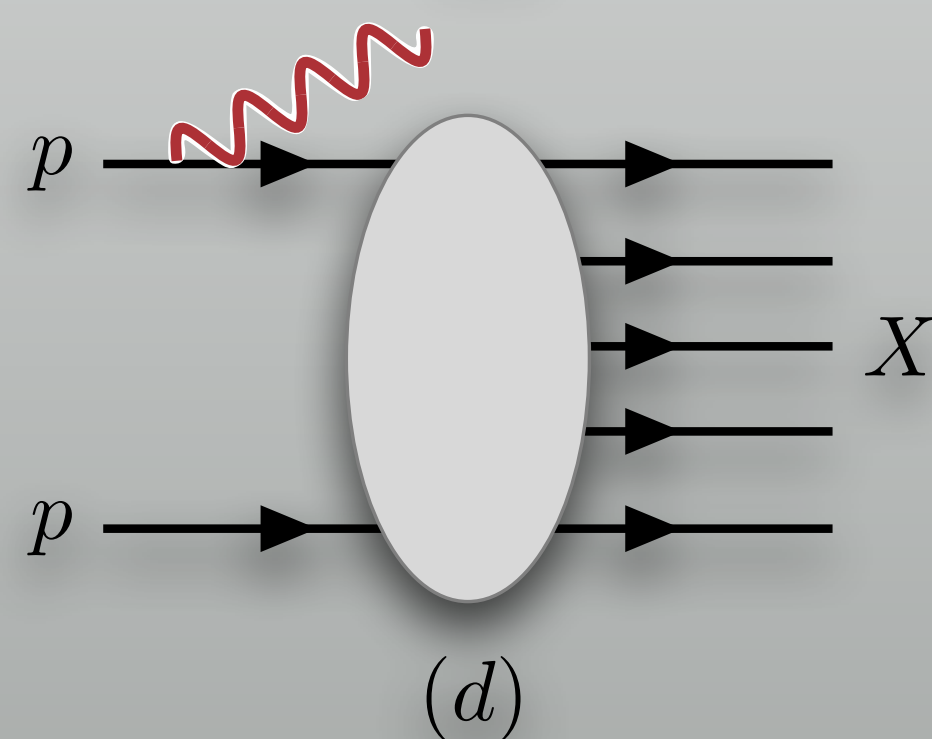
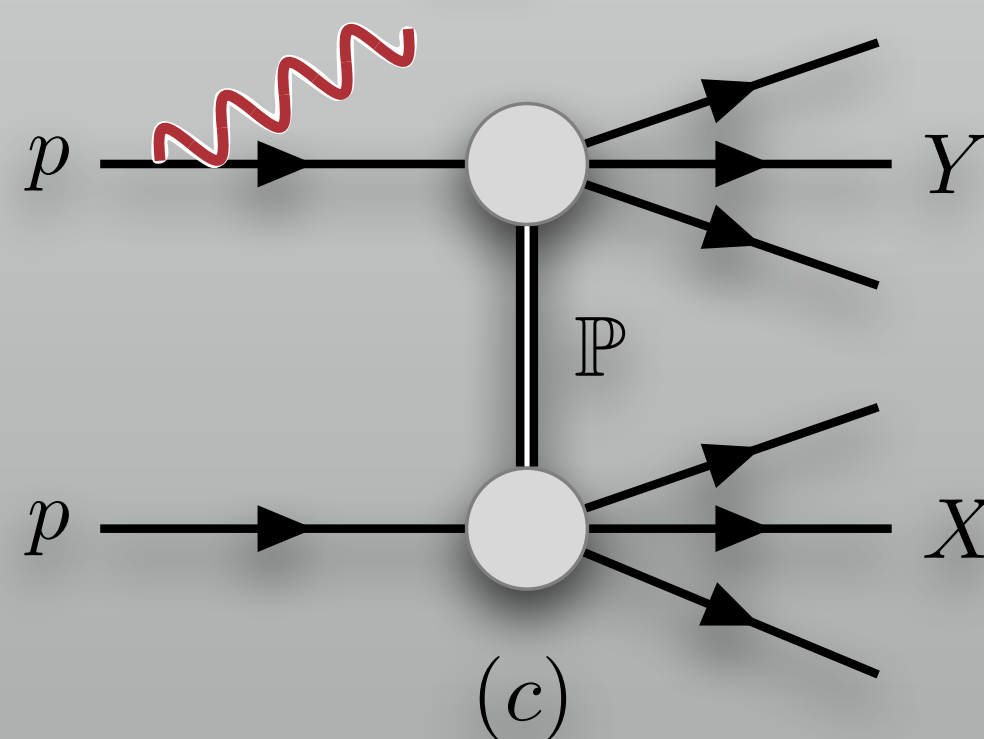


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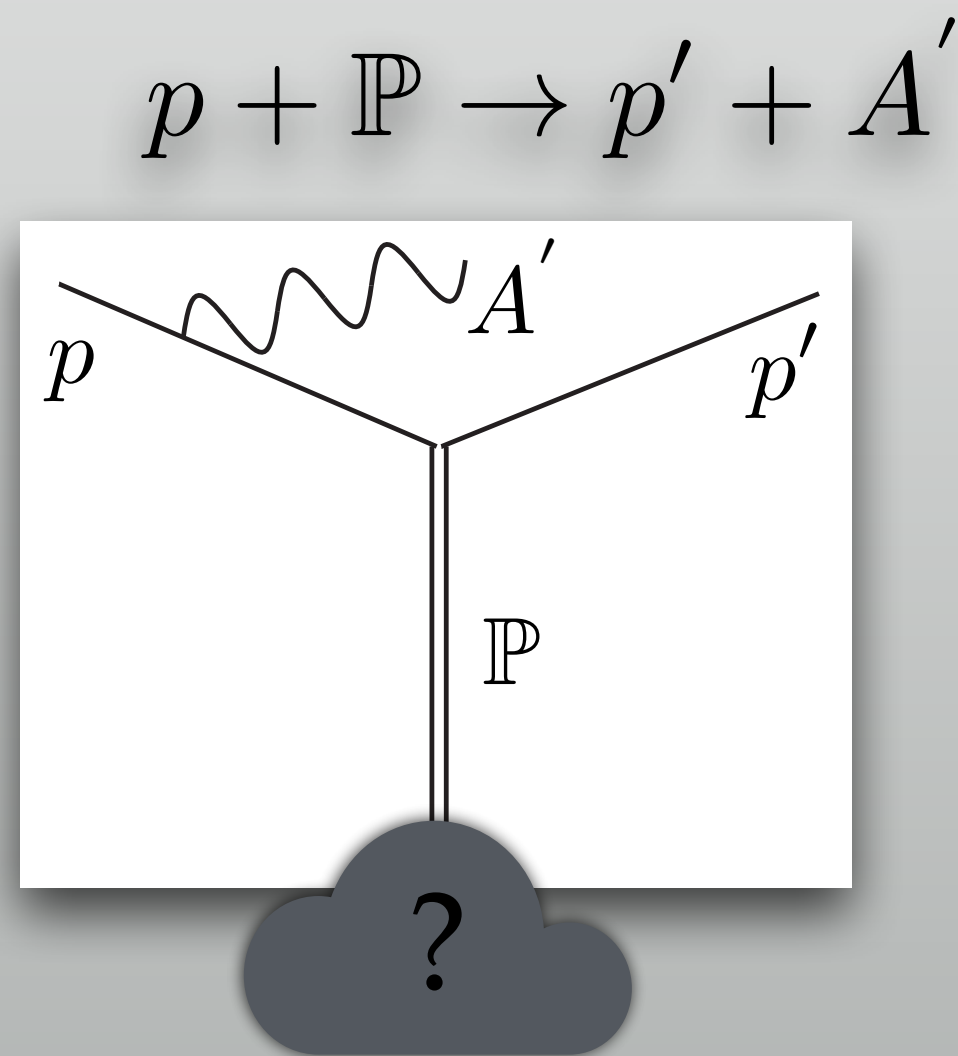
~ 56%

inelastic

Production of BSM particles

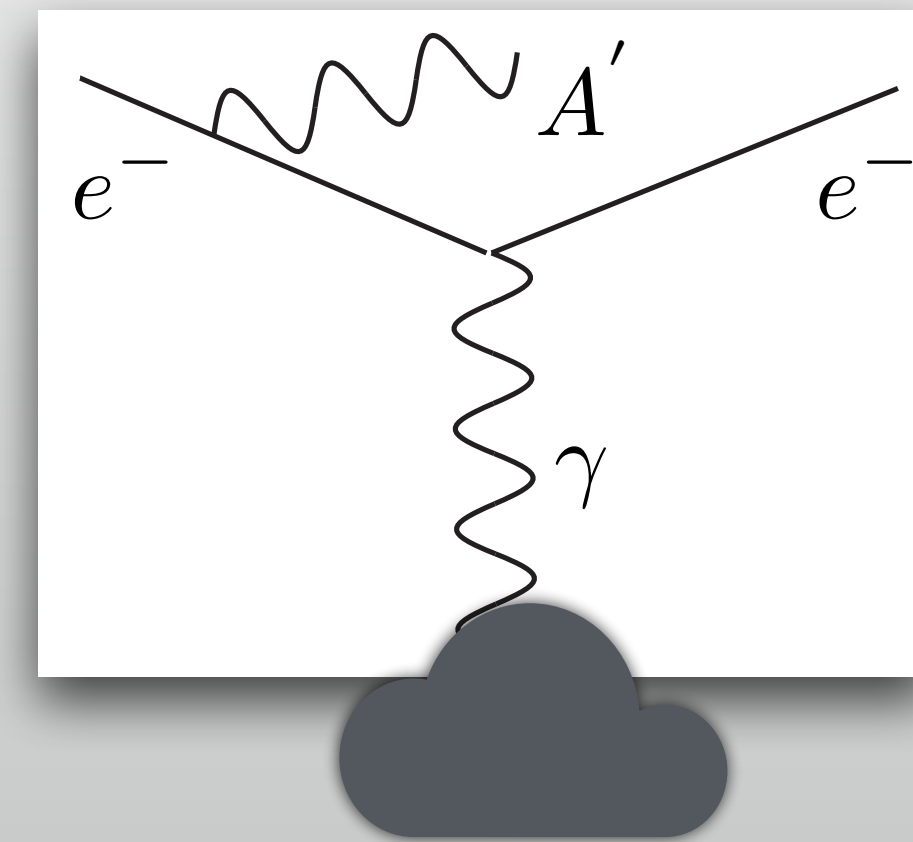
Proton
Bremsstrahlung

→ arXiv 1311.3870 Blümlein, Brunner



like FWW-approach for electron
beams

$$e^- + \gamma \rightarrow e^- + A'$$



But: what is the “pomeron cloud”?
Only valid for quasi-elastic
processes?

Production of BSM particles

Proton
Bremsstrahlung

→ [arXiv 1311.3870](#) Blümlein, Brunner

→ Recent developments:

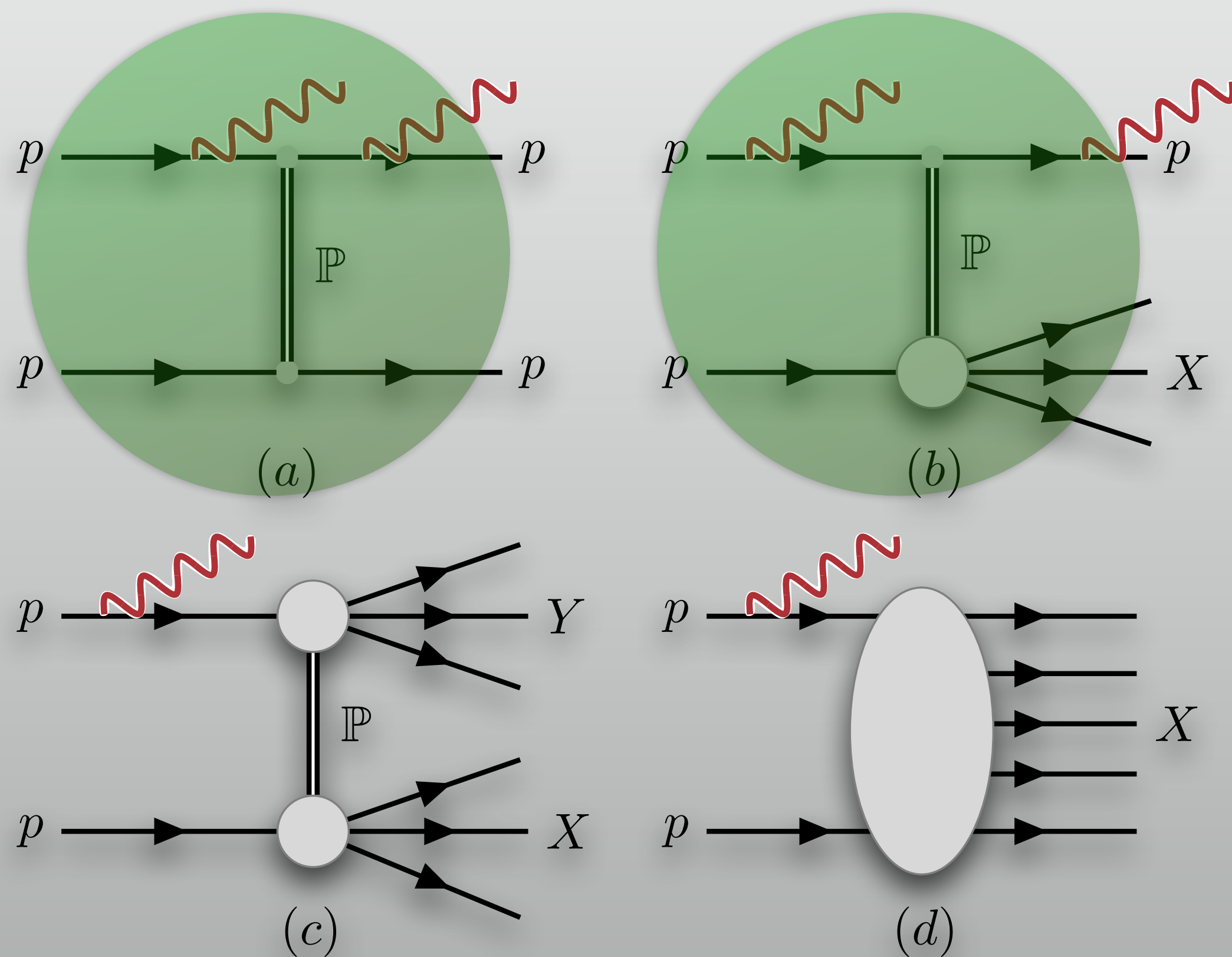
[arXiv 2108.05900](#) Saeid Foroughi-Abari, Adam Ritz

$$d\sigma^{pp \rightarrow A' + X}(s) \approx dP_{p \rightarrow p' A'} \times \sigma_{p' p}(s')$$

Production of BSM particles

Proton
Bremsstrahlung

arXiv 2108.05900 Saeid Foroughi-Abari, Adam Ritz



Cancellation between Initial- and Final-State radiation!

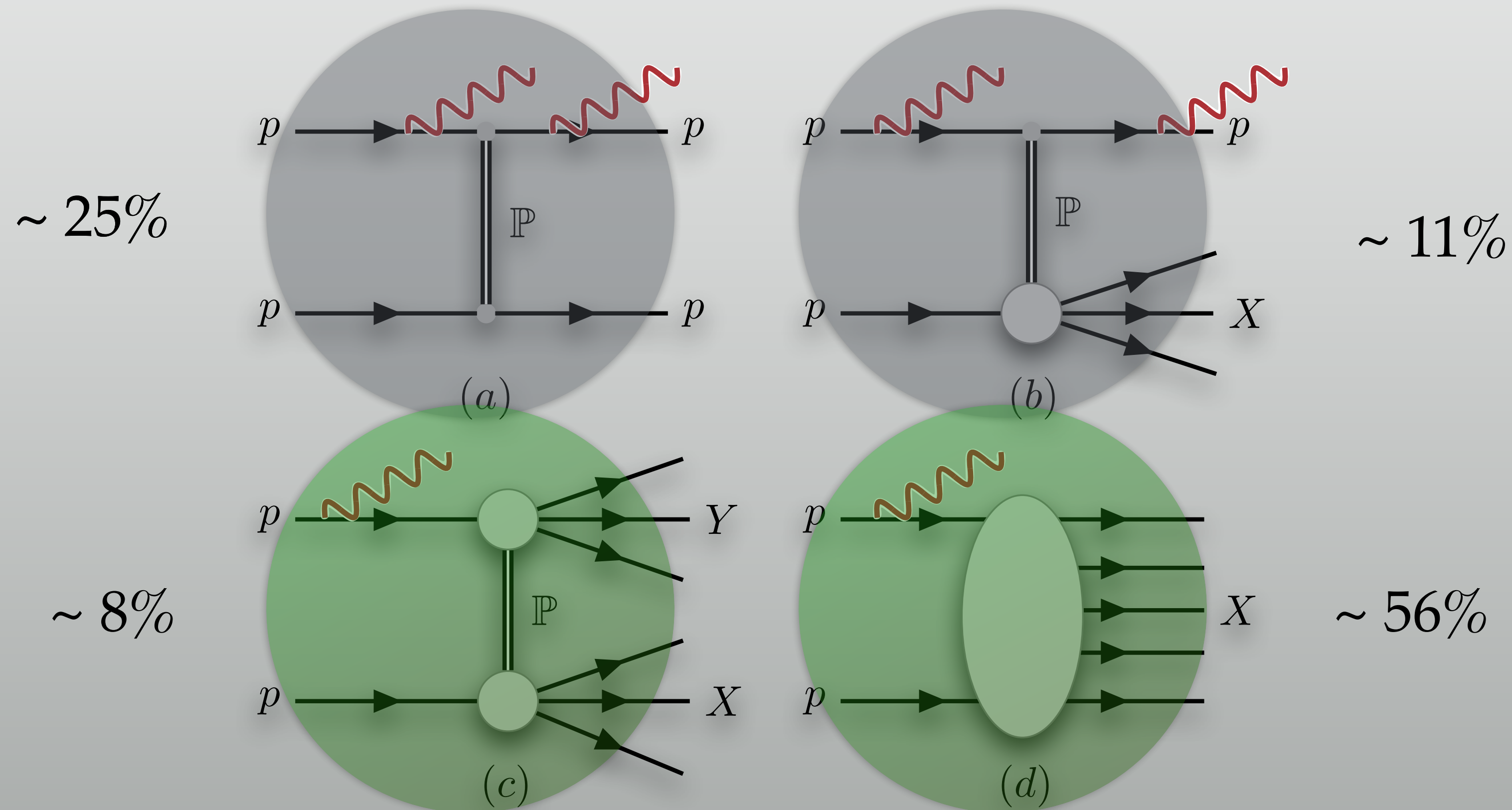
→ Elastic and single-diffractive events suppressed!

Production of BSM particles

Proton
Bremsstrahlung

→ Recent developments:

arXiv 2108.05900 Saeid Foroughi-Abari, Adam Ritz

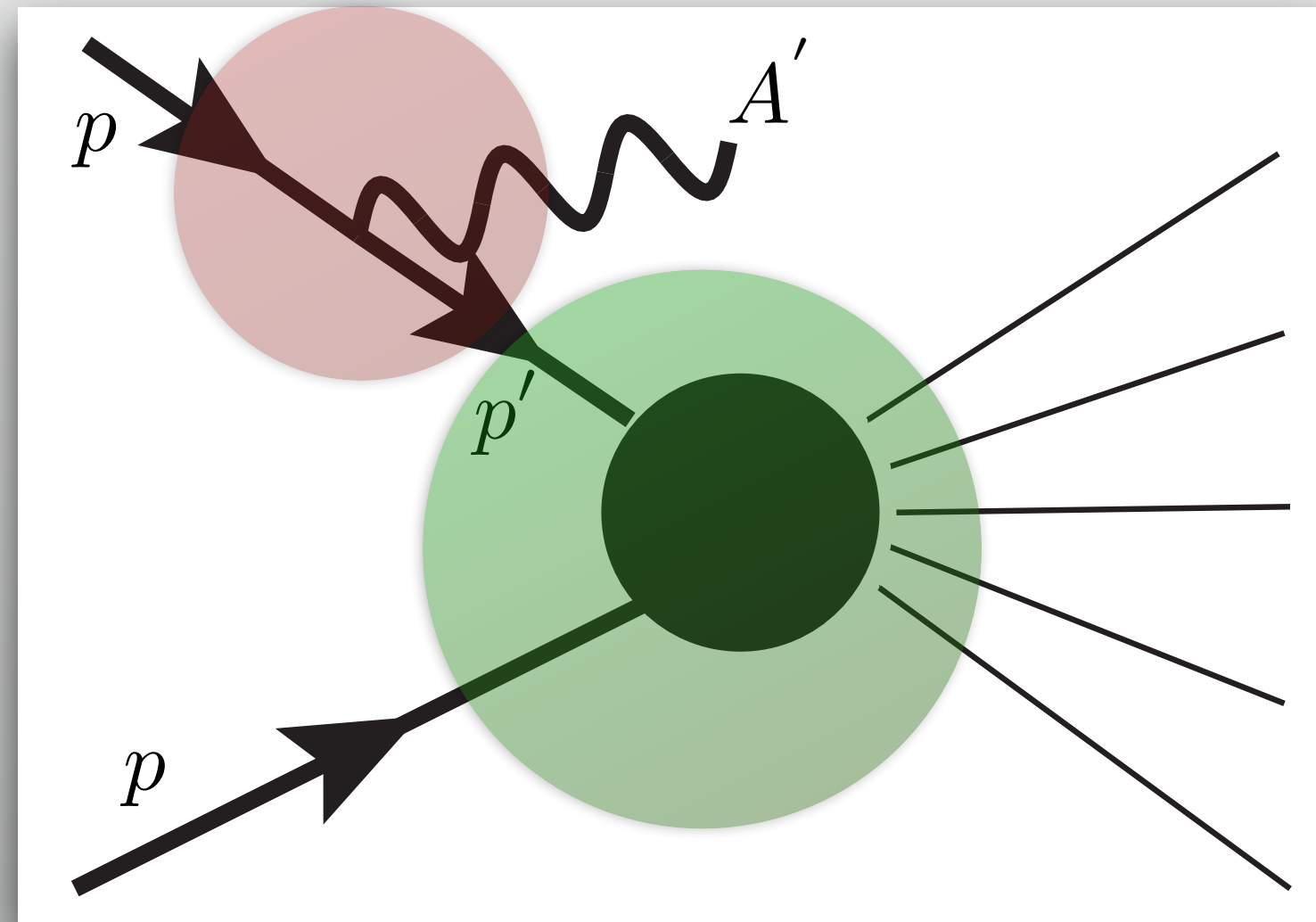


Production of BSM particles

Proton
Bremsstrahlung

→ Work in progress:

Saeid Foroughi-Abari, PR, Adam Ritz



$$d\sigma^{pp \rightarrow A' + X}(s) \approx dP_{p \rightarrow p' A'} \times \sigma_{p'p}(s')$$

X

Improved description of splitting vertex:

- Splitting function
- Form factors

Splitting functions

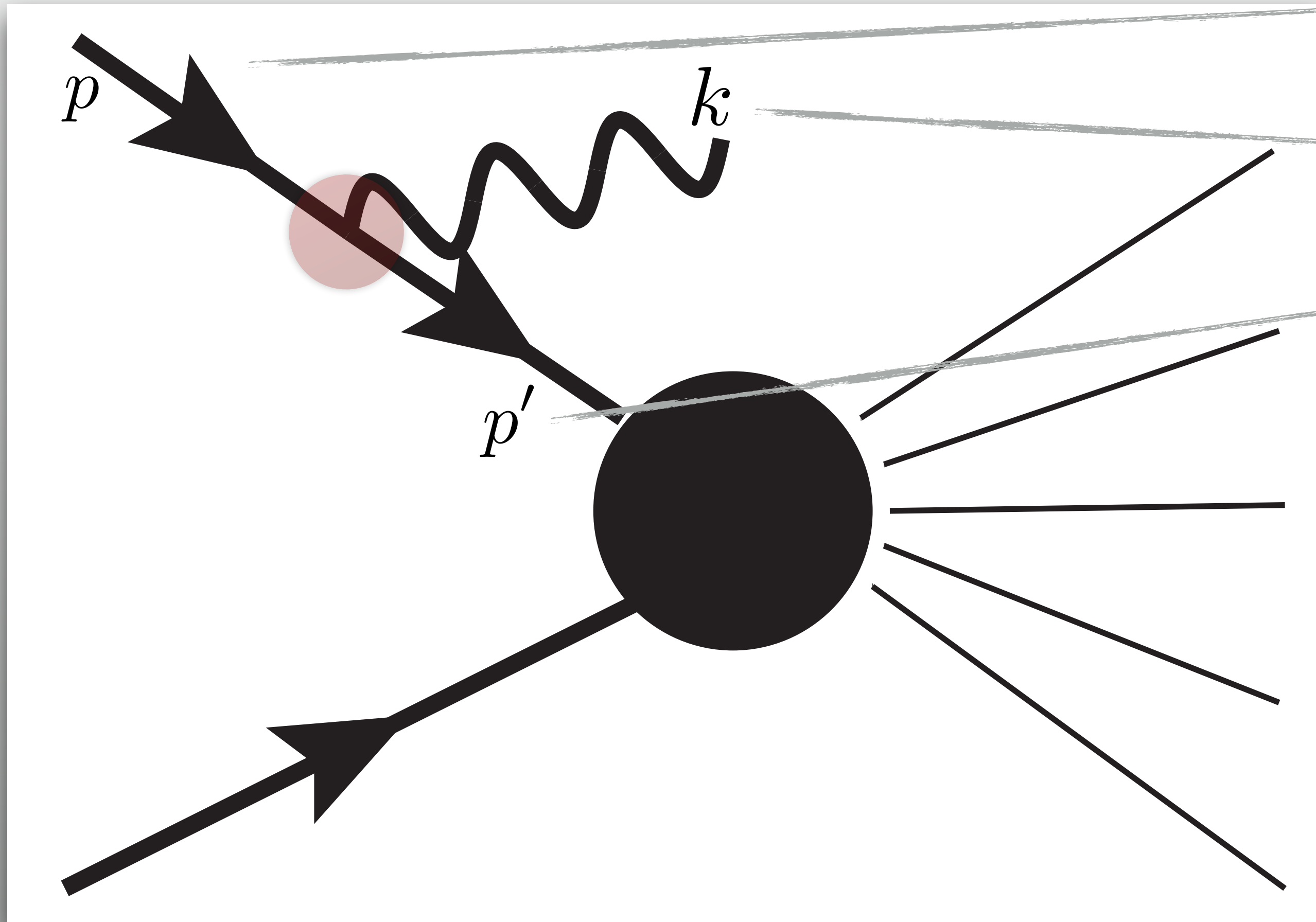
Infinite momentum frame:

$$p^\mu = \left(p_p + \frac{m_p^2}{2p_p}, \mathbf{0}, p_p \right),$$

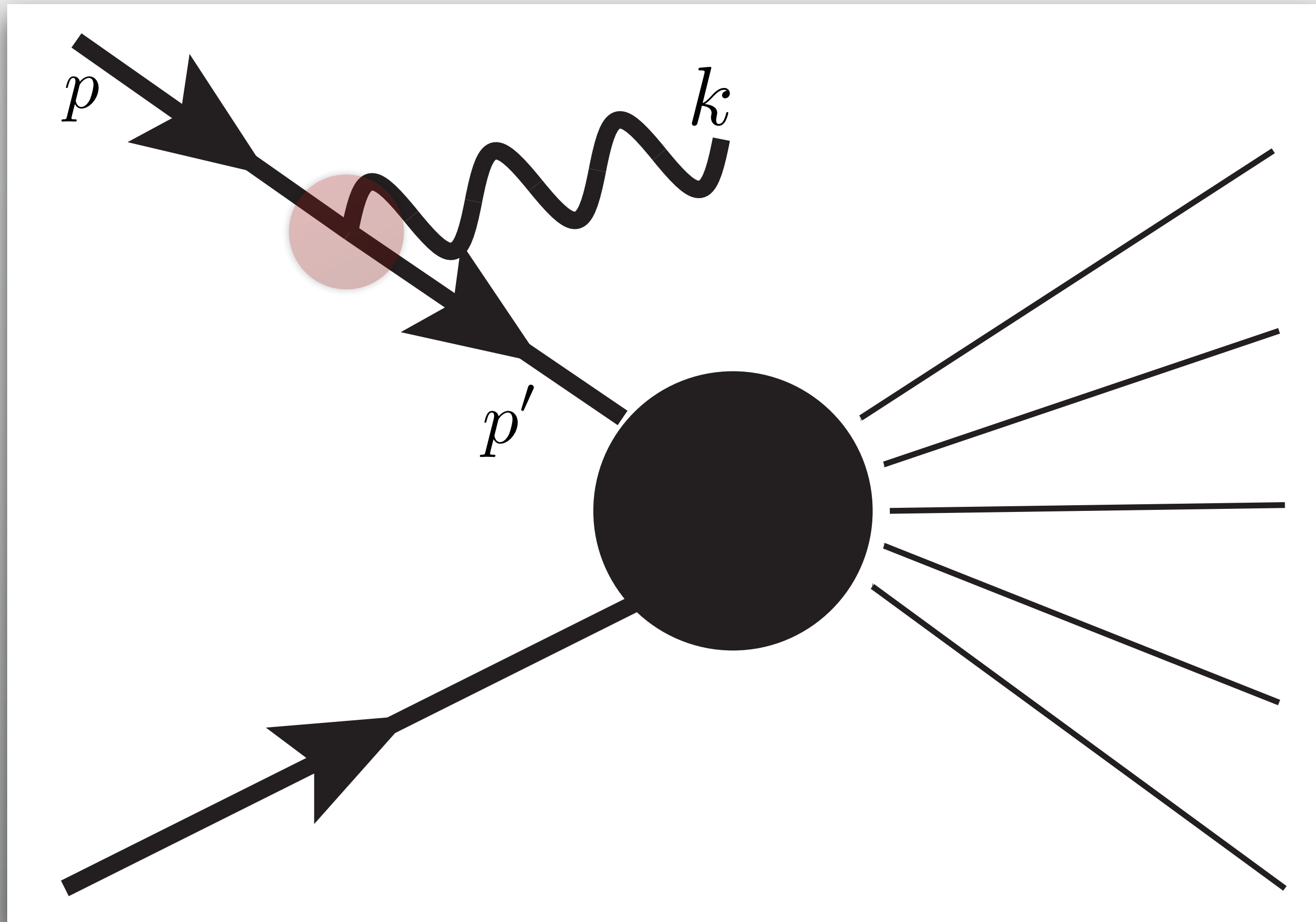
$$k^\mu = \left(zp_p + \frac{p_T^2 + m_{A'}^2}{2zp_p}, \mathbf{p}_T, zp_p \right)$$

$$p'^\mu = \left((1-z)p_p + \frac{p_T^2 + m_p^2}{2p_p(1-z)}, -\mathbf{p}_T, (1-z)p_p \right)$$

$$d\mathcal{P}_{p \rightarrow p' A'} \equiv w(z, p_T^2) dz dp_T^2$$



Splitting functions



$$d\mathcal{P}_{p \rightarrow p' A'} \equiv w(z, p_T^2) dz dp_T^2$$

Straightforward calculation:

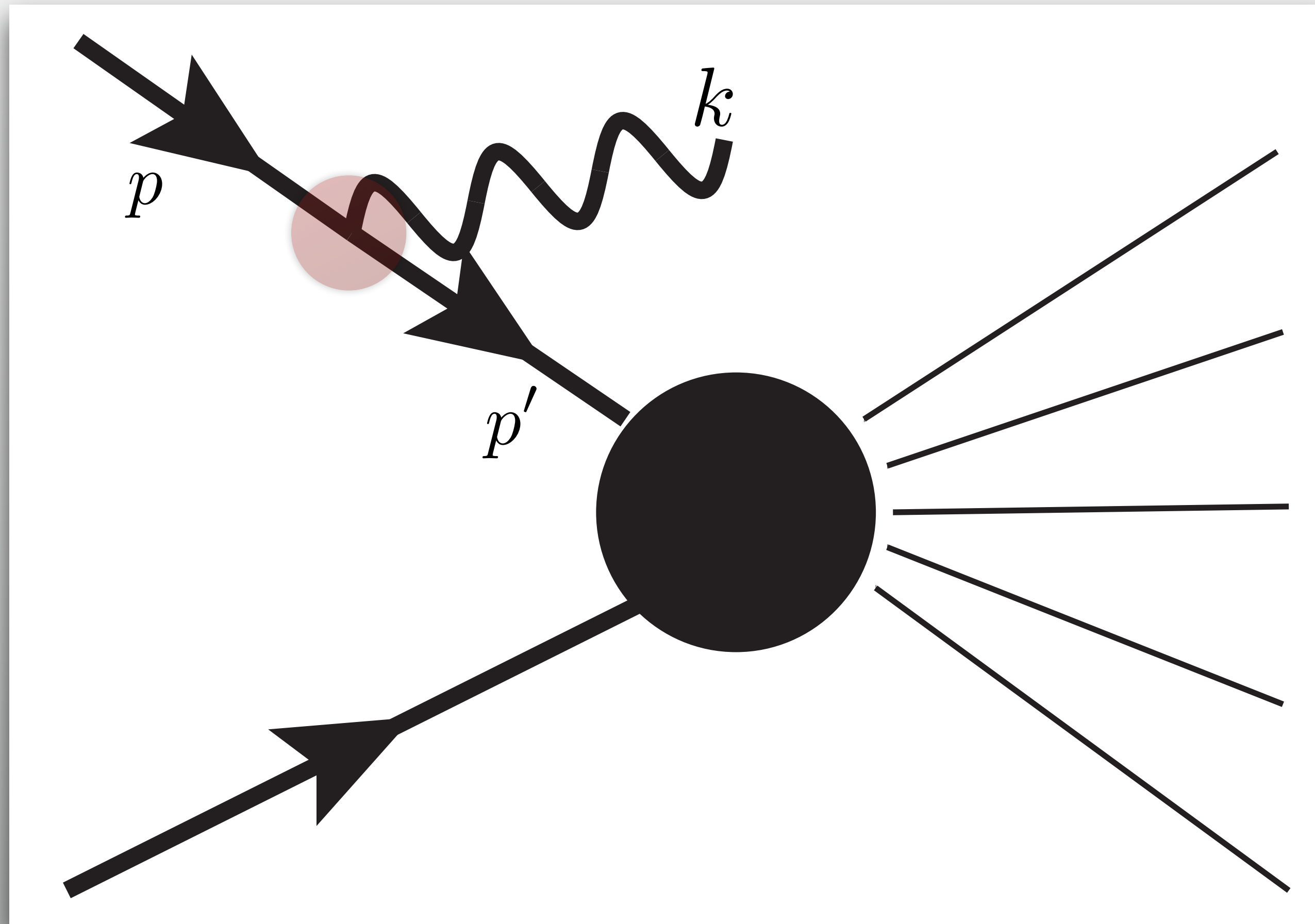
$$w_V(z, p_T^2) = \frac{\alpha\epsilon^2}{2\pi H} \left[z - z(1-z) \left(\frac{2m_p^2 + m_{A'}^2}{H} \right) + \frac{H}{2zm_{A'}^2} \right]$$

with $H(z, p_T^2) \equiv p_T^2 + z^2 m_p^2 + (1-z)m_{A'}^2$

divergent for $m_{A'} \rightarrow 0$

→ Dawson approach

Splitting functions



$$d\mathcal{P}_{p \rightarrow p' A'} \equiv w(z, p_T^2) dz dp_T^2$$

Dawson approach:

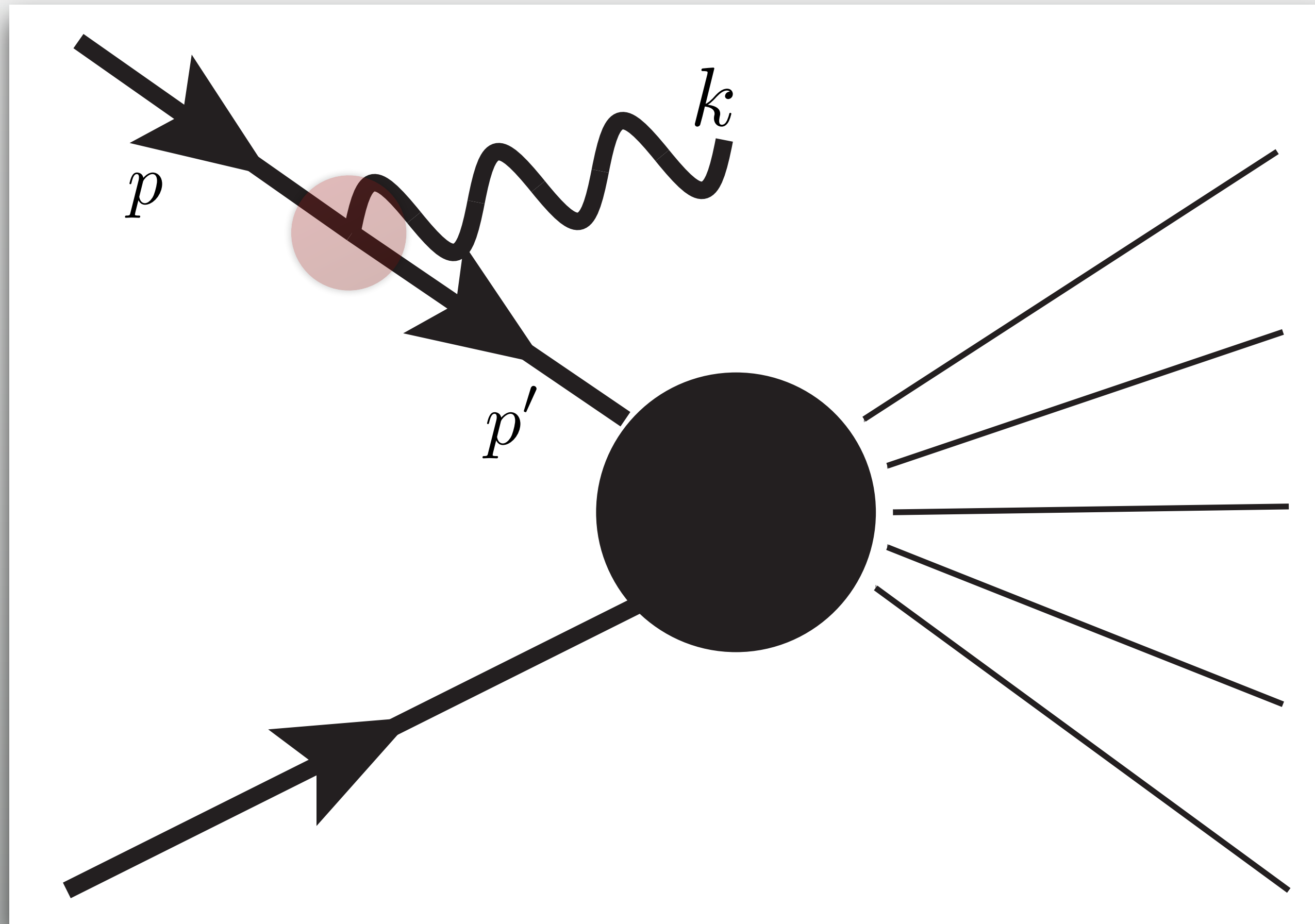
$$\bar{u}^r(p') \not{k} u^s(p) = 0$$

Use

$$\varepsilon_0^\mu(k) \equiv \varepsilon_L^\mu(k) - \frac{k^\mu}{m_{A'}} = \frac{m_{A'}}{(k_0 + |\vec{k}|)} (-1, \vec{k}/|\vec{k}|)$$

$$\rightarrow w_V^{\text{eff}}(z, p_T^2) = \frac{\alpha \epsilon^2}{2\pi H} \left[\frac{1 + (1-z)^2}{z} - z(1-z) \left(\frac{2m_p^2 + m_{A'}^2}{H} \right) \right]$$

Splitting functions



$$d\mathcal{P}_{p \rightarrow p' A'} \equiv w(z, p_T^2) dz dp_T^2$$

Dawson approach:

$$\bar{u}^r(p') \not{k} u^s(p) = 0$$

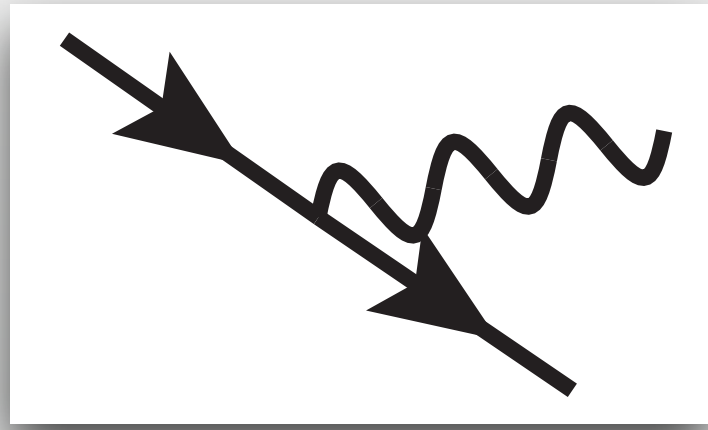
Use

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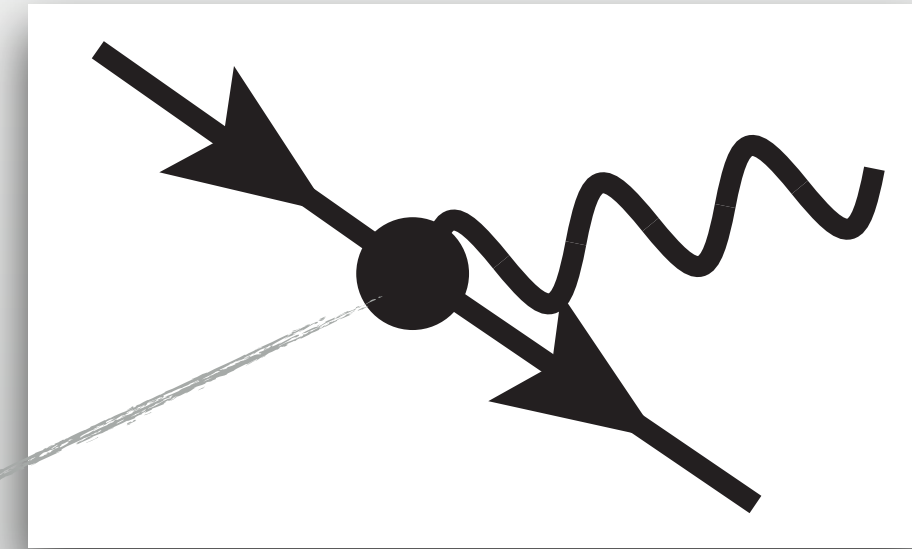
$$m_{A'}, m_p \rightarrow 0$$

$$\rightarrow w_V^{\text{eff}}(z, p_T^2) = \frac{\alpha \epsilon^2}{2\pi H} \left[\frac{1 + (1-z)^2}{z} - z(1-z) \left(\frac{2m_p^2 + m_{A'}^2}{H} \right) \right]$$

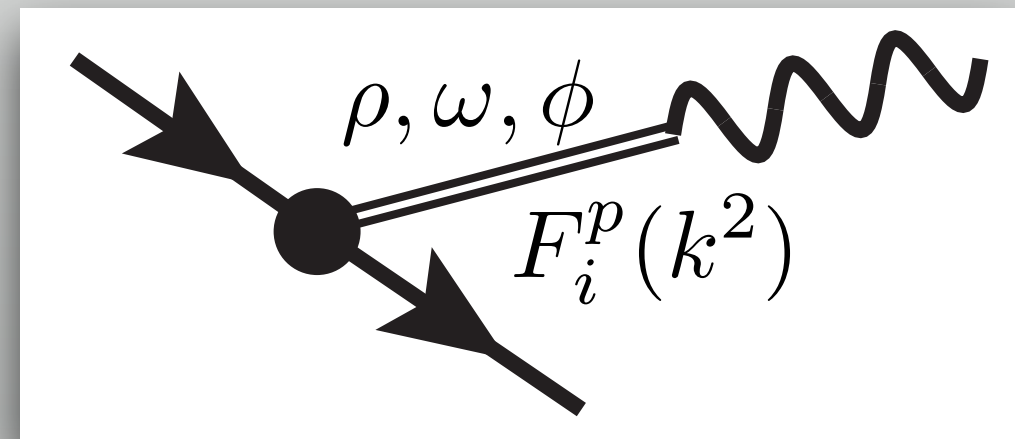
Form factors



$$\mathcal{L}_{\text{eff}} \supset \epsilon e A'_\mu \bar{p} \gamma^\mu p$$



$$\langle N | J_{\text{em}}^\mu | N \rangle = \bar{u}(p') \left[\gamma^\mu F_1^N(k^2) + \frac{i\sigma^{\mu\nu}(p'-p)_\nu}{2m_N} F_2^N(k^2) \right] u(p)$$



→ Hadronic resonant contributions

1601.06190 VMD, Unitarity & Analytic Model
2109.12961 Dispersion relation

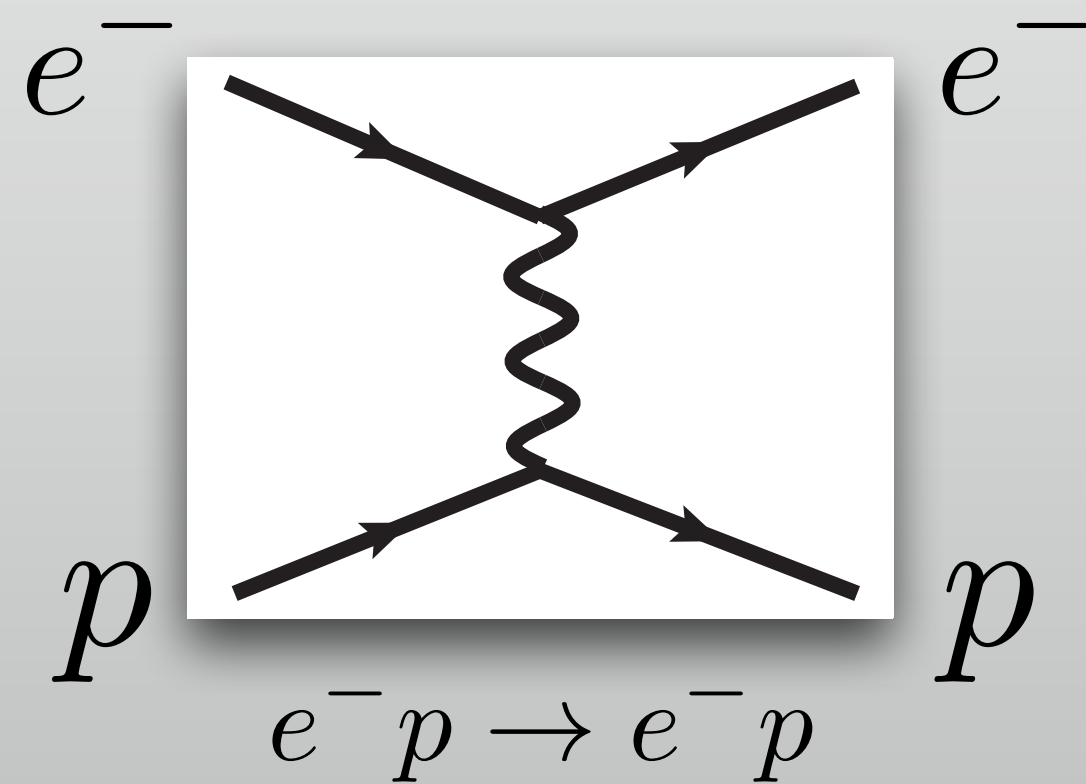
Photon interactions with protons

Sachs Form factors:

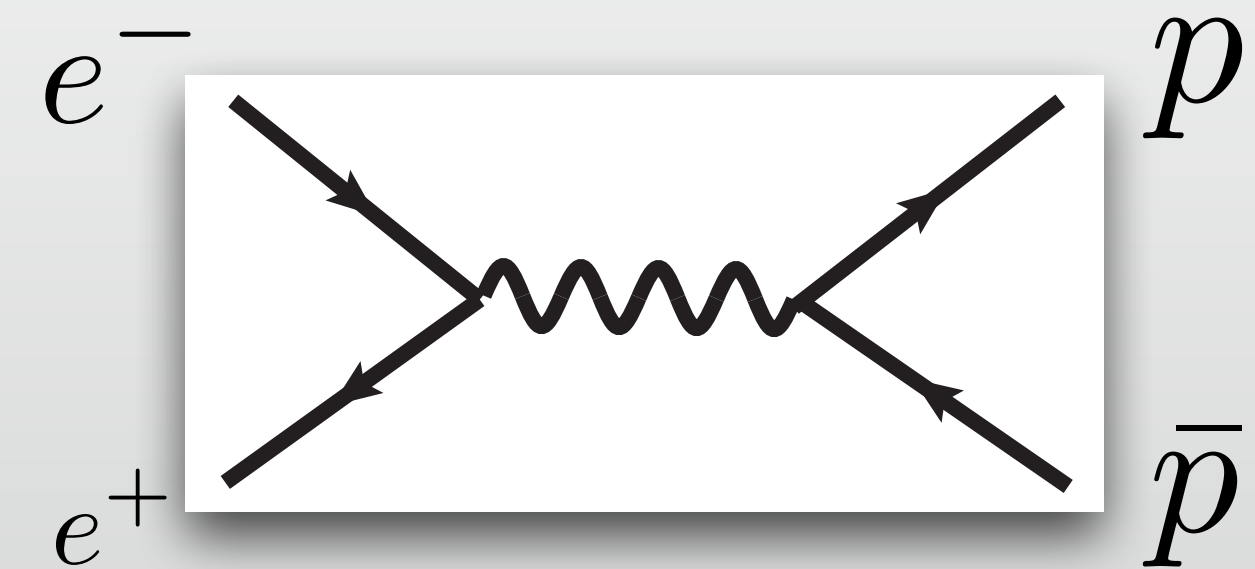
$$G_E(t) = F_1(t) - \tau F_2(t),$$

$$G_M(t) = F_1(t) + F_2(t),$$

Scattering



Production



$$e^+ e^- \rightarrow p \bar{p}$$

$$\sigma_{e^+e^- \rightarrow p\bar{p}}(q^2) = \frac{4\pi\alpha^2\beta}{3q^2} C(q^2) \left[|G_M(q^2)|^2 + \frac{2m_p^2}{q^2} |G_E(q^2)|^2 \right]$$

$$\equiv \frac{4\pi\alpha^2\beta}{3q^2} C(q^2) \left(1 + \frac{2m_p^2}{q^2} \right) |G_{\text{eff}}^p(q^2)|^2.$$

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} \frac{1}{\epsilon(1+\tau)} \underbrace{[\tau G_M^2(Q^2) + \epsilon G_E^2(Q^2)]}_{=\sigma_R},$$

State of the Art Form factors

2109.12961

New insights into the nucleon's electromagnetic structure

Yong-Hui Lin,¹ Hans-Werner Hammer,^{2,3} and Ulf-G. Meißner^{1,4,5}

¹*Helmholtz Institut für Strahlen- und Kernphysik and Bethe Center for Theoretical Physics, Universität Bonn, D-53115 Bonn, Germany*

²*Technische Universität Darmstadt, Department of Physics, 64289 Darmstadt, Germany*

³*ExtreMe Matter Institute EMMI and Helmholtz Forschungsakademie Hessen für FAIR (HFHF), GSI Helmholtzzentrum für Schwerionenforschung GmbH, 64291 Darmstadt, Germany*

⁴*Institute for Advanced Simulation and Institut für Kernphysik, Forschungszentrum Jülich, D-52425 Jülich, Germany*

⁵*Tbilisi State University, 0186 Tbilisi, Georgia*

(Dated: September 28, 2021)

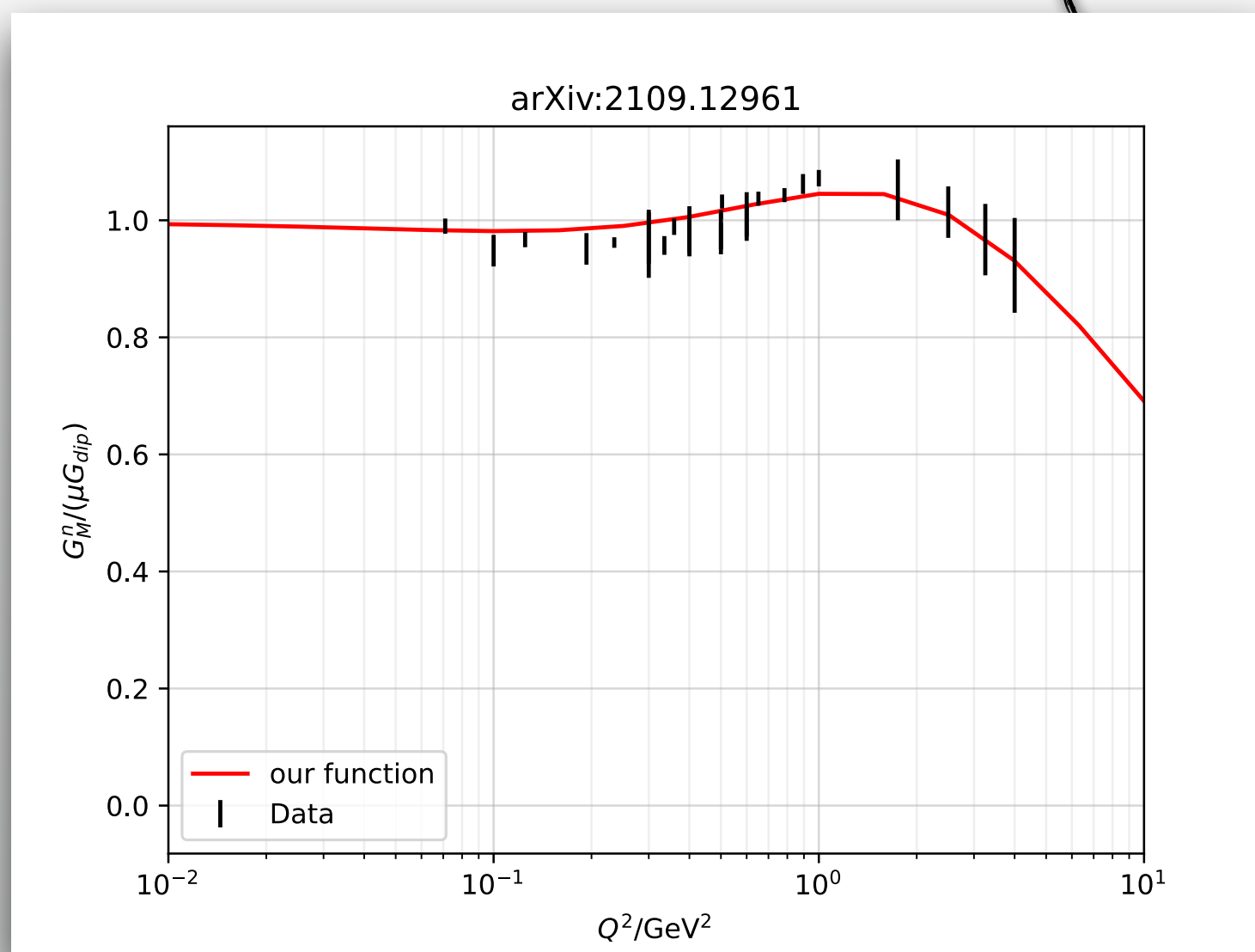
We present a combined analysis of the electromagnetic form factors of the nucleon in the space- and timelike regions using dispersion theory. Our framework provides a consistent description of the experimental data over the full range of momentum transfer, in line with the strictures from analyticity and unitarity. The statistical uncertainties of the extracted form factors are estimated using the bootstrap method, while systematic errors are determined from variations of the spectral functions. We also perform a high-precision extraction of the nucleon radii and find good agreement with previous analyses of spacelike data alone. For the proton charge radius, we find $r_E^p = 0.840^{+0.003+0.002}_{-0.002-0.002}$ fm, where the first error is statistical and the second one is systematic. The Zemach radius and third moment are in agreement with Lamb shift measurements and hyperfine splittings. The combined data set of space- and timelike data disfavors a zero crossing of $\mu_p G_E^p / G_M^p$ in the spacelike region. Finally, we discuss the status and perspectives of modulus and phase of the form factors in the timelike region in the context of future experiments as well as the onset of perturbative QCD.

- Continuum contributions $\pi\pi, K\bar{K}, \rho\pi$
- Vector meson poles
- Discussion of two-photon effect
- Fit to space- and timelike data
- ...

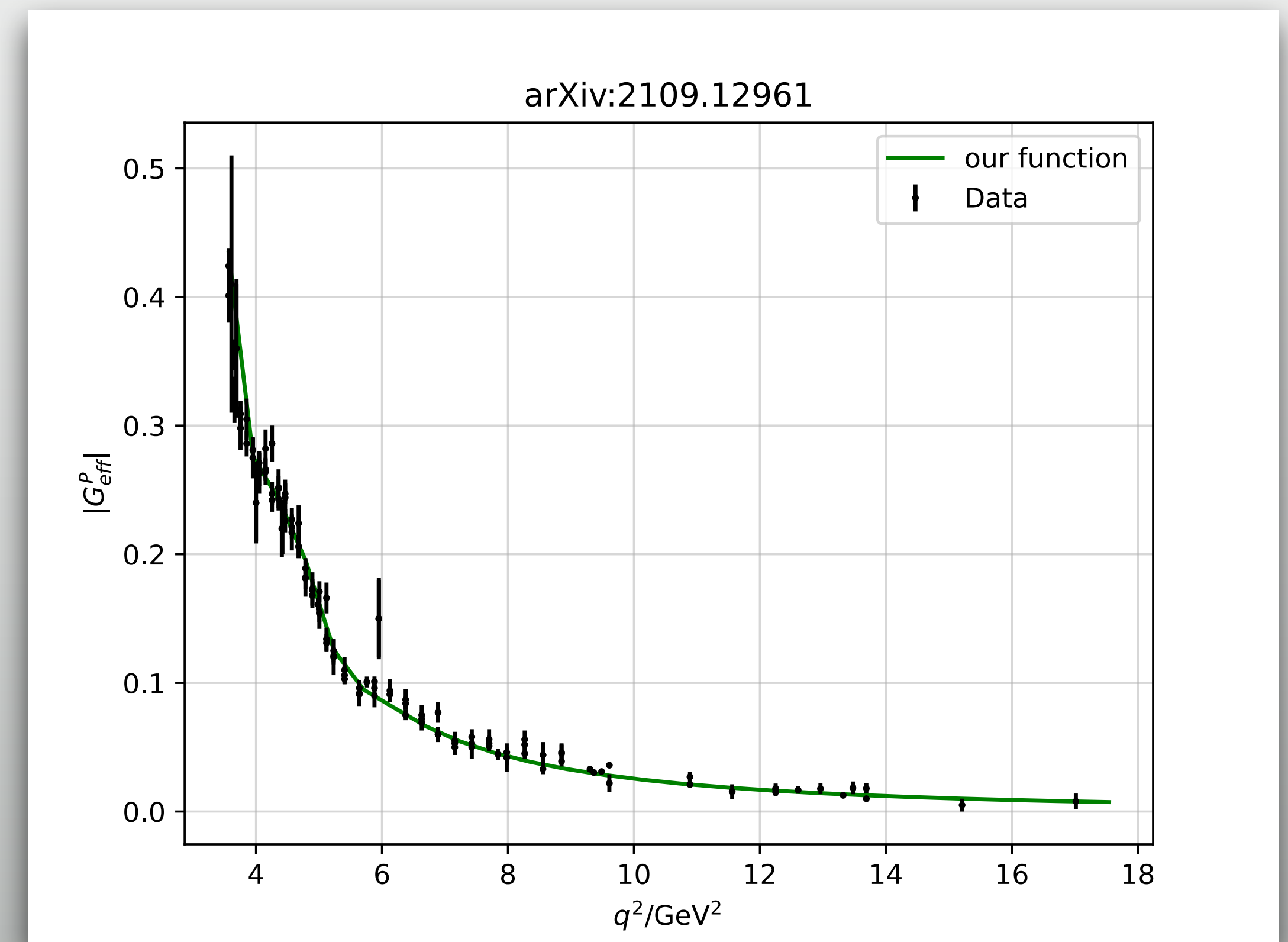
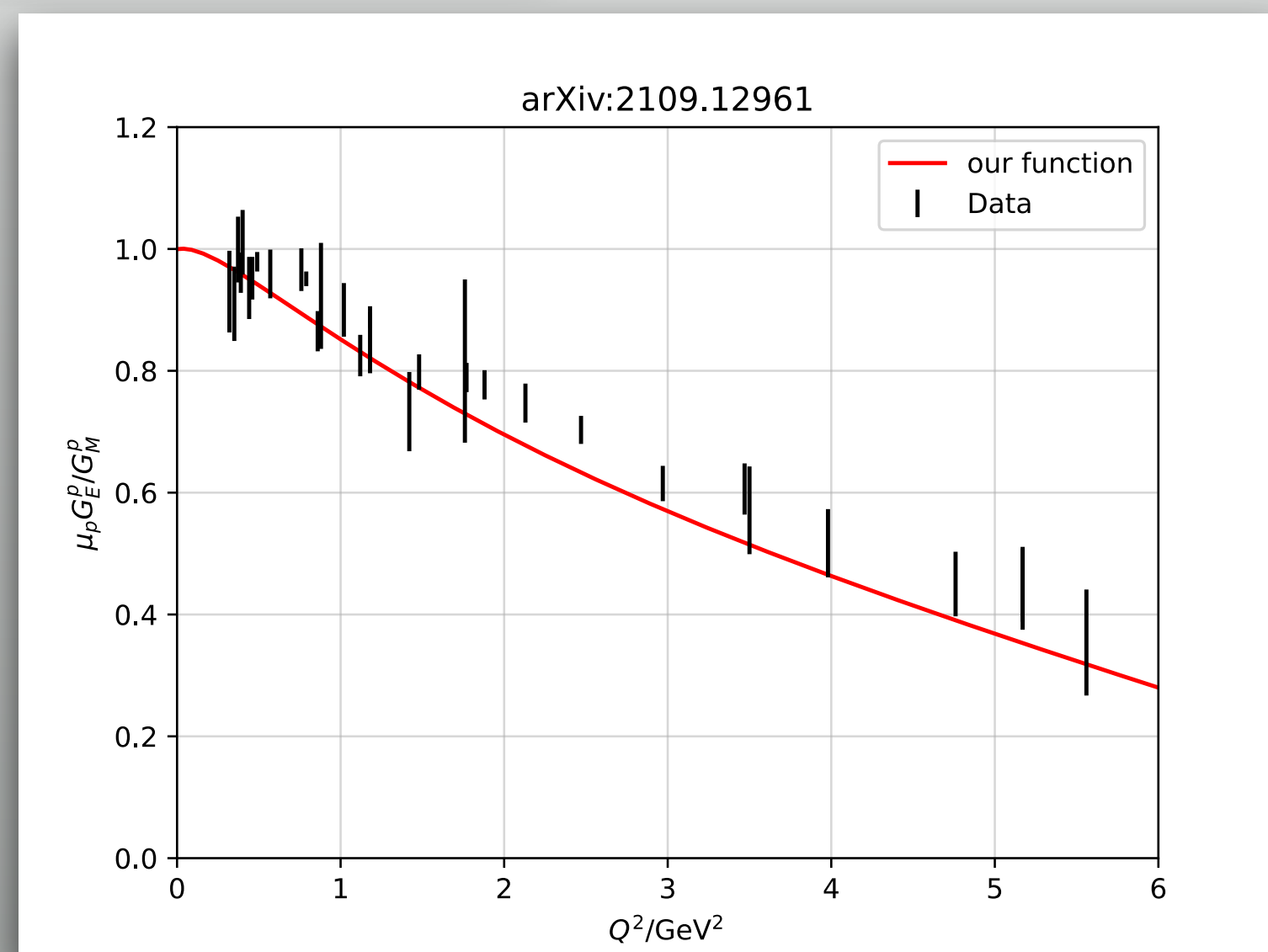
State of the Art Form factors

2109.12961

Production

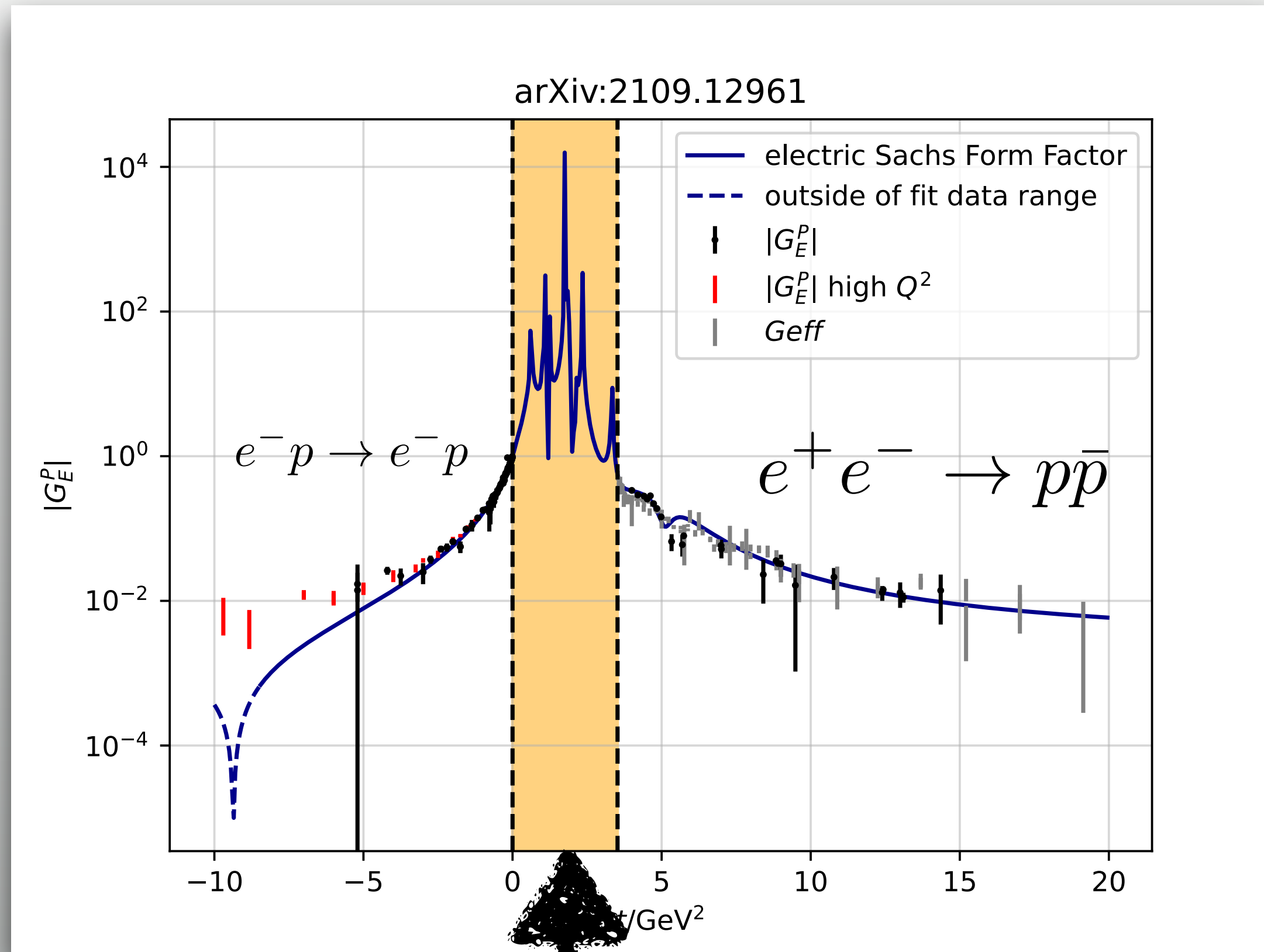


Scattering

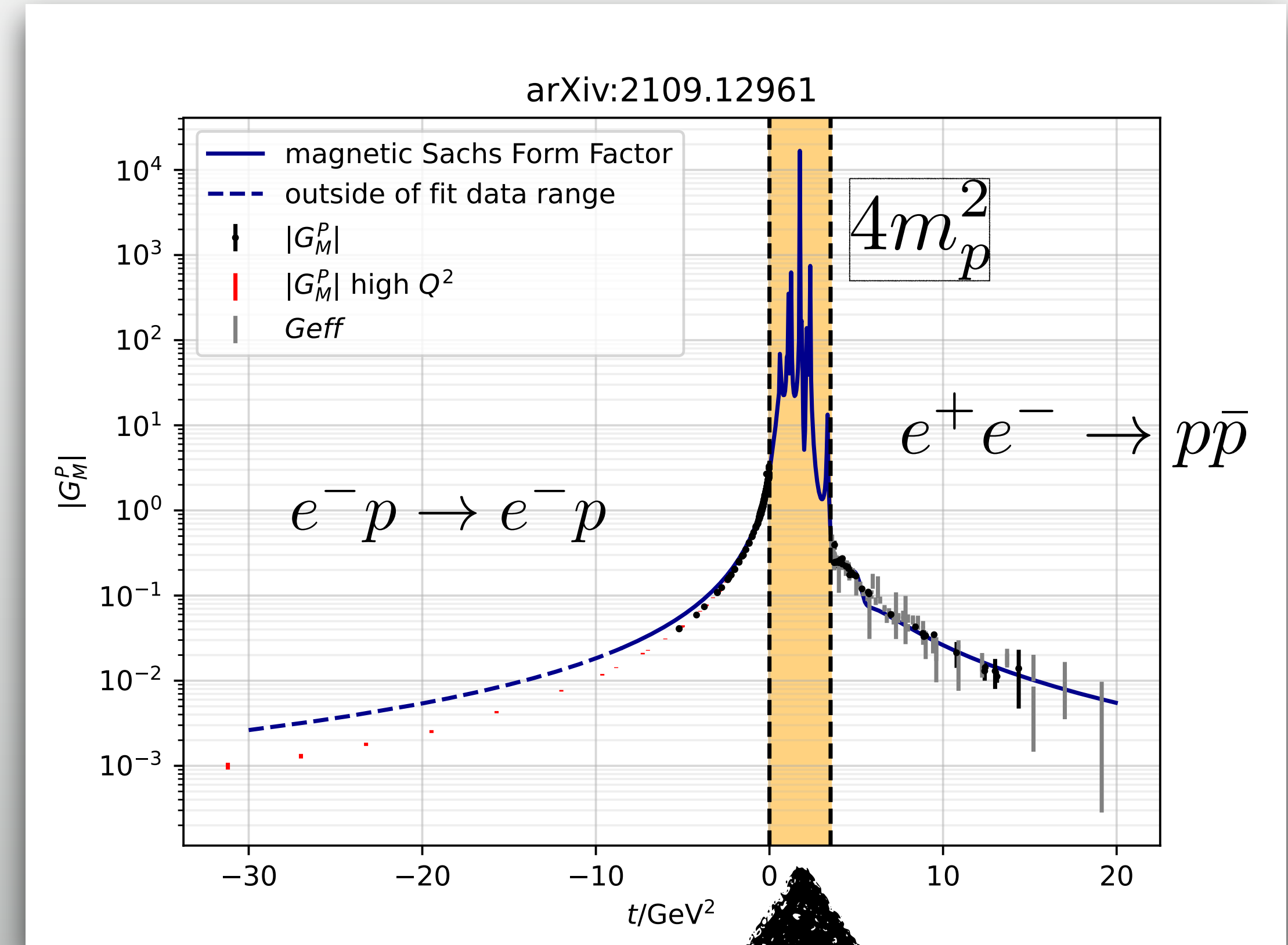


State of the Art Form factors

2109.12961



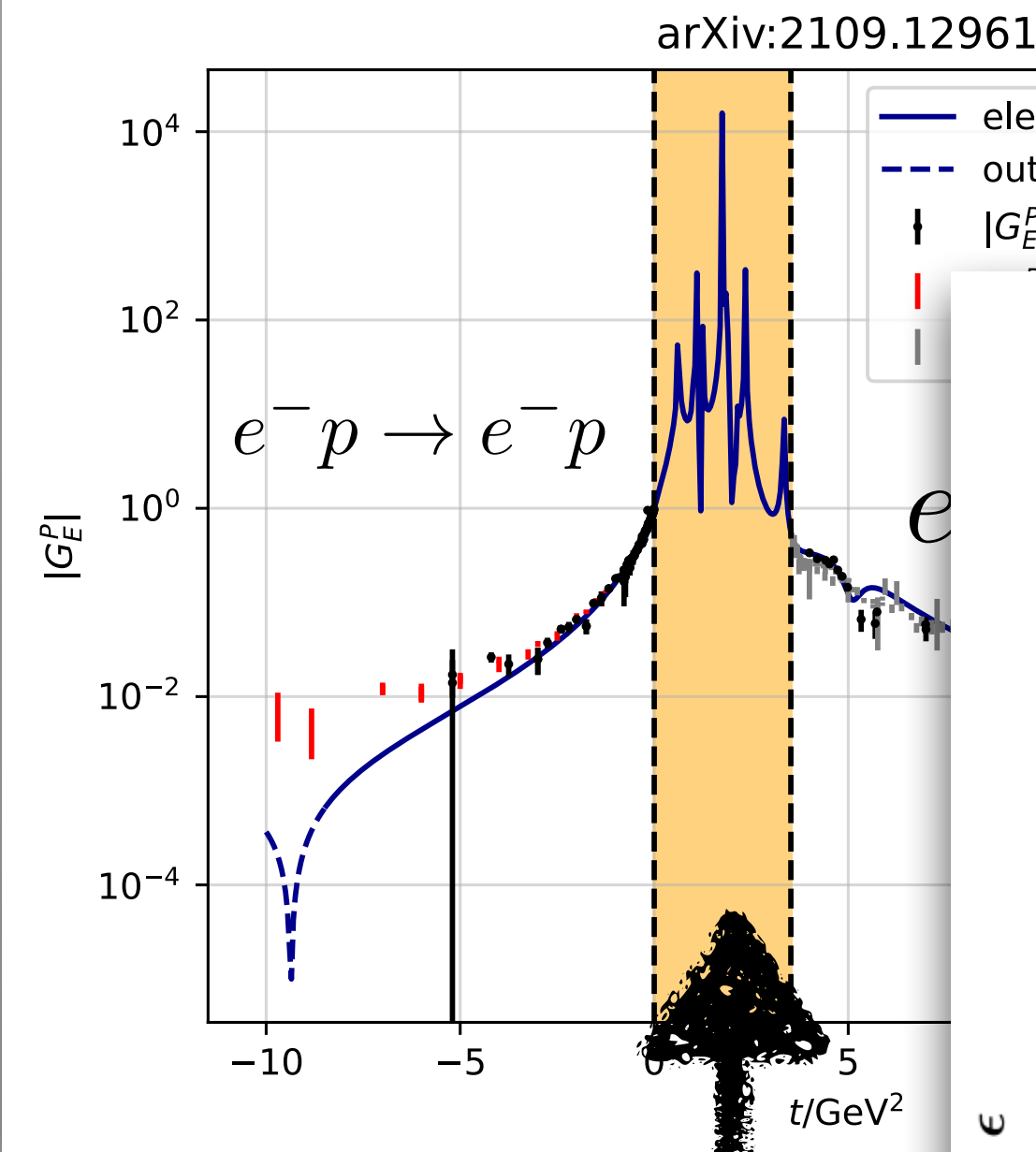
Unphysical region



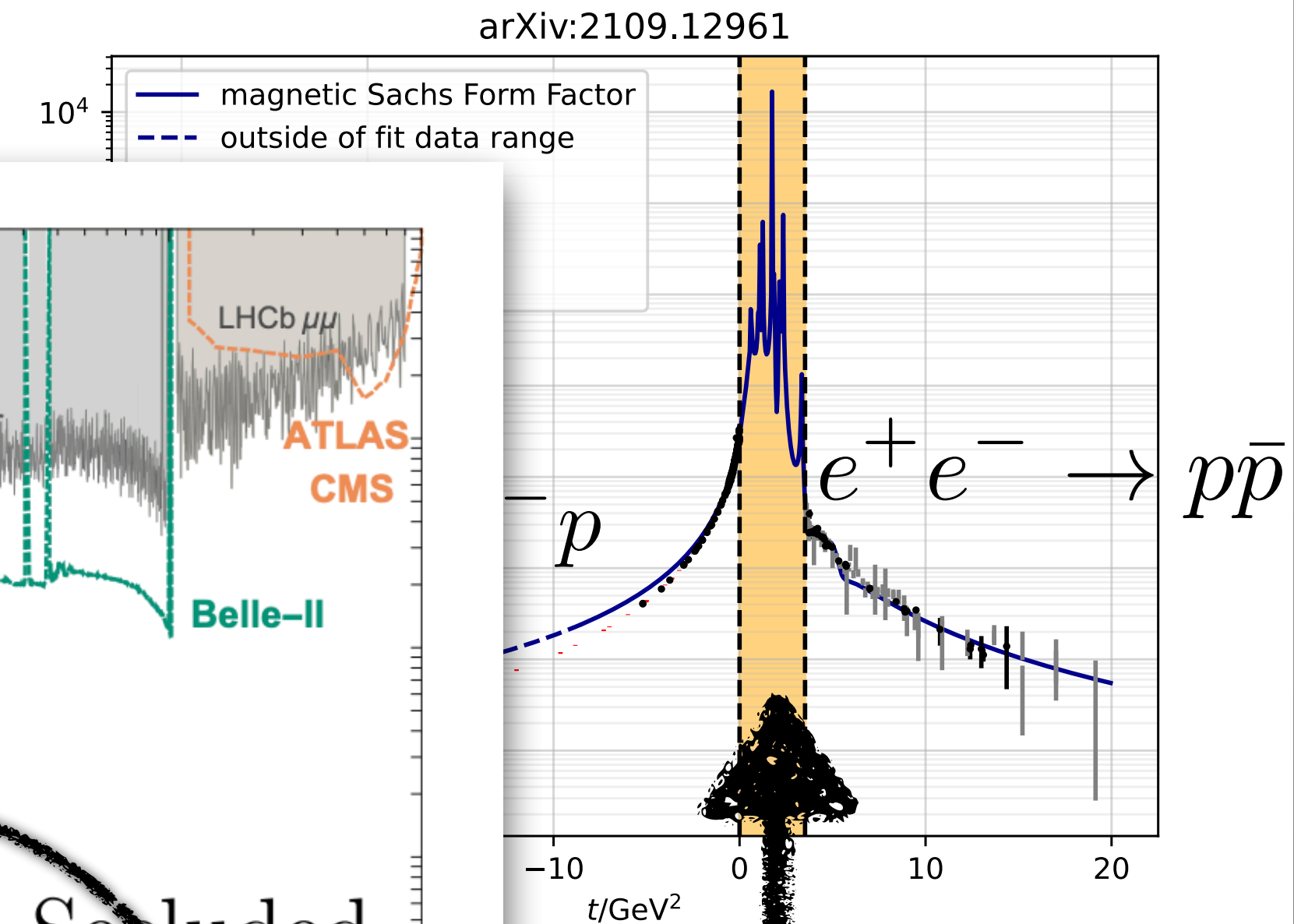
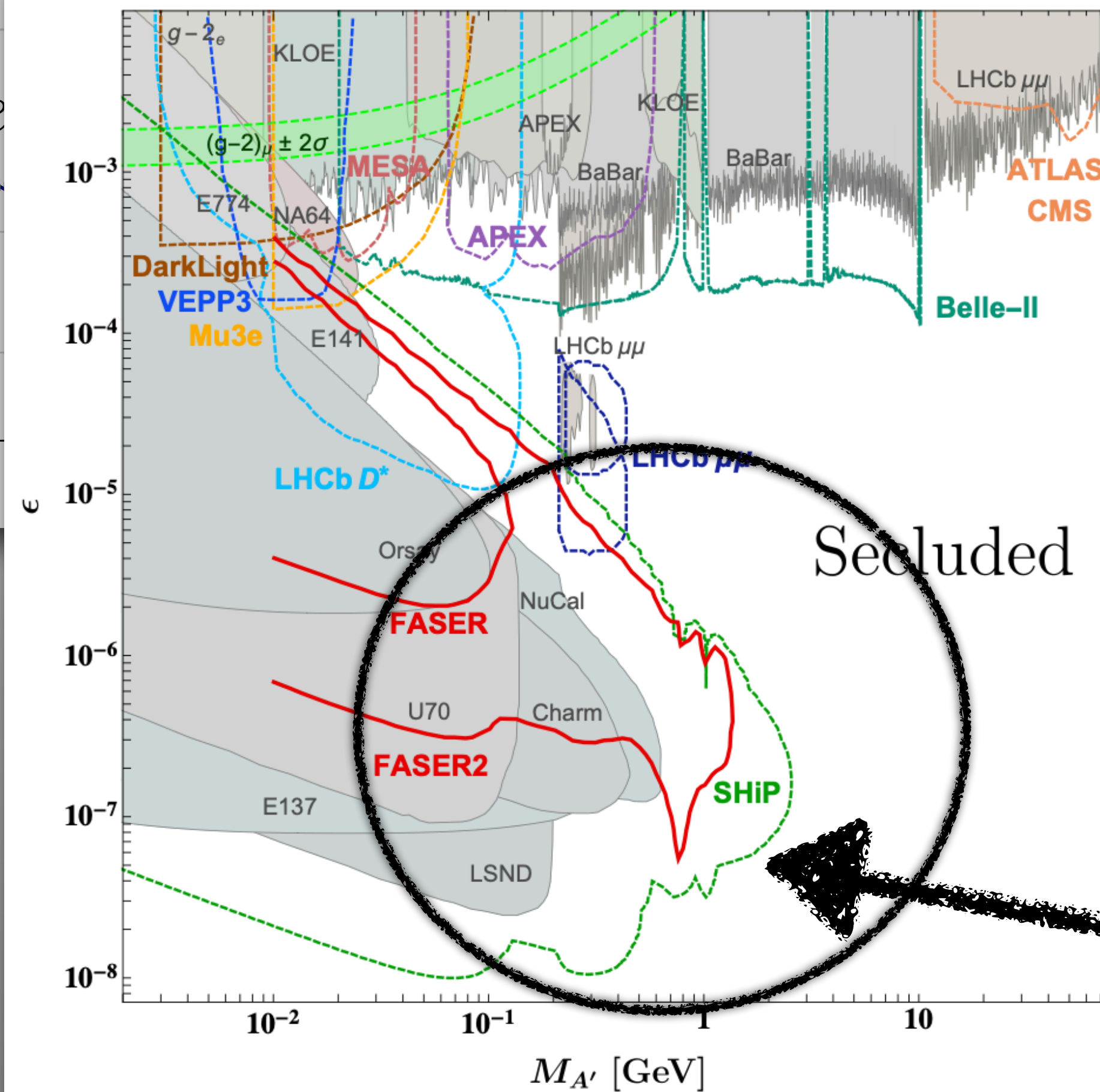
Unphysical region

State of the Art Form factors

2109.12961



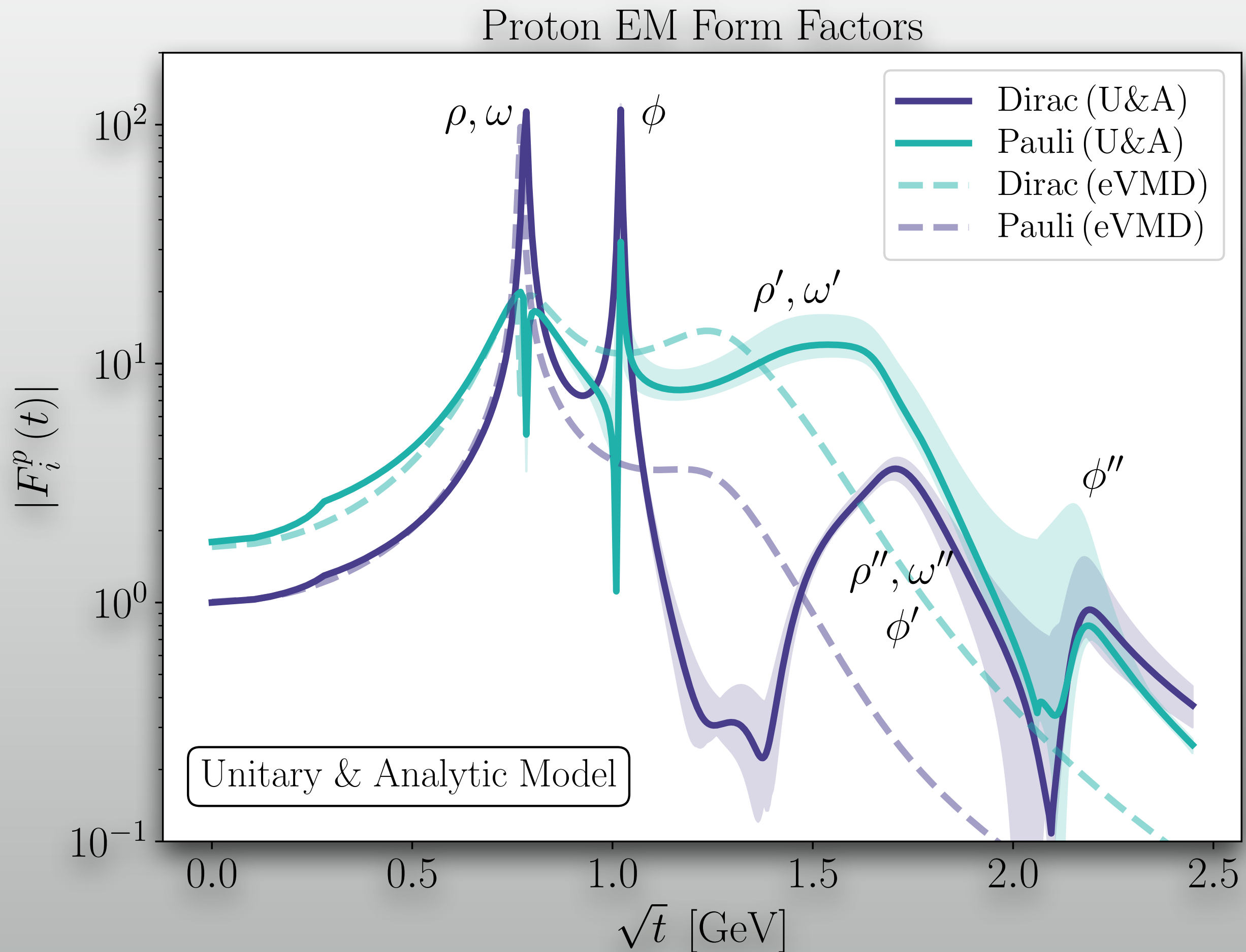
Unphysical region



Unphysical region

Our region of interest!

Form factors



Based on PDG values for
physical vector meson poles

$$m_\omega, m_{\omega'}, m_{\omega''}, \Gamma_\omega, \Gamma_{\omega'}, \Gamma_{\omega''}$$

$$m_\rho, m_{\rho'}, m_{\rho''}, \Gamma_\rho, \Gamma_{\rho'}, \Gamma_{\rho''}$$

$$m_\phi, m_{\phi'}, m_{\phi''}, \Gamma_\phi, \Gamma_{\phi'}, \Gamma_{\phi''}$$

1601.06190 VMD, Unitarity&Analytic Model

Form factors

Off-shell form factor:

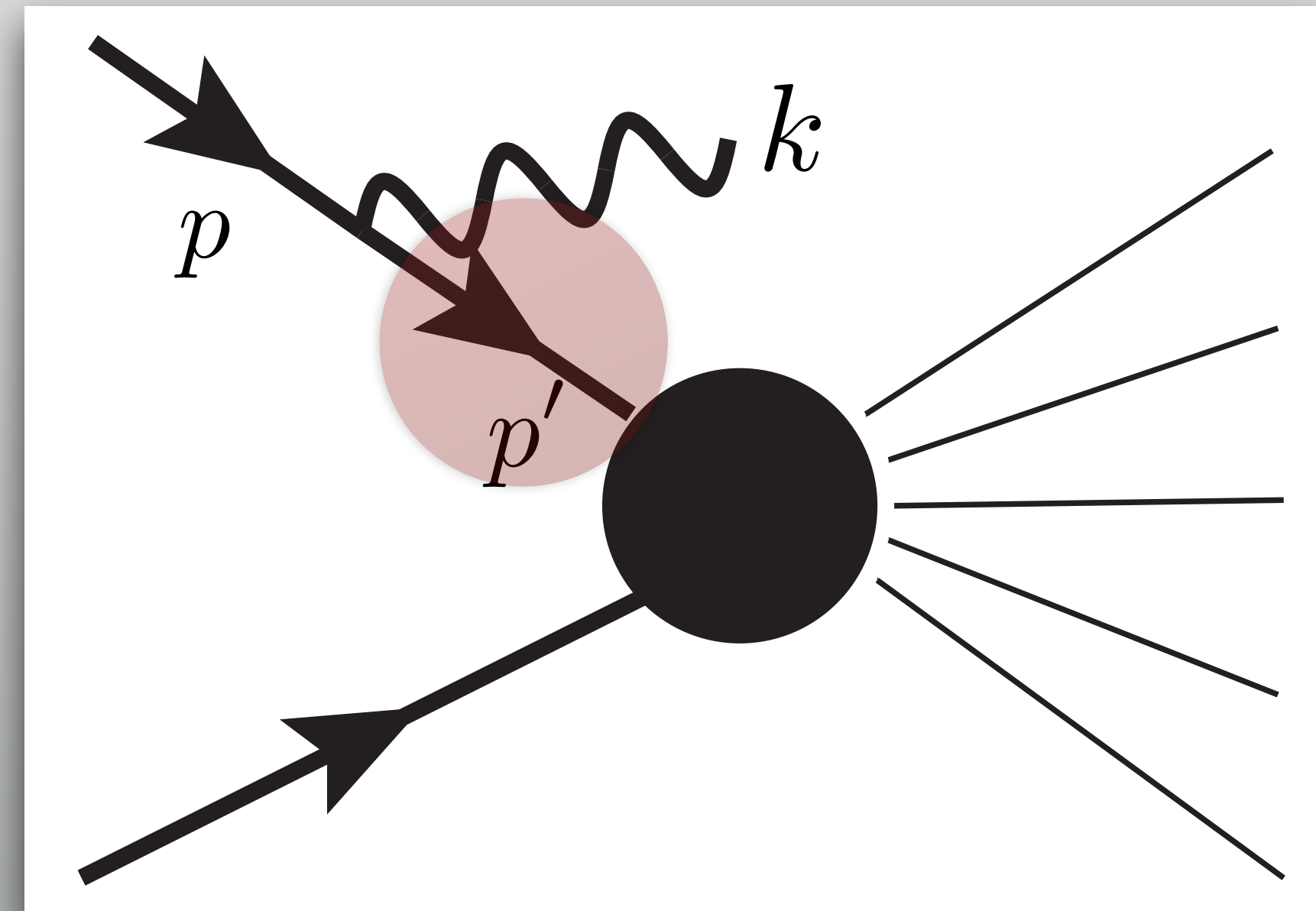
$$F_{pp'V}(p'^2) = \frac{1}{1 + (p'^2 - m_p^2)^2 / \Lambda_p^4}$$

with $1 \text{ GeV} \lesssim \Lambda_p \lesssim 2 \text{ GeV}$

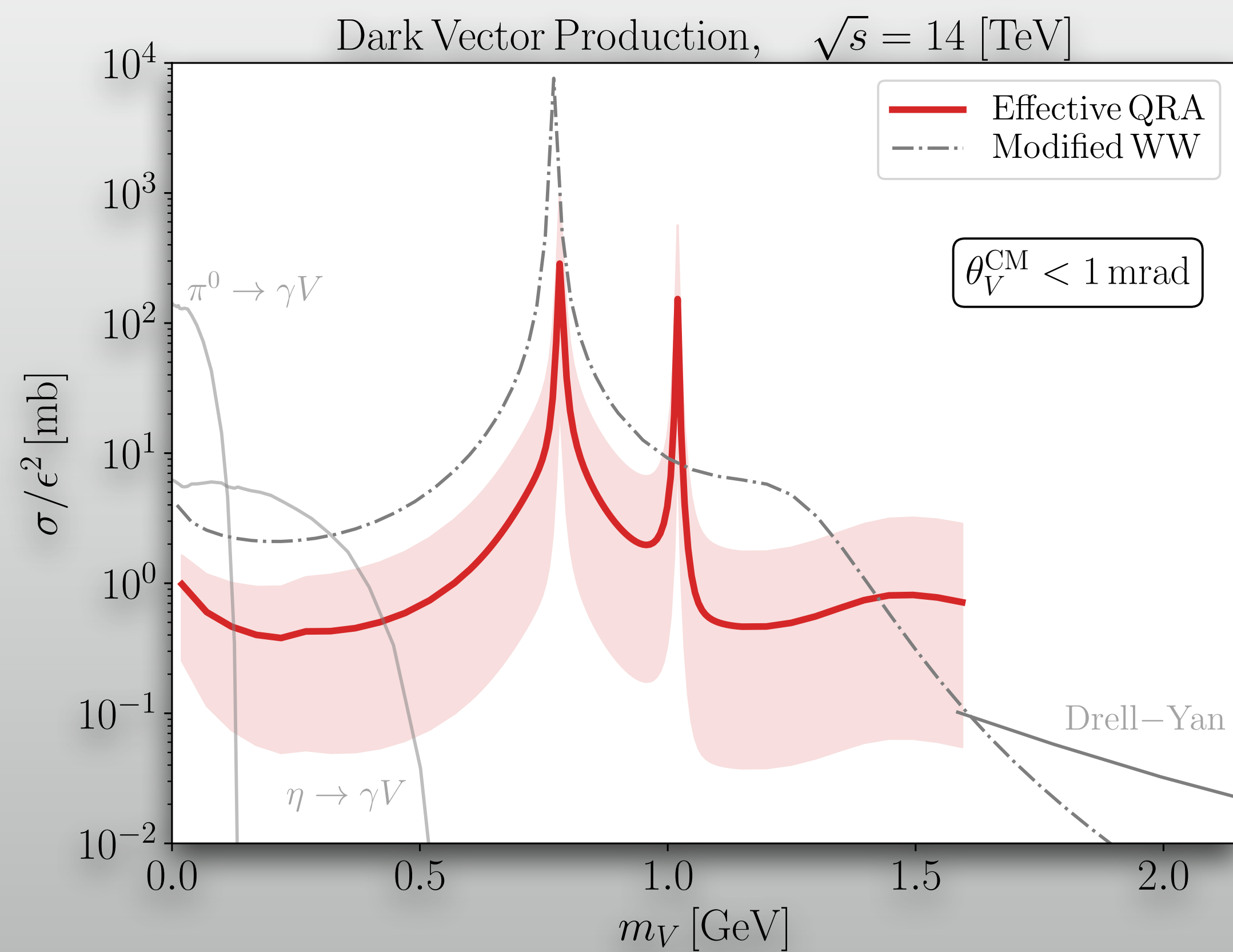
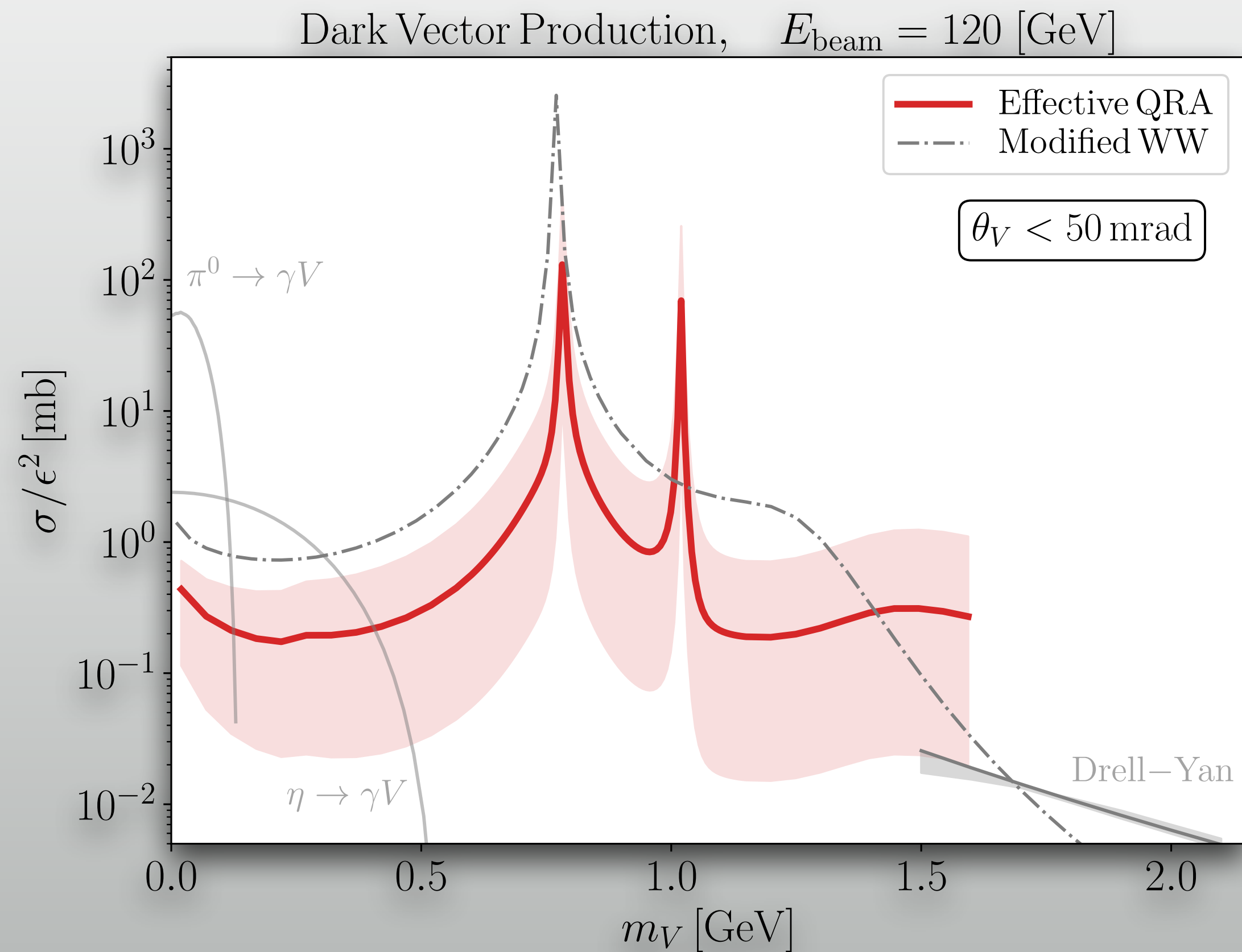
controls off-shell behaviour of intermediate proton



$$F_i(k^2, p'^2) = F_{pp'V}(p'^2) \times F_i(k^2)$$



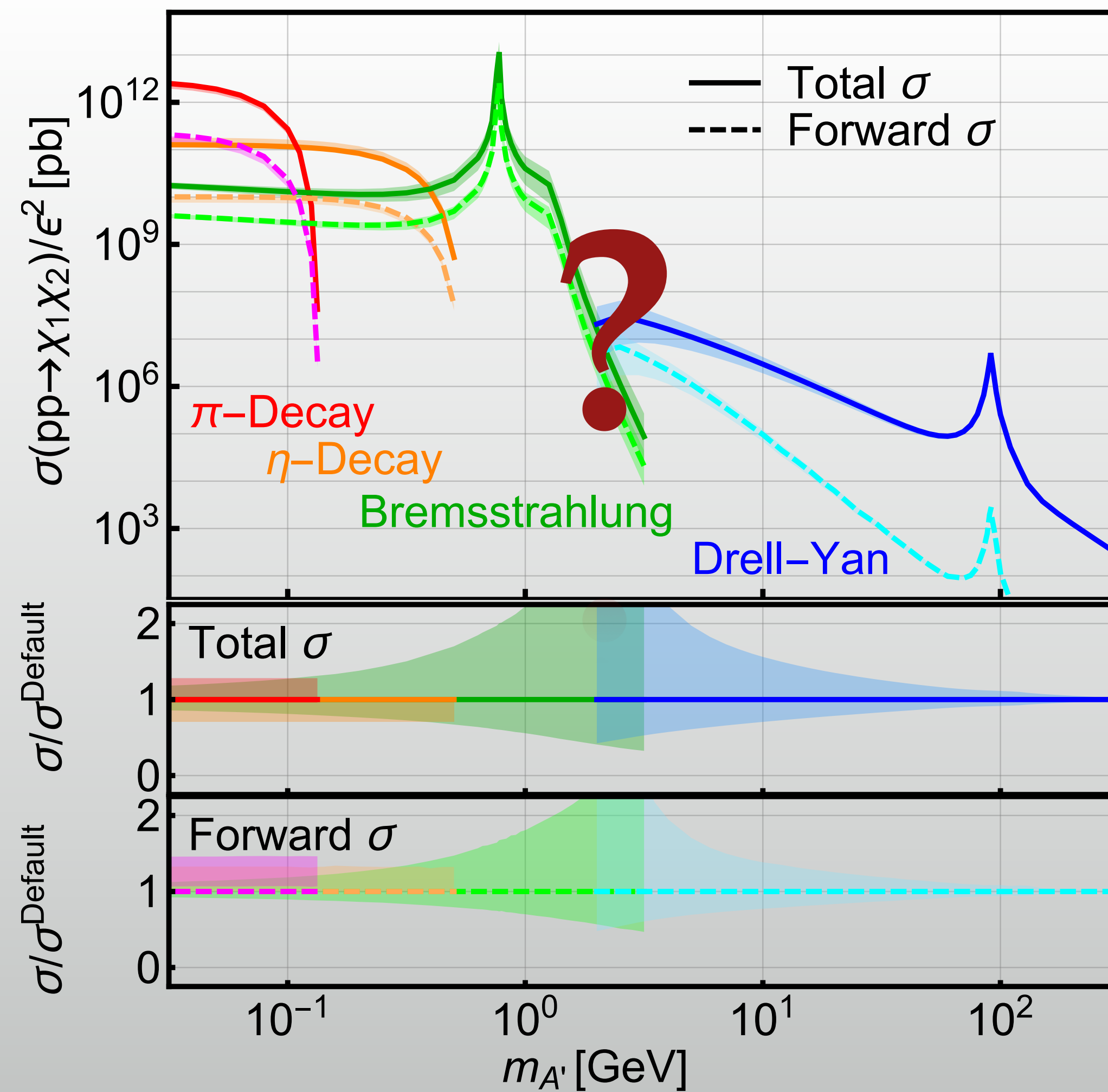
Preliminary results:



$$1 \text{ GeV} \lesssim \Lambda_p \lesssim 2 \text{ GeV}$$

Production of BSM particles

- Drell-Yan → sub-dominant
- Meson decays → in lower mass range
- Bremsstrahlung
- Is there more than this?



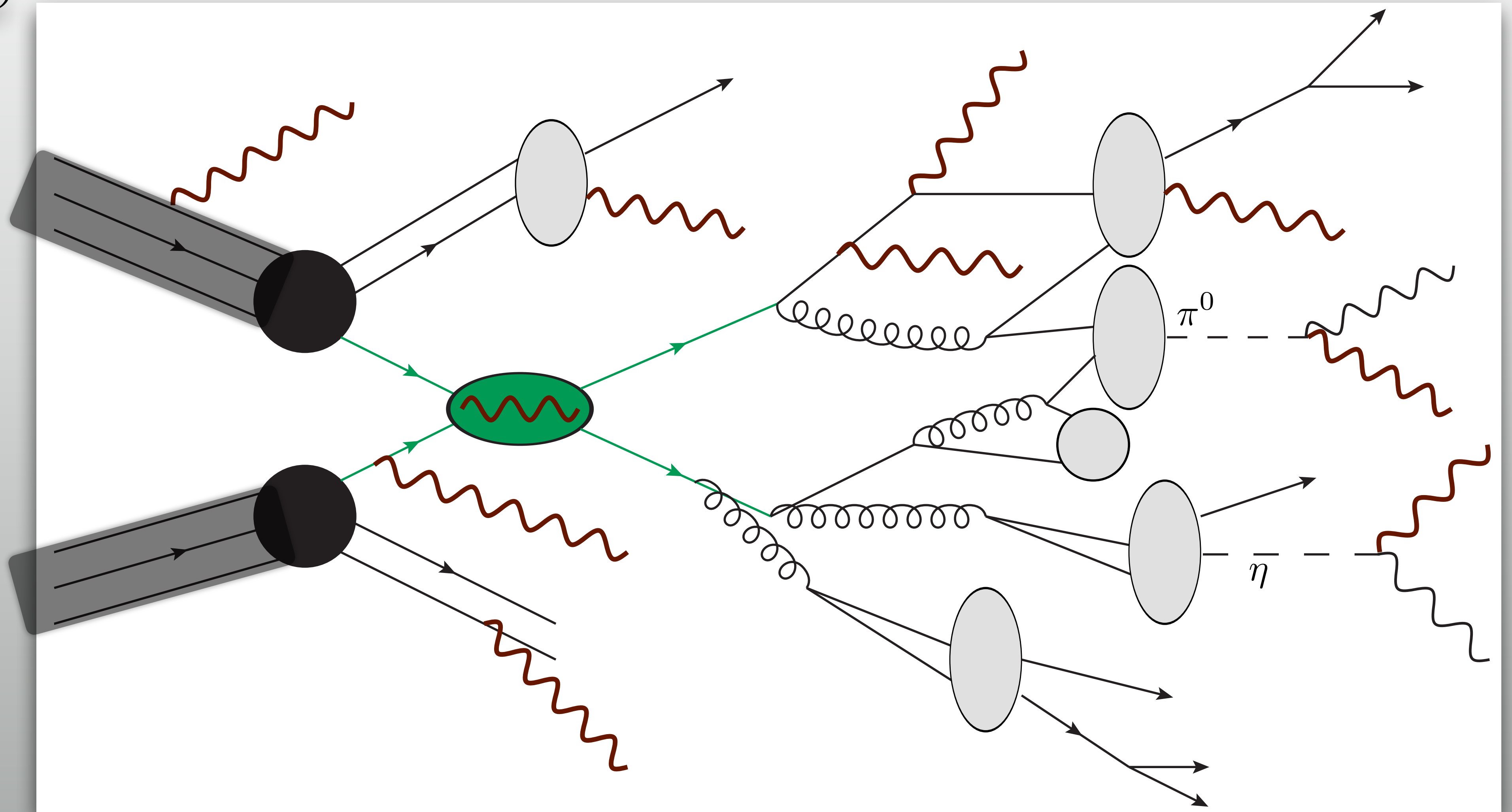
$$\mathcal{L} = -\frac{\epsilon}{2} F_{\mu\nu} X^{\mu\nu}$$

arXiv 1810.01879

A. Berlin, F. Kling

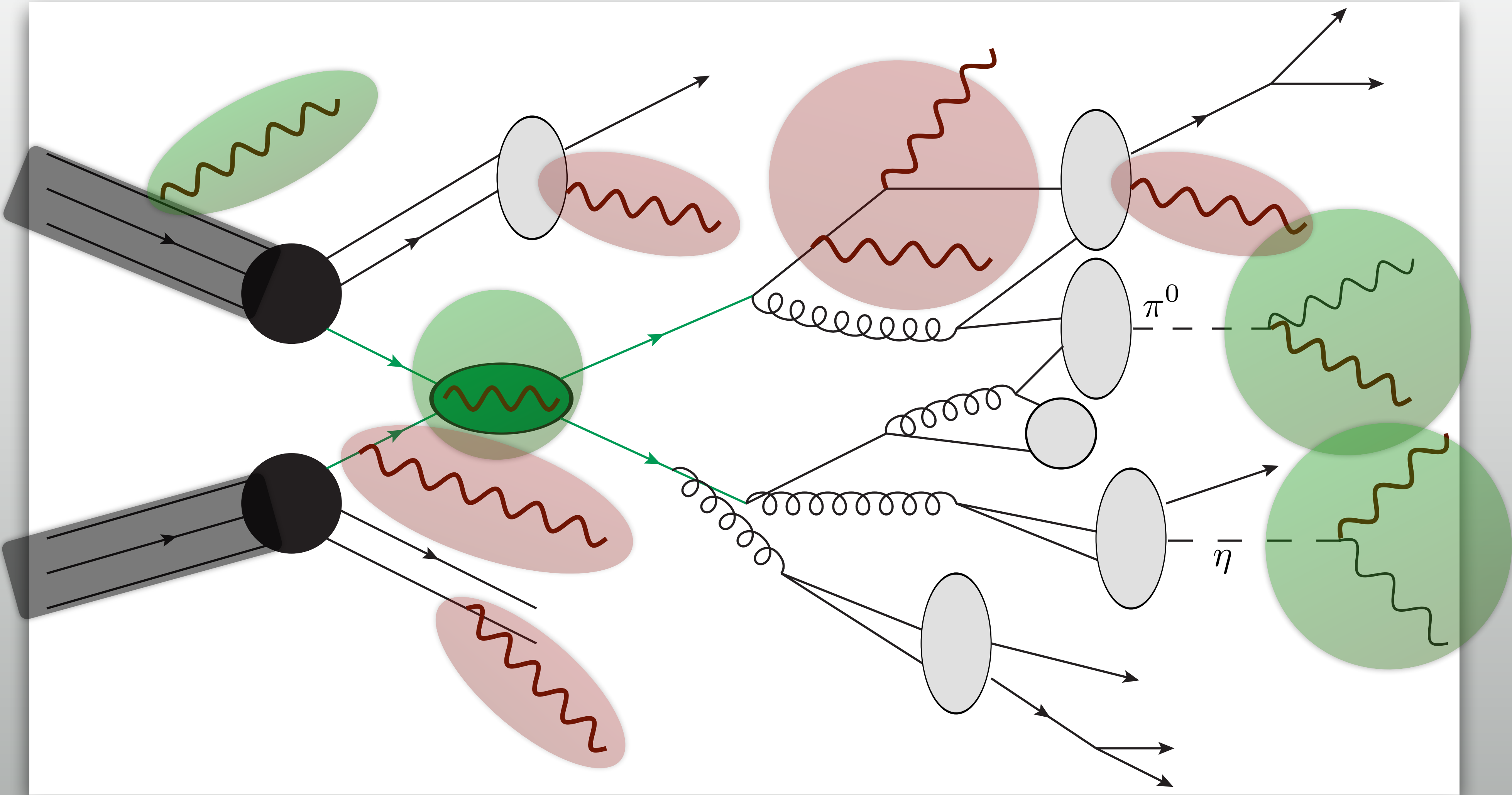
Where/how do we produce dark photons in proton-proton collisions?

$$\mathcal{L} = -\frac{\epsilon}{2} F_{\mu\nu} X^{\mu\nu}$$



Additional Production Modes?

Unknown Sources

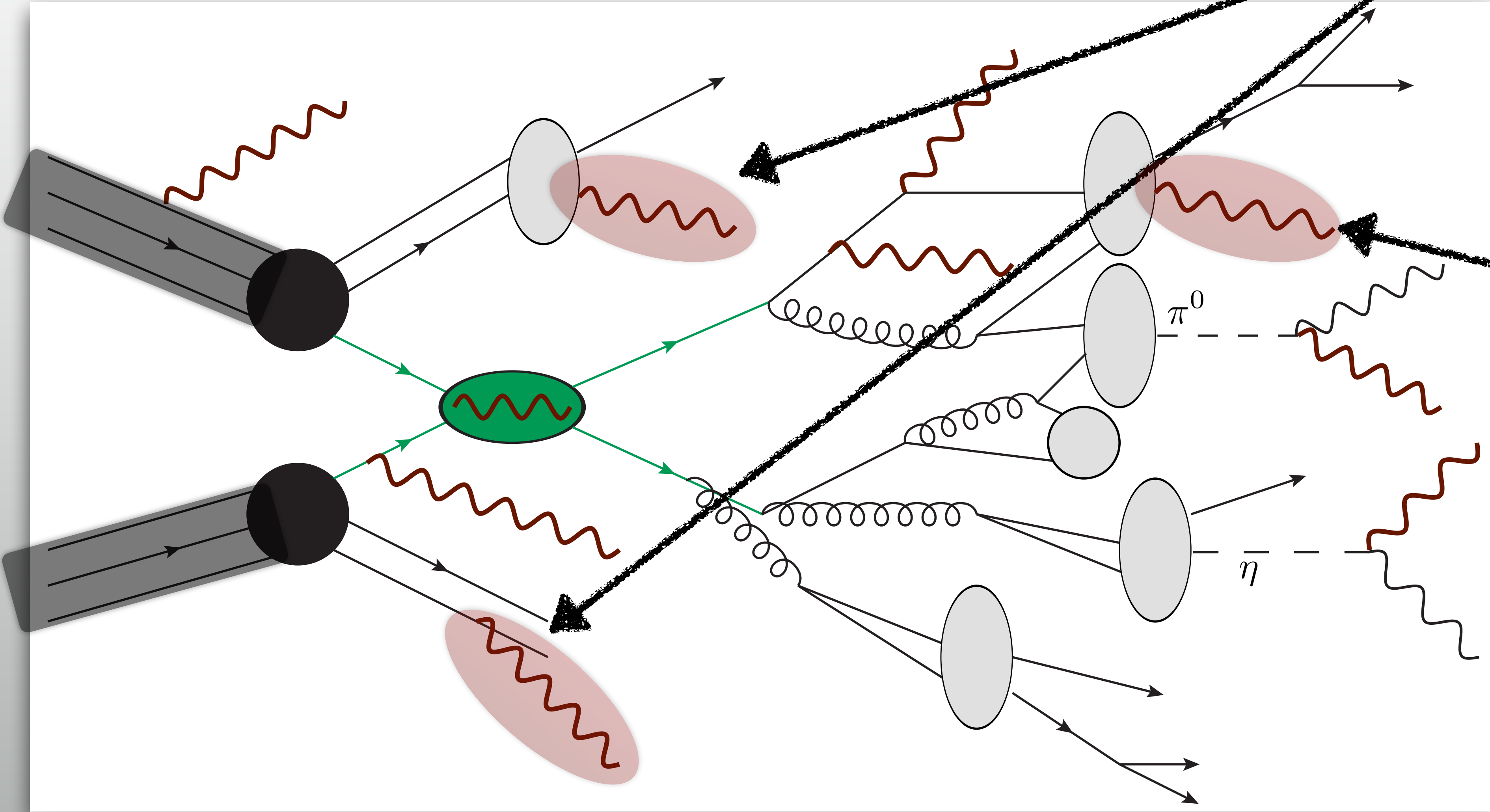


Additional Production Modes?

Beam remnants

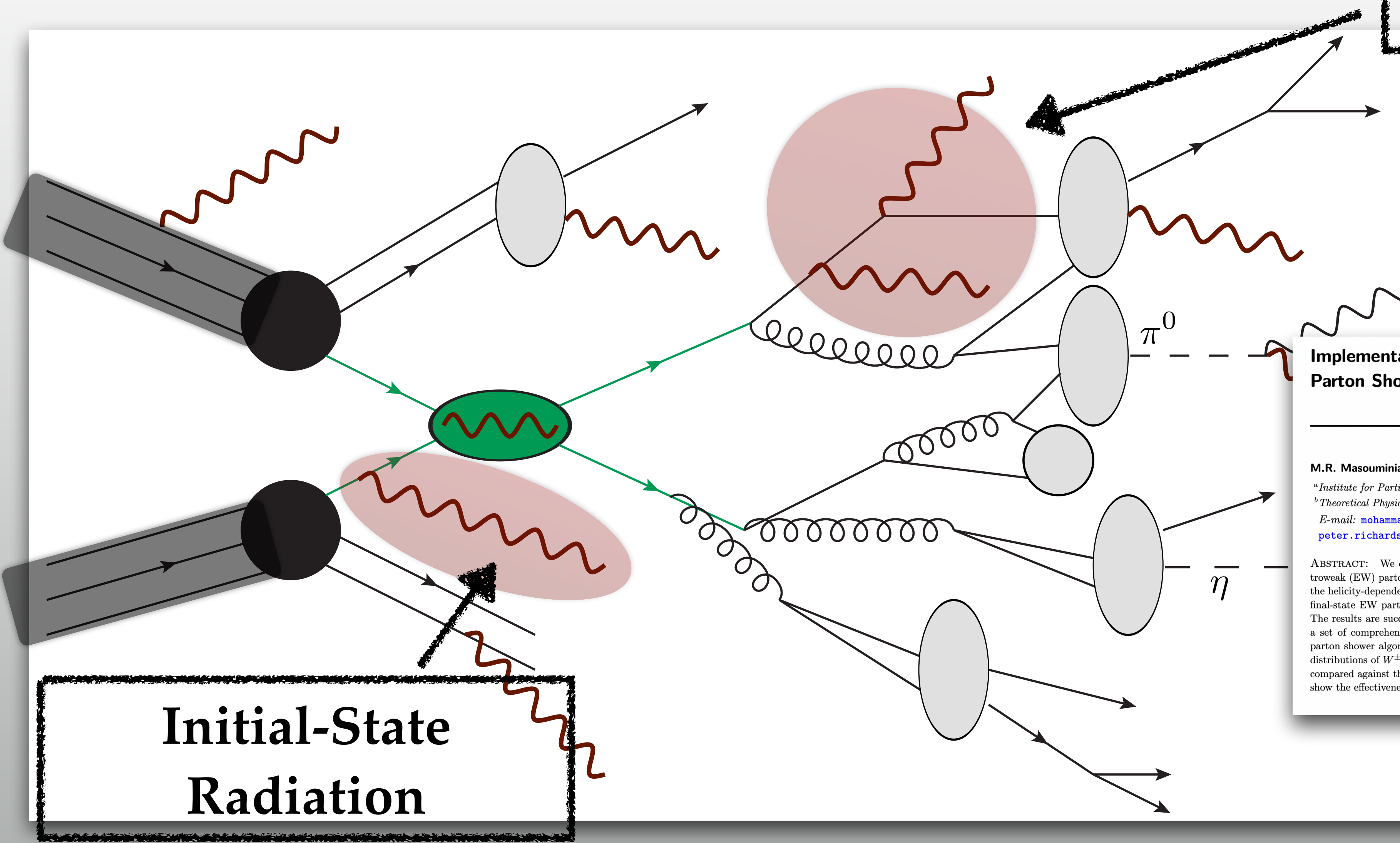
Before or as part of hadronization

Hadronization of final states



Additional Production Modes?

Final-State Radiation



Initial-State Radiation

Implementation of Angularly Ordered Electroweak Parton Shower in Herwig 7

M.R. Masouminia,^a P. Richardson^{a,b}
^a*Institute for Particle Physics Phenomenology, Durham University, Durham, UK*
^b*Theoretical Physics Department, CERN, Switzerland*
E-mail: mohammad.r.masouminia@durham.ac.uk,
peter.richardson@durham.ac.uk



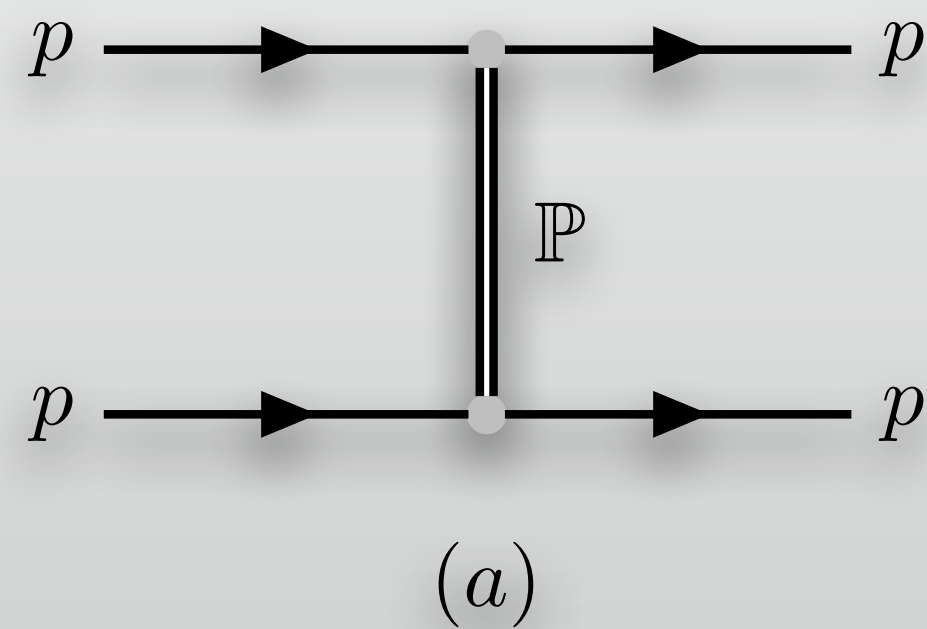
ABSTRACT: We discuss the necessary steps for implementing an angularly ordered (AO) electroweak (EW) parton shower in Herwig 7 multi-purpose event generator. This includes calculating the helicity-dependent *quasi-collinear* EW branching functions that correspond to the full range of final-state EW parton shower, in addition to the initial-state EW gauge vector boson radiations. The results are successfully embedded in the AO Herwig 7 shower algorithm and have undergone a set of comprehensive and conclusive performance tests. Furthermore, we have used this EW parton shower algorithm, alongside the existing *QCD + QED* AO shower, to predict the angular distributions of W^\pm bosons in LHC events with high transverse momentum jets. These results are compared against the explicitly generated underlying events as well as the existing ATLAS data to show the effectiveness of the newly implemented *QCD + QED + EW* AO parton shower scheme.

arXiv 2108.10817

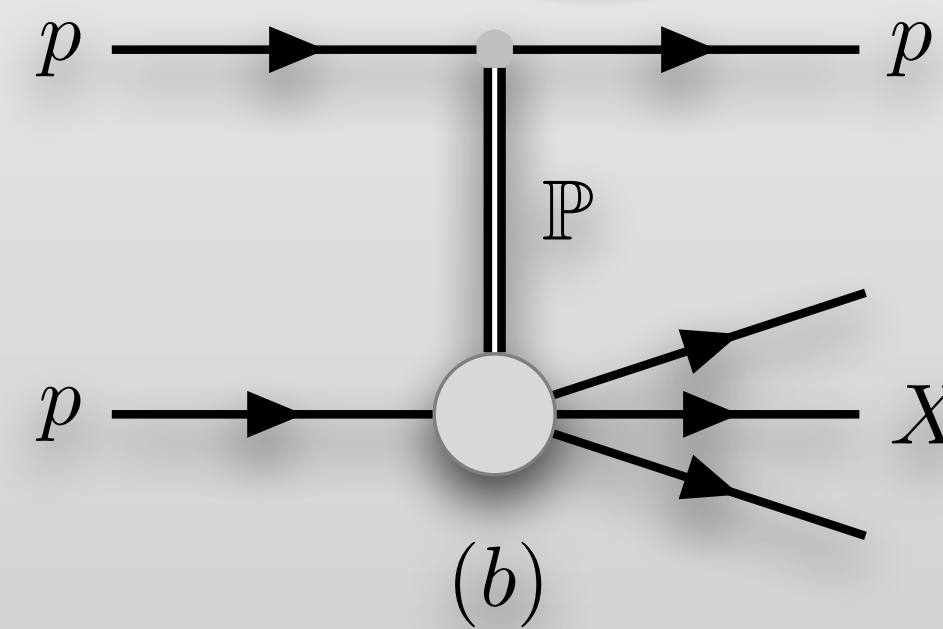
Identify relevant processes

Any process including **partons**

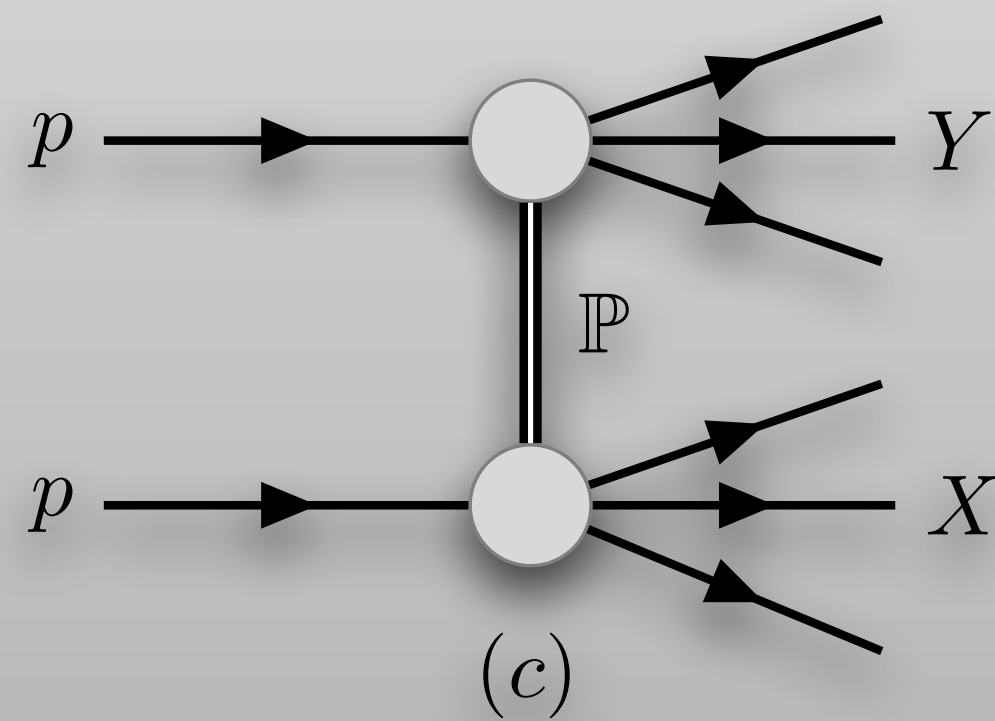
~ 25%



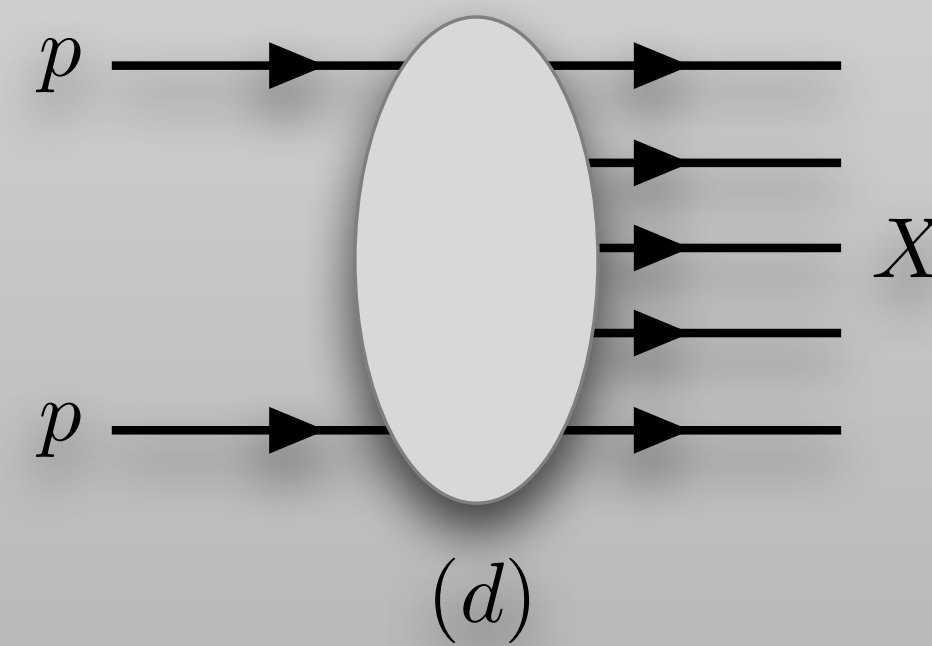
~ 11%



~ 8%

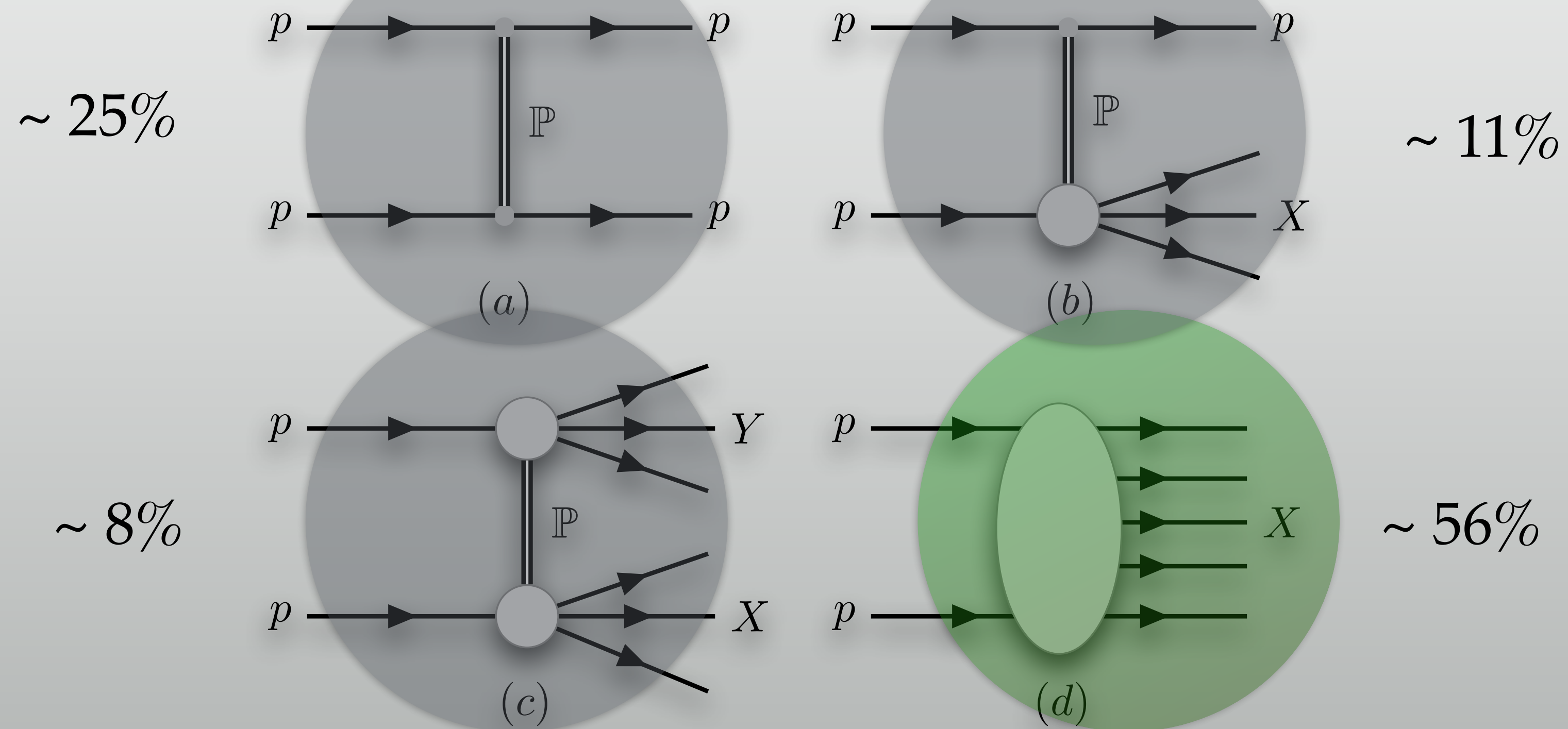


~ 56%



Identify relevant processes

Any process including partons



Modelling in Herwig



Dijet-like Event $qq' \rightarrow qq' +$

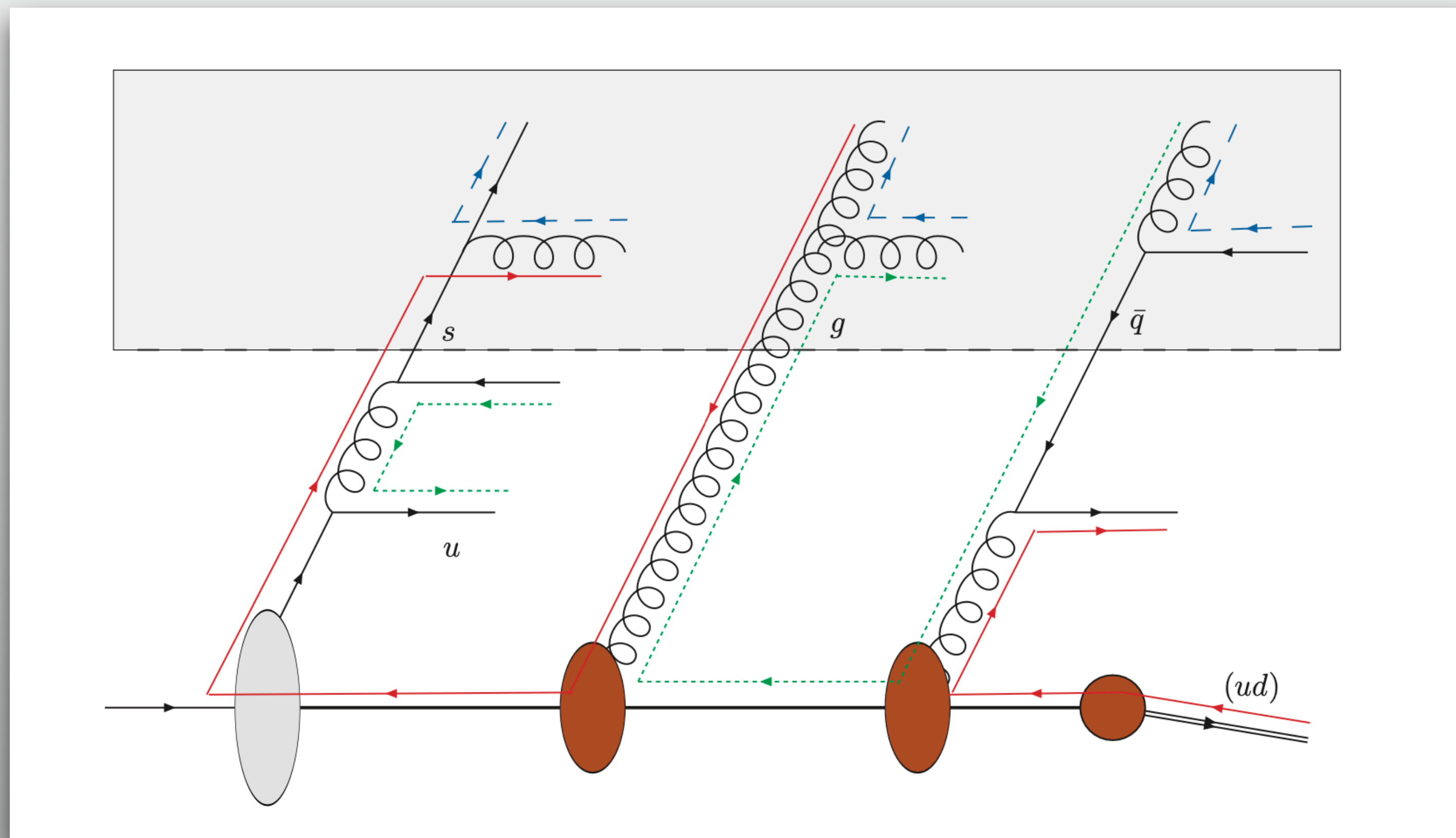
Multiple Partonic Interactions (MPI)

QCD 2-to-2

$$qq \rightarrow qq$$

$$gg \rightarrow gg$$

$$gq \rightarrow gq\dots$$

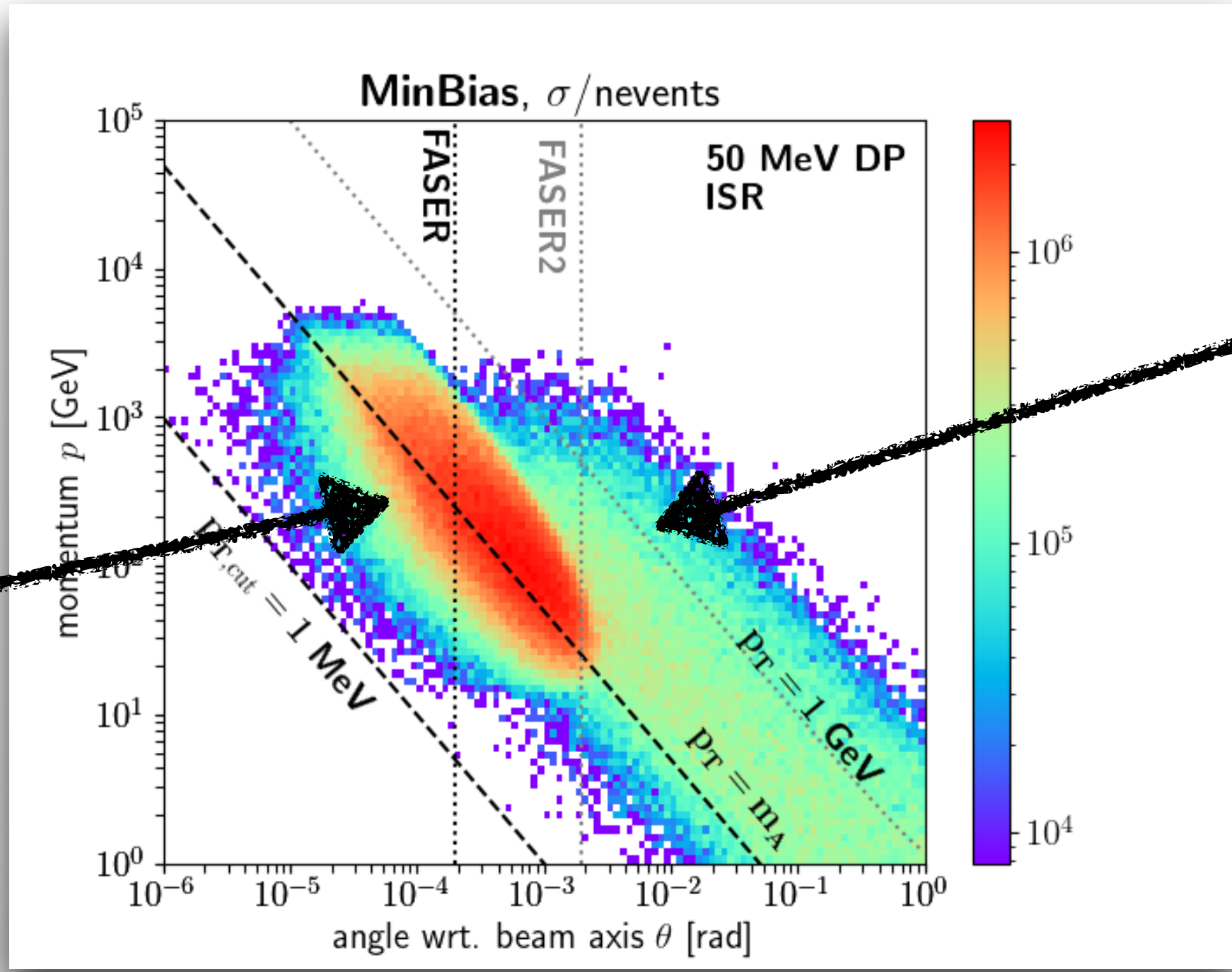


arXiv 0803.3633

Bähr, Gieseke,
Seymour

ISR

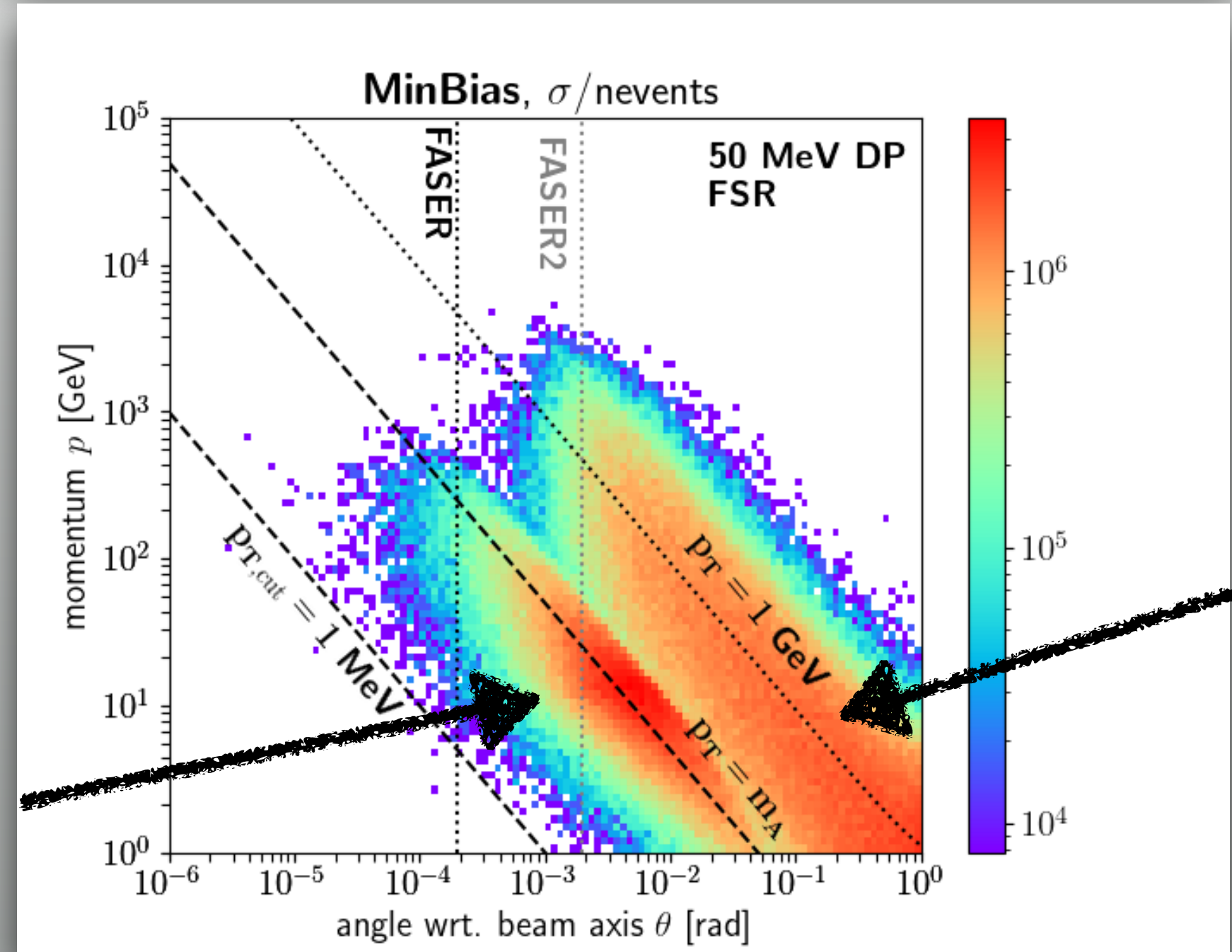
MPI



Dijet-like event

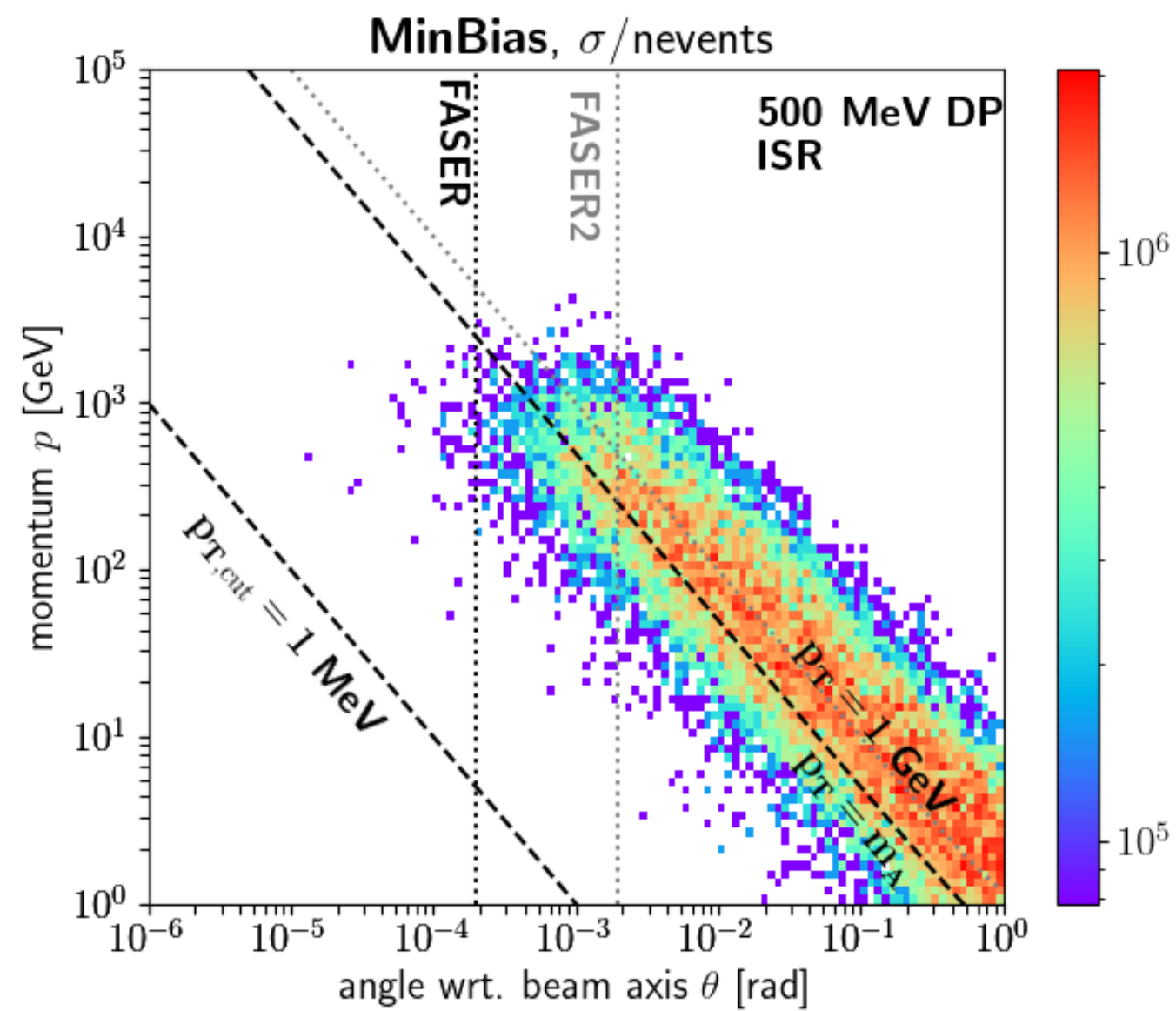
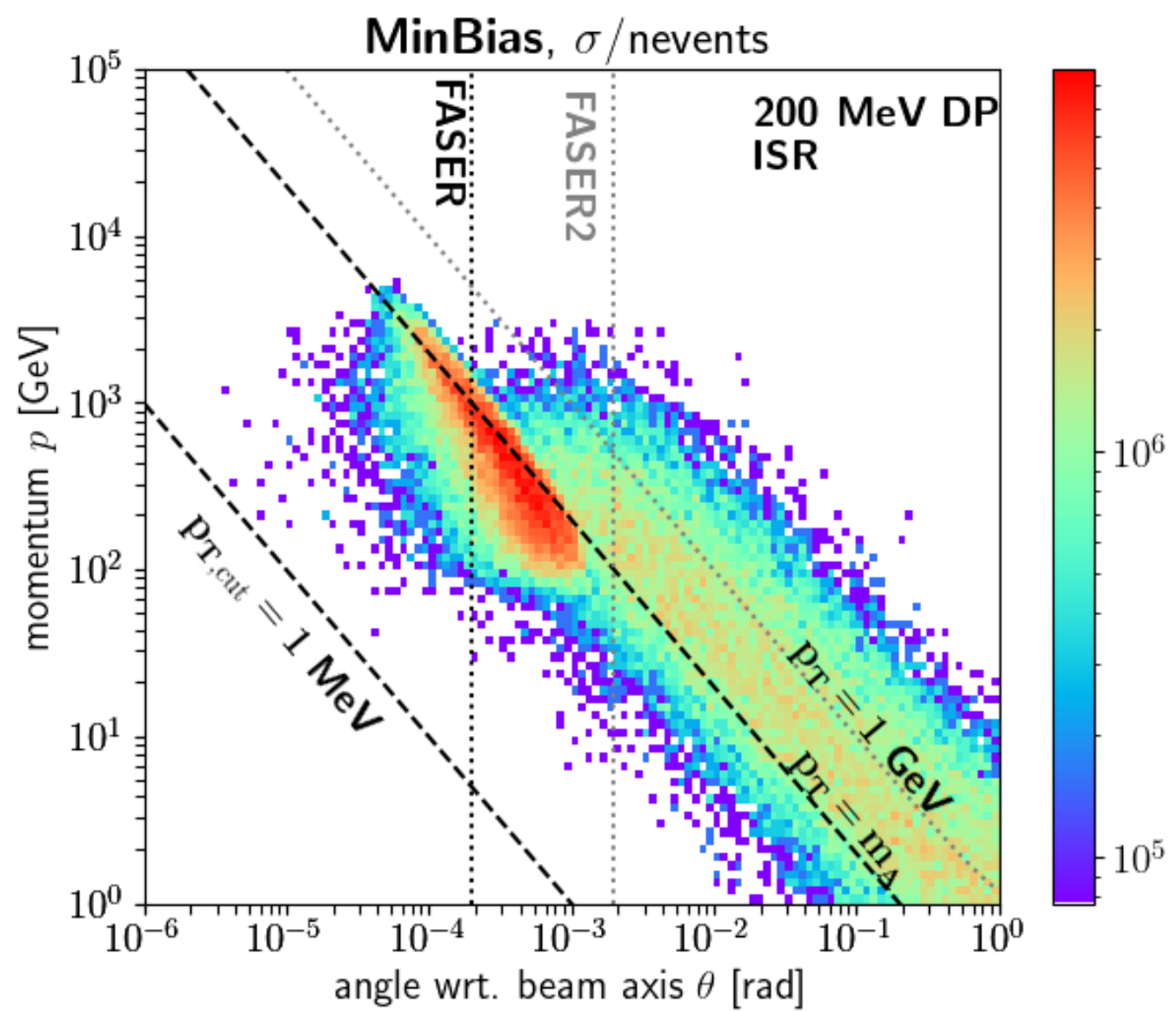
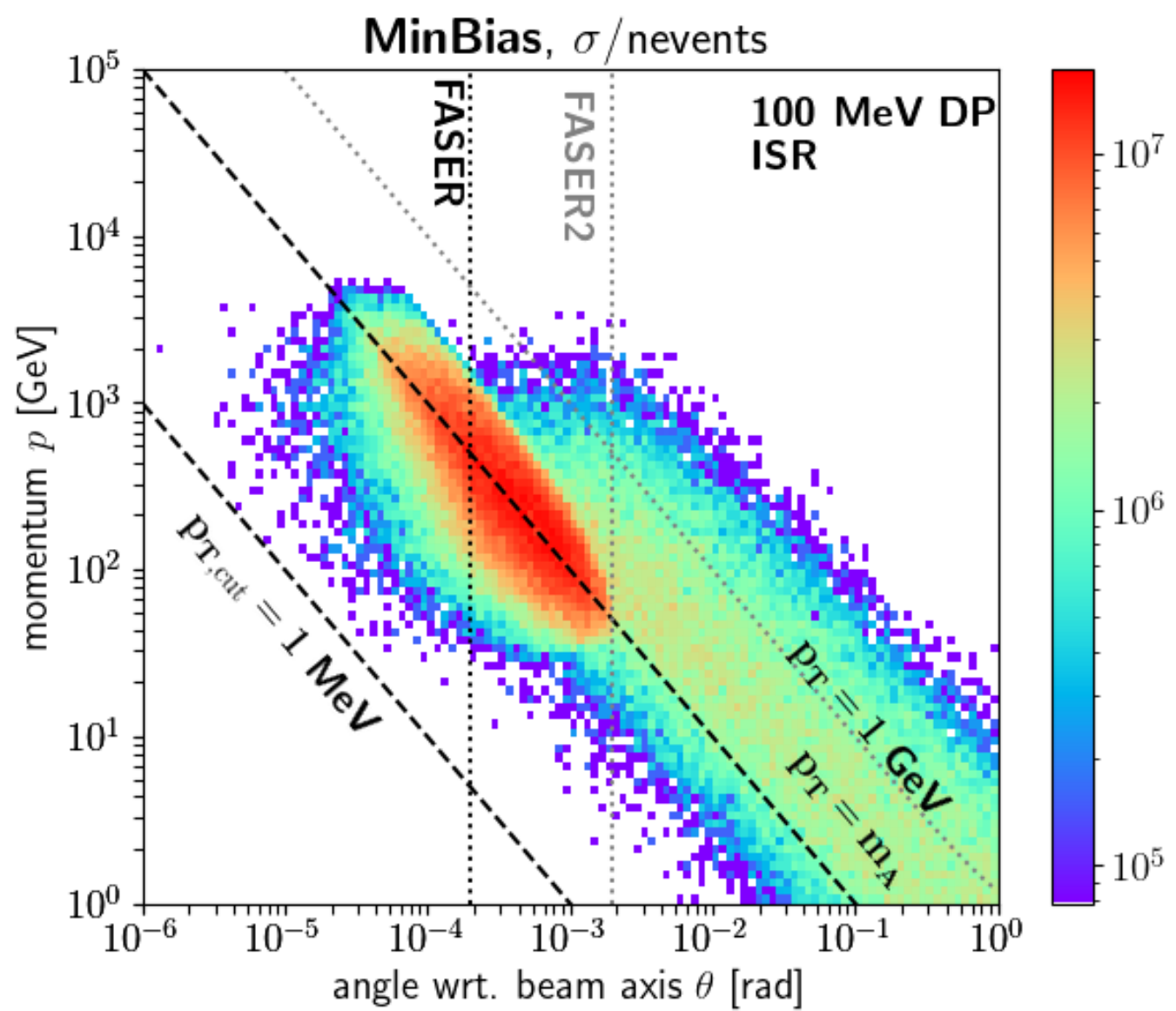
FSR

MPI

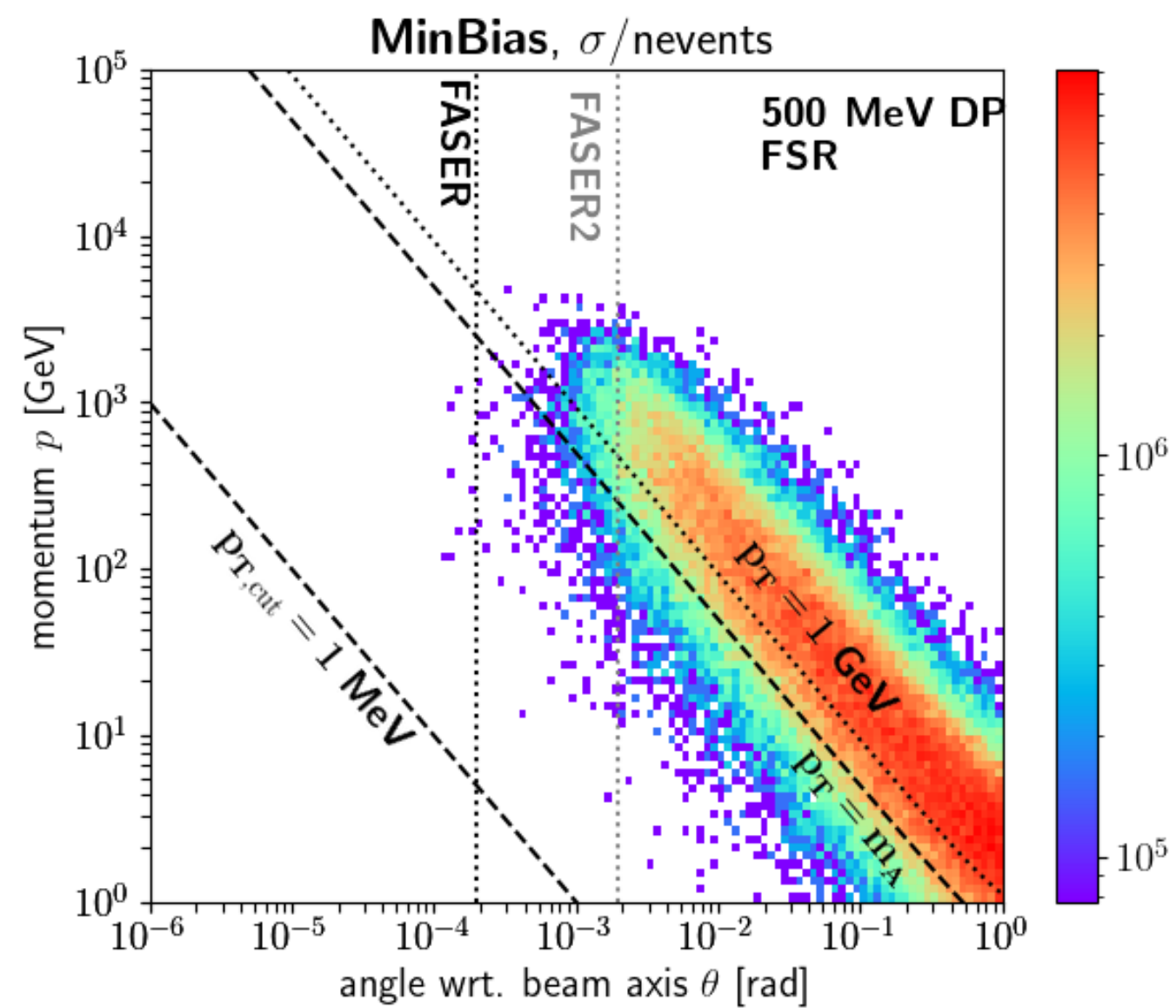
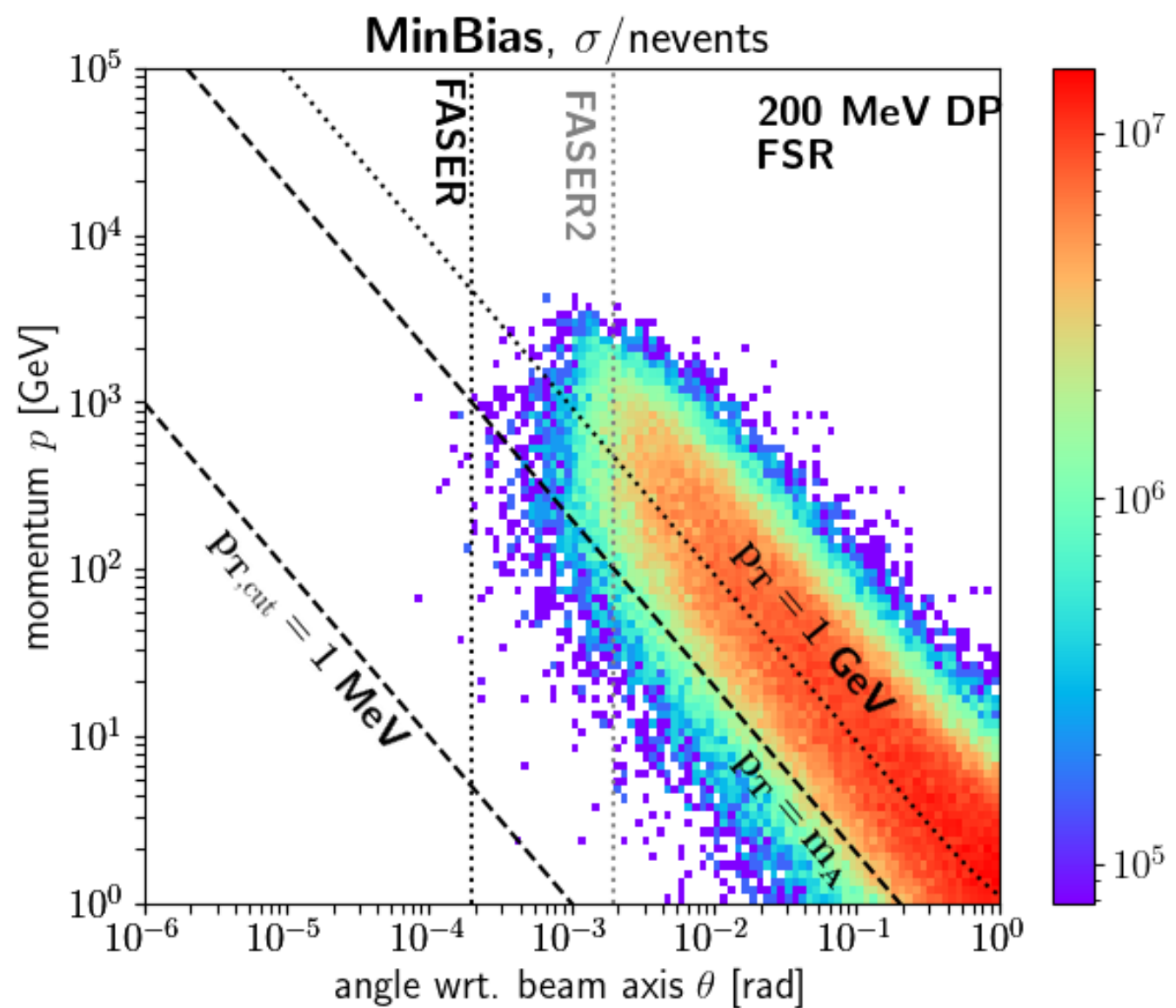
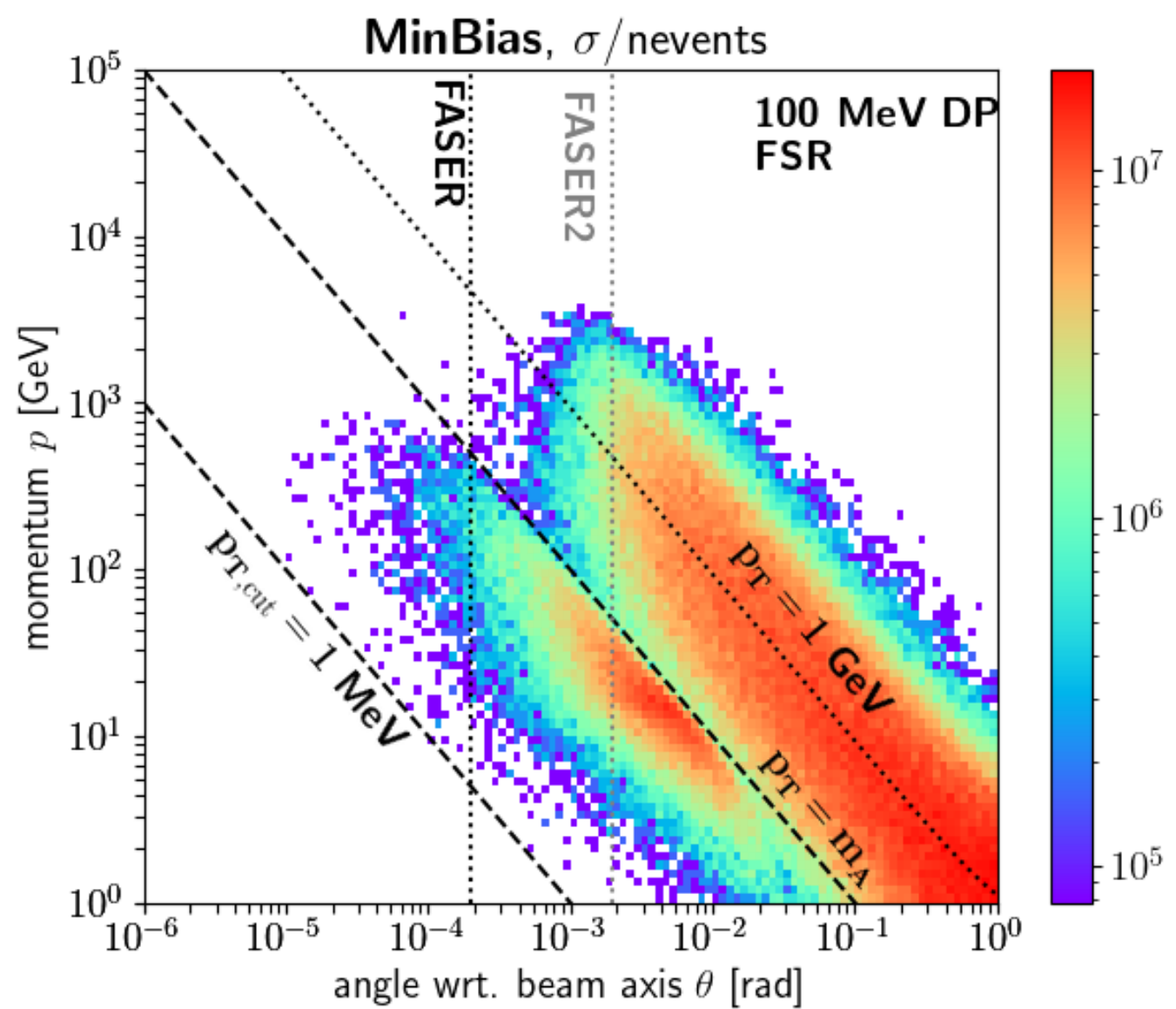


Dijet-like event

ISR

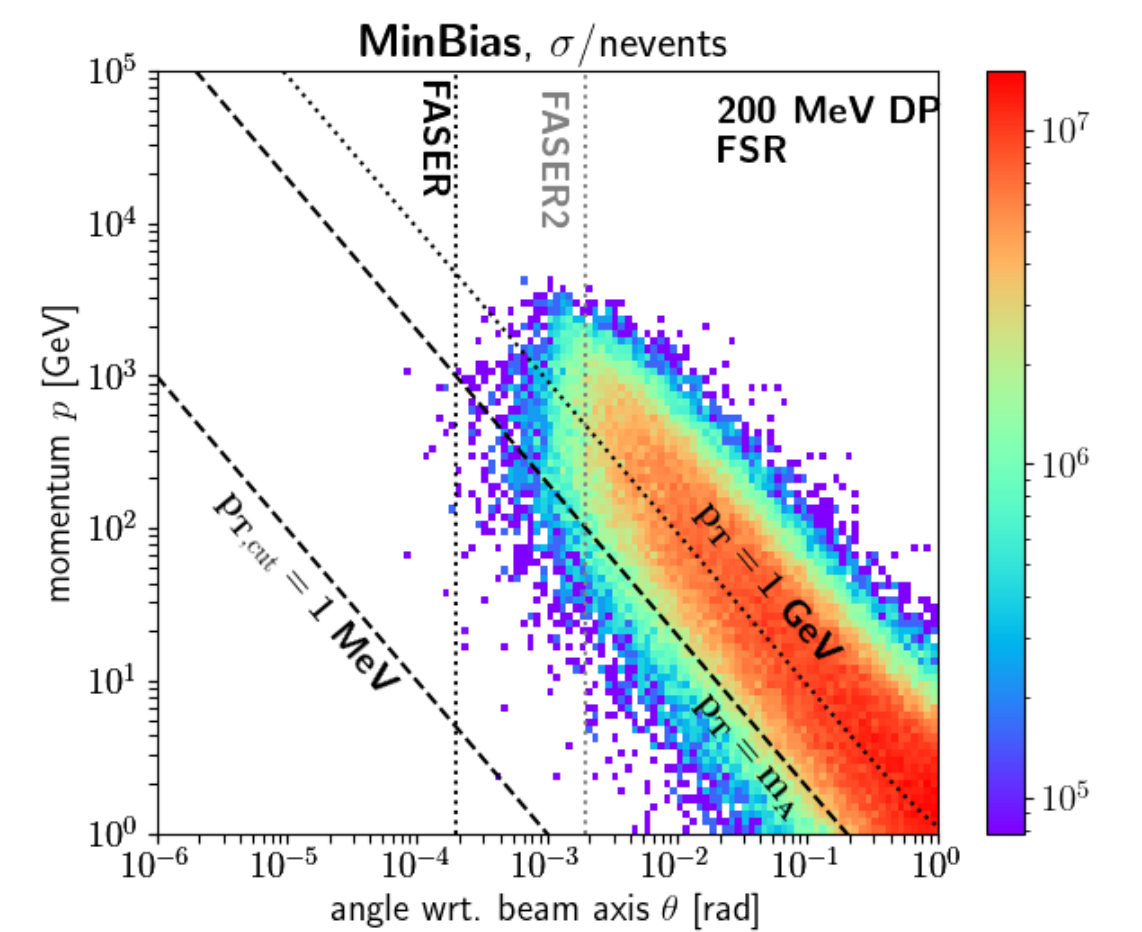
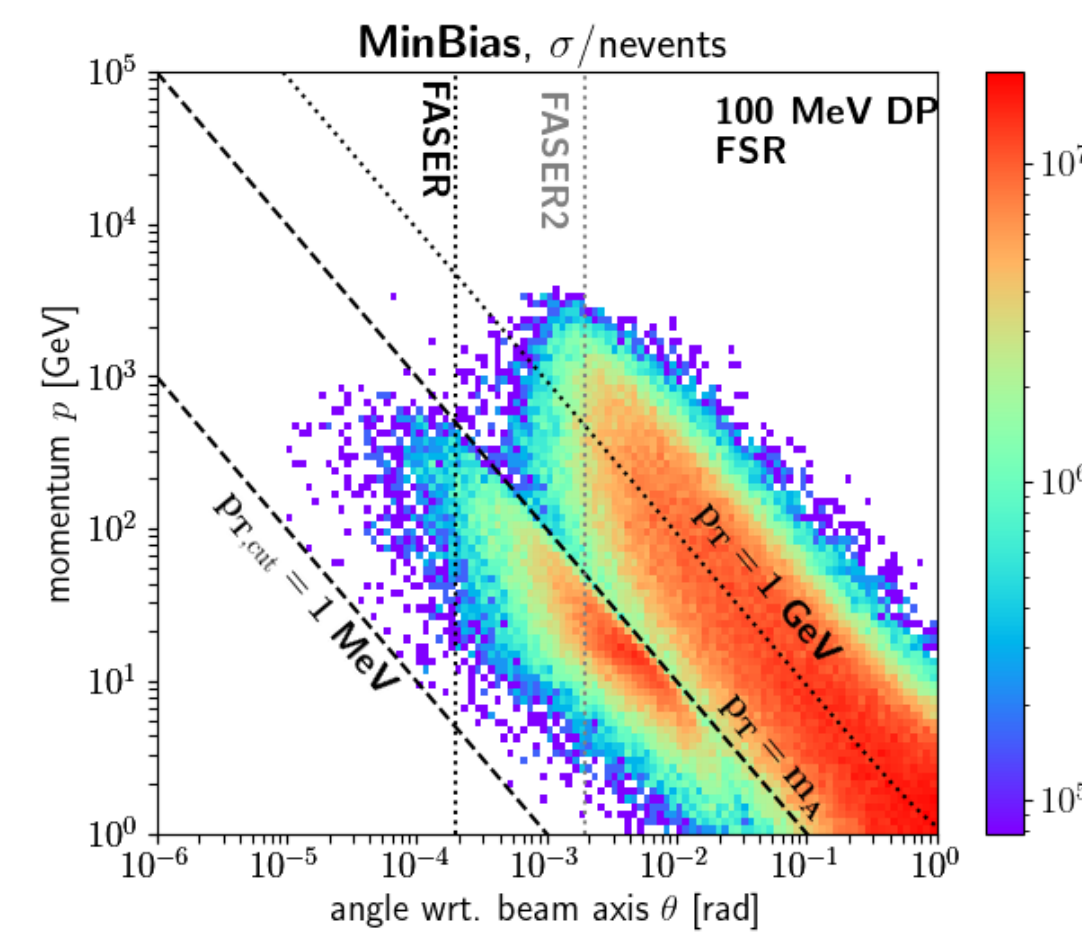
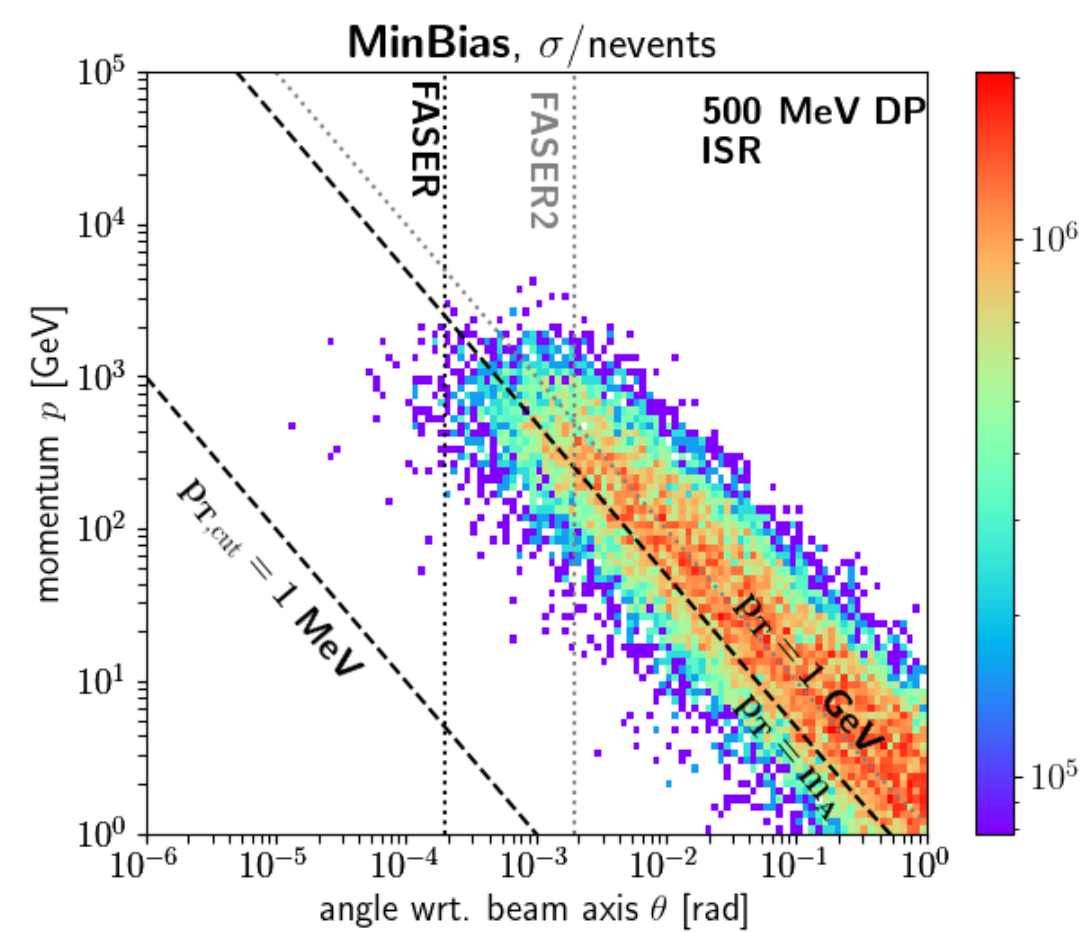
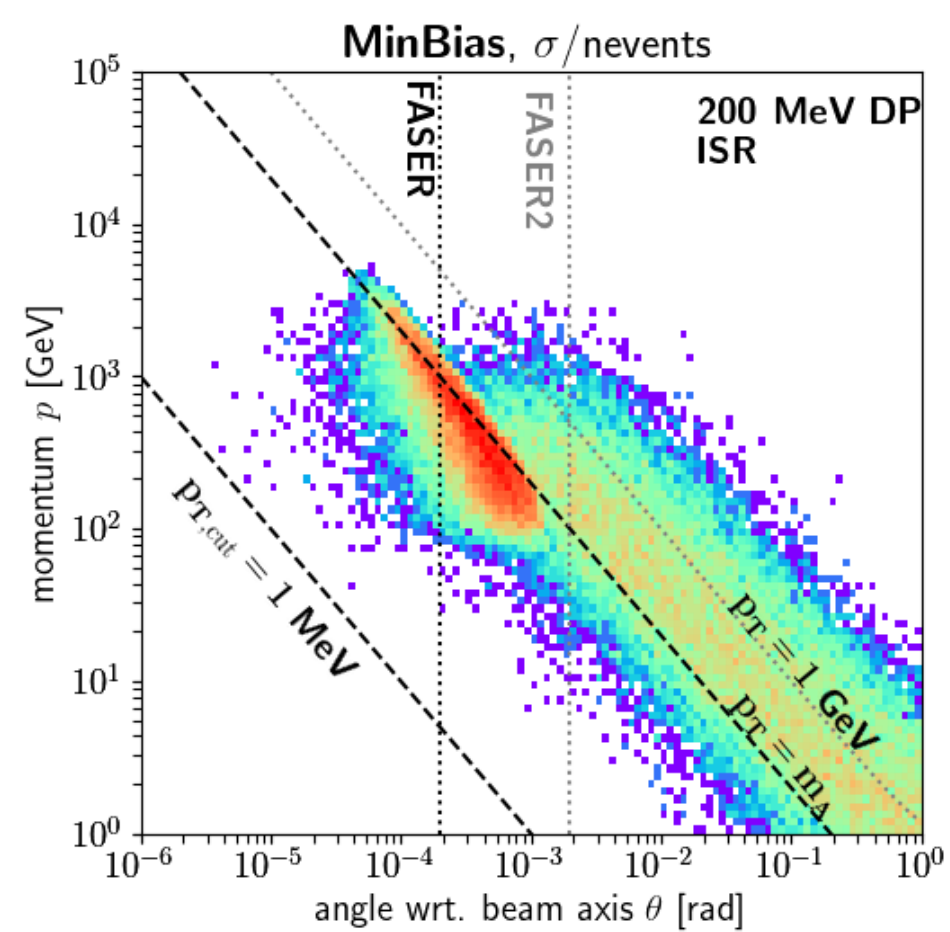


FSR



Validity

- Why does MPI vanish in both cases for higher masses and p_T ?
- Are ISR and Bremsstrahlung related?



ISR

FSR

Conclusions

- Production of dark photons only partially understood!
- future experiments soon will have a more robust description of proton Bremsstrahlung

→ SHiP

→ Forward Physics Facility

+ How relevant are additional production modes?

→ ISR and FSR + are Bremsstrahlung and ISR related?

→ Is there more to do in beam remnants and hadronization?

