

The Dark Universe Introduction

Steven Lowette

Vrije Universiteit Brussel – IIHE

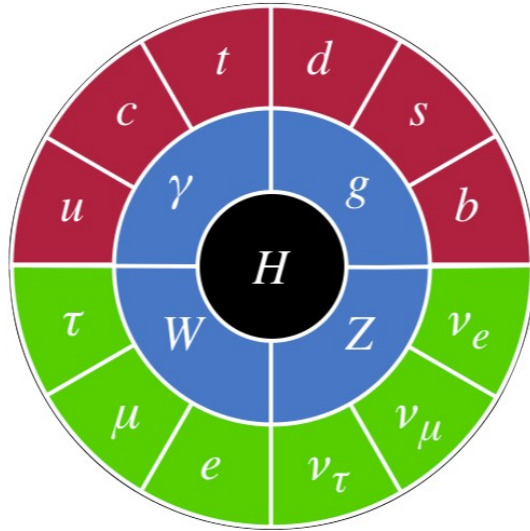
9 March 2023

Symposium – Brussels Town Hall

A decade of discoveries in High Energy Physics



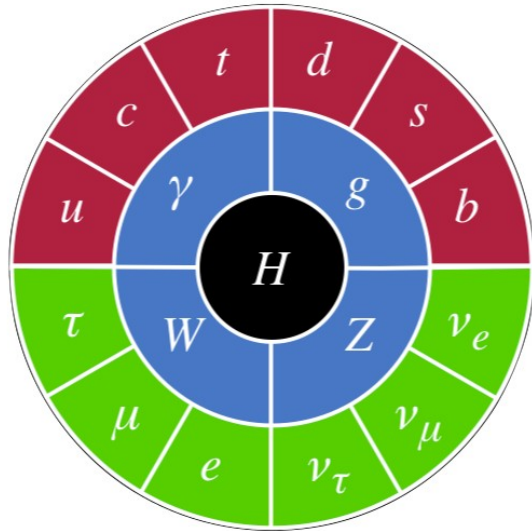
The Dark Sector



~~SM gauge
couplings~~



The Dark Sector



portal coupling

$$H^\dagger H \phi^\dagger \phi$$

H portal

$$F_{\mu\nu} X^{\mu\nu}$$

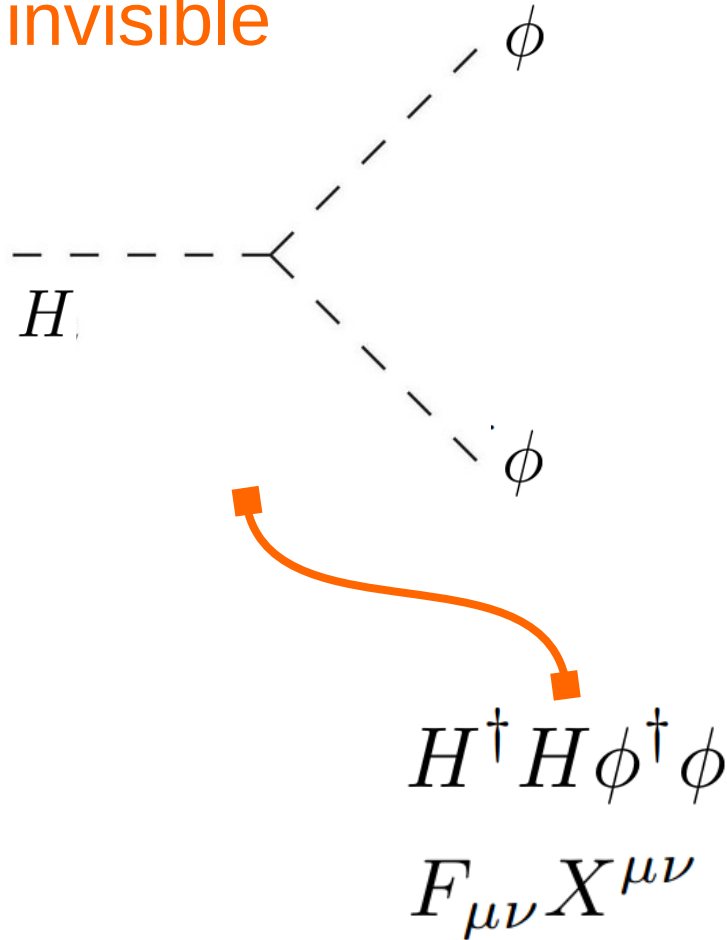
vector portal

$$H L N_R$$

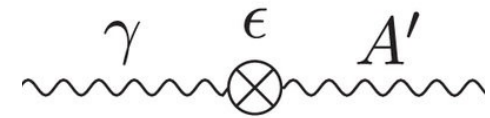
neutrino portal

The Dark Sector: example signatures

H → invisible



dark photon



H portal

vector portal

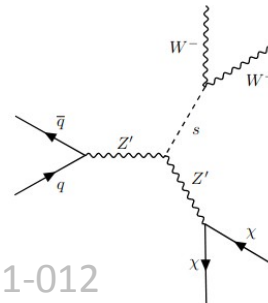
$H^\dagger H \phi^\dagger \phi : H \rightarrow \text{invisible}$

- **$H\phi\phi$ vertex** arises from simplest H portal model
 - H can **decay to invisible particles** ($m_\phi < m_H/2$)
 - would not affect other H property measurements
- **broad LHC search program**
 - in all H production modes
- **incredible precision reached:**
 - **$BR(H \rightarrow \text{invisible}) < 0.15\%$ (0.08%)**

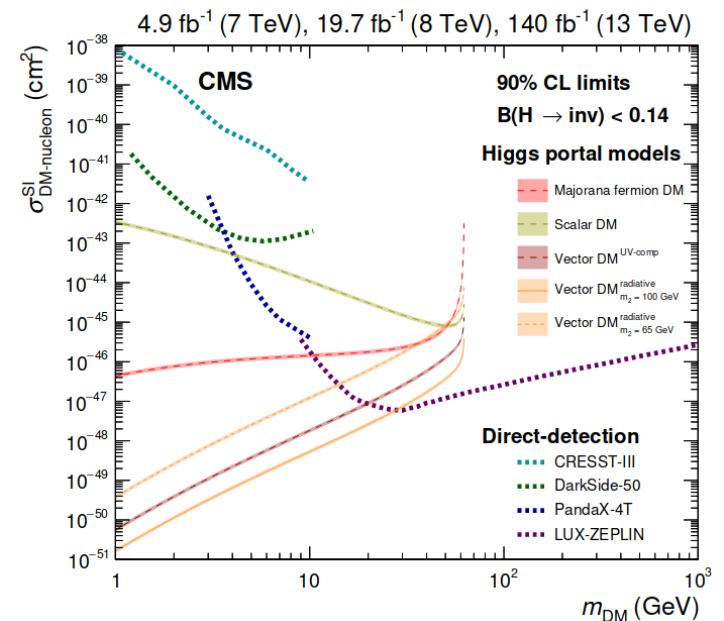
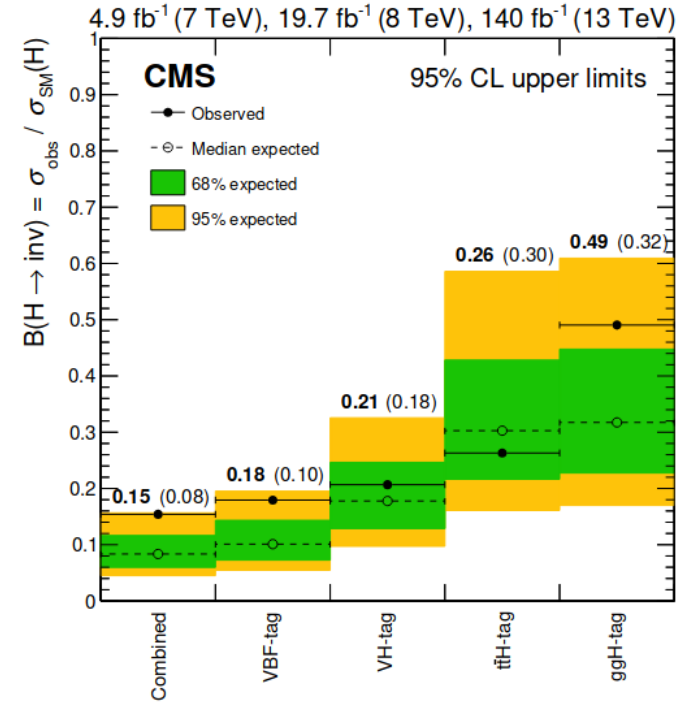
obs. (exp.) 95% C.L. limit

- strong constraints on **dark matter – nucleon** scattering cross sections
- extra scalars may alter the electroweak phase transition
 - **electroweak baryogenesis**

- **also complex dark-H signatures probed**



CMS-PAS-EXO-21-012

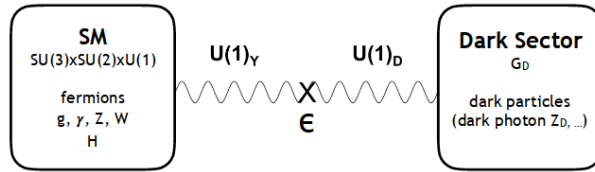


arXiv:2303.01214

arXiv:2303.01214

$F_{\mu\nu} X^{\mu\nu}$: dark photons

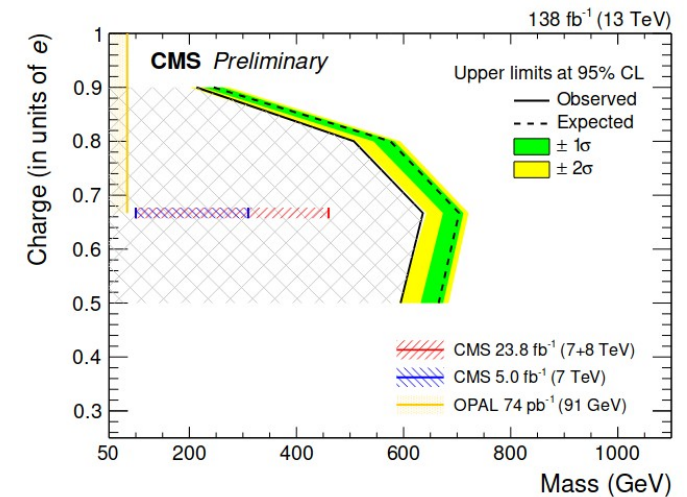
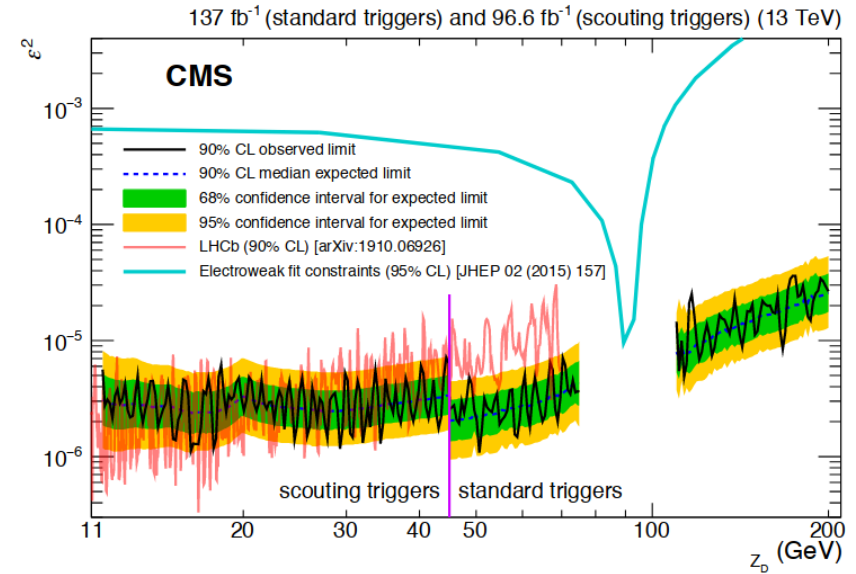
- **dark photon** from new U(1) symmetry
 - vector portal: **kinetic mixing** with regular photon



- **CMS search for $\mu^+\mu^-$ resonance**
 - first time using special triggers to access low mass



- **millicharges** ~ mixing ϵ
 - unique search for **fractionally charged particles**
 - **milliQan**: dedicated new experiment searching for millicharged particles



Next frontiers: long-lived particles

- new **long-lived particles** are a prediction in many dark sector models

arXiv:1903.04497

- **rich experimental frontier**

- new simulations, trigger developments, reconstruction, unusual backgrounds,...

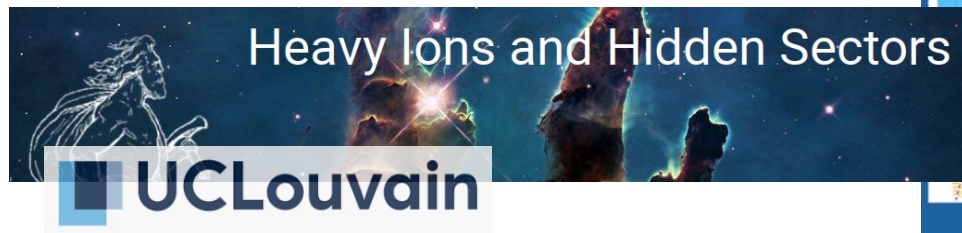
Journal of Physics G: Nuclear and Particle Physics

MAJOR REPORT • OPEN ACCESS

Searching for long-lived particles beyond the Standard Model at the Large Hadron Collider

- **activities**

- displaced leptons from compressed spectrum (JHEP 11 (2020) 112)
- low-mass H-portal scalar in $H \rightarrow ss \rightarrow \mu^+\mu^- K^+K^-$ (JHEP 01 (2020) 115, ongoing)
- strongly interacting dark matter (JHEP 11 (2015) 108, Eur.Phys.J.C 82 (2022) 3, 213)
- 6-quark dark matter (arXiv:1708.08951, ongoing)
- freezin dark matter and cosmology (JHEP 09 (2018) 037)
- new physics with heavy ions (arXiv:1812.07688, PRL 124 (2020) 8, 081801)
- ...



The Dark Universe

11:00

The Dark Universe -- Chair: Steven Lowette

Speakers: Aqeel Ahmed (Vrije Universiteit Brussel) , Iason Baldes (Paris) , Steven Lowette (Vrije Universiteit Brussel)

Introduction

Speaker: Steven Lowette (Vrije Universiteit Brussel)

The BEH mechanism at high temperature: a link to matter production ¶

Speaker: Iason Baldes

Dark Side of the Universe

Speaker: Aqeel Ahmed (Vrije Universiteit Brussel)

Aqeel Ahmed

- Max-Planck-Institute für Kernphysik – Heidelberg
- with the VUB pheno group from 2018-2020
- collider signatures of naturalness → dark matter

Iason Baldes

- Ecole Normale Supérieure – Paris
- with the ULB pheno group from 2018-2022
- link to cosmology → electroweak baryogenesis

DARK SIDE OF THE UNIVERSE

Aqeel Ahmed

Max-Planck-Institut für Kernphysik — Heidelberg



A Decade of Discoveries in High-Energy Physics, Brussels — March 09, 2023

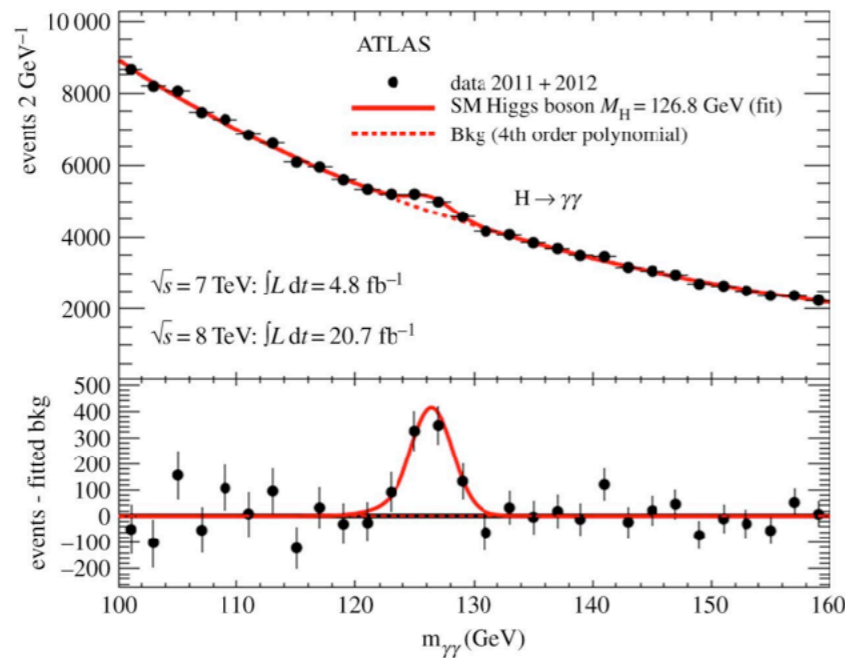
former postdoc (2018–2020) at  **VUB** VRIJE
UNIVERSITEIT
BRUSSEL



A DECADE OF DISCOVERIES

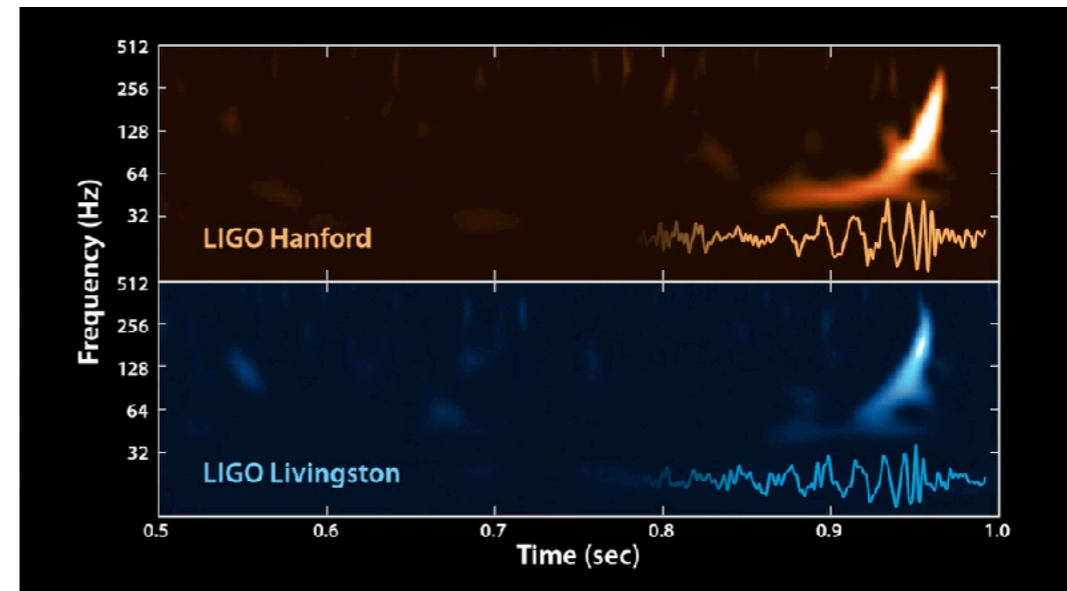
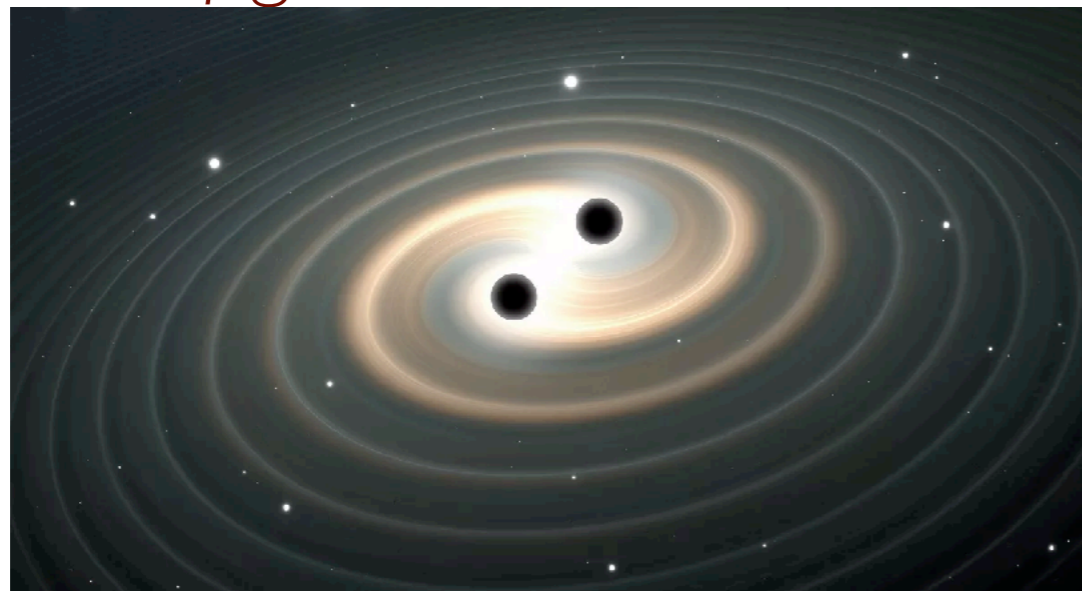
Discovery of the H boson at CERN

morning session



Detection of gravitational waves at LIGO/VIRGO

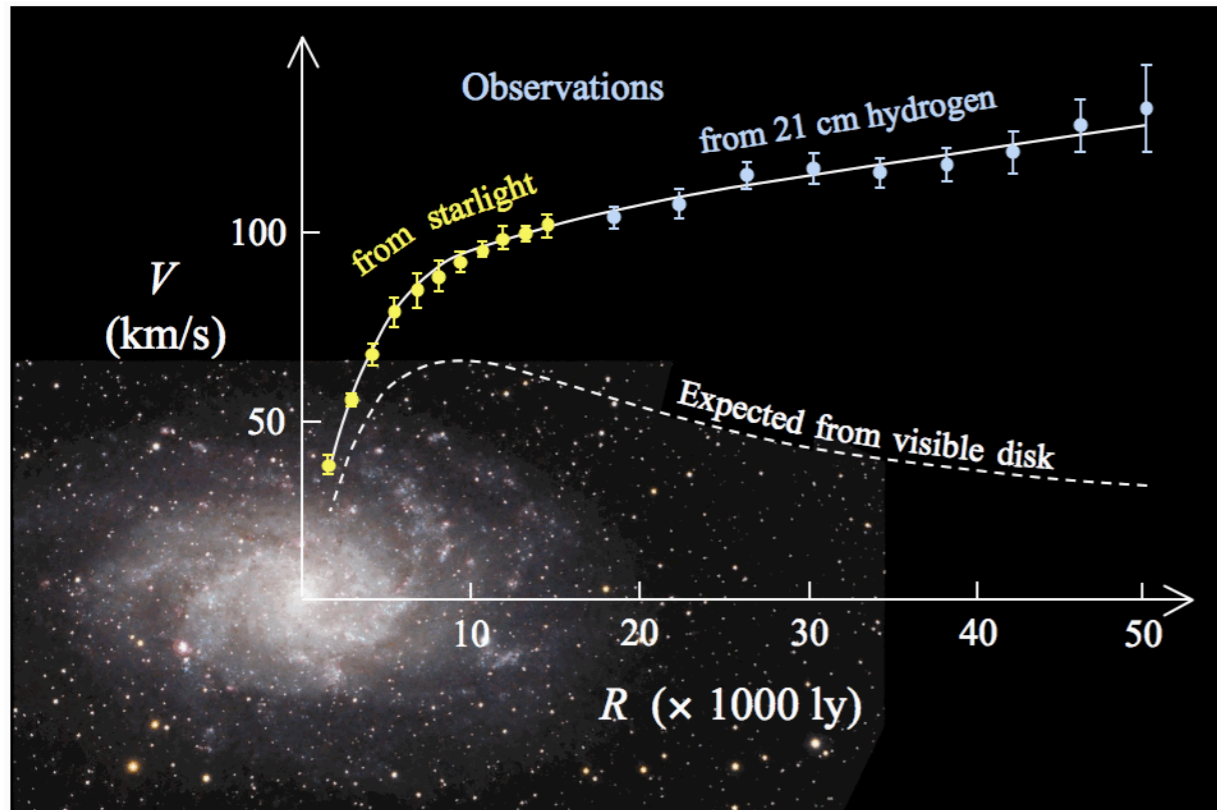
see also Baldes' talk and afternoon session



Dark matter: No breakthrough discovery but enormous progress on theoretical and experimental fronts.

DARK MATTER: EVIDENCE

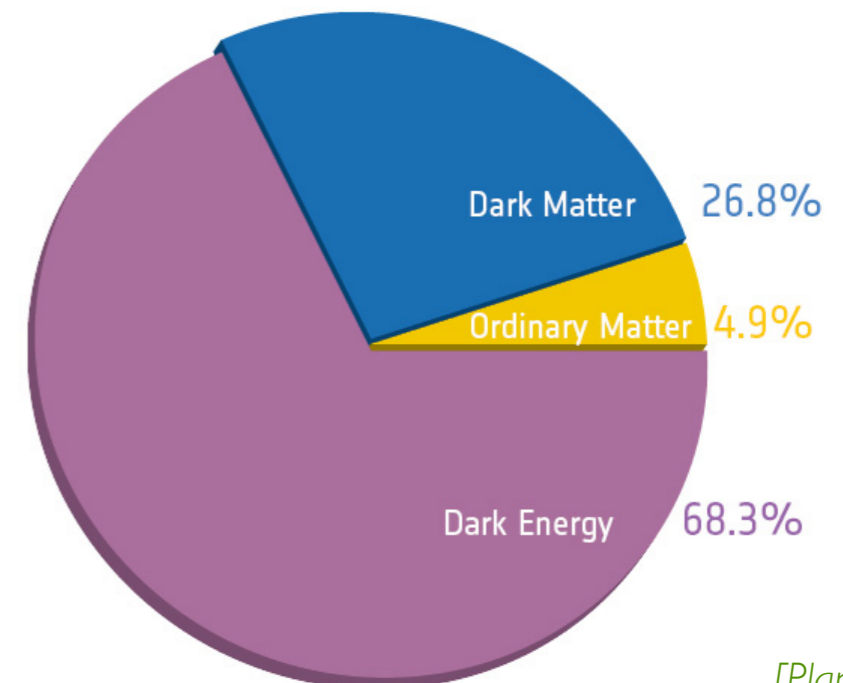
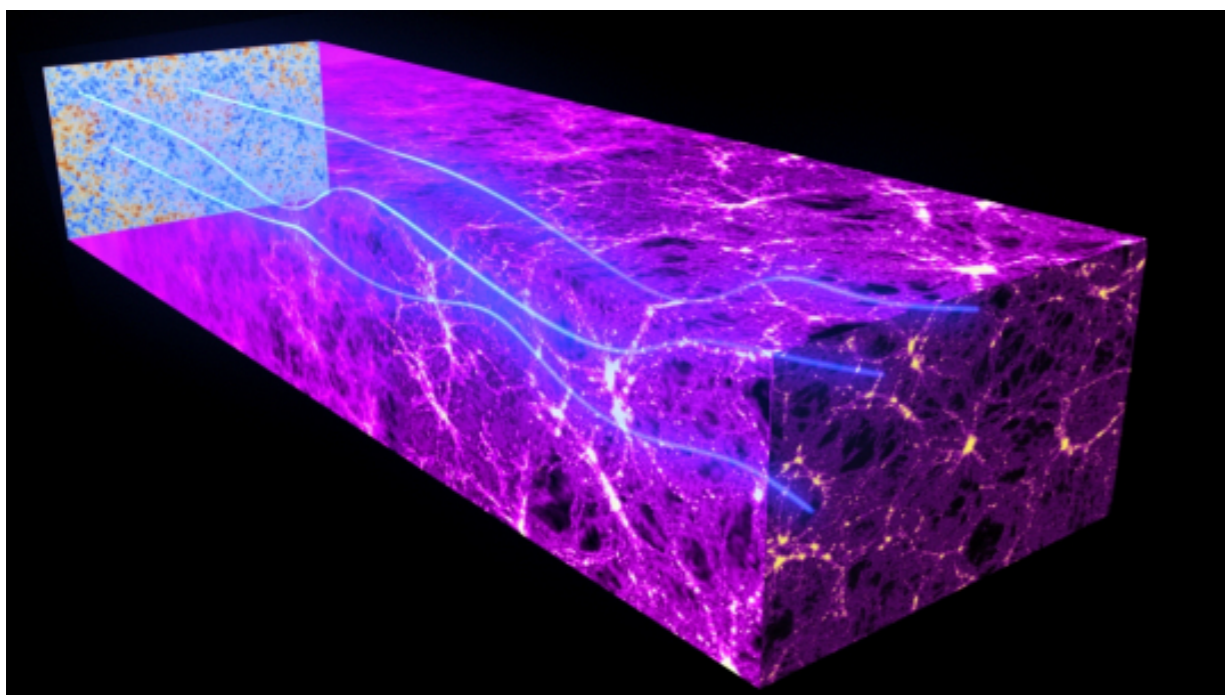
Galactic rotation curves



Cluster collisions



Gravitational lensing of CMB



[Planck Satellite]



DARK MATTER: WHAT WE KNOW

■ Dark matter is:

- *Massive* (as it interacts with gravity)

- *Stable* (at cosmological scales)

- (electrically) *Neutral*

- *Cold* (non-relativistic)

- *Clumps to form halos*

- *Relic abundance* $\Omega_{\text{DM}} h^2 = 0.12(12)$ [Planck, 2018]

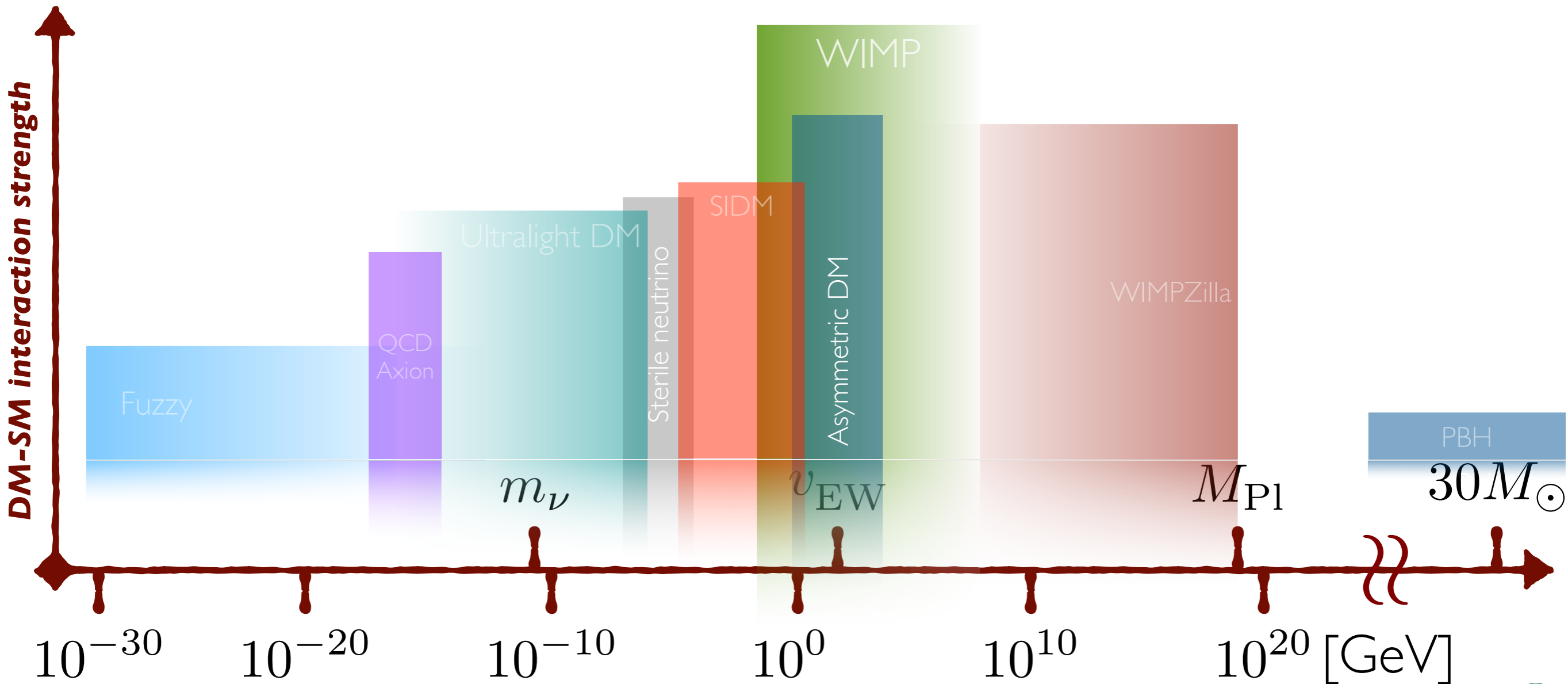
- No evidence of non-gravitational interactions b/w DM and SM

WHAT DARK MATTER CAN BE?

■ DM mass:

■ particle state, $m_{\text{DM}} \in [10^{-30}, 10^{19}] \text{ GeV}$

■ composite state, e.g. Primordial Blackholes, $m_{\text{DM}} \sim \mathcal{O}(10) M_{\odot}$

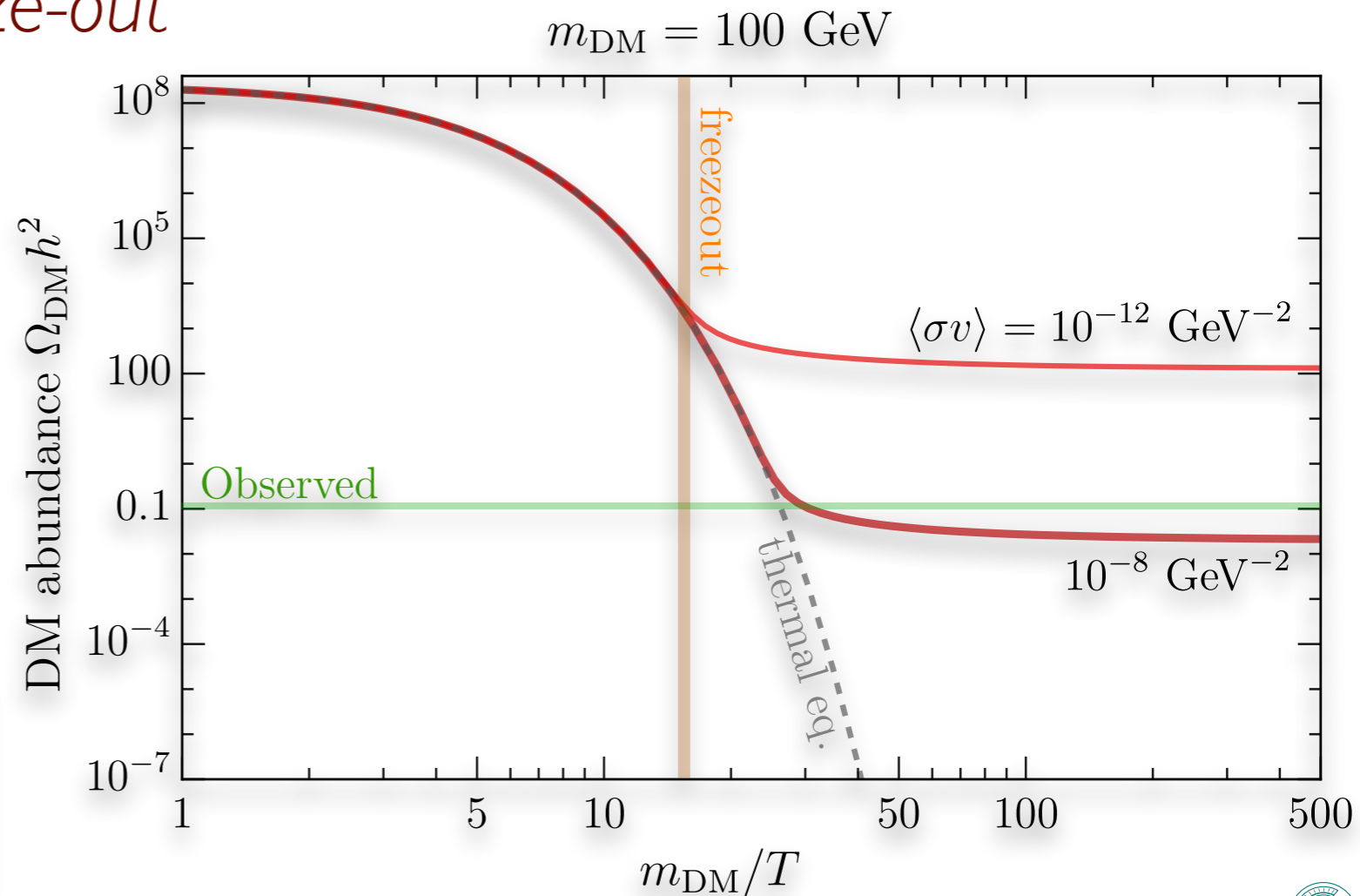
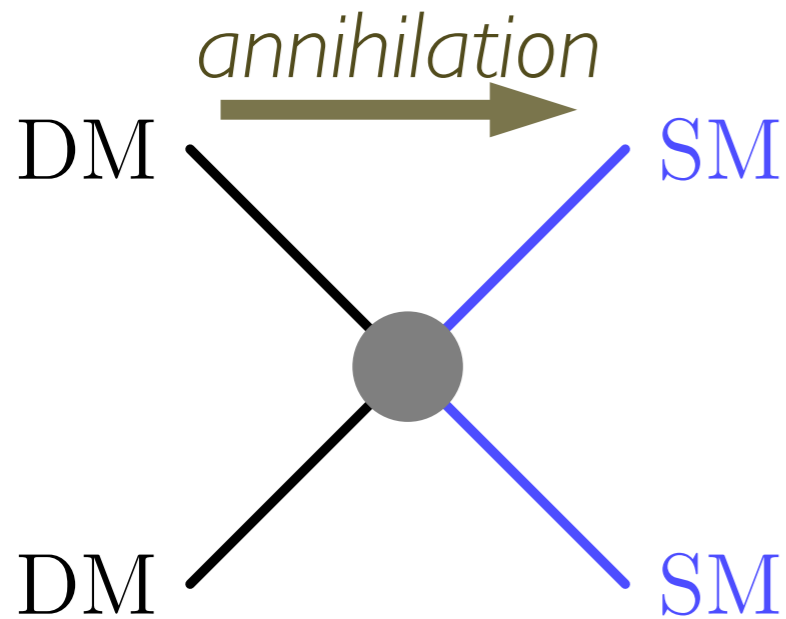


PLAN

- I. WIMP DM and its current status
- II. DM model building beyond the WIMP paradigm:
 - Theoretical ideas
 - Experimental prospects
- III. Conclusions

WIMP DM

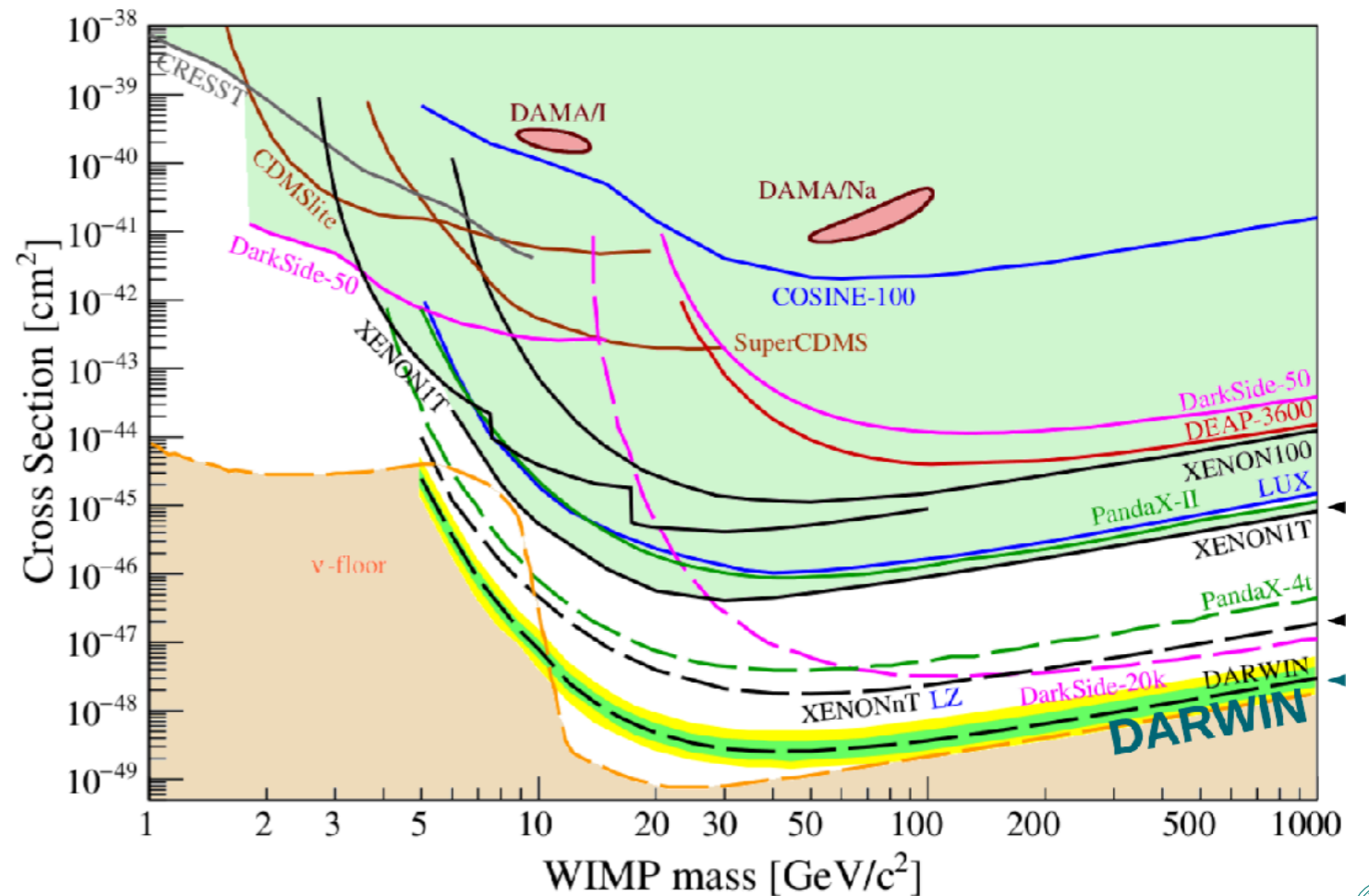
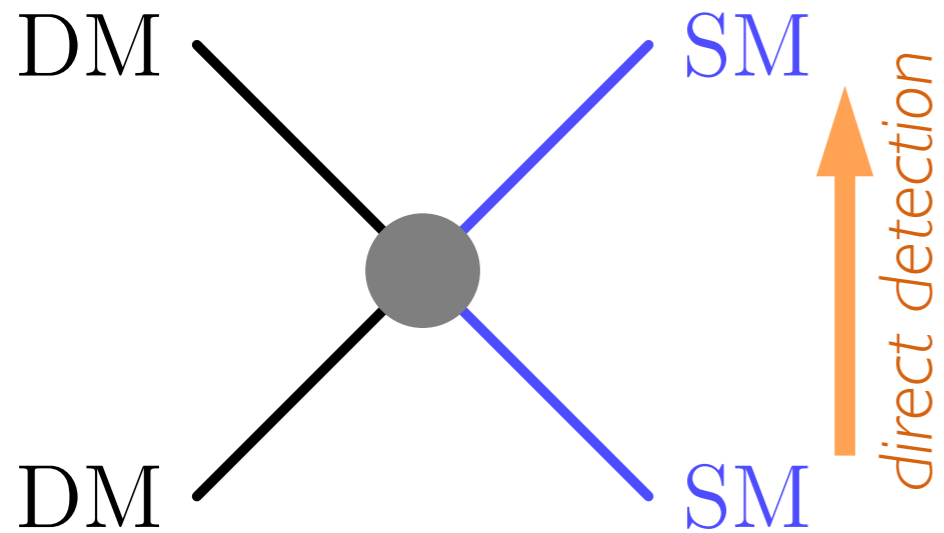
- DM mass $m_{\text{DM}} \sim m_{\text{EW}} \sim \mathcal{O}(100) \text{ GeV}$
- DM-SM interaction strength $\sim G_F$
- DM-SM are in **thermal equilibrium** at high temperatures
- Production via **thermal freeze-out**



$$\Omega_{\text{DM}} h^2 \sim 0.12 \left(\frac{m_{\text{DM}}}{20 T} \right) \frac{10^{-8} \text{ GeV}^{-2}}{\langle \sigma v \rangle}$$

WIMP DETECTION

- WIMP DM at experiments:
 - Direct detection, e.g. *XENON*, *LZ*, ...



WIMP DETECTION

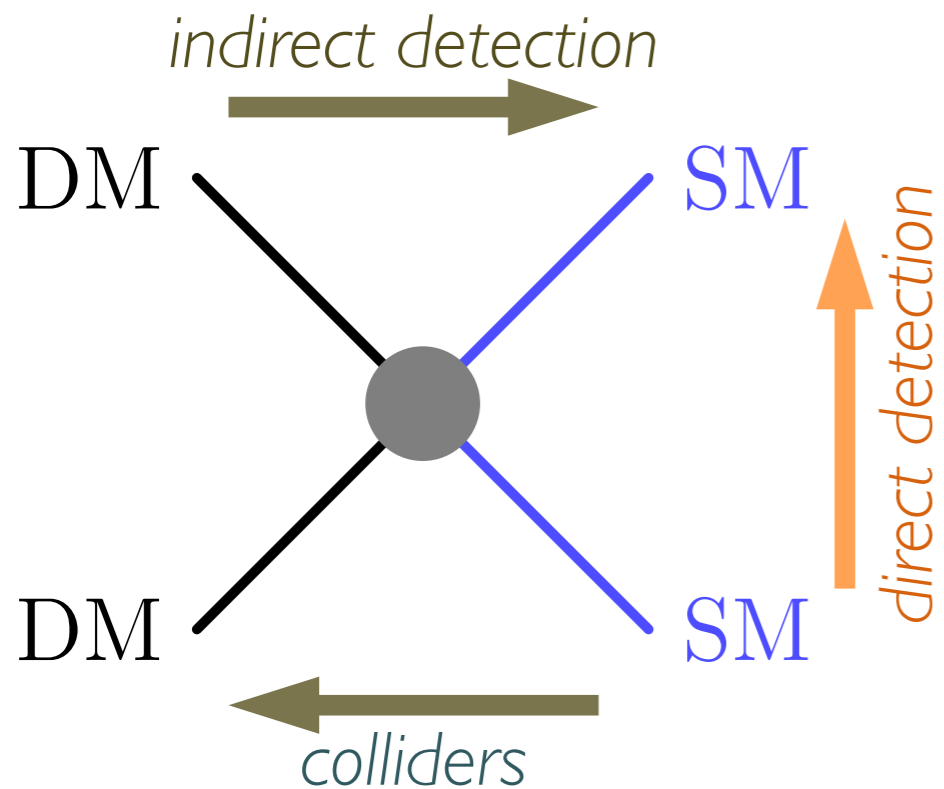
- WIMP DM at experiments:

- Direct detection, e.g. *XENON, LZ, ...*

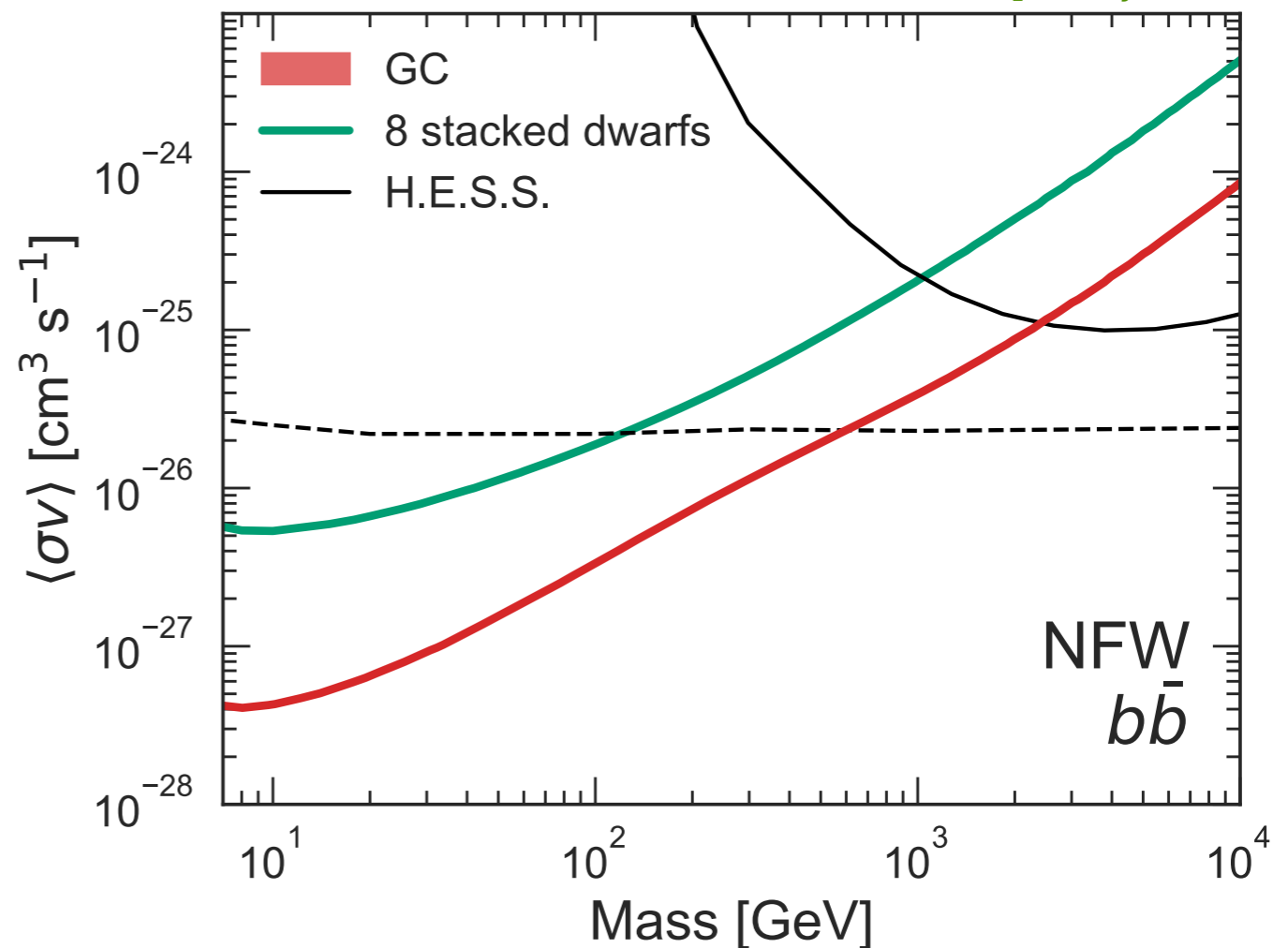
- *Indirect detection, e.g. Fermi gamma-ray space telescope, AMS, ...*

- *Collider experiments, e.g. LHC, future colliders*

see Steven's slides



[Abazajian et al., 2020]



FREEZE-IN DM

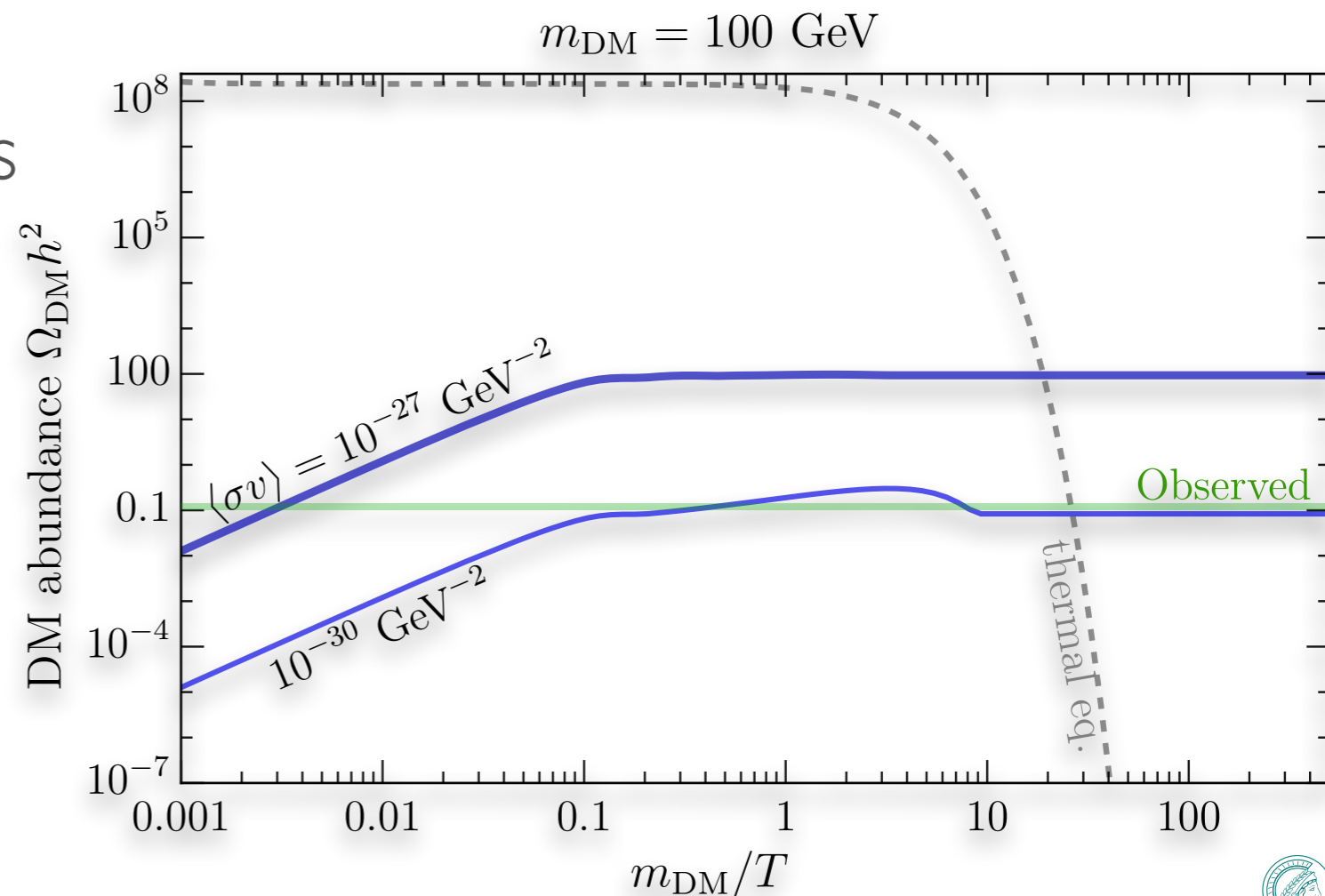
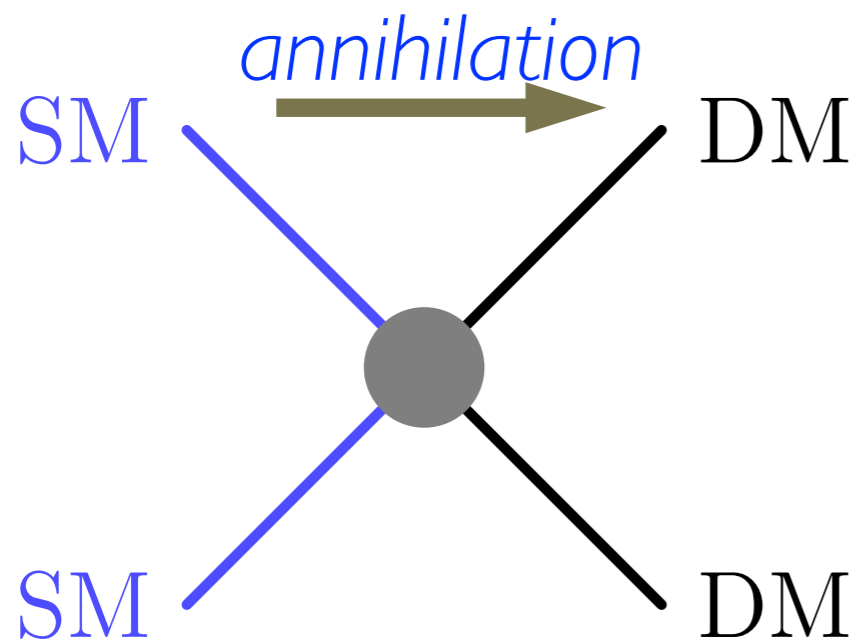
■ Freeze-in DM is one of the main alternatives to WIMP DM:

■ *Out of thermal equilibrium with the SM*

■ Initial abundance is negligible

■ Produced via SM annihilations

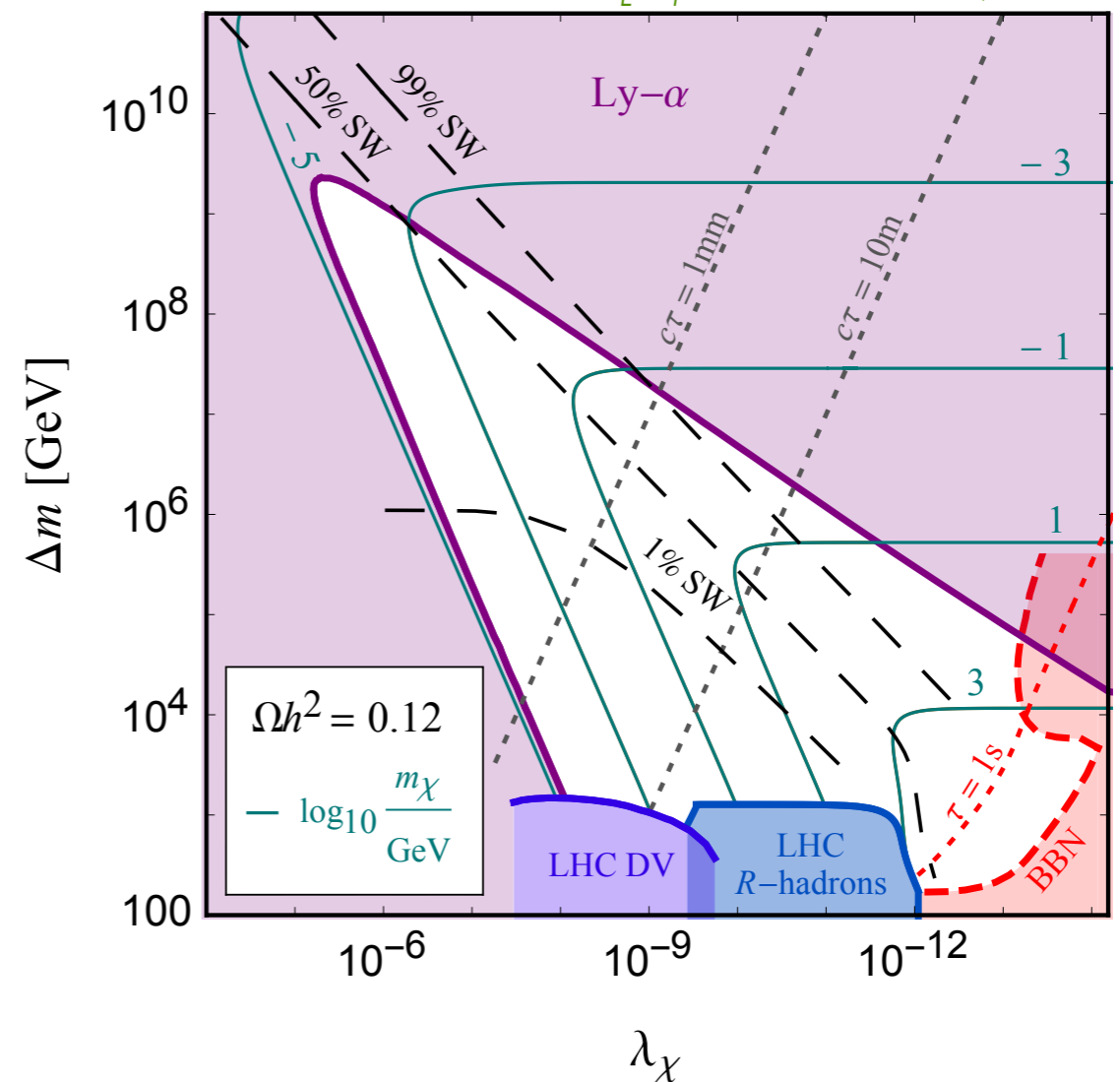
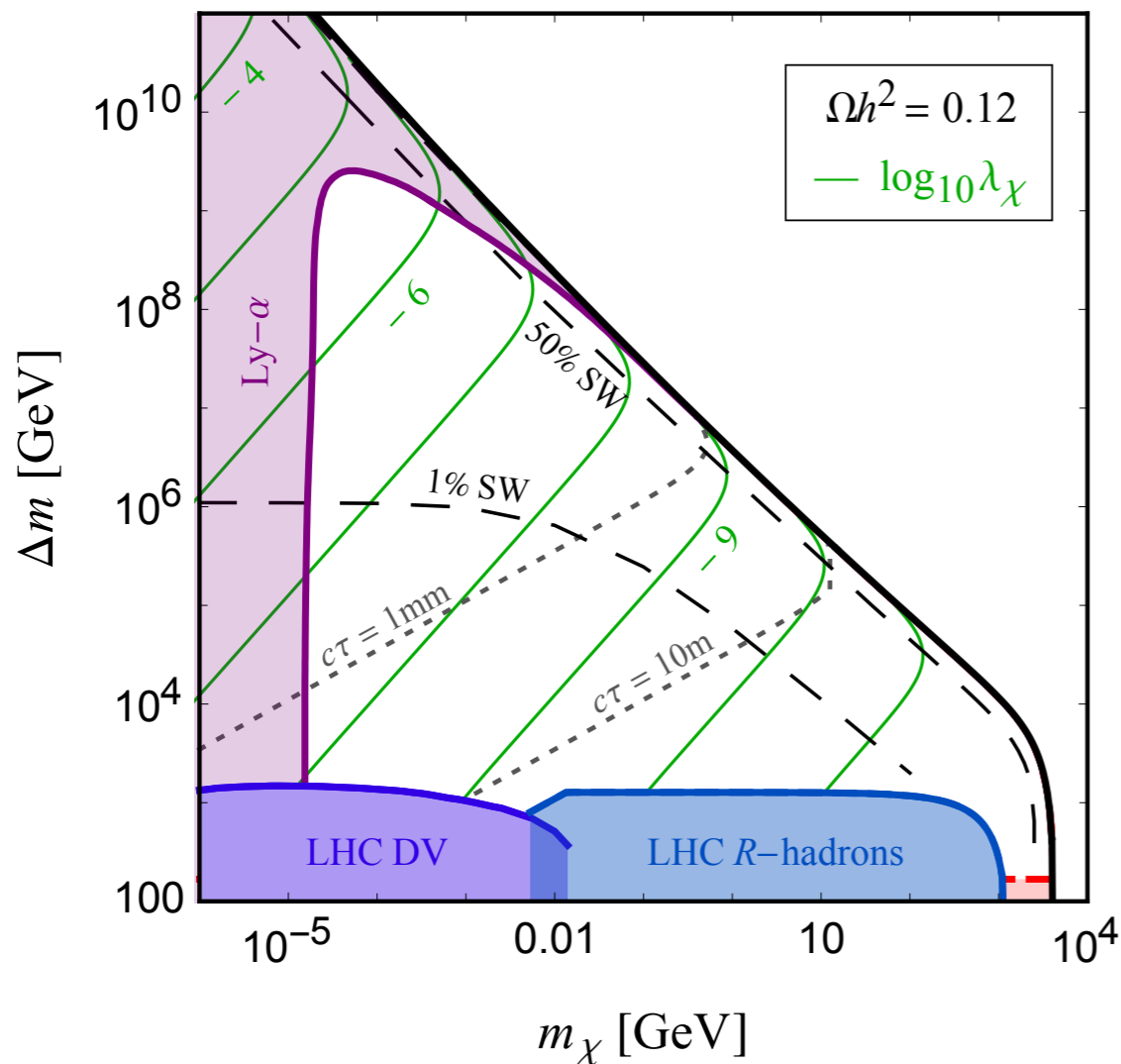
■ DM-SM feeble interactions



FREEZE-IN SIGNATURES

- Freeze-in DM can have its imprints in
 - Astrophysics and *cosmological data*;
 - Collider signals, e.g. long-lived particles through mediators/portals

[Lopez-Honorez et al., 2021-22,+...]



MISALIGNMENT MECHANISM: AXION DM

- Axions solve the strong CP problem

$$\mathcal{L} = \left(\frac{a}{f_a} - \bar{\theta} \right) \frac{\alpha_s}{8\pi} G^{\mu\nu a} \tilde{G}_{\mu\nu}^a$$

$$m_a^2 = \Lambda_{\text{QCD}}^4 / f_a^2$$

- Axion DM can be produced via the misalignment mechanism and behave as cold DM

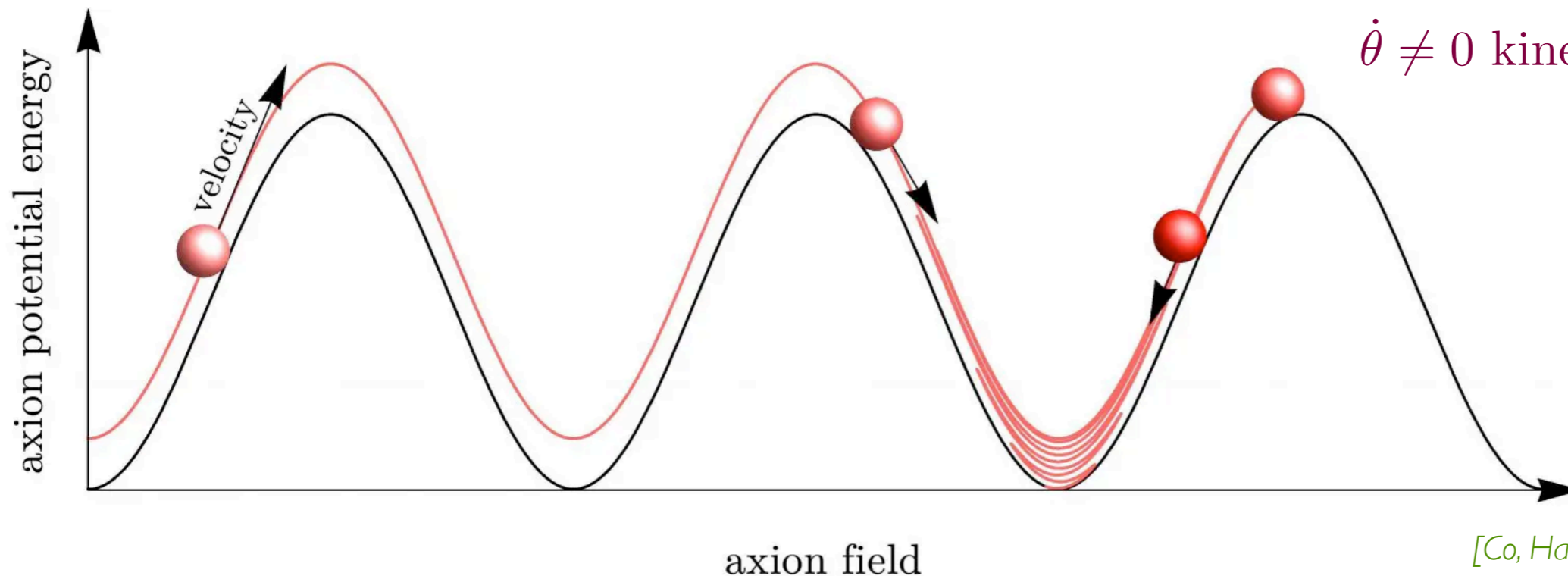
$$\Omega_a h^2 \sim 0.12 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{7/6} \langle \theta_i^2 \rangle$$

$$\theta \equiv a / f_a$$

- Several recent developments for axion DM production, e.g. kinetic misalignment, ...

$\dot{\theta} = 0$ misalignment

$\dot{\theta} \neq 0$ kinetic misalignment

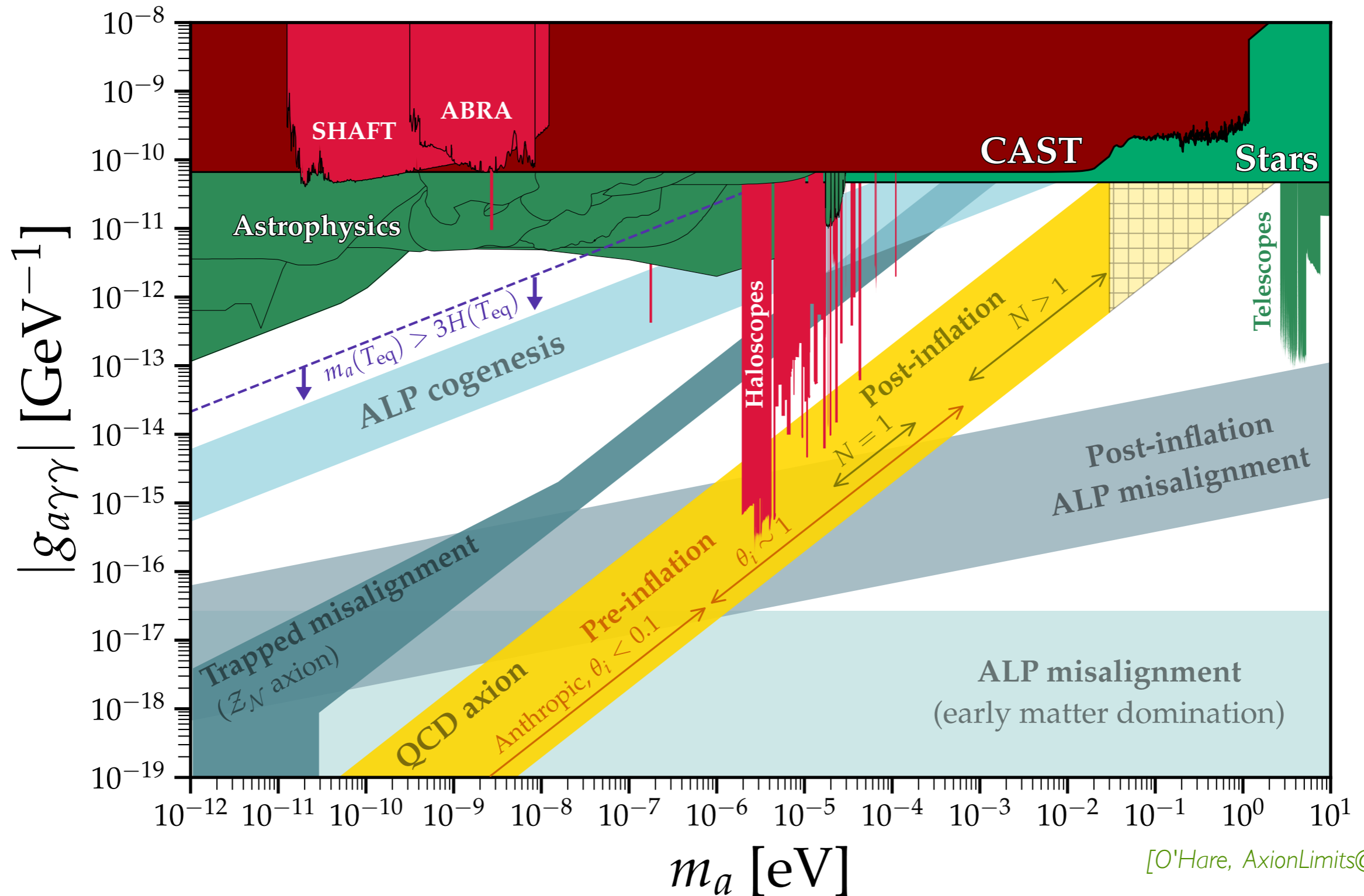


[Co, Hall, Harigaya, 2020,+...]

AXION DM SEARCHES

$$\Omega_a h^2 \sim 0.12 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{7/6} \langle \theta_i^2 \rangle$$

$$g_{a\gamma\gamma} \sim \mathcal{O}(1) \frac{\alpha_{\text{EM}}}{2\pi f_a}$$



[O'Hare, AxionLimits@github]

GRAVITATIONAL DARK MATTER

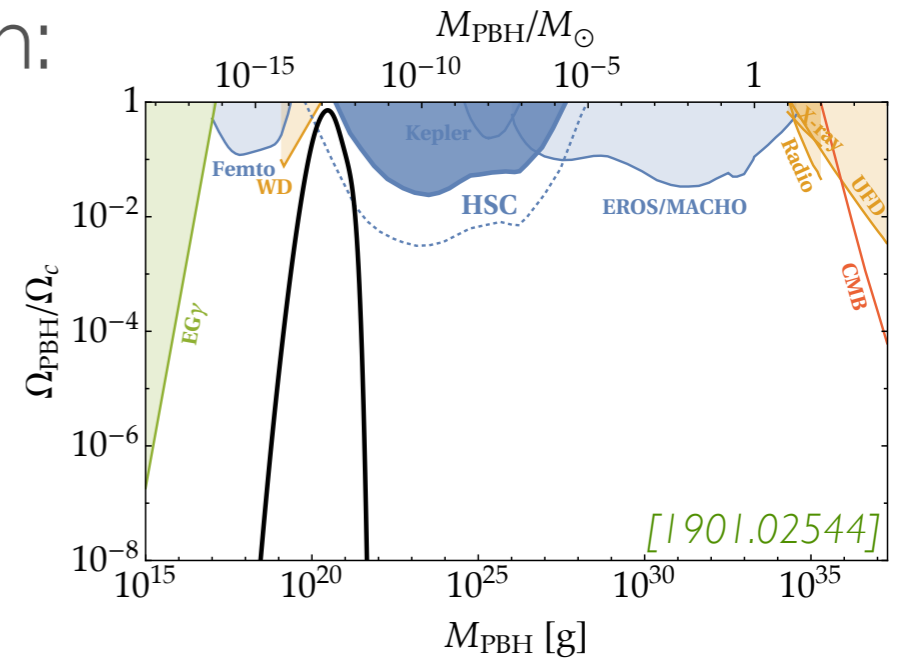
■ Nightmare scenario: DM only interacts gravitationally!

■ Mechanisms for gravitational DM production:

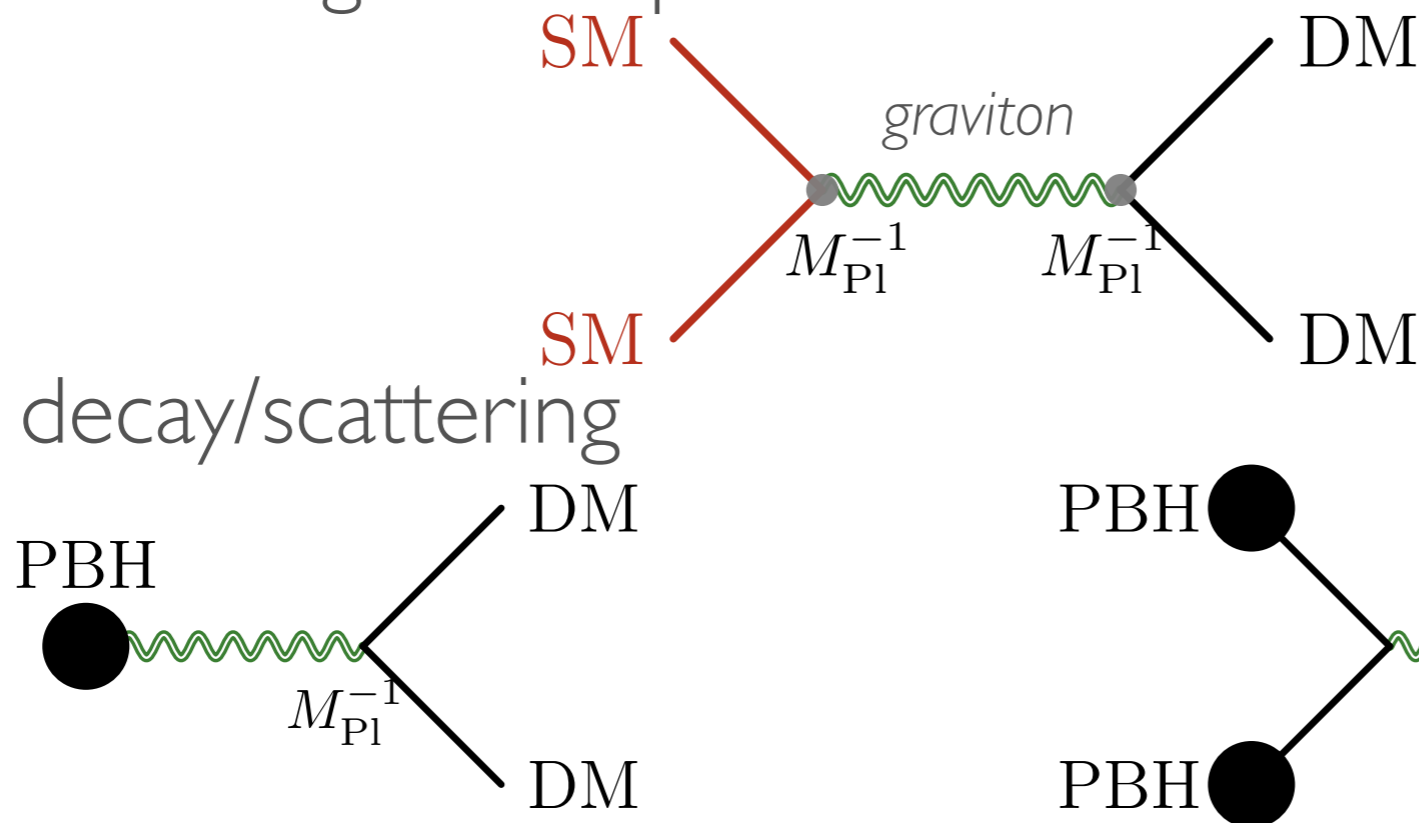
■ Primordial black holes (PBH) as DM

■ Freeze-in via graviton portal

■ PBH decay/scattering



[Garny, Sandora, Sloth 2015; A.A., Grzadkowski, Socha 2019+...]



[Hooper et al, 2019; Baldes et al. 2020+...]

[Bernal, Zapata 2020; +...]



GRAVITATIONAL PARTICLE PRODUCTION

- DM from gravitational particle production due to quantum fluctuations

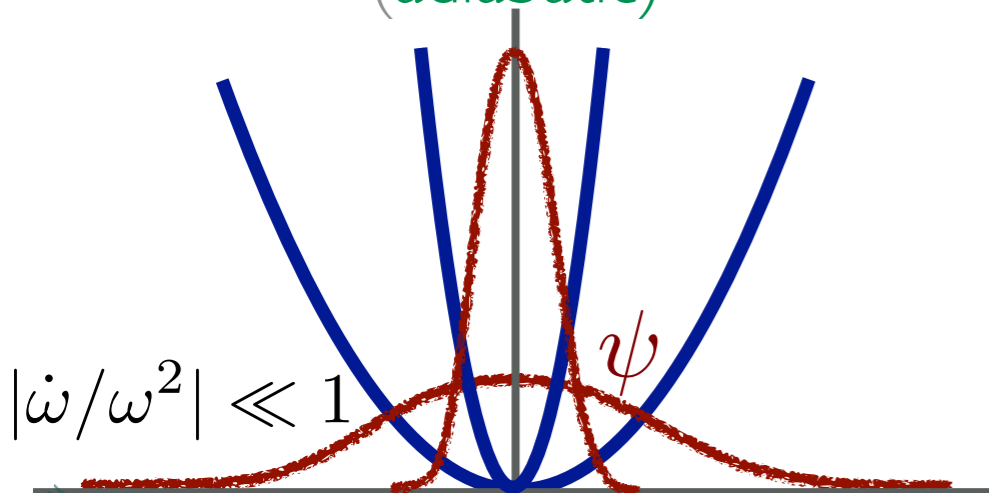
[Schrodinger 1939; Parker 1969; Zeldovich, Starobinsky 1972; Ford 1987; +...]

- Rapid expansion in the early universe (during and after inflation) leads to non-adiabatic production of quantum excitations from the background fields.

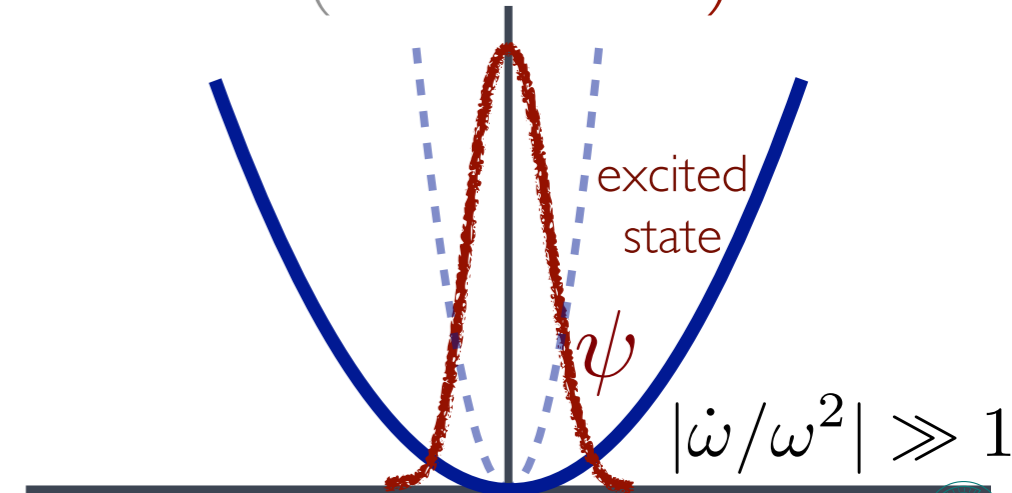
- Analogy: 1D harmonic oscillator with time-dependent frequency

$$\ddot{x}(t) + \omega^2(t)x(t) = 0$$

slow frequency change
(adiabatic)



fast frequency change
(non-adiabatic)



GRAVITATIONAL DARK MATTER

- Why gravitational production of dark matter?
 - An unavoidable consequence of the rapid expansion of the Universe
 - DM number density is tied with inflationary observables; Hubble scale H_I , reheating temperature T_{rh} , inflaton potential $V(\phi)$, etc.
 - Besides gravity, no other interaction is required
 - Gravitational DM production is sensitive to inflationary and reheating dynamics so it may be a probe into early universe cosmology.

*[Graham, Mardon, Rajendran 2015,
A.A., Grzadkowski, Socha, 2020+...]*

EARLY UNIVERSE COSMOLOGY

[A.A., Grzadkowski, Socha, 2020]

■ Inflation: de Sitter solution

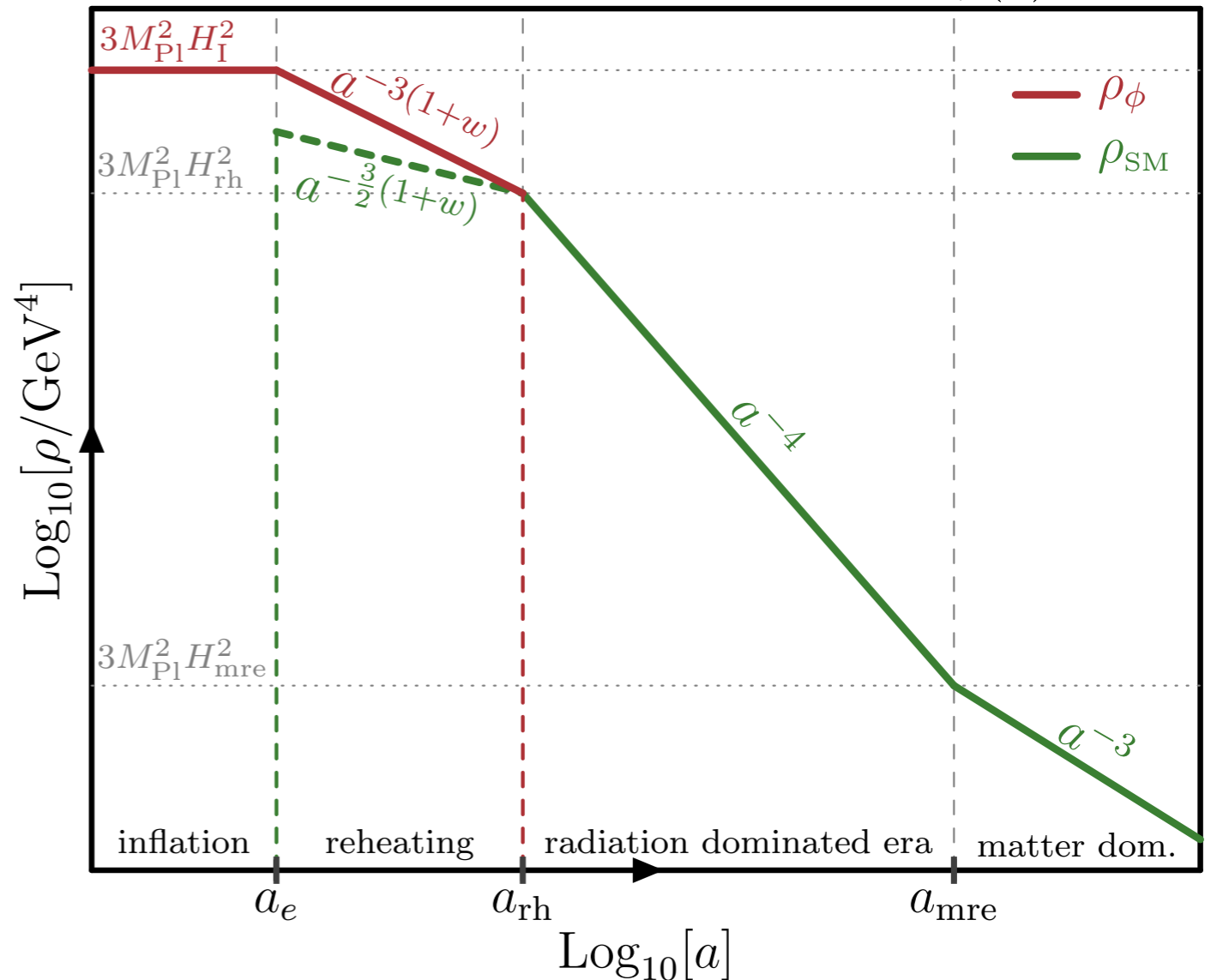
■ Reheating: non-standard cosmology

$$w = \rho_\phi / p_\phi : [-1/3, 1]$$

■ Standard cosmological evolution after reheating

$$\rho(a) = 3M_{\text{Pl}}^2 H^2(a)$$

Evolution of the energy density $\rho(a)$



$$H(a) = \begin{cases} H_I, & a \leq a_e \\ H_I \left(\frac{a_e}{a} \right)^{\frac{3(1+w)}{2}}, & a_e < a \leq a_{\text{rh}} \\ H_{\text{rh}} \left(\frac{a_{\text{rh}}}{a} \right)^2, & a_{\text{rh}} < a \end{cases}$$

$$\gamma \equiv \sqrt{\frac{H_{\text{rh}}}{H_I}}$$

Reheating

$$10^{-18} \leq \gamma \leq 1$$

$T_{\text{rh}} \gtrsim 10 \text{ MeV}$

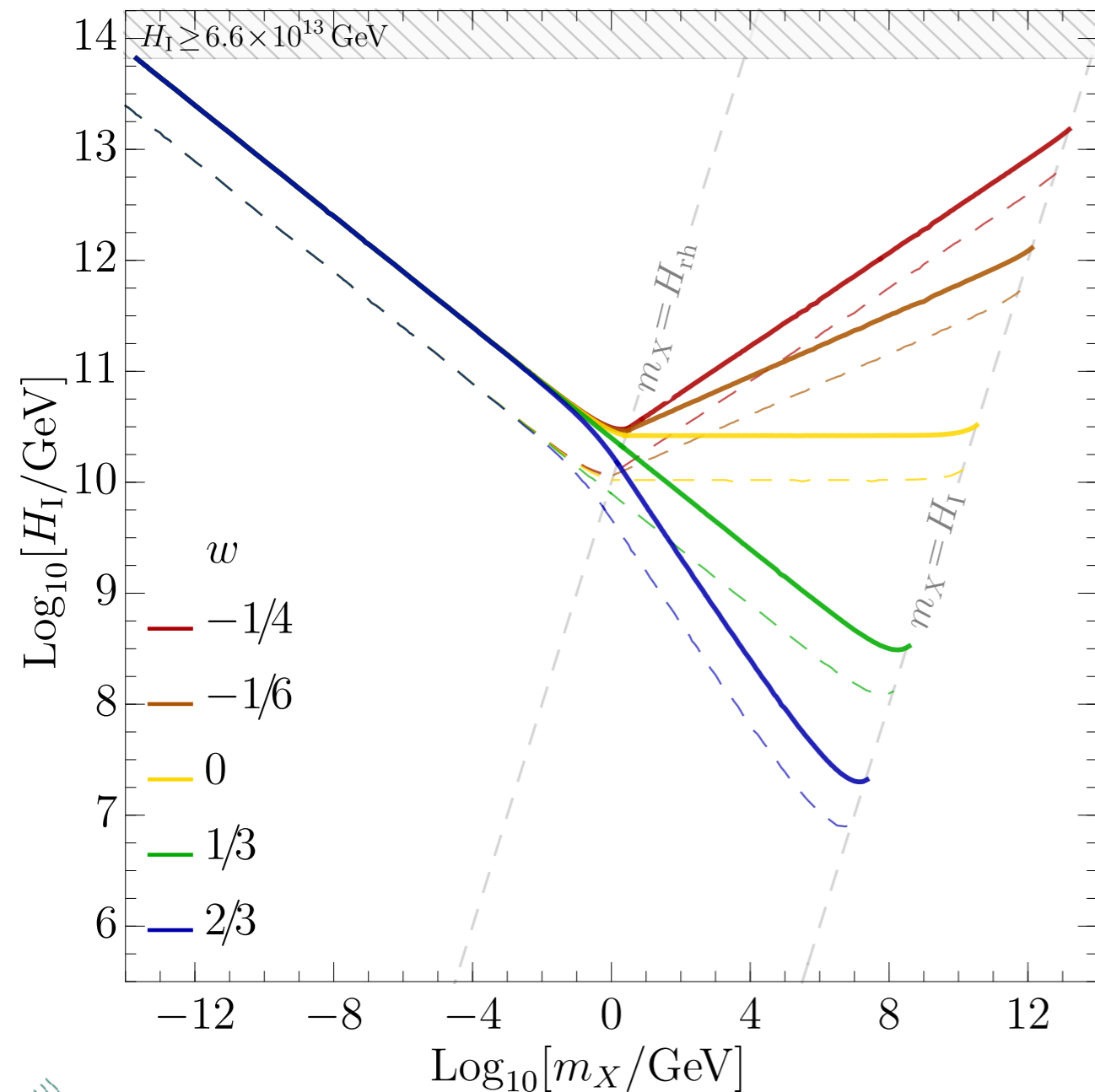


DARK PHOTON DM

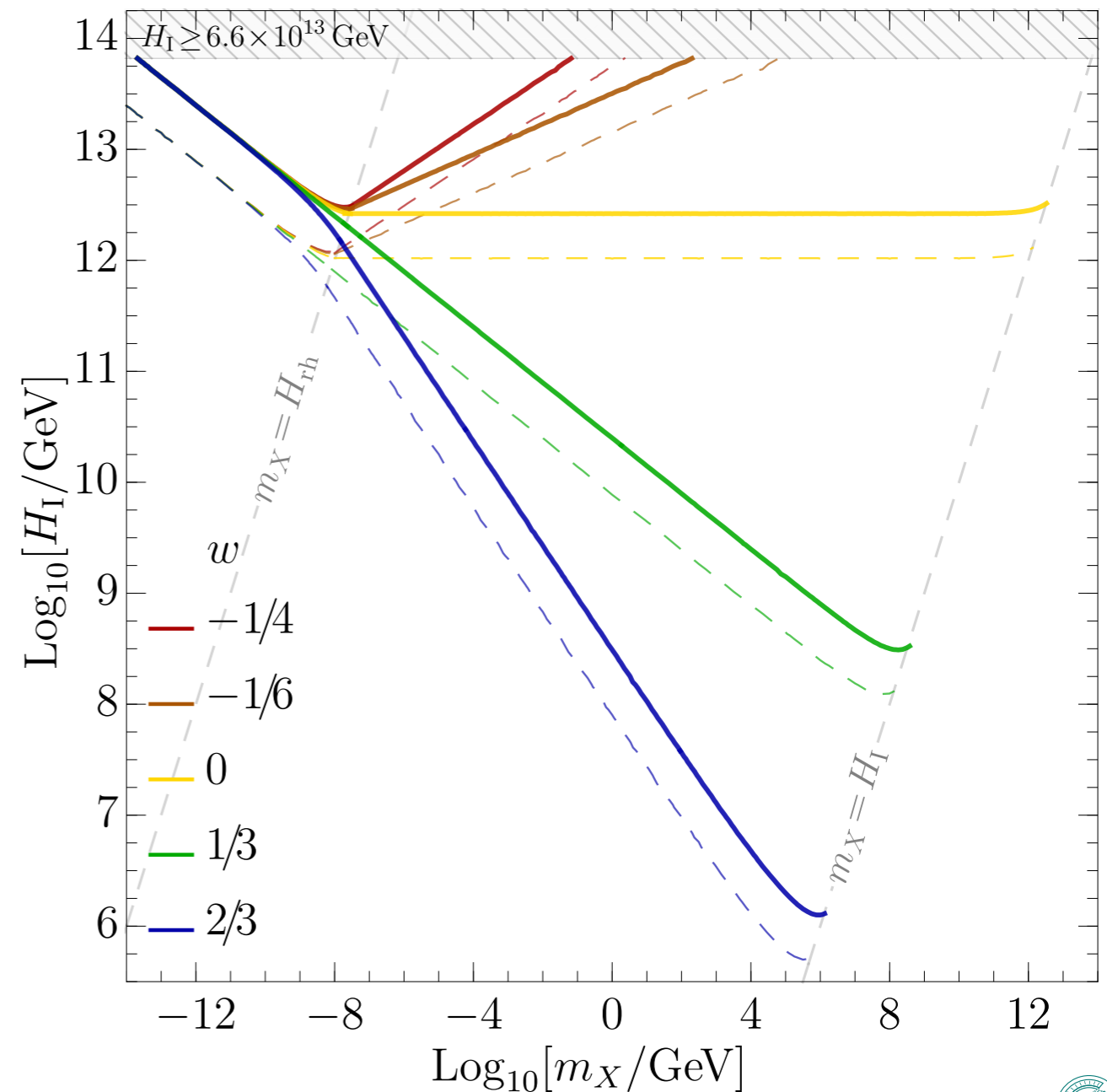
$$\Omega_X h^2 \approx 0.12 \times \begin{cases} \left(\frac{m_X}{0.33 \text{ GeV}} \right)^{\frac{2w}{1+w}} \left(\frac{H_I}{3.3 \times 10^{10} \text{ GeV}} \right)^{\frac{5+w}{2(1+w)}} \left(\frac{10^{-5}}{\gamma} \right)^{\frac{3w-1}{1+w}}, & H_{\text{rh}} \leq m_X < H_I \\ \left(\frac{m_X}{2 \times 10^{-14} \text{ GeV}} \right)^{1/2} \left(\frac{H_I}{6.6 \times 10^{13} \text{ GeV}} \right)^2, & m_X < H_{\text{rh}} \end{cases}$$

[A.A., Grzadkowski, Socha, 2020]

$\Omega_X h^2 = 0.12$ (solid), 0.012 (dashed) for $\gamma = 10^{-5}$

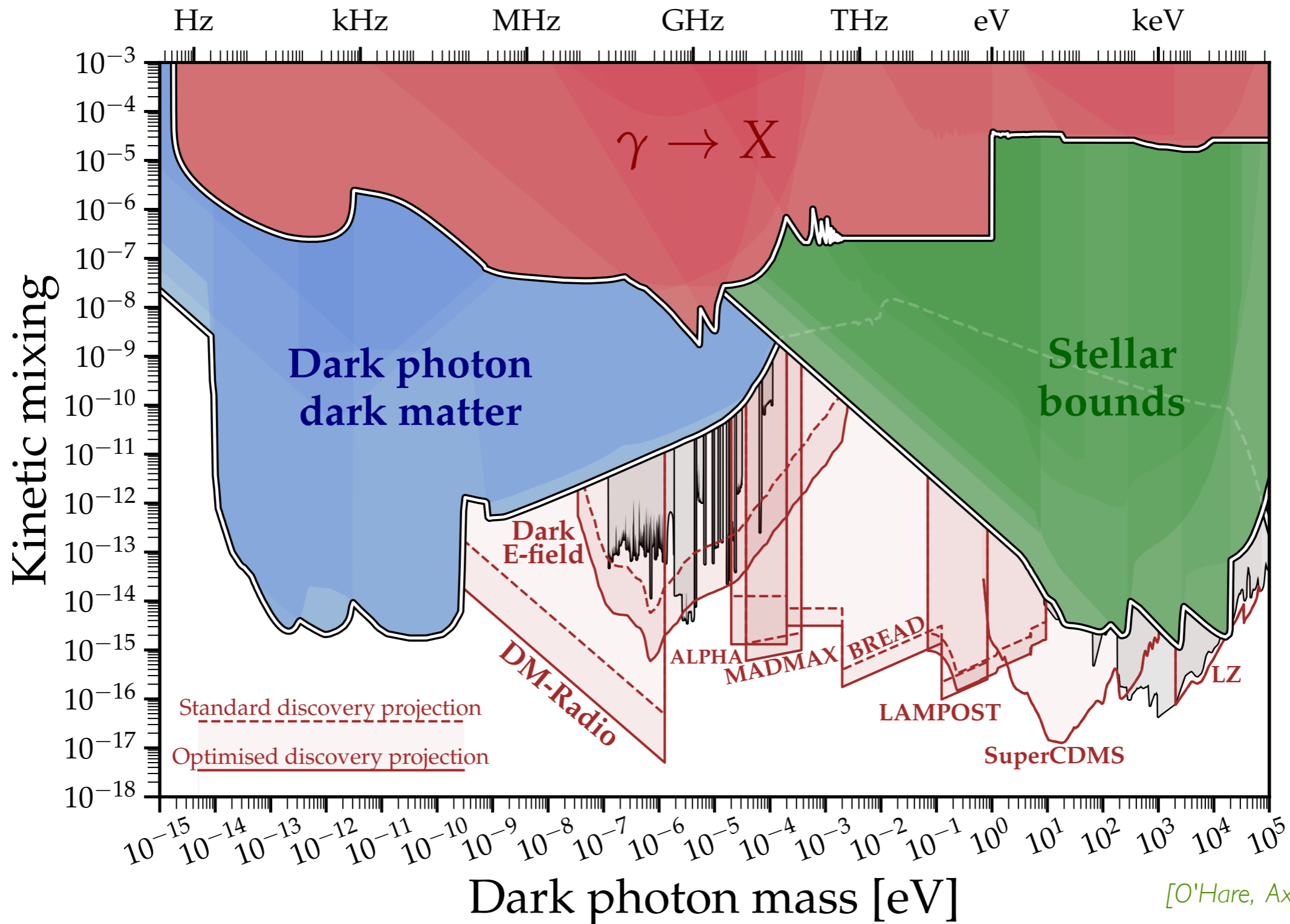


$\Omega_X h^2 = 0.12$ (solid), 0.012 (dashed) for $\gamma = 10^{-10}$



DARK PHOTON SEARCHES

■ Gravitational dark photon DM can be detected if $\mathcal{L} \supset \frac{\epsilon}{2} F_{\mu\nu} X^{\mu\nu}$



[O'Hare, AxionLimits@github]

CONCLUSIONS

- Simplified WIMP DM models are in tension with ever-strengthening experimental bounds.
- Significant progress has been made on freeze-in DM, gravitationally produced DM, axion DM, PBHs,
- Early universe cosmology, e.g. inflation/reheating dynamics, can play a significant role in non-thermal DM production.
- The search for DM remains a well-motivated task for BSM physics for the coming decades.

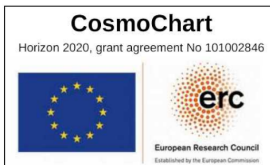
Thank you!

The BEH mechanism at high temperature – a link to matter production –

Iason Baldes

2018-2019: ULB EoS, 2019-2022: ULB F.R.S.-FNRS,

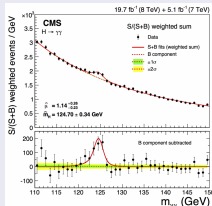
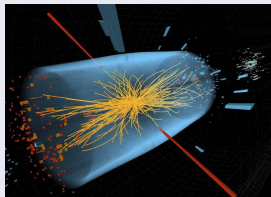
2022-: LPENS



A Decade of Discoveries in High Energy Physics
Brussels, 9 March 2023

The SM H boson and its mass

Following the 2012 discovery

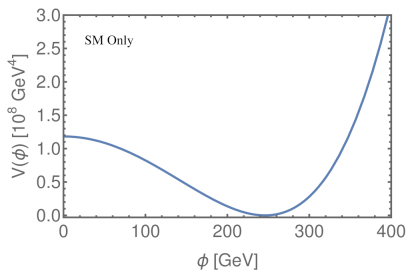


The mass is now also precisely measured:

$$m_h = 125.25 \pm 0.17 \text{ GeV}$$

[PDG Avg. of ATLAS/CMS results]

SM H Potential

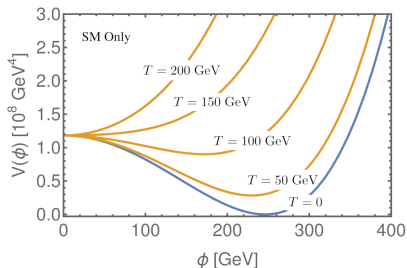


$$\langle H \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ \phi \end{pmatrix} \quad V(\phi) = -\frac{1}{2}\mu_H^2\phi^2 + \frac{1}{4}\lambda_H\phi^4$$

All the parameters are known:

$$\sqrt{2}\mu_H = m_h = 125 \text{ GeV [LHC]} \quad v_{\text{EW}} = \sqrt{\frac{m_h^2}{2\lambda_H}} = 246 \text{ GeV [Muon decay]}$$

At finite temperature



$$V(H) \approx -\frac{1}{2}\mu_H^2\phi^2 + \frac{1}{4}\lambda_H\phi^4 + \frac{1}{2}c_H T^2\phi^2$$

The thermal mass coefficient is related to other SM couplings:

$$c_H \approx \left(\frac{\lambda_H}{2} + \frac{3g_2^2}{16} + \frac{g_Y^2}{16} + \frac{y_t^2}{4} \right) \approx 0.4$$

Cosmological Puzzles

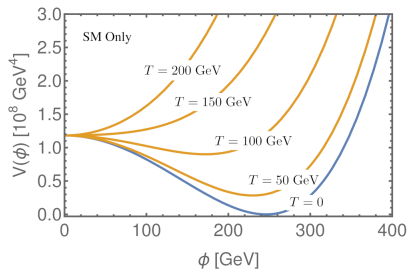
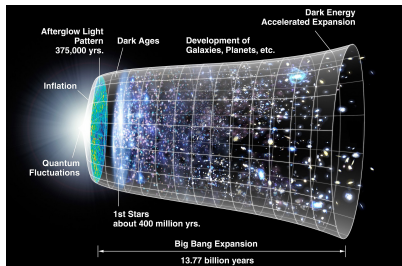
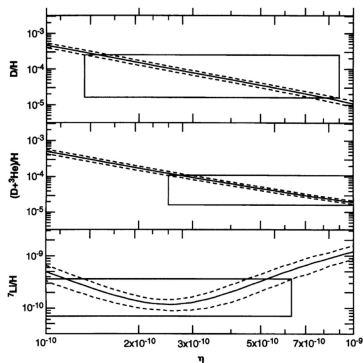
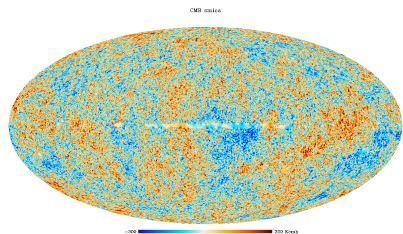


Image: NASA/Wikipedia

Cosmological puzzles which could be related to the SM H boson

- Dark Energy
- Inflation
- Dark Matter
- **Baryon Asymmetry → Ordinary Matter Density**

The matter-antimatter asymmetry



CMB in agreement with BBN:

$$Y_B \equiv \frac{n_b - n_{\bar{b}}}{s} = (0.86 \pm 0.02) \times 10^{-10}$$

Sakharov Conditions

- 1 B violation
- 2 C and CP violation
- 3 Departure from thermal equilibrium (or spontaneously broken CPT)

SM + FLRW

- 1 (B+L) violation present in symmetric phase at $T \gtrsim 100$ GeV from non-perturbative EW sphaleron process.
- 2 CP violation observed in quark sector.
- 3 Can be driven by expansion.

Electroweak baryogenesis - basic picture

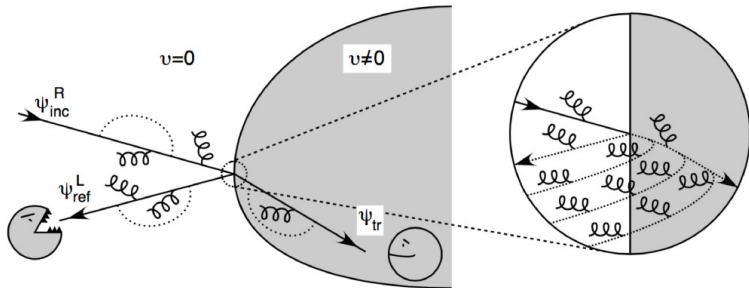
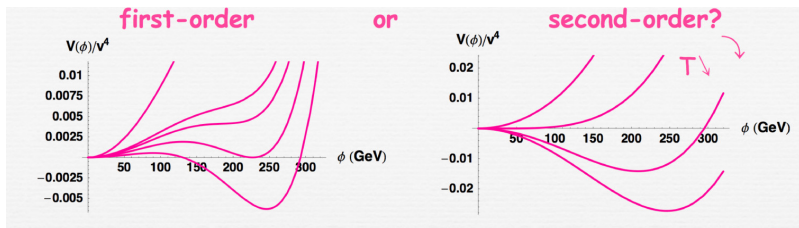


Image from - Gavela, Hernandez, Orloff, Pène, Quimbay [hep-ph/9406289]

- CP violating collisions with the bubble walls lead to a chiral asymmetry.
- Sphalerons convert this to a Baryon Asymmetry.
- This is swept into the expanding bubble where sphalerons are suppressed.

Electroweak baryogenesis - Requirements



Electroweak baryogenesis requires:

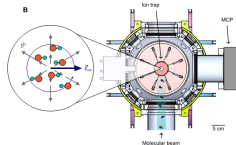
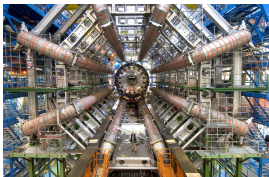
- A strong first order phase transition ($\phi_n/T_n \gtrsim 1$)
- Sufficient CP violation

However in the SM:

- The H boson mass is too large
- Quark masses are too small

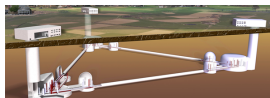
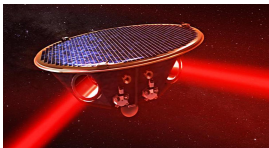
We require new EW-scale physics!

Experimental signatures



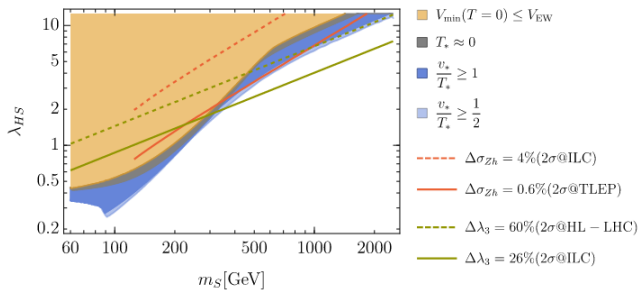
BSM Experimental signatures for EWBG

- 1 Collider signals associated with $V(H)$ modification.
- 2 Electric Dipole Moments associated with low scale CP violation.
- 3 Gravitational waves from the strong FOPT.



Singlet model

First order EW Phase Transition from a singlet - Choi, Volkas '93 + ...



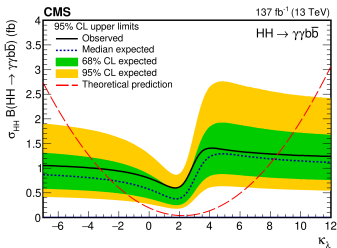
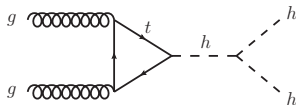
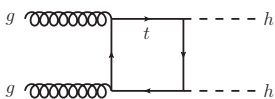
- Beniwal et al, 1702.06124

Modification of h^3 coupling

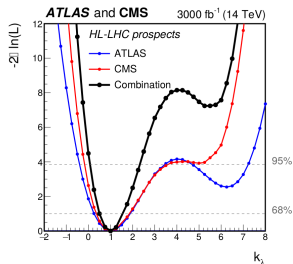
$$\lambda_3 \approx \frac{m_h^2}{2v_{EW}} + \frac{\lambda_{HS}^3 v_{EW}^3}{24\pi^2 m_S^2}$$

Collider signatures - Triple h coupling

SM: $V(h) = \frac{1}{2} m_h^2 h^2 + \lambda_H v_{EW} h^3 + \frac{1}{4} \lambda_H h^4$ with $v_{EW} = \sqrt{\frac{m_h^2}{2\lambda_H}} = 246$ GeV.



- 2011.12373

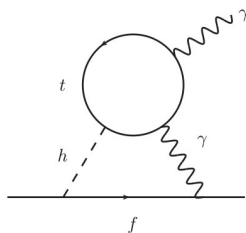


- 1902.00134

Measuring the cubic term is long term challenge.

Some, but not all, singlet models returning a strong FOPT can be excluded by HL-LHC.

Electron EDM constraint



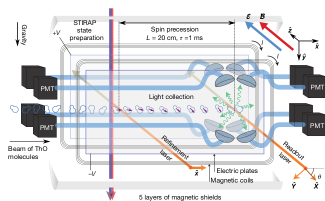
$$\frac{i}{2} d_e (\bar{e} \sigma^{\mu\nu} \gamma_5 e) F_{\mu\nu}$$

Rough estimate of the EDM - Glioti, Rattazzi, Vecchi, 1811.11740.

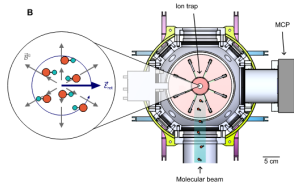
$$|d_e| \sim 10^{-29} \text{ e cm } \theta_{\text{CP}} \left(\frac{50 \text{ TeV}}{\Lambda} \right)^2 \quad 1\text{-loop}$$

$$|d_e| \sim 10^{-29} \text{ e cm } \theta_{\text{CP}} \left(\frac{2.5 \text{ TeV}}{\Lambda} \right)^2 \quad 2\text{-loop}$$

Experimental searches - EDMs

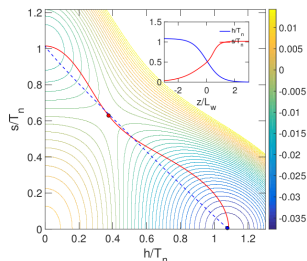


ACME II (ThO): $|d_e| < 1.1 \times 10^{-29} \text{ e cm}$ - Nature 562, 355–360 (2018)



Colorado (HfF^+): $|d_e| < 4.1 \times 10^{-30} \text{ e cm}$ - 2212.11841

Hiding the CP violation



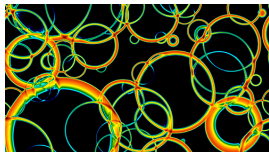
$\mathcal{L} \supset \frac{1}{2} \bar{\chi} ((\eta P_R + \eta^* P_L) S + m_\chi) \chi + y \bar{L}_\tau H_2 P_R \chi + \text{h.c.}$ - from [1] below.

One idea is to hide the CP violation in the dark sector

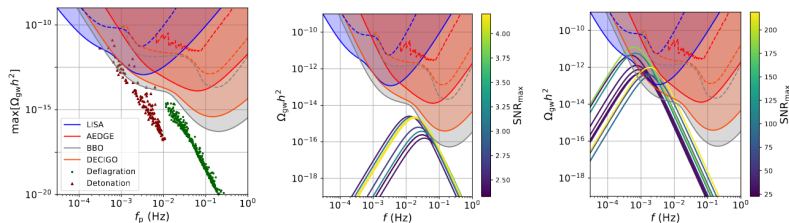
- 1 “Electroweak baryogenesis from a dark sector”,
Cline, Kainulainen, Tucker-Smith, 1702.08909.
- 2 “Electroweak Baryogenesis From Dark CP Violation,”
Carena, Quirós, Zhang, 1811.09719 and 1908.04818.

- eEDM at 3 or 4-loops.

Experimental searches - GWs



From a simulation by Weir et al.



Singlet model - Cline et al. 2102.12490

Strong transitions are detectable by LISA.

LVK searches for PTs have started e.g.:

..., Turbang, ... Mariotti, Sakellariadou, Sevrin et al. 2209.14707

Going beyond EWBG

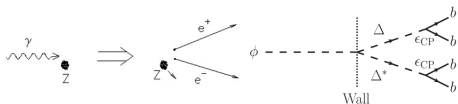
Cosmological PTs are of course of interest beyond EWBG.

- Associated Baryogenesis models not necessarily:
 - 1 Constrained experimentally by EDMs.
 - 2 Limited to small γ_{wall} in order to match Y_B .
- Larger range of GW signals possible.

Say we have a phase transition with $\langle \phi \rangle : 0 \rightarrow v_\phi$

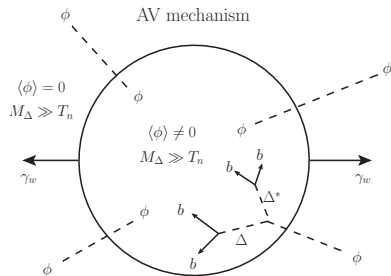
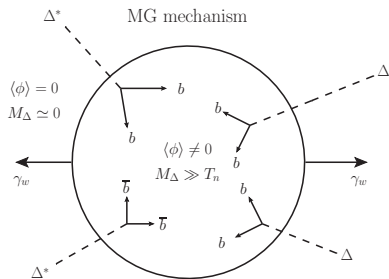
- Incoming quanta $E \sim \gamma_{\text{wall}} T \gg v_\phi$: can produce massive states
- Non-adiabatic transitions $\sigma \rightarrow S$, with $M_S \gg m_\sigma, v_\phi$
- Pair production of heavy particles $\sigma \rightarrow S + S$, with $M_S \gg m_\sigma, v_\phi$

- Azatov, Vanvlasselaer, 2010.02590 Azatov, Vanvlasselaer, Yin 2101.05721



The wall breaks translational symmetry

Going beyond EWBG - Ultra relativistic walls



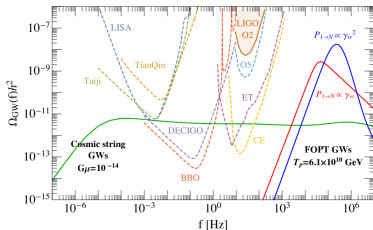
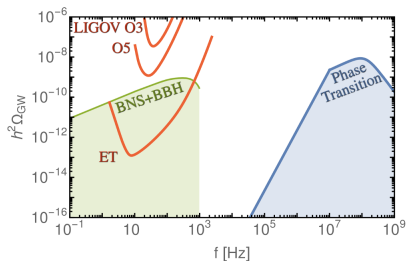
IB, Blasi, Turbang, Mariotti, Sevrin 2106.15602

For strong phase transitions with $\langle \phi \rangle \gg T_n$.

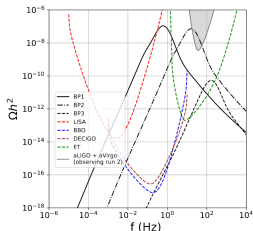
The final states are far from equilibrium and can source the baryon asymmetry through their decays.

Closely related work: Azatov, Vanvlasselaer, Yin 2106.14913

Going beyond EWBG - GW signals



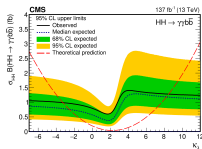
-IB, Blasi, Turbang, Mariotti, Sevrin 2106.15602 - Peisi Huang, Ke-Pan Xie 2206.04691



-Dasgupta et al. 2206.07032

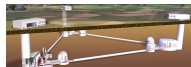
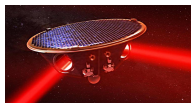
Larger range of GW signals possible.

Conclusions



Conclusion

- SM H boson could play a leading role in the early universe
→ possible connections to Inflation, DM, Baryon asymmetry...
- EWBG well motivated, testable
→ now rather constrained baryogenesis mechanism (EDMs).
- But: Beyond the SM EW or other first order PT hardly constrained.
- Exciting interplay of future collider tests and Gravitational waves.

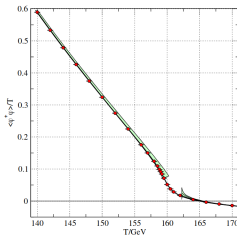
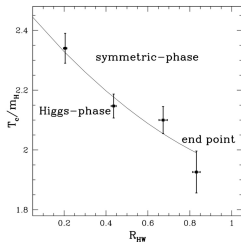


The matter-antimatter asymmetry

Textbook Argument for Baryogenesis

- In a symmetric universe $n_b/s = n_{\bar{b}}/s \approx 10^{-20}$
- The post-inflation causal volume is too small for baryons/antibaryons to be sufficiently separated
- $n_b/s = n_{\bar{b}}/s \approx 10^{-10}$ would be reached at $T \approx 40$ MeV when $M_{H-3} \approx 10^{-7} M_{\odot}$
- Need a mechanism to generate the asymmetry

Electroweak phase transition - Lattice Studies



- Csikor, Fodor, Heitger, hep-ph/9809291,

D'Onofrio, Rummukainen 1508.07161

SM with $m_h = 125$ GeV predicts a crossover.

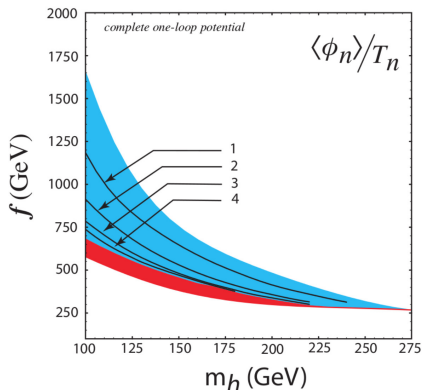
Nevertheless, only the minimum (VEV) of the potential, and the 2nd derivative there (m_h), is known if we allow for BSM physics.

The SM scalar potential can be modified.

Require a modification of the SM Scalar potential

Successful electroweak baryogenesis requires suppressed washout:

$$\frac{\Gamma_{\text{sph}}}{V} \sim 10^{1\div 4} \left(\frac{\alpha_W T}{4\pi} \right)^4 \left(\frac{2M_W(\phi)}{\alpha_W T} \right)^7 \text{Exp} \left[-\frac{3.2M_W(\phi)}{\alpha_W T} \right] \Rightarrow \frac{\phi_n}{T_n} \gtrsim 1$$



$$V(\Phi) = m^2 |\Phi|^2 + \lambda |\Phi|^4 + \frac{1}{f^2} |\Phi|^6$$

Other options:

- Singlet models/tree level barriers
- Multi-step transitions
- Thermal barriers from bosonic loops

- Delaunay, Grojean, Wells [0711.2511]

CPV and The Baryonic Yield

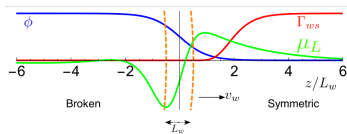


Image from 1706.08534 - Bruggisser, Konstandin, Servant

Diffusion equation

$$\partial_z n_B = \frac{3}{2} v_w^{-1} \Gamma_{ws} (N_c \mu_L T^2 - \mathcal{A} n_B), \quad \Gamma_{ws} = 10^{-6} T \exp(-a\phi(z)/T)$$

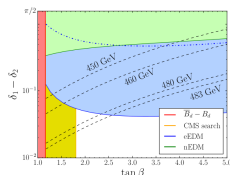
$$\eta_B = \frac{n_B(-\infty)}{s} = \frac{135 N_c}{4\pi^2 v_w g_* T} \int_{-\infty}^{+\infty} dz \Gamma_{ws} \mu_L e^{-\frac{3}{2} \mathcal{A} \frac{1}{v_w} \int_{-\infty}^z dz_0 \Gamma_{ws}}$$

$$\eta_B \sim \frac{\Gamma_{ws} \mu_L L_w}{g_* T} \sim \frac{10^{-8} \mu_L}{T} \quad \text{for} \quad L_w \sim \frac{1}{T}$$

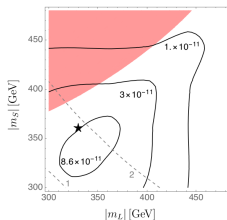
ACME: $|d_e| < 8.7 \times 10^{-29}$ e cm (2013) $|d_e| < 9.4 \times 10^{-29}$ e cm (2017)

Is electroweak baryogenesis dead?

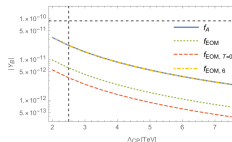
James M. Cline^{1,2}



1.



2.

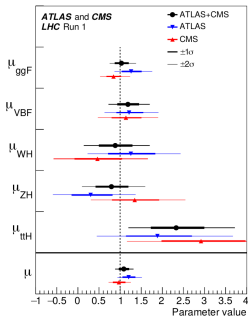
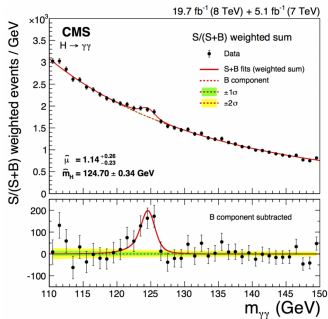


3.

- 1 1611.05874 - Dorsch, Huber, Konstandin, No
- 2 1707.02306 - Egana-Ugrinovic
- 3 1710.04061 - de Vries, Postma, van de Vis, White

Severe constraint on EWBG!

LHC constraints - Limit on Mixing



$$\mu = 1.09 \pm 0.11$$

LHC Run 1

7 + 8 TeV

1606.02266

$$\mu = 1.10 \pm 0.06$$

LHC Run 2

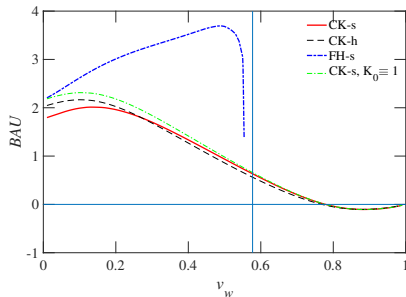
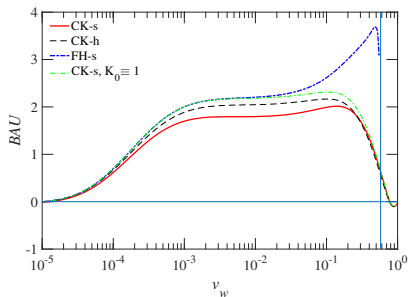
13 TeV

1810.02521

$$\theta \lesssim \mathcal{O}(0.1)$$

But: problem if $v_{\text{wall}} \simeq 1$.

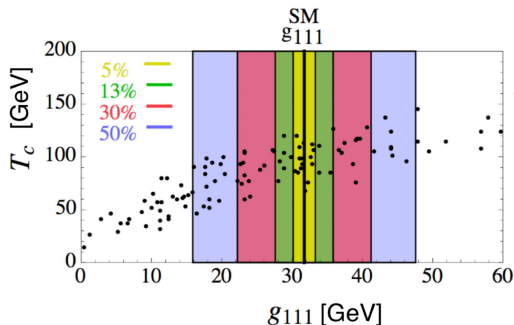
- Less of the plasma is pushed by the wall at high v_{wall} .
- This suppresses the BAU.



From: Cline, Kainulainen 2001.00568

Also see: Dorsch, Huber, Konstandin 2106.06547

Collider signatures - Singlet models difficult to detect



Somewhat optimistically:

$\sim 30 - 50\%$ HL-LHC or TLEP

$\sim 13\%$ ILC

$\sim 3 - 8\%$ 100 TeV pp

- Correlation between T_c and triple Higgs couplings $g_{111}h^3$ in a singlet model. - Profumo, Ramsey-Musolf, Wainwright, Winslow [1407.5342]
- And/or: mixing reducing the signal strength.
Currently LHC: $\theta \lesssim \mathcal{O}(0.1)$ compatible with singlet models of EWBG.
- And/or: direct searches for heavy singlet states.