[R. Alemany]
[CERN BE/OP]
Lectures UCLouvain (30.03.2022)

#### Colliders and The Large Hadron Collider

- I. LHC layout
- II. LHC Operational cycle
- III. Beam measurements



#### Why colliders?



What is the centre of mass energy → the interesting energy because it is the only one available for the collision products → when two particles of masses m1 and m2 collide? 

Square of the four-momentum:

$$|\mathbf{p}| = (E, \vec{p})$$

$$\mathbf{p}^2 = E^2 - \vec{p}^2 = m^2$$

$$(\mathbf{p}_1 + \mathbf{p}_2)^2 = E_{cm}^2 = (E_1 + E_2)^2 - (\vec{p}_1 + \vec{p}_2)^2$$

$$E_{\rm cm}^2 = (m_1^2 + m_2^2 + 2m_2 E_{1,\rm lab})$$

(m2,-p1)

$$E_{\rm cm}^2 = (E_1 + E_2)^2$$

#### Why colliders?



$$E_{\rm cm}^2 = (m_1^2 + m_2^2 + 2m_2 E_{1,\rm lab})$$

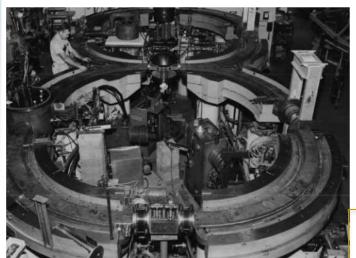
To get a Ecm = 630 GeV to do a SppS collider (SPS in the 80's discovered the weak neutral currents with this collider), what should be the value of E1,lab?

E1,lab ~ 212 TeV

To get a Ecm = 14 TeV to do an LHC collider (LHC discovered the Higgs in 2012) what should be the value of E1,lab?

E1,lab ~ 100 PeV

If colliders would not have been invented we would not have discovered yet many particles



#### First collider

double-ring electron-electron collider, built by a small group of Princeton and Stanford physicists. (Courtesy Stanford University)

reference to their possibility stems from a Russian publication of the 1920s; it would not be surprising if the same idea occurred independently to many people. The first collider actually used for particle-physics experiments, built at Stanford in the late 1950s, produced electron-electron collisions (see photograph on the left). Other early machines, generating electron-positron collisions, were built in Italy, Siberia and France. Since then there has been

#### Why colliders?

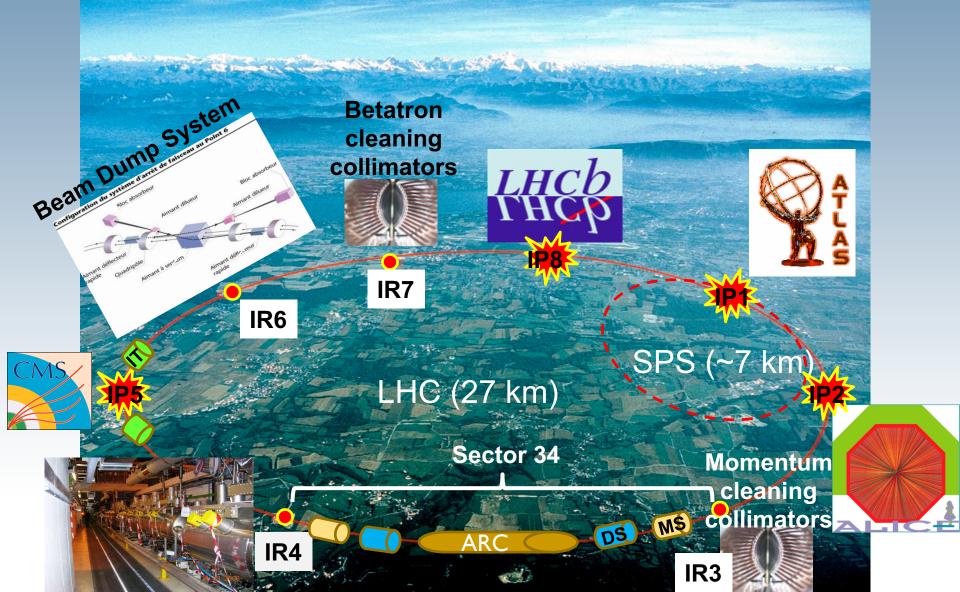


From the point of view of the energy, a collider is very very interesting However it has some disadvantages:

- The beam is several orders of magnitude less dense than a fixed target → less number of interactions per second
  - How do we overcome this issue?
     We collide for many hours → e.g. LHC collisions ~ 12 15 hours
- 2. The collision products go out in all directions → experiment has to cover all the space around the beam pipe
  - Experiments located around the beam pipe are difficult to take in/out to install a new experiment.
  - Fixed target experiments are more flexible, e.g. SPS north area: tens
    of experiments per year.

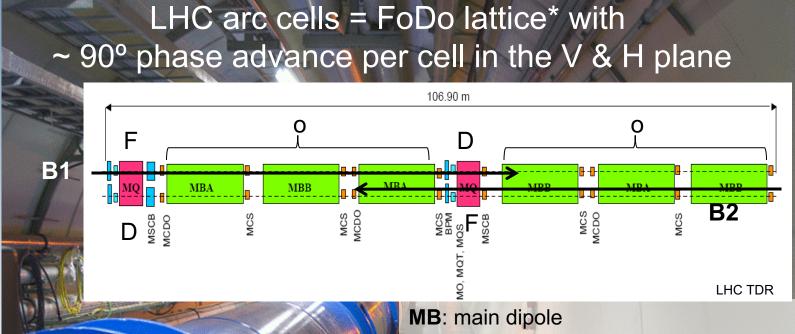
#### I. Basic layout of the machine





### I. Basic layout of the machine: the arc





The FoDo-Lattice

A magnet structure consisting of focusing and defocusing quadrupole lenses in alternating order with nothing in between.

(Nothing = elements that can be neglected on first sight: drift, bending magnets, RF structures ... and especially experiments...)

MQ: main quadrupole MQT: Trim quadrupole

MQS: Skew trim quadrupole

MO: Lattice octupole (Landau damping)

MSCB: Skew sextupole +

Orbit corrector (lattice chroma+orbit)

MCDO: Spool piece sextupole

MCDO: Spool piece octupole +

Decapole

BPM: Beam position monitor





Golden formula (you should know by heart) Circumference → FIXED!!! by LEP

$$B\rho = \frac{p}{Ze}$$

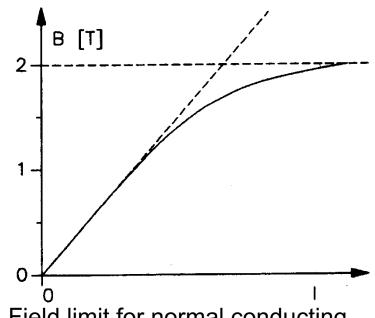
FIXED, no choice 
$$\rho \approx \frac{26658.9 \text{ m}}{2\pi} \cdot 66\% \approx 2780 \text{ m}$$

~ 66% of the lattice elements are dipoles

p = nucleon momentum → defined by the physics case → TeV range → 7 TeV

$$B = \frac{p}{\rho Ze} \approx 3.33 \frac{p\left(\frac{GeV}{c}\right)}{\rho(m)} \neq 8.39 T$$

We need SUPERCONDUCTING technology

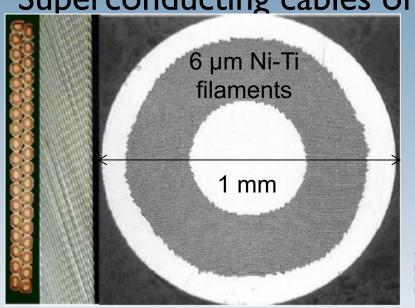


Field limit for normal conducting magnets due to saturation

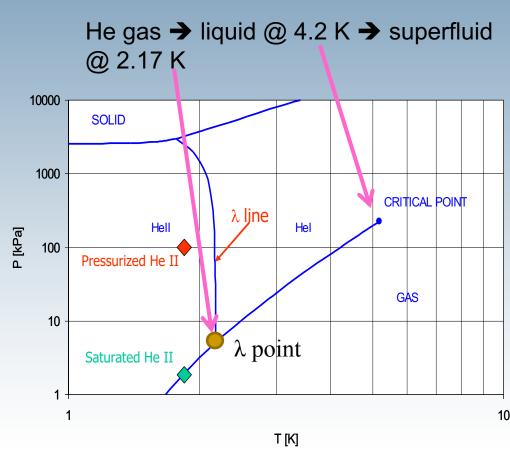
## I. Basic layout of the machine: Superconducting magnets



Superconducting cables of Nb-Ti

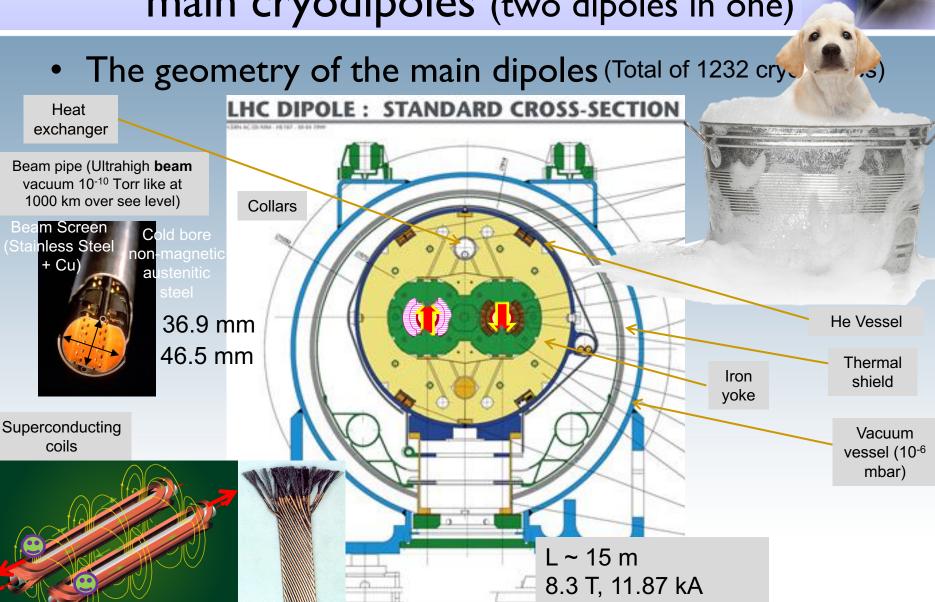


LHC ~ 27 km circumf. with 20 km of superconducting magnets operating @8.3 T. An equivalent machine with normal conducting magnets would have a circumference of 100 km and would consume 1000 MW of power → we would need a dedicated nuclear power station for such a machine. LHC consumes ~ 10% nuclear power station



Total amount of He used @LHC ~500-700 T

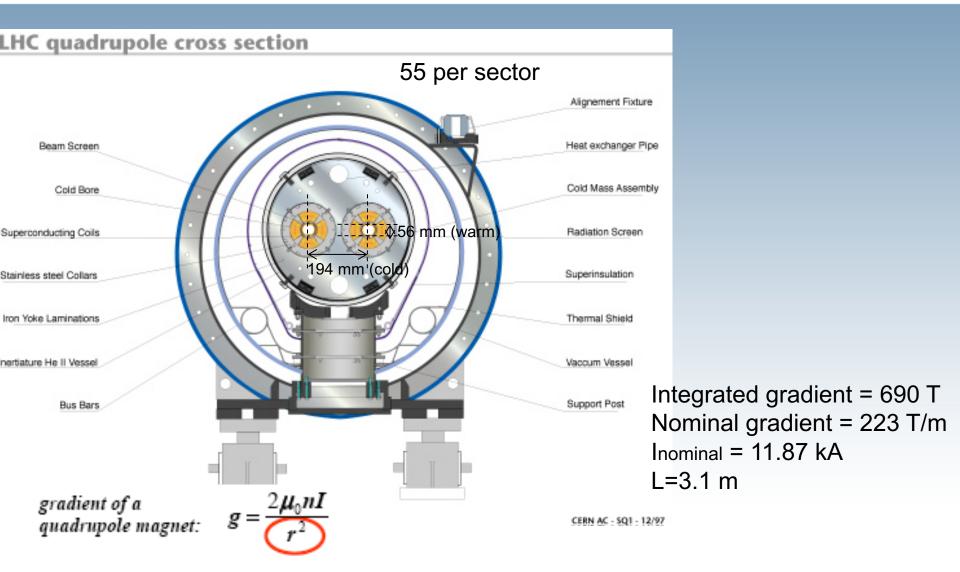
I. Basic layout of the machine: main cryodipoles (two dipoles in one)



T = 1.9 K, ~27.5 ton

### I. Basic layout of the machine: main quadrupoles





I. Basic layout of the machine:
 main dipoles → Field quality

#### The magnetic field of the main dipoles:

The **stability** of the geometry of the superconducting coils is essential to the field quality.

Mechanical stress during coil assembly

Thermal stresses during cool-down

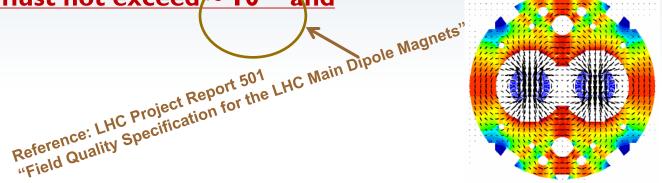
**Electromagnetic stresses** during operation

= sources of deformations of the coil geometry

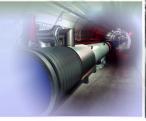
Additional sources of field errors are the <u>dimensional tolerances</u> of the magnet components and of the manufacturing and assembling tooling.

The relative variations of the integrated field and of the field shape

imperfections must not exceed ~ 10 -4 and



## I. Basic layout of the machine: main dipoles Field quality



#### Why the tolerances are so tight?

→ Because the field quality determines how long the particles can circulate in the accelerator

$$B_{y} + iB_{x} = B_{ref} \sum_{n=1}^{\infty} (b_{n} + ia_{n}) \left(\frac{x + iy}{R_{ref}}\right)^{n-1}$$

$$n = 1 \rightarrow \text{dipole}$$

$$n = 2 \rightarrow \text{quadrupole}$$

$$n = 3 \rightarrow \text{sextupole}$$

$$n = 4 \rightarrow \text{octupole}$$

$$n = 5 \rightarrow \text{decapole}$$

Up to which order do we care?

$$\rightarrow$$
 Up to n = 7 at least



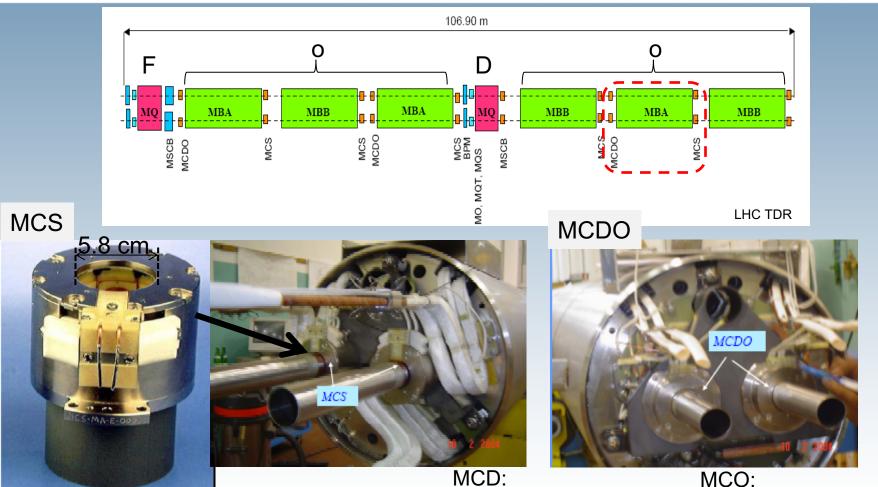
## I. Basic layout of the machine: main dipoles → Field quality



Error	Side effects if error is > few 10 <sup>-4</sup>	Corrected with	
al	Closed orbit perturbations and thus feed-down	Dipole correctors (MCB)	
Ы	contributions from higher order multiple errors		
a2	Linear coupling and vertical dispersion	Skew quadrupoles (MQS)	
b2	Tune change, $\beta$ and dispersion beating	Trim quadrupoles (MQT)	
a3	Chromatic coupling and Q"	Skew sextupoles (MSS)	
b3	b2 feed-down at injection (persistent current effect) and off-momentum $\beta$ -beat	Sextupole (MCS)	
a4	Dynamic aperture (DA) at injection		
b4	DA and Q" at injection	Octupole (MCO)	
a5	DA for off-momentum particles at injection		
b5	DA and Q" at injection	Decapole (MCD)	
a6 to b7	DA at injection  Sextupole  1.5e-06 1e-06		

## I. Basic layout of the machine: dipole corrector magnets

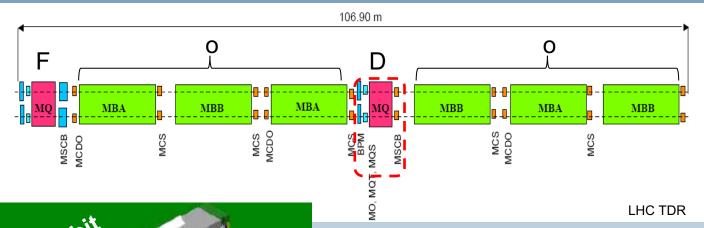




Nominal main field strength = 1630 T/m<sup>2</sup> Inominal = 550 A, 1.9 K, L=15.5 cm, ~10 kg Nominal main field strength ~ 120 T/m<sup>4</sup> Inominal = 550 A, 1.9 K, L=11 cm, ~6 kg Nominal main field strength = 8200 T/m<sup>3</sup> Inominal = 100 A, 1.9 K, L=11 cm, ~6 kg

### I. Basic layout of the machine: quadrupole corrector magnets





MSCB
Orbit
Chromaticity

Courtesy of Tesla Eng

376 tw/naperture assemblies supplied by Tesla Eng.

MSM (sextupole):

MCBM (dipole):

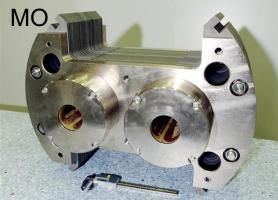
Nominal main field

strength = 2.93 T

Inominal = 55 A, 1.9 K,

L=78.5 cm, ~143 kg

Nominal main field strength = 4430 T/m<sup>2</sup> Inominal = 550 A, 1.9 K, L=45.5 cm, ~83 kg Landau damping



Nominal main field strength = 63100 T/m<sup>3</sup> Inominal = 550 A, 1.9 K L=38 cm, ~8 kg

#### MQT/MQS

MQT: Trim quadrupole

**MQS**: Skew trim quadrupole

#### MQT/MQS:

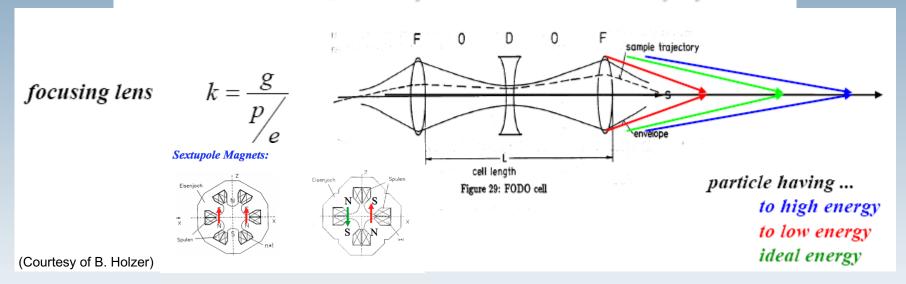
Nominal main field strength = 123 T/m Inominal = 550 A, 1.9 K

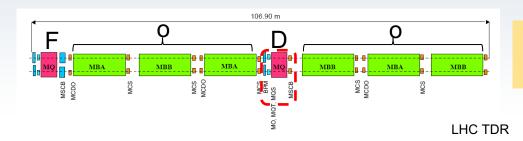
L=38 cm, ~250 kg

### I. Basic layout of the machine: quadrupole corrector magnets



### 20.) Chromaticity: A Quadrupole Error for ∆p/p ≠ 0

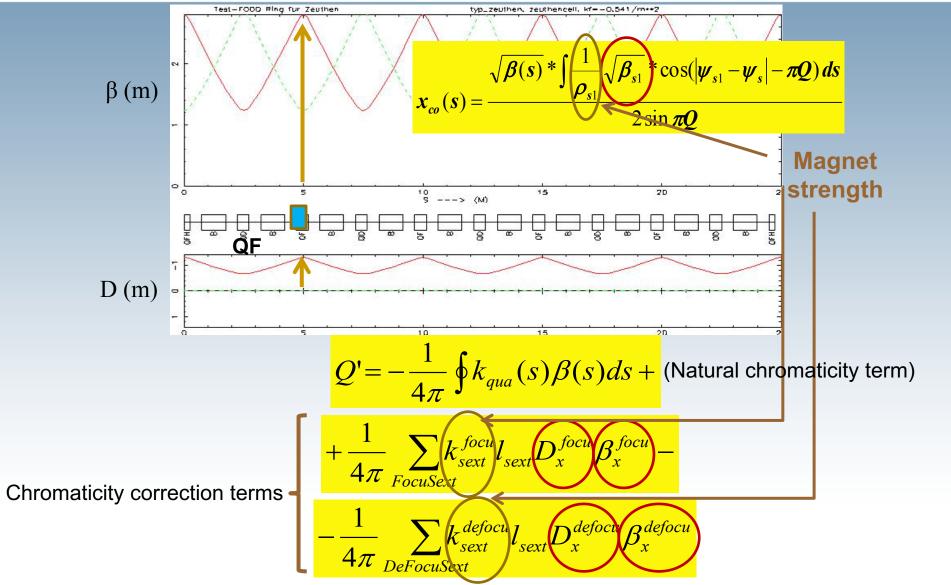




Why the orbit and sextupole correctors are placed close to a quadrupole?

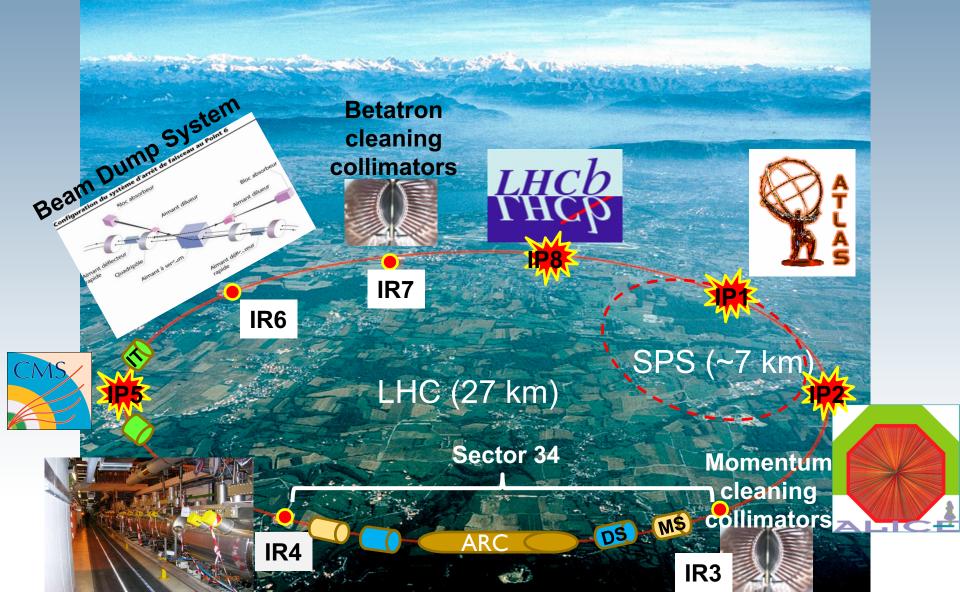
## I. Basic layout of the machine: quadrupole corrector magnets





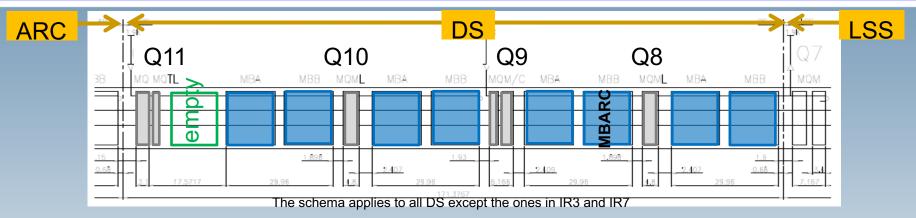
#### I. Basic layout of the machine





### I. Basic layout of the machine: Dispersion suppression

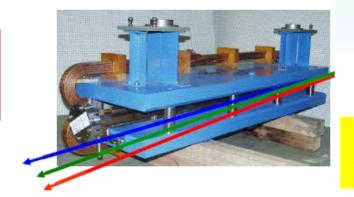




- Cancels the horizontal dispersion generated on one side by the arc dipoles and on the other by the separation/recombination dipoles and the crossing angle bumps
- Helps in matching the insertion optics to the periodic solution of the arc
- If only dipoles are used they cannot fully cancel the dispersion, just by a factor 2.5.
   Therefore individual powered quadrupoles are required (Q8-Q11 with I ~ 6000 A)

dipole magnet

$$\alpha = \frac{\int B \ dl}{p e}$$



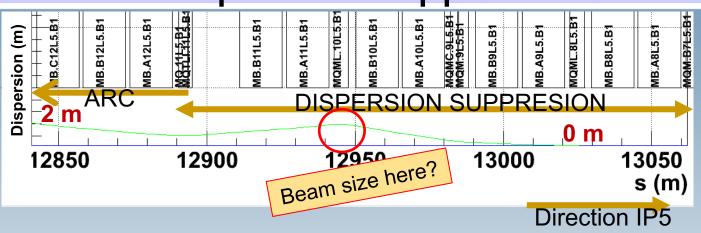
$$x'' + K(s)x = \frac{1}{\rho} \frac{\Delta p}{p}$$



$$x(s) = x_{\beta}(s) + D(s) \frac{\Delta p}{p}$$

### I. Basic layout of the machine: Dispersion suppression





D(s) is created by the dipole magnets and afterwards focused by the quadrupoles

The inhomogeneous solution changes the beam size ...

$$\varepsilon = \frac{\varepsilon_n}{\beta_{rel}\gamma_{rel}}$$

$$\sigma = \sqrt{\sigma_{\beta}^{2\Box} + \sigma_{D}^{2\Box}}$$
$$= \sqrt{\varepsilon\beta + D_{\Box}^{2} \left(\frac{\Delta p}{p}\right)^{2}}$$

$$\sigma = \sqrt{3.5 \cdot 10^{-6} + 4 \cdot 10^{-6}}$$

What is the beam size at 450 GeV?

$$σ = \sqrt{\frac{\varepsilon_n}{p}}β + D_{...}^2 \left(\frac{\Delta p}{p}\right)^2$$
 $\Rightarrow γ(@7 \text{ TeV})~7463 γ(@450 \text{ GeV})~480 β≈1$ 
 $\Rightarrow σ = \sqrt{1.3 \cdot 10^{-6} + 4 \cdot 10^{-6}} ≅ σ_D + 30% σ_β$ 

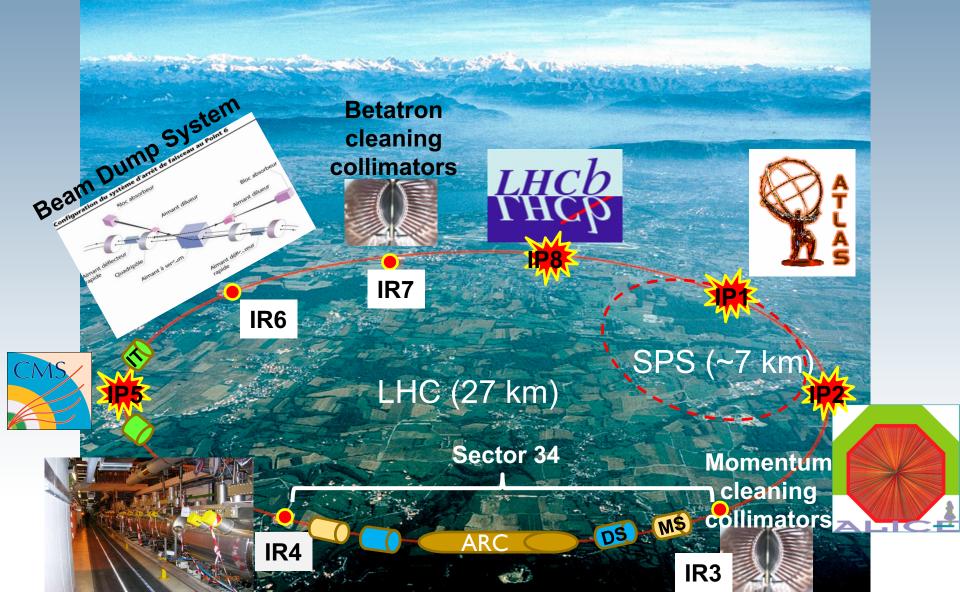
At 7 TeV in LHC:  $\epsilon n = 3.5 \,\mu m$  rad,  $\beta = 180 \, m, D = 2 \, m$ ,  $\Delta p/p \sim 10^{-3}$ 

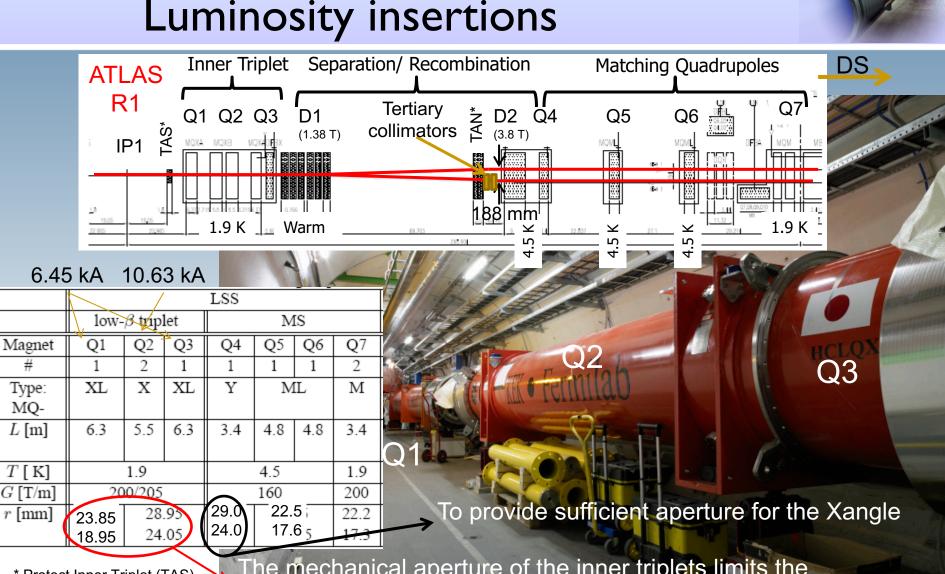
When you design your beam pipe you have to take into account the contribution of D(s)

Why do we suppress dispersion before reaching the IP?

#### I. Basic layout of the machine







\* Protect Inner Triplet (TAS) and D2 (TAN) from particles coming from the IP The mechanical aperture of the inner triplets limits the maximum  $\beta^*$  @IPs and the maximum Xangle  $\rightarrow$  limit peak lumi

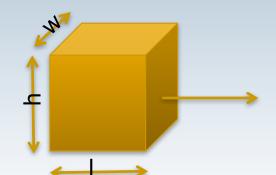
### Luminosity

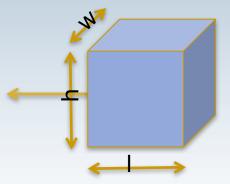
### Two figures of merit in a collider: energy and luminosity



#### Luminosity = number of interactions per unit of time

$$\frac{dR}{dt} = L\sigma$$
cross-section at the given energy of the interested physics
$$cm^{-2}. s^{-1}$$





Number of particles of the first beam in the cube:

$$n_1 = \frac{I_1}{qs}$$
I/s=intensity per cm<sup>2</sup>
with  $I = q/t$ 

The orange beam is crossed by n2 particles per unit of time:

$$n_2 = \frac{I_2}{qs}$$

#### Luminosity



The number of events per unit of time for two beams of 1 cm2 is:

$$\frac{dn}{dt} = \frac{I_1}{qs} \cdot \frac{l}{c} \cdot \frac{I_2}{qs} \cdot \sigma$$

and for beams of a general area s = wh

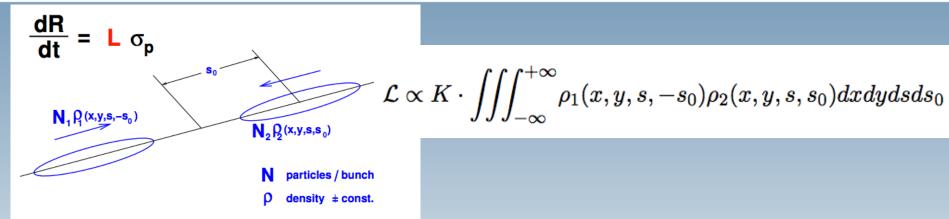
$$L = \frac{I_1 I_2}{q^2 c} \cdot \frac{l}{wh}$$

If we replace q = 1.6e-19 coulomb and c = 3e10 cm/s:

$$L \approx 1.3 \times 10^{27} \cdot I_1 I_2 \cdot \frac{l}{wh} \text{ [cm}^{-2}. \text{ s}^{-1}]$$



#### Luminosity for bunched beams



For two bunches colliding HEAD ON with Gaussian beam density distribution and equal beam sizes:

$$\mathcal{L} = rac{N_1 N_2 f N_b}{2\pi \sqrt{\sigma_{1x}^2 + \sigma_{2x}^2} \sqrt{\sigma_{2y}^2 + \sigma_{2y}^2}}$$

Ni: number of particles in one bunch

f: revolution frequency

Nb: number of bunches in one beam

sigma: beam size

Reference: Concept of luminosity

Werner Herr and Bruno Muratori\* CERN. Geneva. Switzerland



### Luminosity for bunched beams

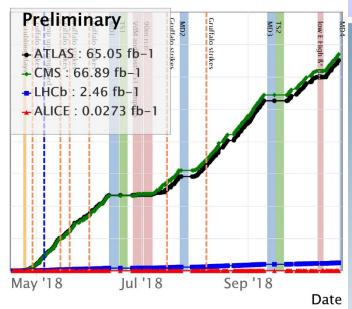
#### Examples:

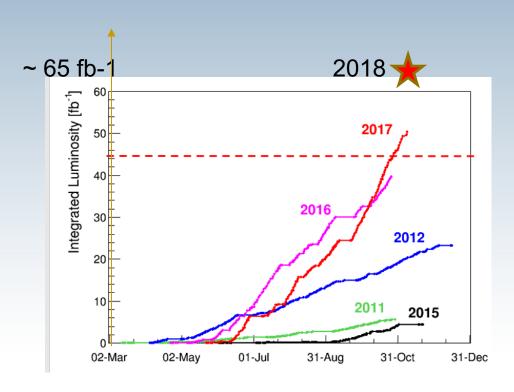
	Energy	$\mathcal{L}$	rate	$\sigma x/\sigma_y$	Particles
	(GeV)	$\mathrm{cm^{-2}s^{-1}}$	$s^{-1}$	$\mu$ m/ $\mu$ m	per bunch
SPS $(p\bar{p})$	315x315	6 10 <sup>30</sup>	$4 \ 10^5$	60/30	$\approx 10 \ 10^{10}$
Tevatron (p $\bar{p}$ )	1000x1000	50 10 <sup>30</sup>	$4 \ 10^6$	30/30	$\approx 30/8 \ 10^{10}$
HERA (e <sup>+</sup> p)	30x920	40 10 <sup>30</sup>	40	250/50	$\approx 3/7 \ 10^{10}$
LHC (pp)	7000x7000	$10000 \ 10^{30}$	$10^{9}$	17/17	11 10 <sup>10</sup>
$LEP (e^+e^-)$	105x105	100 10 <sup>30</sup>	≤ 1	200/2	$\approx 5 \ 10^{11}$
$PEP(e^+e^-)$	9x3	3000 10 <sup>30</sup>	NA	150/5	$\approx 2/6 \ 10^{10}$
KEKB ( $e^+e^-$ )	8x3.5	$10000\ 10^{30}$	NA	77/2	$\approx 1.3/1.6 \ 10^{10}$

### LHC Integrated performance



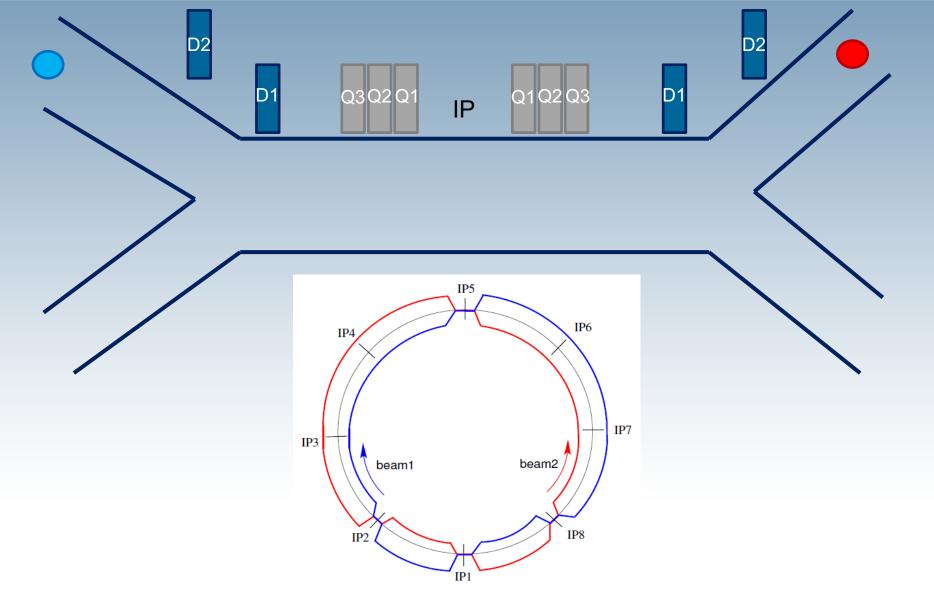


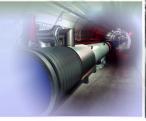


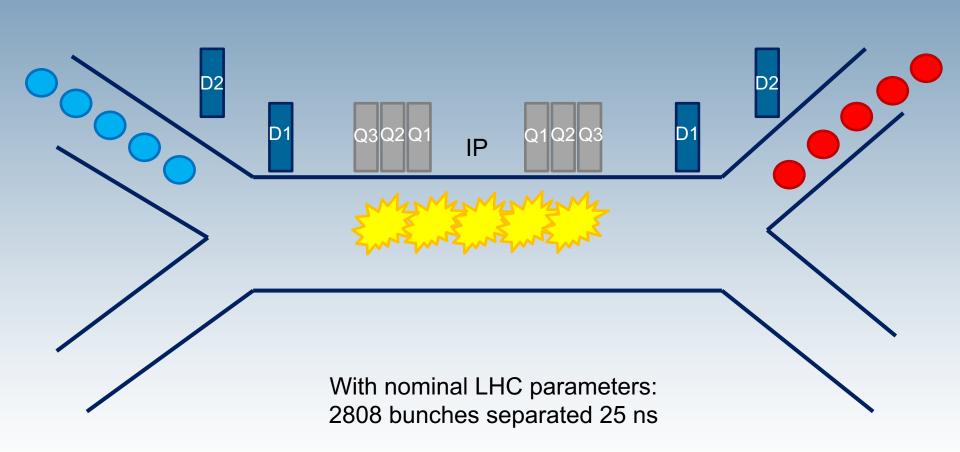


Delivered integ. luminosity [fb-1]



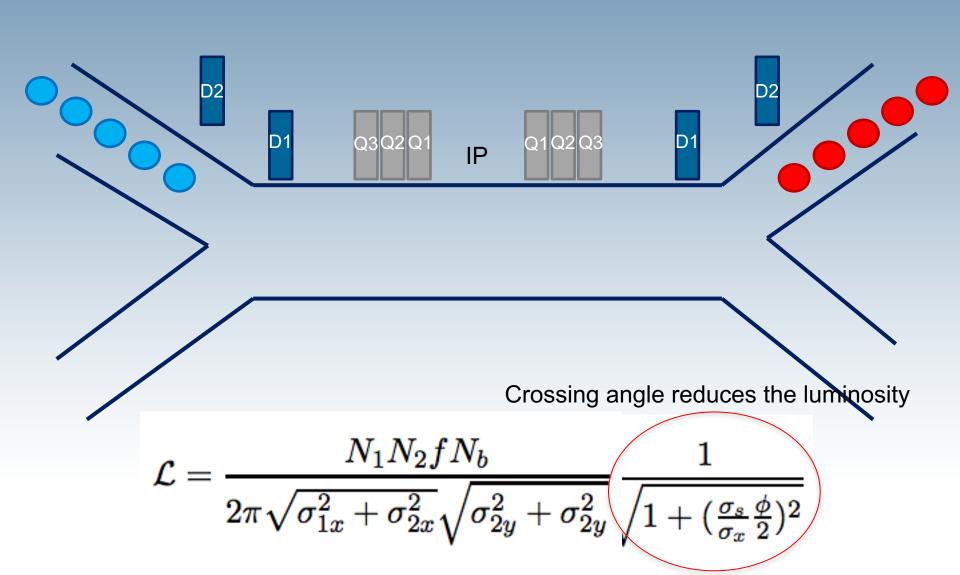




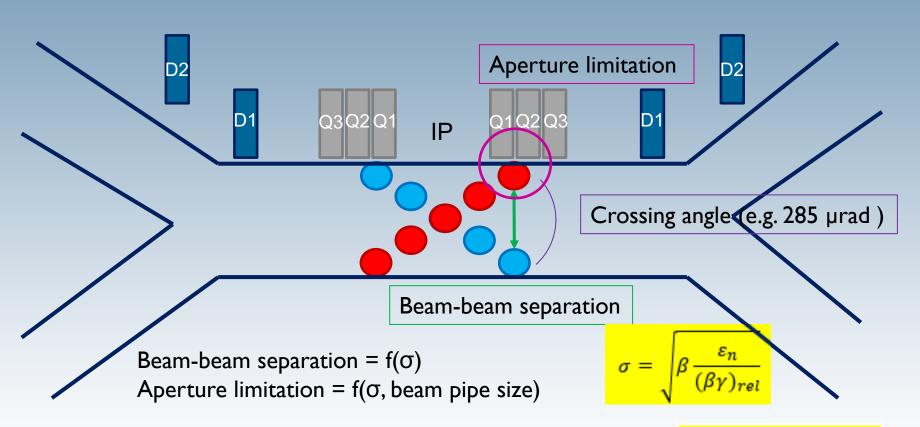


We can have up to 30 parasitic interactions around the IP









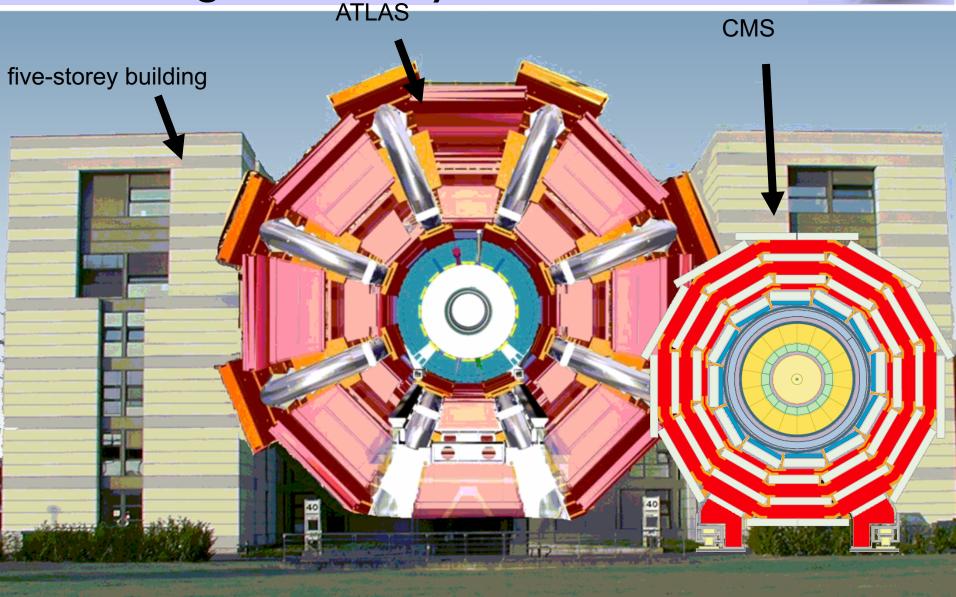
 $\beta$  is very small at the IP, but very big at the Inner triplets  $\Rightarrow$ 

$$\beta(s) = \beta^* + \frac{s^2}{\beta^*}$$

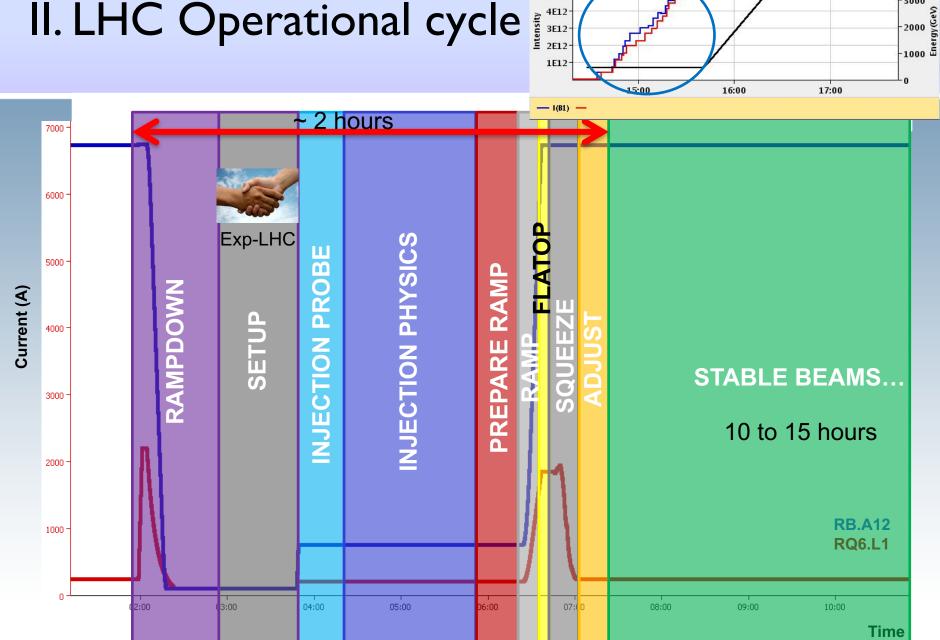
Therefore, the quadrupoles around the IP have such a big apertures

# I. Basic layout of the machine: High luminosity insertions ATLAS





### II. LHC Operational cycle

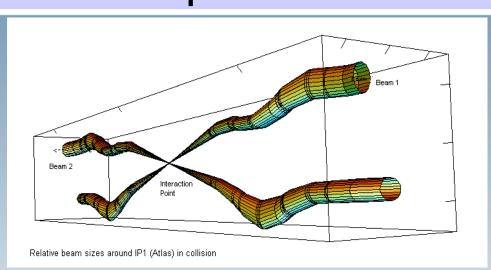


5E12

3000

### II. LHC Operational cycle:Squeeze → reduce β\*





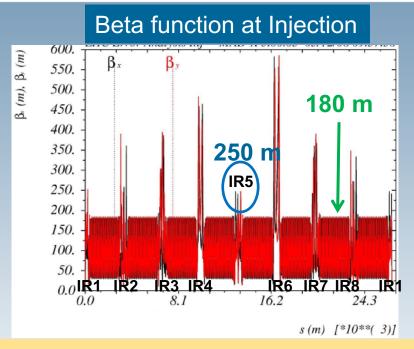
Squeeze the beam size down as much as possible at the collision point to increase the chances of a collision

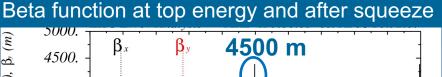
$$L \approx \frac{N_1 N_2 f_{rev} N_b}{4\pi \sigma^2} \quad \sigma = \sqrt{\beta \frac{\varepsilon_n}{(\beta \gamma)_{rel}}}$$

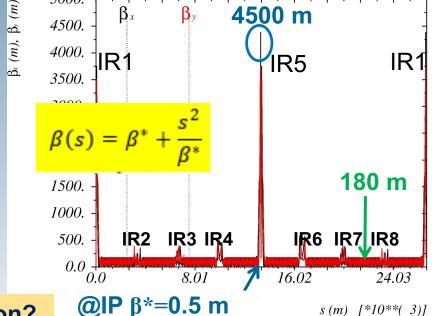
- So even tough we squeeze our N<sub>1,2</sub>=100,000 million protons per bunch down to 16 microns (1/5 the width of a human hair) at the interaction point. We get only around 20 collisions per crossing with nominal beam currents.
- The bunches cross (every 25 ns) so often we end up with around 600 million collisions per second at the start of a fill with nominal current.
- Most protons miss each other and carry on around the ring. The beams are kept circulating for hours → 10 hours

### II. LHC Operational cycle: Squeeze $\rightarrow$ reduce $\beta^*$ ( $\beta$ @IP)

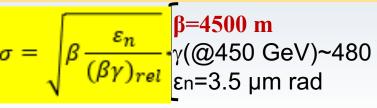








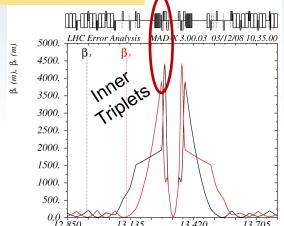
Why we cannot have  $\beta^*=0.5$  m at injection?



#### Remember:

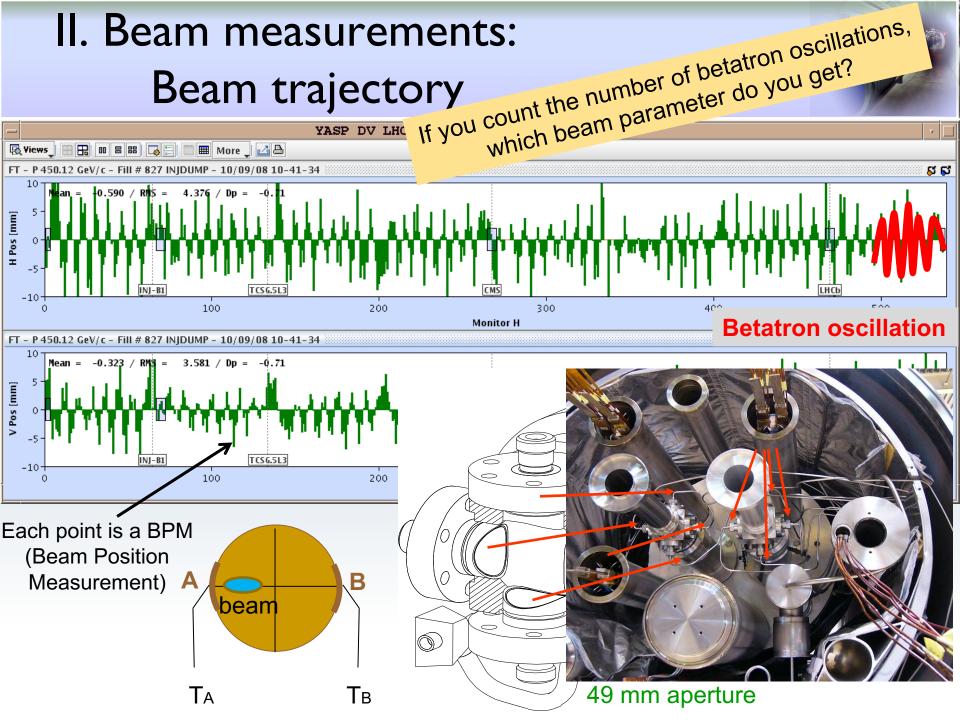
there is no D(s) here

 $\sigma \sim 6 \text{ mm} !!$ 



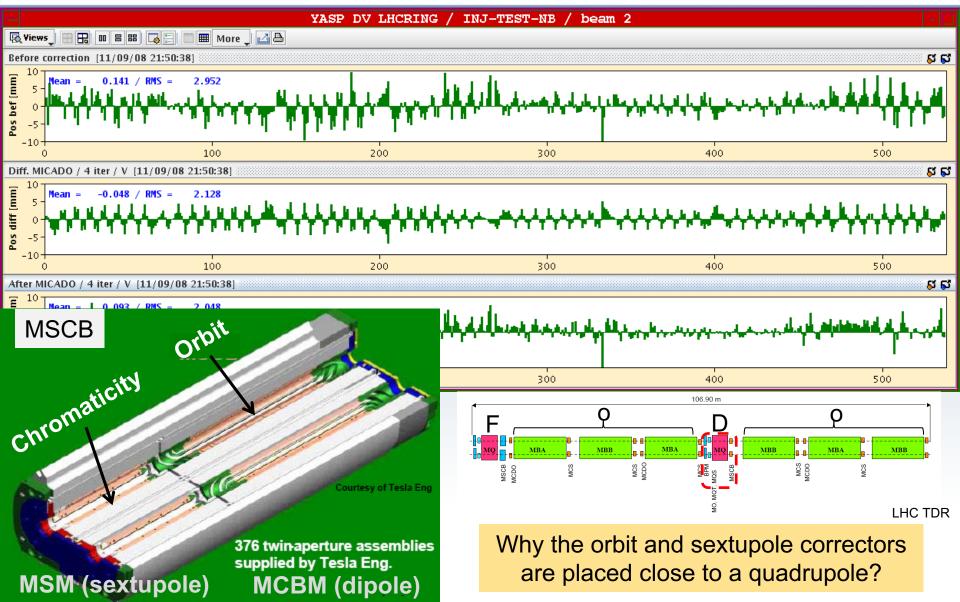
Rbeampipe~29/24 mm we could only accommodate ~ 4 times the beam size and we need at least 7σ clearance

> @ 7 TeV  $\sigma_{\rm IT} \sim 1.2 \, \rm mm$

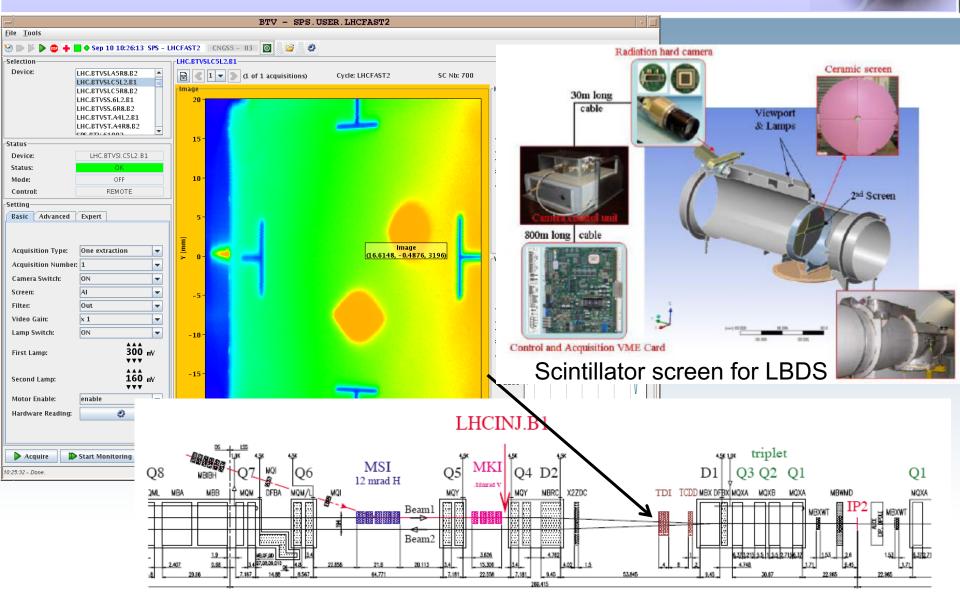


## II. Beam measurements: Beam trajectory correction





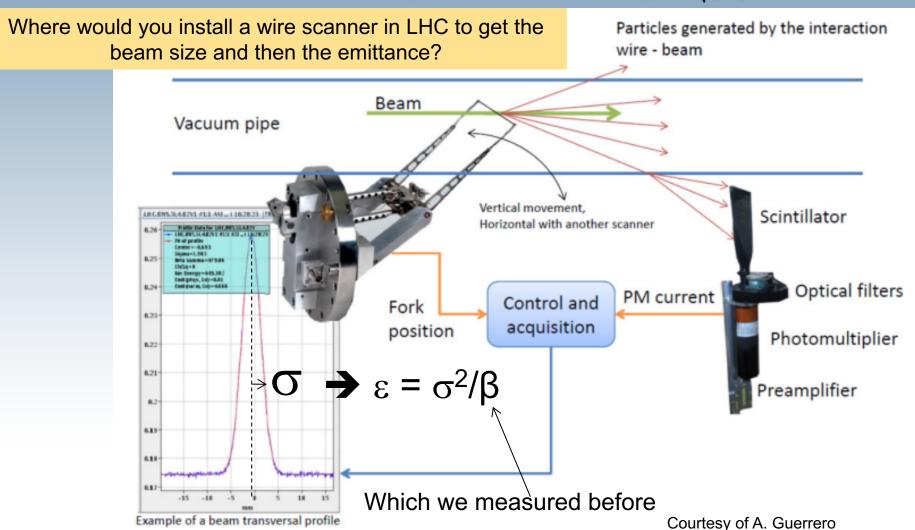
### II. Beam profile measurements: Beam I on TDI screen – Ist and 2<sup>nd</sup> turns



### II. Beam profile measurements: Emittance measurement - Wire scanner

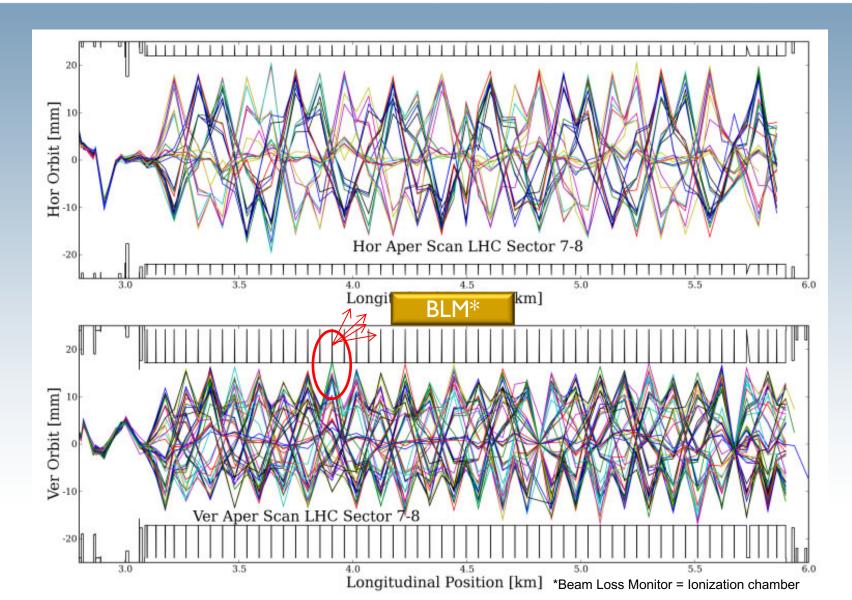


Emittance is the figure of merit for profile measurements. But it is a derived quantity from the beam size  $(\sigma = \sqrt{\beta \varepsilon})$ 



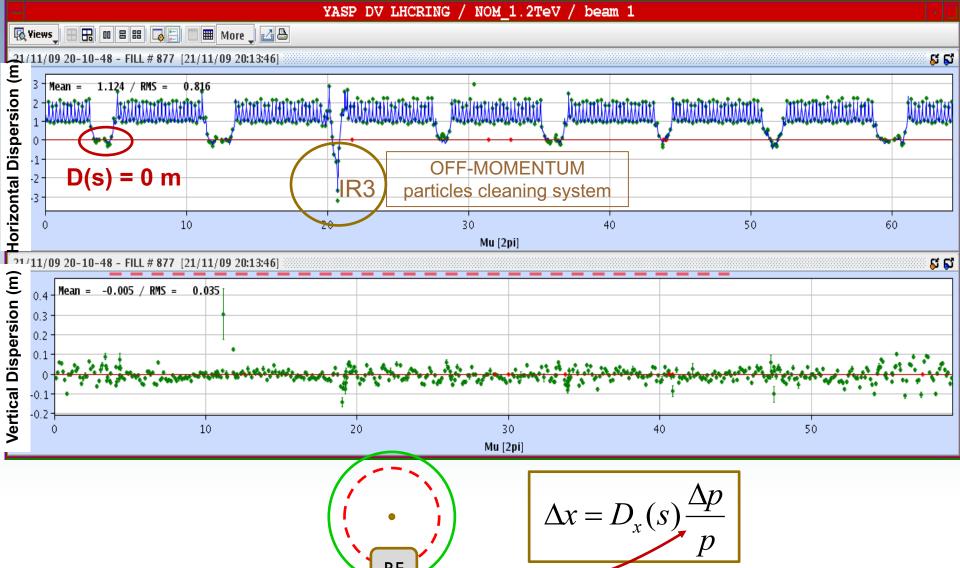
## II. Beam measurements: Aperture scan





### II. Beam measurements: Dispersion measurement



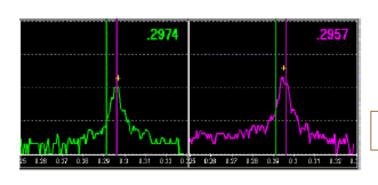


### II. Beam measurements: Beta measurement



a quadrupol error leads to a shift of the tune:

1<sup>st</sup> Change quadrupole strength in steps

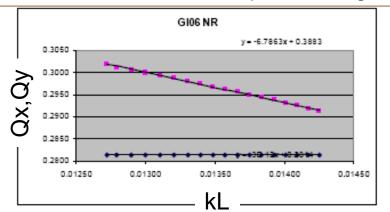


$$\Delta Q = \int_{s0}^{s0+l} \frac{\Delta k \beta(s)}{4\pi} ds \approx \frac{\Delta k l_{quad} \overline{\beta}}{4\pi}$$

2<sup>nd</sup> Measure Tune

Example: measurement of \beta in a storage ring: tune spectrum

#### 3<sup>rd</sup> Plot Tune vs Quadrupole strength

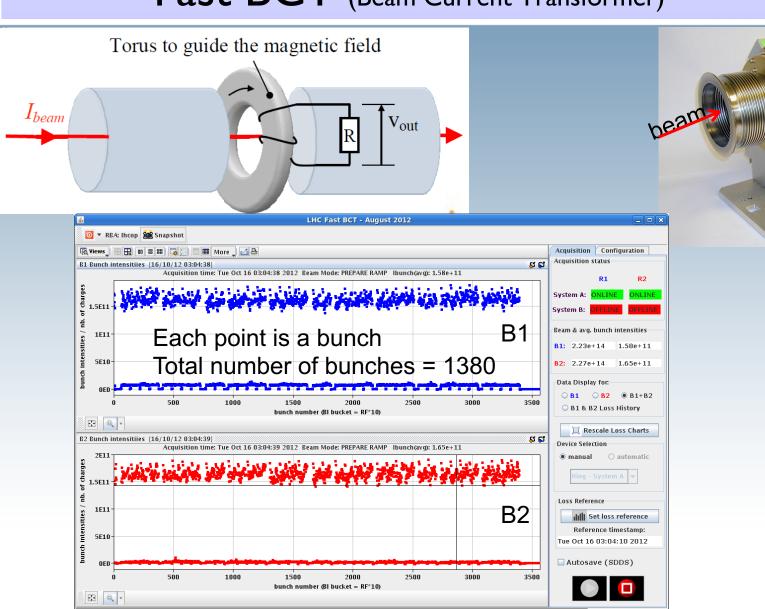


Courtesy of B. Holzer (lectures)

## II. Beam measurements: Fast BCT (Beam Current Transformer)



**FBCT** 



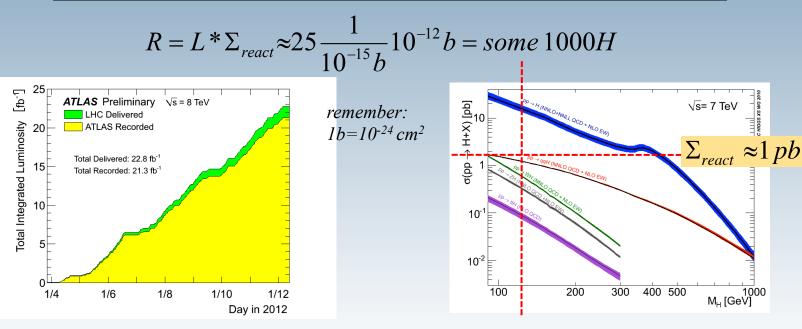


### **SPARE SLIDES**

### High Light of the LHC



production rate of events is determined by the cross section  $\Sigma_{react}$  and a parameter L that is given by the design of the accelerator: ... the luminosity



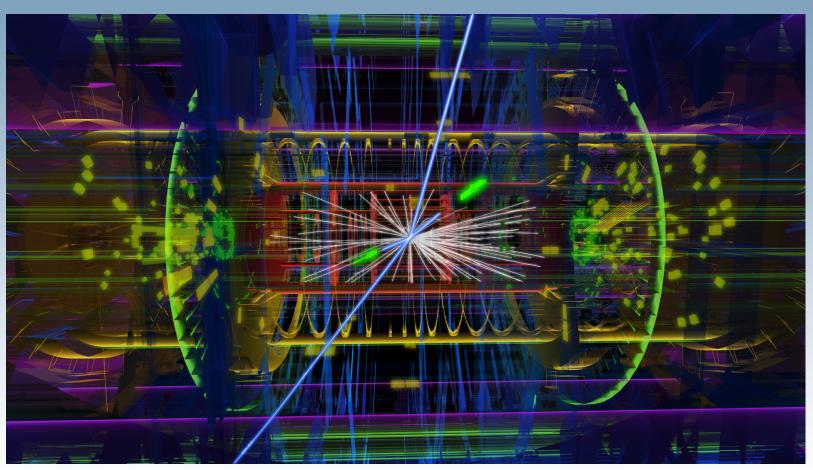
Integrated luminosity during RUN I

$$\int Ldt \approx 25 \, fb^{-1}$$

Official number: 1400 clearly identified Higgs particles "on-tape"

### High Light of the HEP year





ATLAS event display: Higgs => two electrons & two muons

## I. Basic layout of the machine: Superconducting magnets



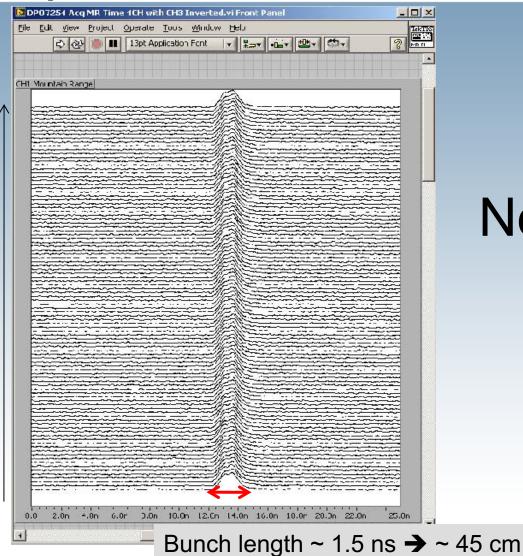
- Superfluid helium 
   Why is it so great?!!
  - very high thermal conductivity 

     is able to conduct away heat a thousand times better than a metallic conductor like copper
  - very low viscosity coefficient → can penetrate tiny cracks, deep inside the magnet coils to absorb any generated heat
  - very high heat capacity prevents small transient temperature fluctuations

## XIV. Beam captured – mountain range display







#### Now RF ON

### Beam parameters (nominal)

 $\beta^* = 0.55 \text{ m}$ 

Transverse normalized emittance

RMS bunch length  $\varepsilon = 0.5 \text{ nm rad}$ 

RMS beam size @IPI & IP5  $\rightarrow \sigma_{x,y} = \sqrt{\epsilon \beta}$ 

RMS beam size @IP2 & IP8  $\rightarrow \sigma_{x,y} = \sqrt{\epsilon \beta}$ 

Geometric luminosity reduction factor (F)

Instantaneous lumi @IPI & IP5 (IP2Pb-Pb, IP8)

Beam current

Stored energy/beam

		Injection	Collision	2012
Proton energy	GeV	450	7000	4000
Particles/bunch		$1.15 \times 10^{11}$		1.6 × 10 <sup>11</sup>
Num. bunches		2808		1380
Longitudinal emittance (40)	eVs	1.0	2.5	

µm rad

Α

MI

cm

μm

μm

 $cm^{-2}s^{-1}$ 

3.5

23.3

11.24

375.2

279.6

3.75

362

7.55

16.7

70.9

0.836

 $10^{32}$ )

 $10^{34}(10^{27},$ 

 $\beta^* = 0.6 \text{ m}$ 

7 1033

rad

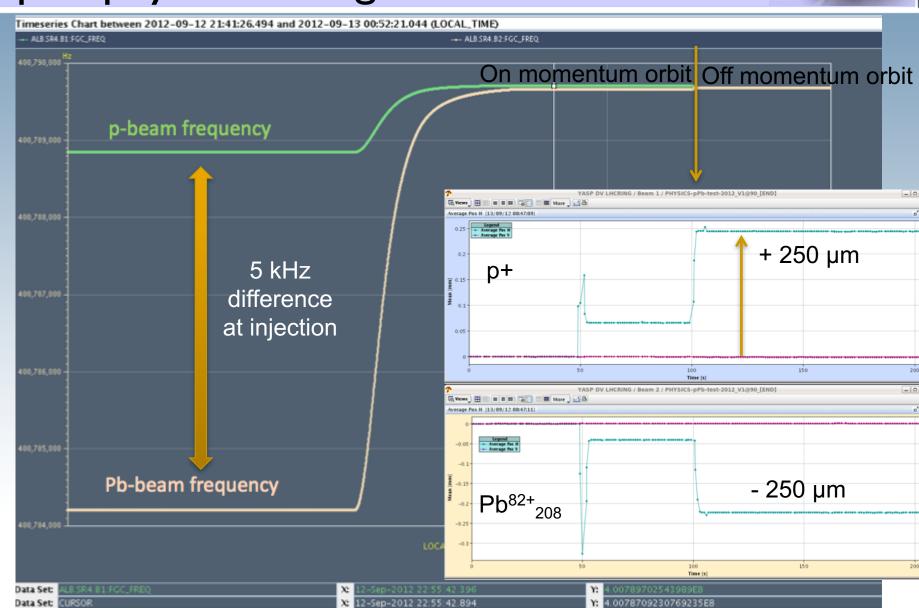
 $\varepsilon n = 2.5 \mu m$ 

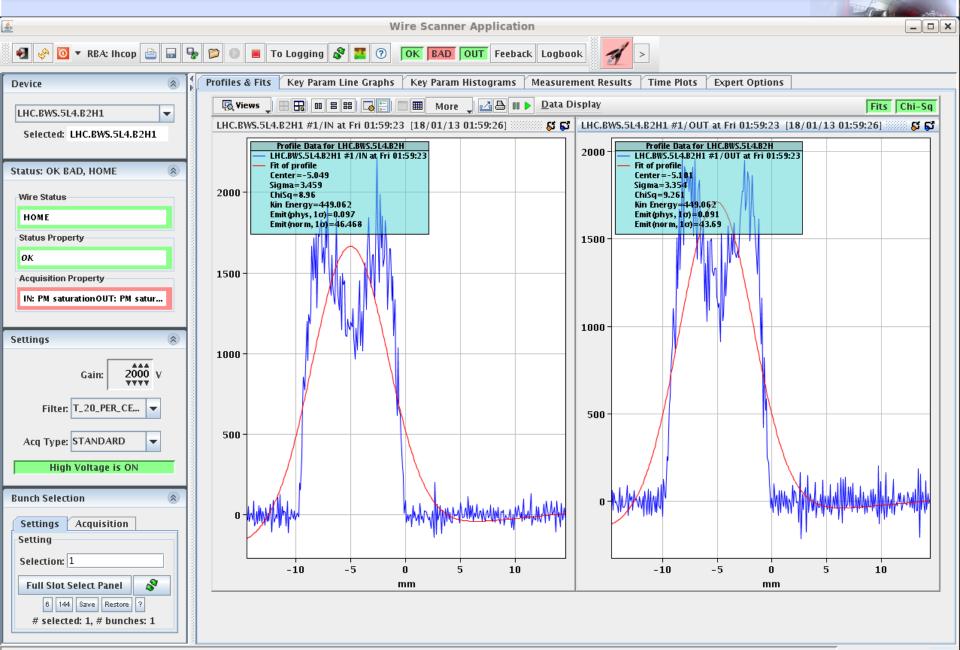
0.582

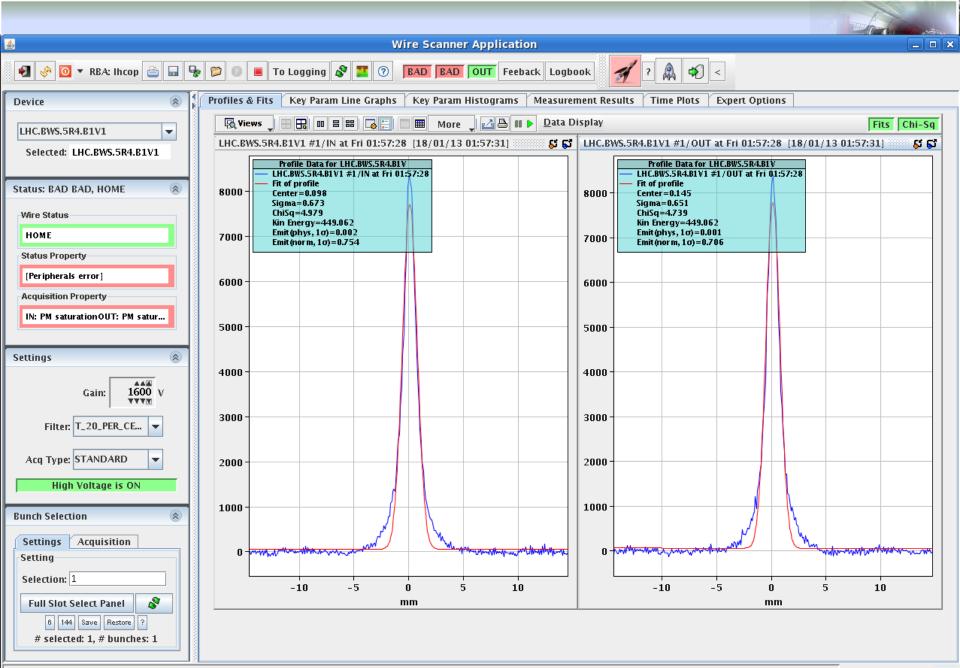
Peak luminosity related data



### pPb physics during 2013







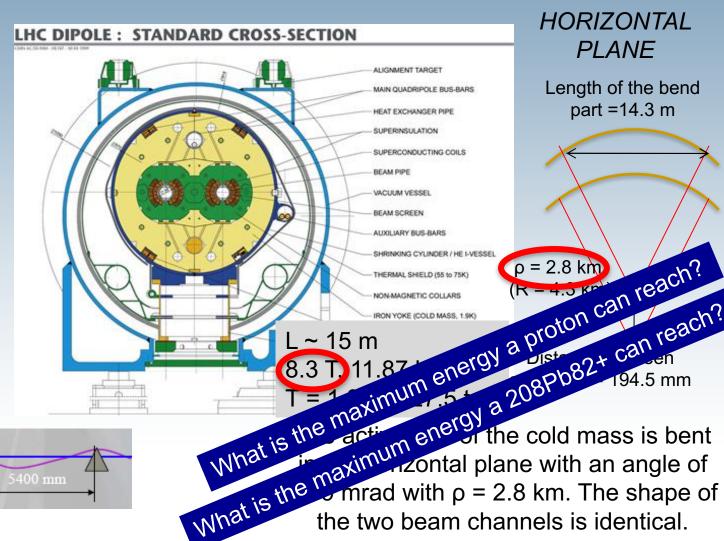
### I. Basic layout of the machine: main cryodipoles (two dipoles in one)

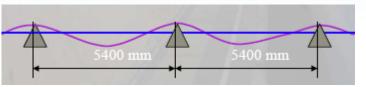


• The geometry of the main dipoles (Total of 1232 cryodipoles)

**VERTICAL** PLANE

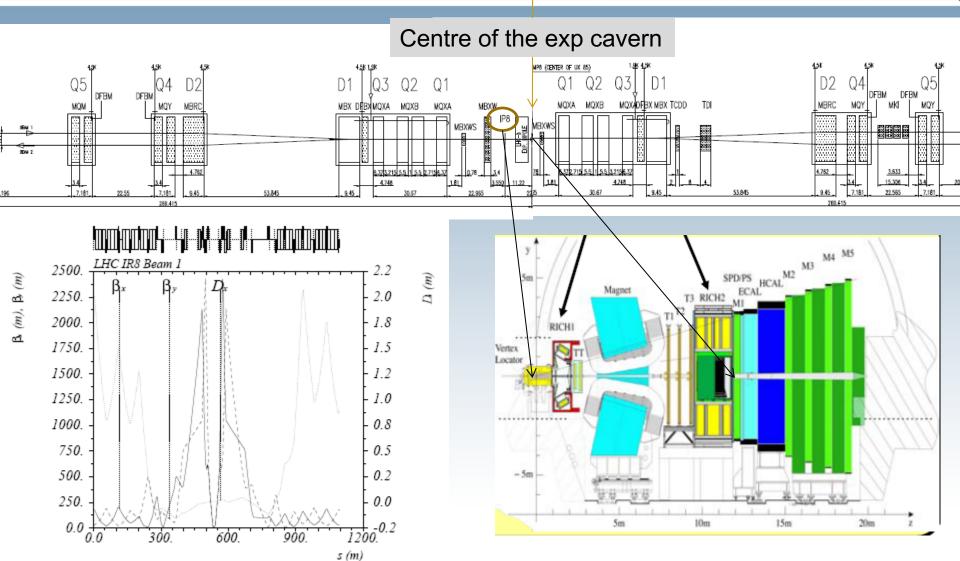
The theoretical shape of the beam channels is a straight line, while the natural shape has ~ 0.3 mm deflection between two supports at 5.4 m distance





## II. The experiments: Low luminosity insertions: LHCb

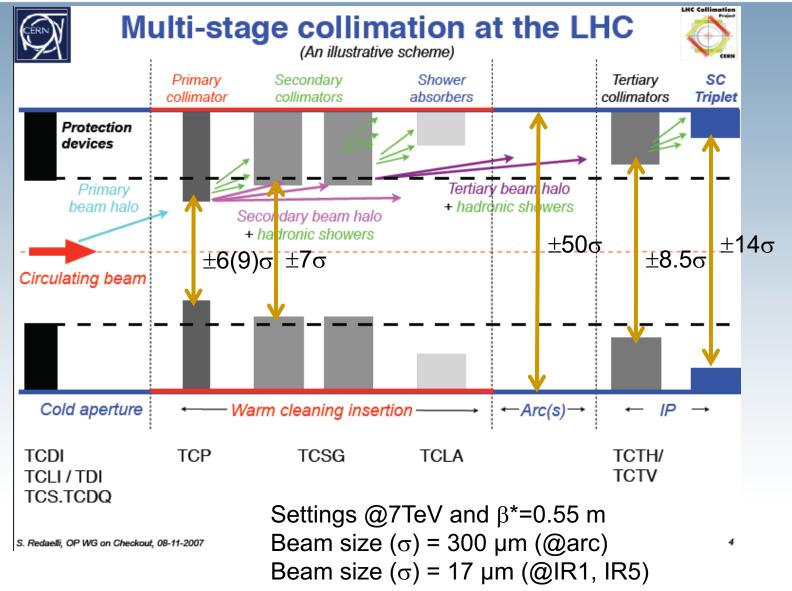




(c) Beam 1, collision optics

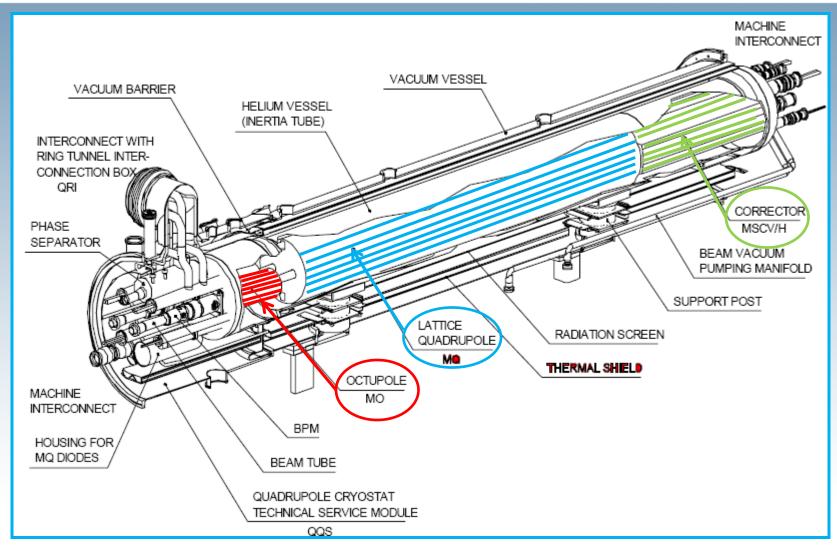
### IV. Momentum and betatron cleaning insertions (IR3, IR7)





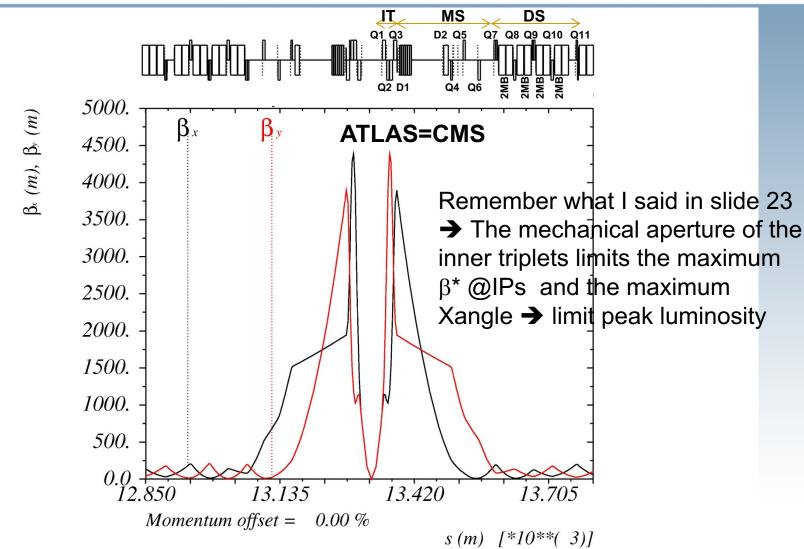
### I. Basic layout of the machine: quadrupole corrector magnets





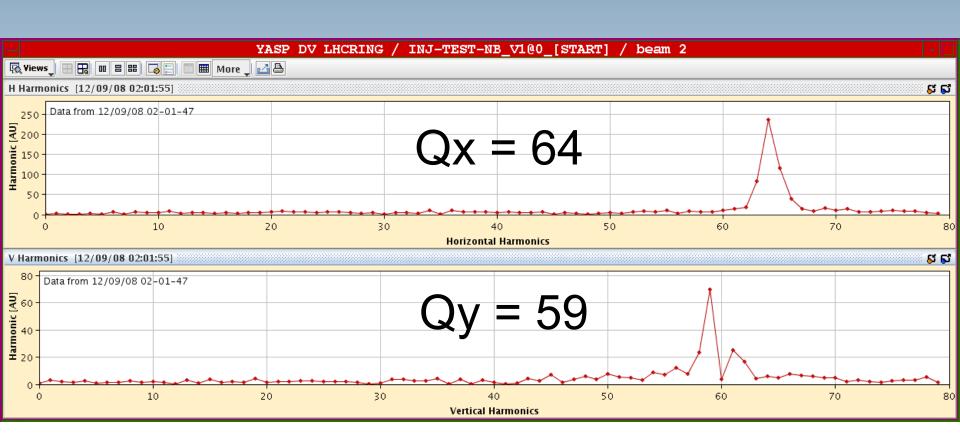
## II. LHC Operational cycle:Squeeze → reduce β\*





# II. Beam measurements: Integer tunes

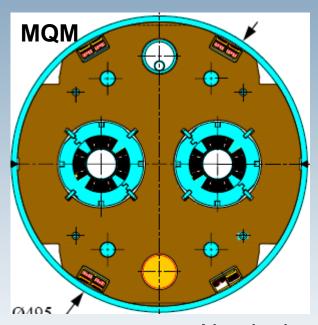




## I. Basic layout of the machine: Dispersion suppression



Quadrupole types: MQ, MQM, MQTL





Nominal gradient = 200/160 T/m

Inominal = 5.4/4.3 kA

Lmag=2.4/3.4/4.8 m

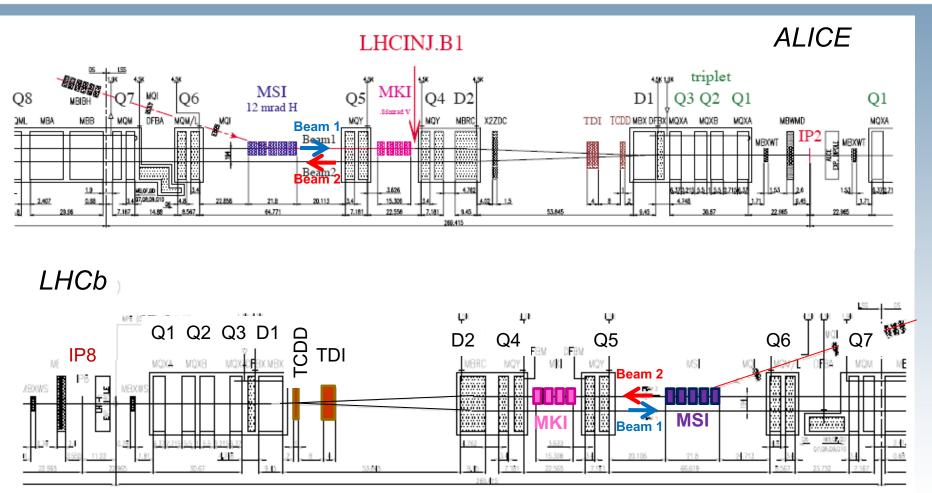
T=1.9/4.5 K

Cold bore  $\varnothing$  = 53/50 mm

**Individual powered apertures** 

## I. Basic layout of the machine: Low luminosity insertions



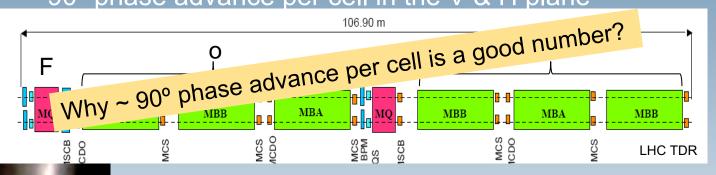


Extra challenge 

the lattice has to accommodate the injection region



LHC arc cells = FoDo lattice\* with ~ 90° phase advance per cell in the V & H plane



Which parameter determines the beam size (ignoring D(s))?

$$r^2 = \varepsilon_x \beta_x + \varepsilon_y \beta_y$$

(In general in proton machines εx≈εy → beams are round)

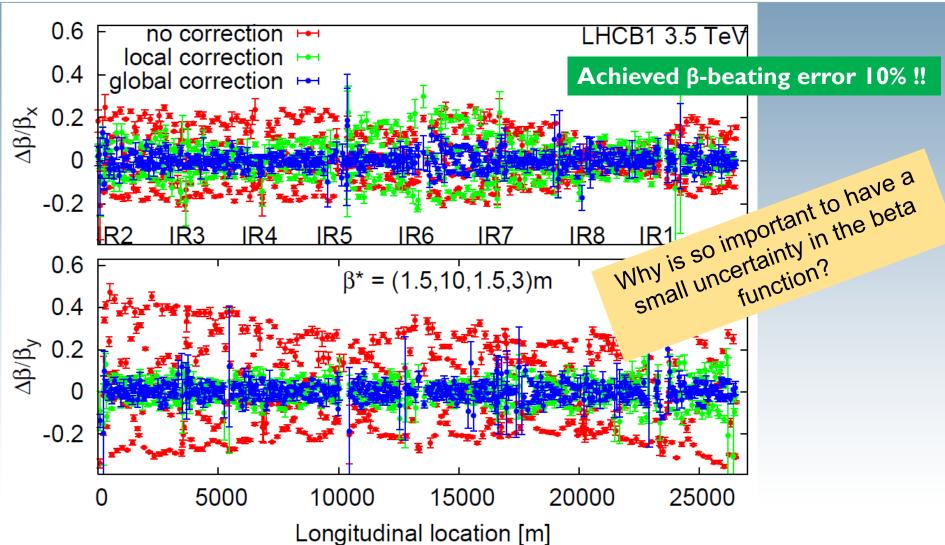
In order to get the maximum aperture possible the  $\beta(s)$  in both planes have to be minimized:

$$\frac{d}{d\mu}(\beta_{max} + \beta_{min}) = \frac{d}{d\mu}\left(\frac{\left(1 + sin\frac{\mu}{2}\right)L}{sin\mu} + \frac{\left(1 - sin\frac{\mu}{2}\right)L}{sin\mu}\right) = \frac{d}{d\mu}\left(\frac{2L}{sin\mu}\right) = 0$$

$$\frac{L}{\sin^2 u} \cos \mu = 0 \quad \Rightarrow \quad \mu = 90^{\circ}$$

### II. Beam measurements: Beta measurement





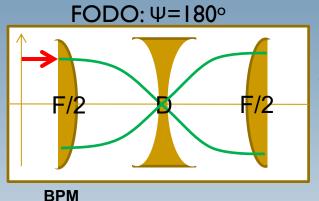
Reference: Record low beta beating in the LHC

Tomas, R. et al. Phy. Rev. Special Topics - Accelerators and Beams15(9).

### II. Beam measurements:



### Non-integer tunes

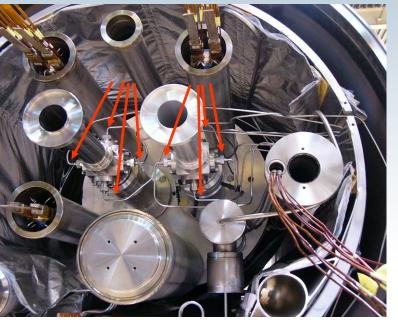


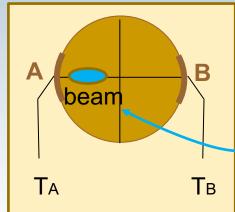
Every 2xT we measure the same position:

$$\frac{1}{2T} = f_Q = Q \cdot \frac{1}{T} = Q \cdot f_{rev}$$

$$\rightarrow Q = 0.5$$

BPM pos → FFT → tune





With a frequency of ~3 kHz we measure the same beam position



$$Qx = .279$$
 <sub>@LHC</sub>  $Qy = .310$