

# **$tWZ$ associated production at the LHC in the SMEFT** [to appear soon]

in collaboration with  
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# Outline

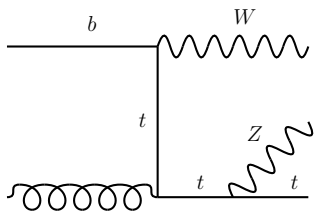
- Introduction
- The  $tWZ$  challenge at NLO
- $tWZ$  in the SM
- $tWZ$  in the SMEFT
- Summary

# Introduction

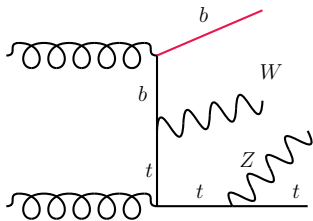
- The SM can not answer most of the pressing questions related to the top quark → **indications of answers might be manifested in new physics phenomena seen in the top-quark sector**
- $tWZ$  is a rare EW process → **potential probe of EW couplings that are not well measured, so far**
- $tWZ$  is sensitive to **unitarity-violating behaviour induced in its sub-amplitudes** via modified EW interactions [[1904.05637](#)] → can potentially serve as a test for the SM hypothesis
- New physics phenomena can be encapsulated in a higher dimensional operators → **SMEFT**
- Accurate theoretical predictions → **study  $tWZ$  at NLO**
- **$tWZ$  at NLO is non-trivial due to its overlap with other processes**

## $tWZ$ from LO to NLO in QCD

At NLO in QCD, the  $tWZb$  appears as a real emission final state

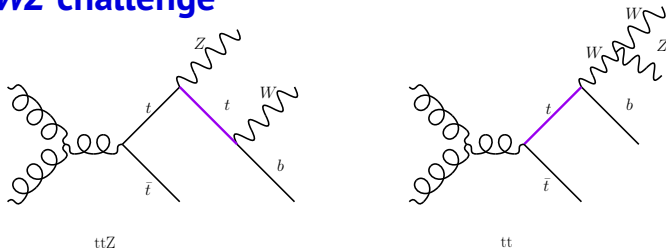


LO



NLO

## The $tWZ$ challenge



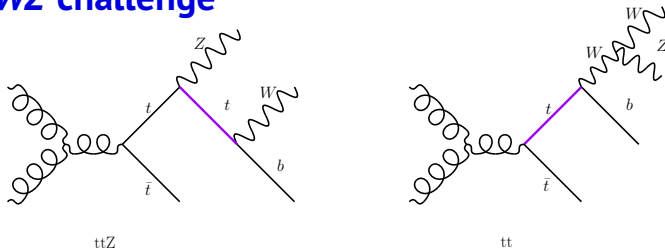
The  $tWZb$  final state can also have resonant contributions, from

$$ttZ, t \rightarrow Wb, \text{ or } tt, t \rightarrow WZb$$

and not necessarily the non-resonant  $tWZb$

- These topologies **do not belong to the genuine  $tWZ$  final state process** but to the leading order  $ttZ$  and  $tt$  processes
- The underlying resonant structure can spoil the perturbative behaviour of the NLO expansion

## The $tWZ$ challenge



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**Suppress contributions from the resonant amplitude!**

## The $tWZ$ challenge cont'd

The amplitude associated to the  $tWZ$  process can be expressed as

$$\mathcal{A}_{tWZ} = \mathcal{A}_{tWZ}^{res} + \mathcal{A}_{tWZ}^{res} \quad (1)$$

and thus the matrix element,

$$|\mathcal{A}_{tWZ}|^2 = |\mathcal{A}_{tWZ}^{res}|^2 + 2\Re\left(\mathcal{A}_{tWZ}^{res} \mathcal{A}_{tWZ}^{\dagger res}\right) + |\mathcal{A}_{tWZ}^{res}|^2 \quad (2)$$

Two Diagram Removal (DR) schemes to handle the resonant part of the matrix element:

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Two Diagram Removal (DR) schemes to handle the resonant part of the matrix element:

- Keep ONLY the non resonant contribution  $\rightarrow$  **DR1**



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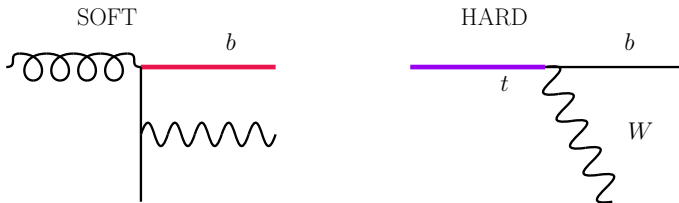
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Two Diagram Removal (DR) schemes to handle the resonant part of the matrix element:

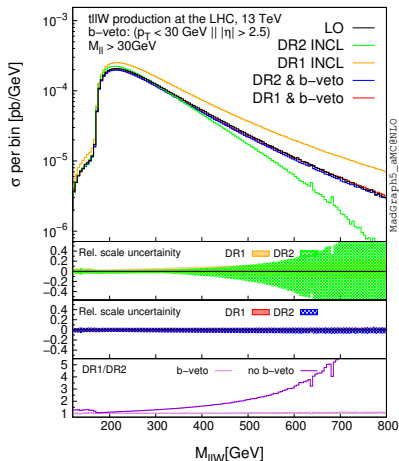
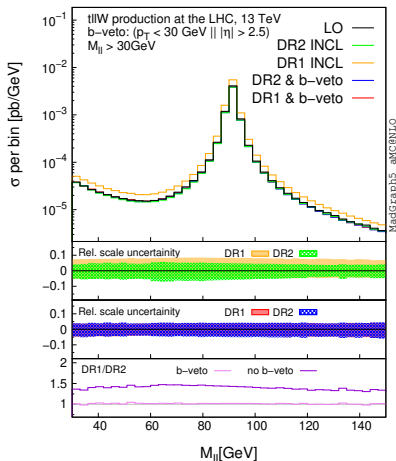
- Keep ONLY the non resonant contribution  $\rightarrow$  **DR1**
- Keep ALSO interference term  $\rightarrow$  **DR2**

## $tWZ$ in the SM

- The current MG5 diagram removal plugin (MadSTR) does not handle  $1 \rightarrow N$  decays where  $N > 2 \rightarrow$  **DR1 and DR2 schemes are implemented by hand**
- The **resonant part of the phase space is suppressed by vetoing hard  $b$ -quarks** as they tend to have come from the decay of a top
- For SM predictions  $\rightarrow$  both the  **$ttZ$**  and the  **$tt$**  overlaps are removed
- A good agreement between the DR1 and the DR2 schemes  $\rightarrow$  **the non-resonant part of  $tWZ$  dominates the phase space**



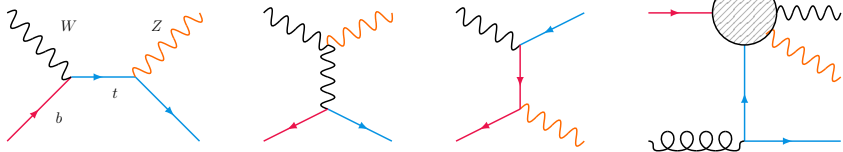
# $tWZ$ in the SM: differential results



The  $b$ -veto significantly improves the DR1-DR2 agreement

# $tWZ$ in the SMEFT

The  $bW \rightarrow tZ$  sub-amplitude:



and therefore the operators included in the analysis are

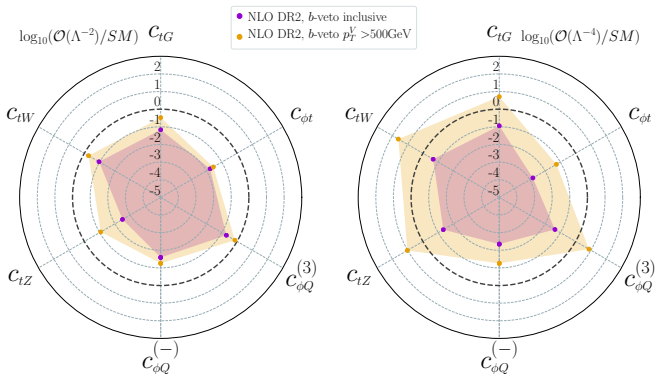
- $\mathcal{O}_{\varphi Q}^{(3)}$ ,  $\mathcal{O}_{\varphi Q}^{(-)}$  → interaction of two fermions fields with the gauge bosons
- $\mathcal{O}_{\varphi t}$  → the right handed  $ttZ$  interaction
- $\mathcal{O}_{tW}$ ,  $\mathcal{O}_{tZ}$  → interaction of the top with the weak isospin and the weak hypercharge gauge fields
- ( $\mathcal{O}_{tG}$  → the gluon-top interaction)

# Generic SMEFT predictions

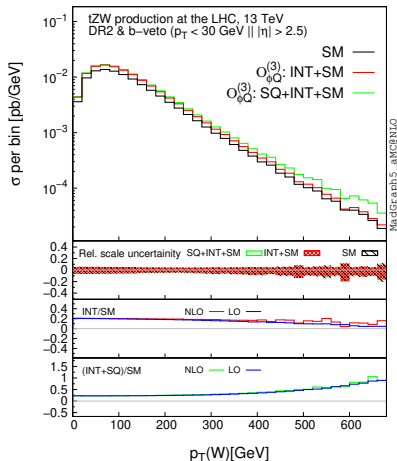
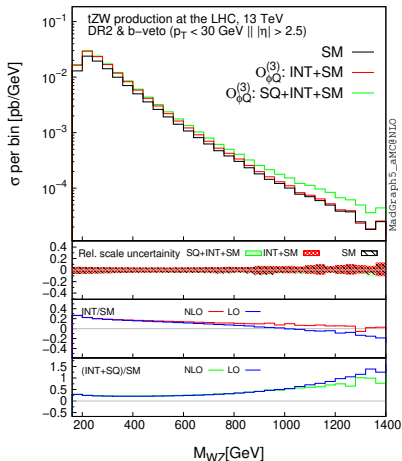
A generic observable in SMEFT can be expressed as

$$\sigma = \sigma_{SM} + \sum_i \frac{C_i^{(6)}}{\Lambda^2} \sigma_i + \sum_{ij} \frac{C_i^{(6)} C_j^{(6)}}{\Lambda^4} \sigma_{ij} \quad (3)$$

- second term → interference contributions of dim-6 operators
- last term → squared contributions of dim-6 operators



# $tWZ$ in the SMEFT: differential results



**SMEFT impacts from LO to NLO accuracy are stable suggesting the DR treatment is correctly identifying the phase space of the  $tWZ$  process**

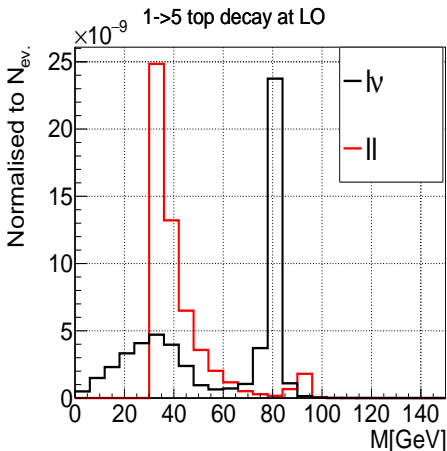
## Summary

- A study of  $tWZ$  production at NLO in QCD is presented in the context of the SM and SMEFT
- The overlap from  $tt$  and  $ttZ$  renders the  $tWZ$  at NLO a non-trivial process to study
- The resonant overlap is handled using the DR schemes at the amplitude-level
- Vetoing hard  $b$ -quarks ensures the  $tWZ$  process dominates the phase space after the diagram removal
- The differential results presented suggest the DR treatment correctly identifies the phase space of the  $tWZ$  process
- The work presented lays the foundation for precision LHC interpretations of  $tWZ$  data in the SMEFT framework

# BACKUP SLIDES

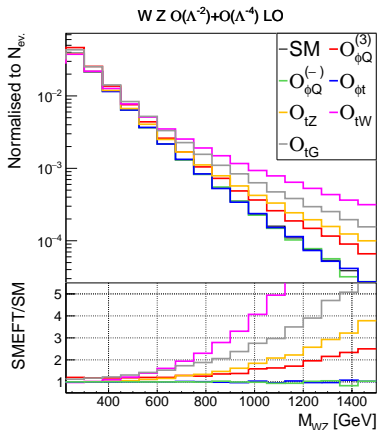
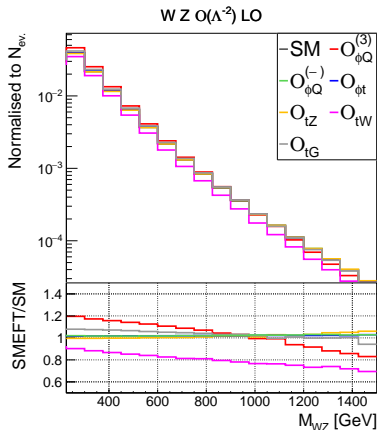


## Why decay the $Z$ in the SM case?



- The  $t \rightarrow ll\nu b$  decay is shown at LO
- This decay is what causes the overlap with  $tt$
- The  $Z$  ‘likes’ to be off-shell more than the  $W$
- To fairly treat the  $tt$  overlap, the  $Z$  should decay (keeping it on-shell suppresses the  $tt$  overlap)
- $Z$  is kept stable in the SMEFT study as the overlap from  $tt$  there is irrelevant

# $tWZ$ in the SMEFT: differential results[2]



$M_{WZ}$  is a proxy for the  $bW \rightarrow tZ$  sub-amplitude scattering energy

