

# The ScotoSinglet Model

Based on: *AB, J. H-García, N. Leerdam, M. White & A. G. Williams*  
[\[arXiv:2010.05937\]](https://arxiv.org/abs/2010.05937)

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When the M meets the P  
(20-21 January, 2021)

# Overview

- Motivation
- The ScotoSinglet Model (ScSM)
  - Model details
  - Analysis
  - Results
- Summary

# Motivation

- Two unsolved puzzles of the Standard Model (SM):
  - 1) Neutrino masses;
  - 2) Dark Matter (DM).

- Possible beyond the SM extensions:

- Radiative neutrino mass models;

[S. M. Boucenna, S. Morisi and J. W. Valle, *Adv. High Energy Phys.* (2014), [arXiv:1404.3751](#)]

[Y. Cai et al., *Front. In Phys.* (2017), [arXiv:1706.08524](#)]

- Weakly Interacting Massive Particles (WIMPs).

[J. L. Feng, *Ann. Rev. Astron. Astrophys.* (2010)]

- The Scotogenic Model (ScM) can address both puzzles *simultaneously*.

[E. Ma, *Phys. Rev. D.* (2006), [hep-ph/0601225](#)]

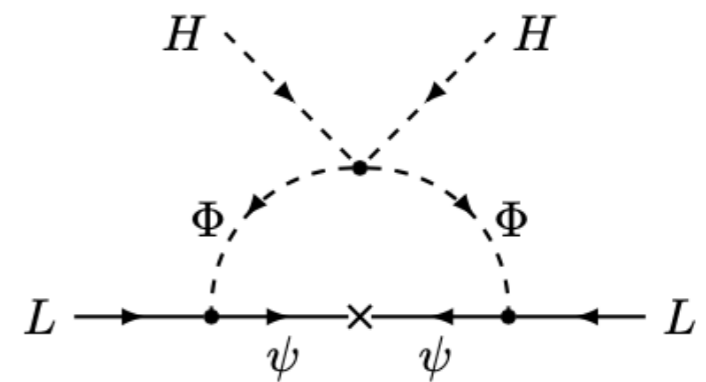
[J. Kubo, E. Ma and D. Suematsu, *Phys. Lett. B.* (20016), [hep-ph/0604114](#)]

# The Scotogenic Model (ScM)

Add a scalar doublet  $\Phi$  and new Majorana singlets  $\psi_i$  ( $i \geq 2$ ).

All fields charged under  $\mathbb{Z}_2$  symmetry:  $\Phi \rightarrow -\Phi$  and  $\psi_i \rightarrow -\psi_i$ .

- Neutrino masses generated at 1-loop level.
- DM can be either the lightest  $\psi_i$ , or lightest component of  $\Phi$  (CP-odd or CP-even).



Neutrino mass generation in the Scotogenic model at 1-loop level.

Strong constraints from  $\nu$  masses, **electroweak precision tests**, **direct detection** & **lepton flavour violation (LFV)** processes  $\implies$  non-trivial to saturate the observed DM abundance.

[C. Hagedorn et al., *JHEP* (2018), [arXiv:1804.04117](https://arxiv.org/abs/1804.04117)]

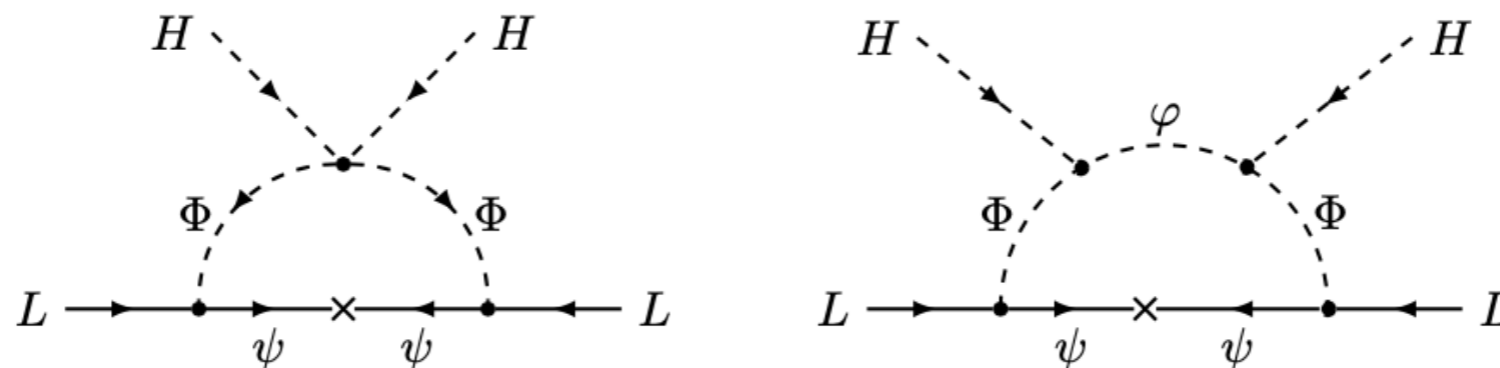
# The ScotoSinglet Model (ScSM)

**ScotoSinglet:** Scotogenic model with an extra scalar singlet.

Simplest extension of the Scotogenic Model.

Interesting features:

- New contribution to neutrino masses;
- 3 potential DM candidates;
- Possibility of maintaining  $\mathbb{Z}_2$  symmetry up to high-energy scales.



Neutrino masses in the ScSM at the 1-loop level. *Left panel:* contribution in the ScM. *Right panel:* new contribution in the ScSM.

# The ScotoSinglet Model (ScSM)

## Model details

Fields	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	$\mathbb{Z}_2$
Real $\varphi$	1	1	0	–
$\Phi$	1	2	1/2	–
$\psi_k$	1	1	0	–

Particle contents of the ScSM. Here  $k = 1, 2$ .

See [arXiv:2010.05937](https://arxiv.org/abs/2010.05937)  
for more details.

## Majorana fermion Lagrangian

$$\mathcal{L}_\Psi = \frac{1}{2} \bar{\Psi} (i\not{\partial} - M_\Psi) \Psi - \bar{\Psi} \mathbf{y}_\Psi \tilde{\Phi}^\dagger L + \text{H.c.},$$

$\mathbf{y}_\Psi = 2 \times 3$  matrix of complex Yukawa couplings

## Scalar potential

$$\begin{aligned} V = & -\mu_H^2 H^\dagger H + \lambda_H (H^\dagger H)^2 + m_\Phi^2 \Phi^\dagger \Phi + \lambda_\Phi (\Phi^\dagger \Phi)^2 + \frac{1}{2} m_\varphi^2 \varphi^2 + \frac{1}{4} \lambda_\varphi \varphi^4 \\ & + \lambda_{H\Phi,1} (H^\dagger H) (\Phi^\dagger \Phi) + \lambda_{H\Phi,2} (H^\dagger \Phi) (\Phi^\dagger H) + \frac{1}{2} [\lambda_{H\Phi,3} (H^\dagger \Phi)^2 + \text{H.c.}] \\ & + \frac{1}{2} \lambda_{H\varphi} H^\dagger H \varphi^2 + \frac{1}{2} \lambda_{\Phi\varphi} \Phi^\dagger \Phi \varphi^2 + [\kappa \Phi^\dagger H \varphi + \text{H.c.}]. \end{aligned}$$

$$\Phi \equiv \begin{pmatrix} \phi^+ \\ \frac{1}{\sqrt{2}} (\phi_R + iA) \end{pmatrix}.$$

# The ScotoSinglet Model (ScSM)

## Model details

After electroweak symmetry breaking & mass-diagonalisation in  $(\phi_R, \varphi)$  basis, namely

$$\begin{pmatrix} \eta_1 \\ \eta_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \phi_R \\ \varphi \end{pmatrix}$$

$\Rightarrow$  **21 parameters** (12 free + 9 nuisance parameters):

$$5 \text{ masses (1 mixing angle)} : \{m_{\psi_1}, m_{\psi_2}, m_{\eta_1}, m_{\eta_2}, m_A, \theta\}, \quad (3.1a)$$

$$6 \text{ couplings: } \{\lambda_\Phi, \lambda_\varphi, \lambda_{H\Phi,1}, \lambda_{H\Phi,2}, \lambda_{H\varphi}, \lambda_{\Phi\varphi}\}. \quad (3.1b)$$

6 complex Yukawas  $\mathbf{y}_\Psi$  (9 nuisance parameters) fixed using Casas-Ibarra parametrisation.

[J. A. Casas & A. Ibarra, *Nucl. Phys. B.* (2001), [hep-ph/0103065](#)]

# The ScotoSinglet Model (ScSM)

## Analysis

- Use MultiNest 3.10.0 to sample the 21-dimensional parameter space.

[F. Feroz, M. Hobson & M. Bridges, *Mon. Not. Roy. Astron. (2009)*, [arXiv:0809.3437](#)]

- Perform a frequentist analysis, specifically *Profile Likelihood Ratio (PLR)*.

$$\ln \mathcal{L}_{\text{total}}(\boldsymbol{\theta}) = \ln \mathcal{L}_{\kappa}(\boldsymbol{\theta}) + \ln \mathcal{L}_{\text{EWPT}}(\boldsymbol{\theta}) + \ln \mathcal{L}_{\mathcal{R}_{\gamma\gamma}}(\boldsymbol{\theta}) \\ + \ln \mathcal{L}_{\text{LFV}}(\boldsymbol{\theta}) + \ln \mathcal{L}_{\Omega h^2}(\boldsymbol{\theta}) + \ln \mathcal{L}_{\text{DD}}(\boldsymbol{\theta})$$

Constraints include:

- 1) Perturbativity and naturalness;
- 2) Electroweak Precision Tests (EWPT);
- 3)  $\mathcal{R}_{\gamma\gamma} \equiv \Gamma(h \rightarrow \gamma\gamma)_{\text{ScSM}} / \Gamma(h \rightarrow \gamma\gamma)_{\text{SM}}$ ;
- 4) Lepton Flavour Violation (LFV) processes;
- 5) DM Relic density ( $\Omega h^2$ );
- 6) Direct Detection (DD).

	Ranges	Priors
<b>Model parameters</b>		
$\{m_{\psi_1}, m_{\psi_2}\}$ (GeV)	[10, 10 <sup>6</sup> ]	Log
$m_{\eta_2}$ (GeV)	[100, 10 <sup>4</sup> ]	Log
$\delta_1 \equiv m_{\eta_1} - m_{\eta_2}$ (GeV)	[10 <sup>-3</sup> , 10 <sup>4</sup> ]	Log
$\delta_A \equiv m_A - m_{\eta_2}$ (GeV)	$[-10^4, -10^{-3}] \cup [10^{-3}, 10^4]$	Log  value
$\theta$ (rad.)	[0, $\pi$ ]	Flat
$\{\lambda_{\Phi}, \lambda_{\varphi}\}$	[10 <sup>-3</sup> , 4 $\pi$ ]	Log
$\{\lambda_{H\Phi,1}, \lambda_{H\Phi,2}, \lambda_{H\varphi}, \lambda_{\Phi\varphi}\}$	$[-4\pi, -10^{-3}] \cup [10^{-3}, 4\pi]$	Log  value
<b>Nuisance parameters</b>		
$\sin^2 \theta_{12}$	[0.275, 0.350]	Flat
$\sin^2 \theta_{13}$	[0.02044, 0.02435] (NO) [0.02064, 0.02457] (IO)	Flat
$\sin^2 \theta_{23}$	[0.433, 0.609] (NO) [0.436, 0.610] (IO)	Flat
$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$	[6.79, 8.01]	Flat
$\frac{\Delta m_{3l}^2}{10^{-3} \text{ eV}^2}$	[2.436, 2.618] (NO) [-2.601, -2.419] (IO)	Flat
$\delta_{\text{CP}}$ (°)	[144, 357] (NO) [205, 348] (IO)	Flat
$\alpha$ (rad.)	[0, 2 $\pi$ ]	Flat
$\{\zeta_1, \zeta_2\}$	[10 <sup>-3</sup> , 10 <sup>3</sup> ]	Log

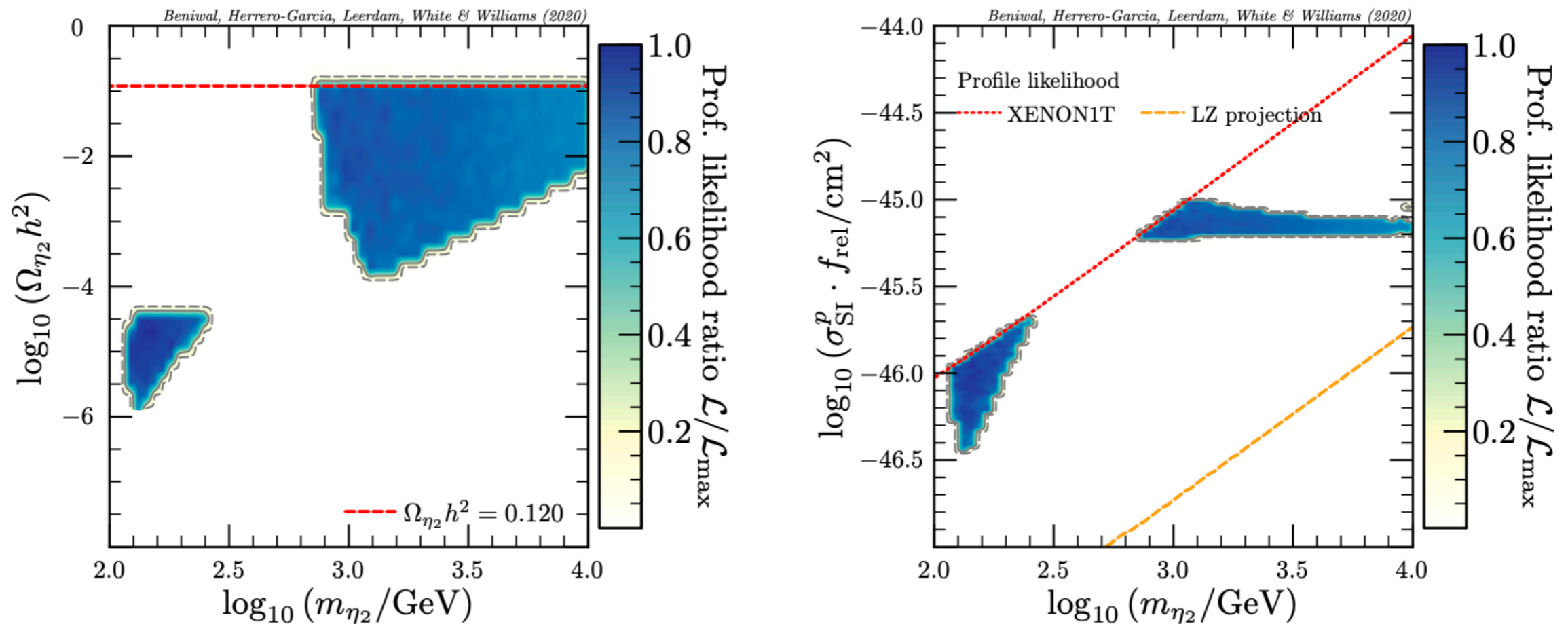
Parameter ranges and priors used in our analysis.



# The ScotoSinglet Model (ScSM)

## Results

Case 1:  $\theta = 0 \implies \eta_2 = \text{scalar singlet}$



Left panel: relic density vs  $\eta_2$  mass. Right panel: effective spin-independent  $\eta_2$ -proton cross-section vs  $\eta_2$  mass.

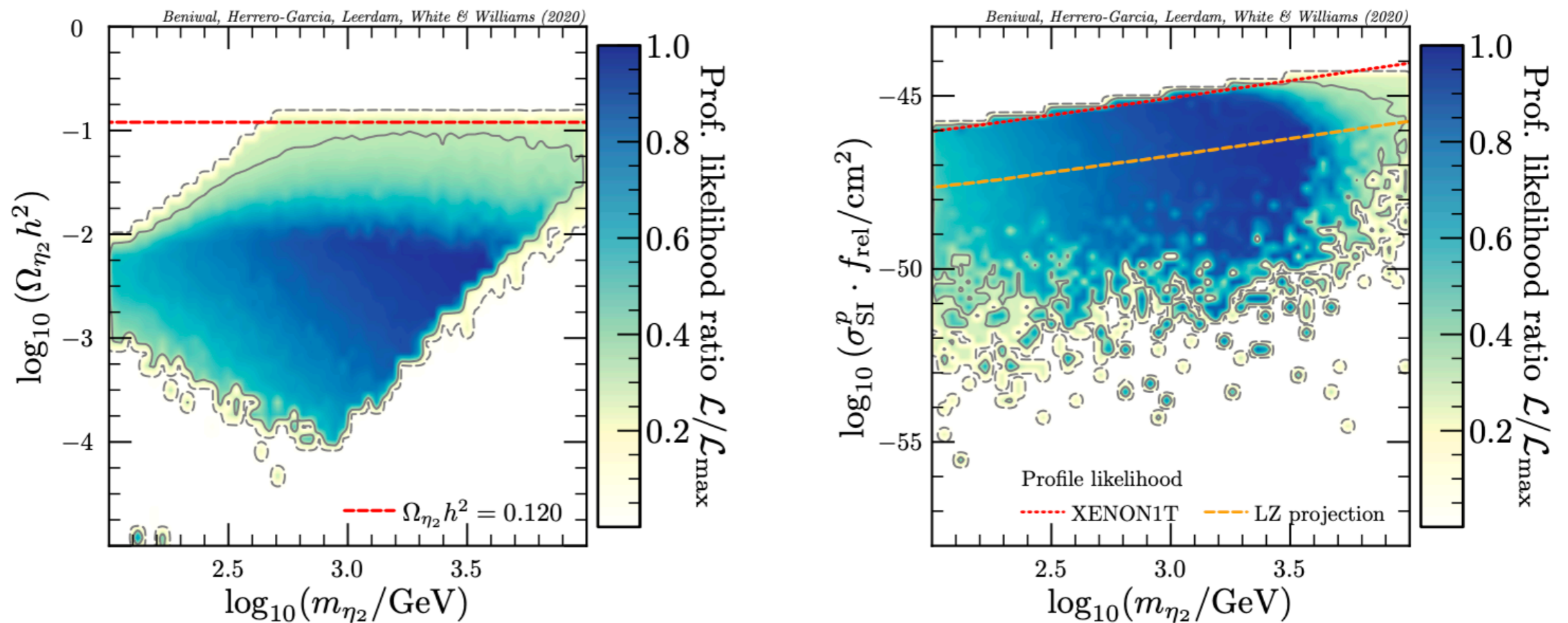
[P. Athron et al., *Eur. Phys. J. C.* (2018), [arXiv:1806.11281](#)]

[P. Athron et al., *Eur. Phys. J. C.* (2019), [arXiv:1808.10465](#)]

# The ScotoSinglet Model (ScSM)

## Results

**Case 2:**  $\theta = \pi/2 \implies \eta_2 = \text{scalar doublet}$

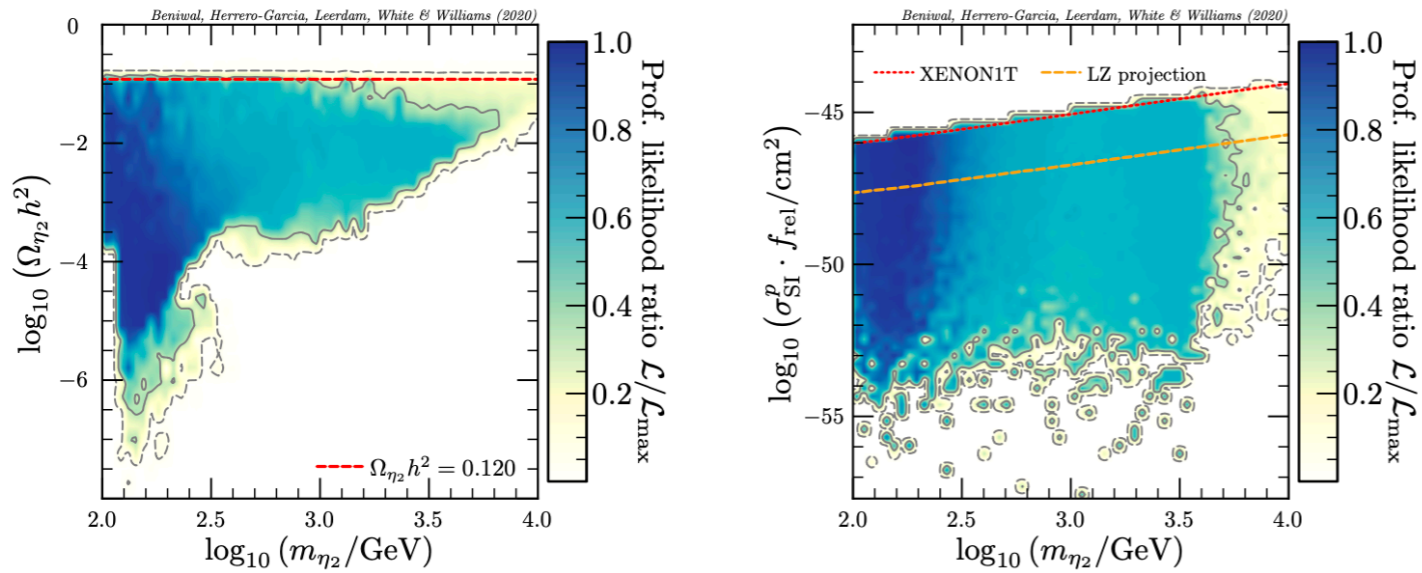


*Left panel:* relic density vs  $\eta_2$  mass. *Right panel:* effective spin-independent  $\eta_2$ -proton cross-section vs  $\eta_2$  mass.

[L. L. Honorez et al., *JCAP* (2007), [hep-ph/0612275](#)]

# The ScotoSinglet Model (ScSM)

## Results

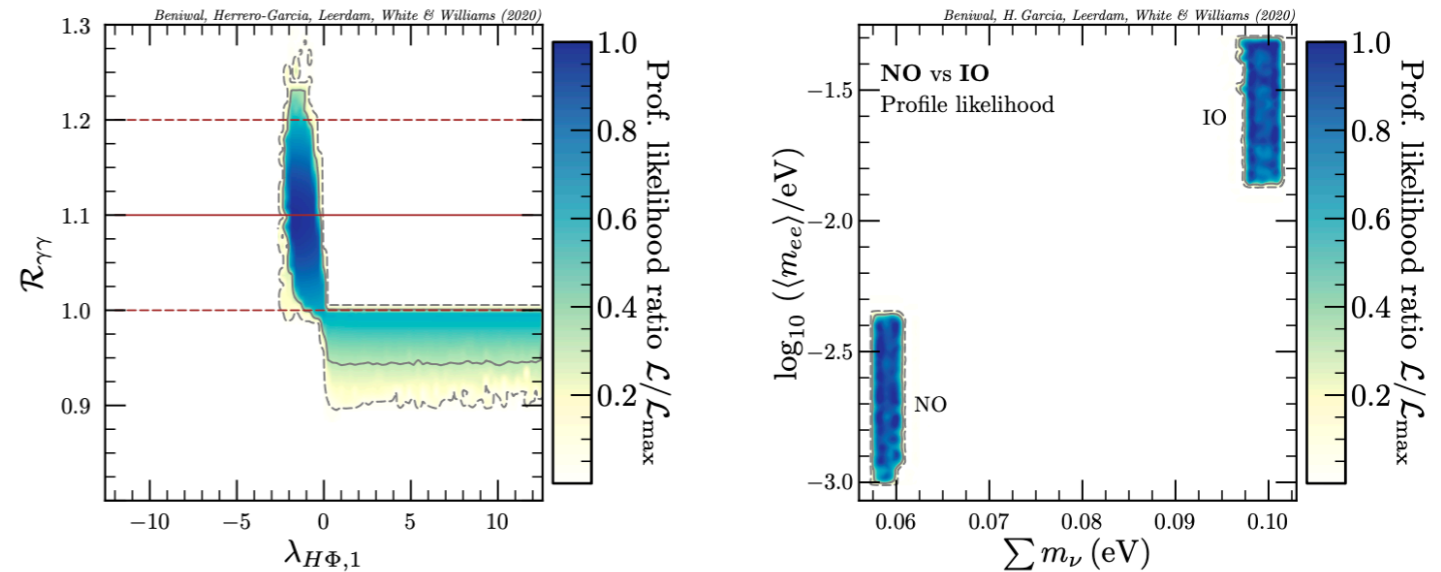


**Case 3: Non-zero mixing**  
( $\theta \neq 0, \pi/2$ )

Left panel: relic density vs  $\eta_2$  mass. Right panel: effective spin-independent  $\eta_2$ -proton cross-section vs  $\eta_2$  mass.

$$\mathcal{R}_{\gamma\gamma} \equiv \frac{\Gamma(h \rightarrow \gamma\gamma)_{\text{ScSM}}}{\Gamma(h \rightarrow \gamma\gamma)_{\text{SM}}} = \left| 1 + \frac{\lambda_{H\Phi,1} v^2}{2m_{\phi^+}^2} \frac{A_0(\tau_{\phi^+})}{A_1(\tau_W) + \frac{4}{3}A_{1/2}(\tau_t)} \right|^2$$

$$\langle m_{ee} \rangle \equiv \sum |U_{ei}^2 m_i|$$



Left panel:  $\mathcal{R}_{\gamma\gamma}$  vs  $\lambda_{H\Phi,1}$ . Right panel:  $\langle m_{ee} \rangle$  vs  $\sum m_\nu$

# Summary

- Proposed a simple variant of the ScM, namely ScSM.
- Interesting features of ScSM:
  - Scalar DM (CP-even or CP-odd) with suppressed direct detection rates **and** correct DM abundance;
  - Two contributions to  $\nu$  masses: usual scotogenic one ( $\propto \lambda_{H\Phi,3}$ ) + new singlet one ( $\propto \kappa^2$ ).
  - Singlet improves the  $\mathbb{Z}_2$  symmetry at high-energy scales.

For more details, see [arXiv:2010.05937](https://arxiv.org/abs/2010.05937).

**Thank you!**

# **Backup slides**

# Dependent parameters

$$m_{\Phi}^2 = m_{\eta_1}^2 \cos^2 \theta + m_{\eta_2}^2 \sin^2 \theta - \frac{1}{2} (\lambda_{H\Phi,1} + \lambda_{H\Phi,2} + \lambda_{H\Phi,3}) v^2, \quad (3.2a)$$

$$m_{\varphi}^2 = m_{\eta_1}^2 \sin^2 \theta + m_{\eta_2}^2 \cos^2 \theta - \frac{1}{2} \lambda_{H\varphi} v^2, \quad (3.2b)$$

$$m_{\phi^+}^2 = m_{\Phi}^2 + \frac{1}{2} \lambda_{H\Phi,1} v^2, \quad (3.2c)$$

$$\lambda_{H\Phi,3} = \frac{1}{v^2} (m_{\eta_1}^2 \cos^2 \theta + m_{\eta_2}^2 \sin^2 \theta - m_A^2), \quad (3.2d)$$

$$\kappa = \frac{1}{v} (m_{\eta_1}^2 - m_{\eta_2}^2) \sin \theta \cos \theta. \quad (3.2e)$$