

# Multiple Emission Kernels for Parton Showers

Emma Simpson Dore<sup>1</sup>

MCnet meeting, Louvain | January 25, 2019

---

<sup>1</sup>In collaboration with Simon Plätzer and Stefan Gieseke

# What does a parton shower do?

- Bridge between hard matrix element and non-perturbative physics
- Facilitates resummation of leading logarithmic contributions
- Uses emission kernels to describe rate of emission for a propagating quark/gluon to emit an additional quark/gluon
- Contains mapping to factorise emissions within the phase space

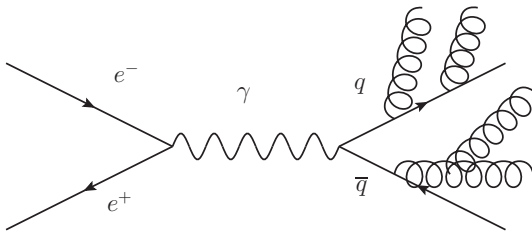


Diagram of an  $e^+e^-$  collision and possible parton shower

# The future of parton showers

Push to higher orders and improved accuracy :

- Design new parton shower, beyond  $1 \rightarrow 2$  branching
- Global recoil scheme for one and more emissions
- Include more spin and colour interferences
- Use higher-order emission kernels

Will be central to address:

- Singularity structure for NNLO matching, increasingly important
- Lack of systematic expansion of uncertainties at higher orders
- Issues with local recoil schemes<sup>2</sup>

---

<sup>2</sup>Dasgupta et al. 2018.

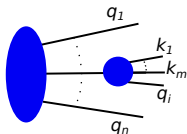
# What do we need to build the parton shower?

- Start from hard process with  $n$  external legs, momenta  $q_i$  is emitter with  $m$  emissions  $k_l$ :

Diagrams with emission from external leg  $q_i$ :

$$\mathbf{Sp}(q_i, k_1 \dots k_m) |\mathcal{M}(q_1, \dots, q_{i-1}, q_i + \sum_{l=1}^m k_l, q_{i+1}, \dots, q_n)\rangle$$

- Where  $\mathbf{Sp}$  is an operator in spin and colour space
- For full matrix element need to sum over all emitters  $\sum_i q_i$
- Develop a new mapping in order to:
  - Distribute recoils globally
  - Factorise without explicitly taking soft or collinear limits



# Mapping for multiple emissions with global recoil

Mapping with both soft and collinear cases included:

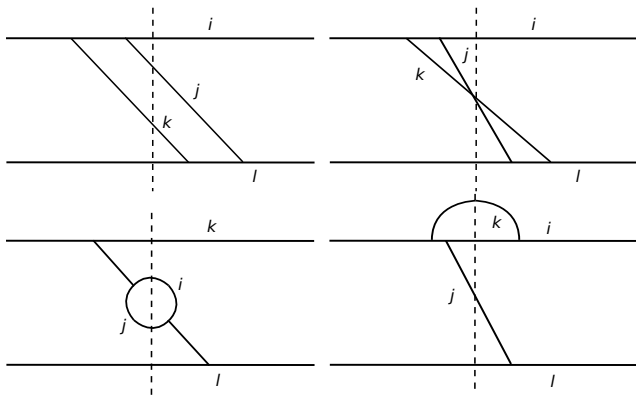
$$q_i^\mu = (1 - \sum_{l=1}^m \alpha_l) \alpha \Lambda^\mu{}_\nu p_i^\nu + y(1 - \sum_{l=1}^m \beta_l) n^\mu - \sum_{l=1}^m \sqrt{y \alpha_l \beta_l} n_{\perp, l}^\mu,$$

$$k_j^\mu = \alpha_l \alpha \Lambda^\mu{}_\nu p_i^\nu + y \beta_l n^\mu + \sqrt{y \alpha_l \beta_l} n_{\perp, l}^\mu,$$

$$q_k^\mu = \alpha \Lambda^\mu{}_\nu p_k^\nu, \quad (k = 1, \dots, n \quad k \neq i)$$

- Above case is single emitter  $q_i$  with  $m$  emissions  $k_l$ , multiple emitter case possible
- Includes soft limit ( $\alpha_l, y \rightarrow 0$ ) and collinear limit ( $y \rightarrow 0$ )
- $n_\perp$  represents the transverse component
- Includes global treatment of recoil via Lorentz transformation  $\Lambda^\mu{}_\nu$

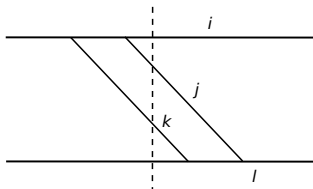
## Focus on two emission case



Examples of 2 emission diagrams

- Complicated interaction of soft and collinear singularities for two emissions
- Devise framework for separating singularities

## Example: Separating singularities

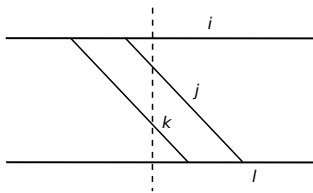


Singularities in the above diagram:

$$\frac{1}{S_{ijk}} \frac{1}{S_{ij}} \frac{1}{S_{kl}} \frac{1}{S_{jkl}}$$

- Where  $S_{ijk} = S_{ij} + S_{ik} + S_{jk}$  and  $S_{ij} = (q_i + q_j)^2$
- Aim to partition singularities so can be expressed as a sum

## Example: Separating singularities



- Factors used for this case, can be generalised to other diagrams:

$$\frac{S_{ijk} + S_{kl}}{S_{ijk} + S_{kl}} \frac{S_{jkl} + S_{ij}}{S_{jkl} + S_{ij}} = \frac{S_{ijk}}{S_{ijk} + S_{kl}} \left[ \frac{S_{jkl}}{S_{jkl} + S_{ij}} + \frac{S_{ij}}{S_{jkl} + S_{ij}} \right] + \frac{S_{kl}}{S_{ijk} + S_{kl}} \left[ \frac{S_{jkl}}{S_{jkl} + S_{ij}} + \frac{S_{ij}}{S_{jkl} + S_{ij}} \right]$$

- The following limits are contained:

- Triple collinear  $S_{ijk} \rightarrow 0$ ,  $S_{kl}, S_{jl} \not\rightarrow 0$
- Triple collinear  $S_{jkl} \rightarrow 0$ ,  $S_{ij}, S_{ik} \not\rightarrow 0$
- Two collinear pairs  $S_{ij}S_{kl} \rightarrow 0$ ,  $S_{ik}, S_{jk}, S_{jl} \not\rightarrow 0$

- The factorisation is safe, only one term contains triple collinear limit



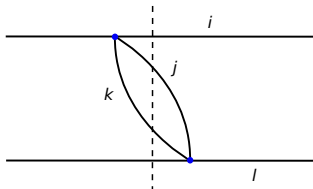
## Partitioning framework

- For the two emissions from one parton case with one spectator, we have four possible triplets:

$$\sum_{i < j < k < l} (V_{ijk,l} + V_{ijl,k} + V_{ikl,j} + V_{jkl,i})$$

- Each term contains sum of diagrams for that triplet with only the relevant singularities i.e.  $V_{ijk,l}$  only contains  $S_{ijk}$  singular contributions
- Use partitioning factors, each term contributes to a different triplet:

$$1 = w_{ijk,jkl} + w_{jkl,ijk} = \frac{S_{jkl}}{S_{ijk} + S_{jkl}} + \frac{S_{ijk}}{S_{ijk} + S_{jkl}}$$



# Fitting all the pieces together

Bulding blocks:

- Mapping for n-emission case
- Combinatorics to give all possible diagrams
- Partitioning framework to add together singular contributions
- Expansion to relevant order in  $\alpha_s$  of matrix elements squared

Example: Triplet spitting for 2 emission case,  $\mathcal{O}(\alpha_s)$

$$\langle n | \mathcal{M}(n) \rangle = g_s^2 \sum_{i < j < k}^n \sum_{\beta} \langle i, j, k | \mathbf{Sp}(i, j, k, i + j + k) \otimes \langle [n]_{ijk} | \mathcal{M}([n]_{ijk}, (i + j + k)) \rangle$$

- First test, emission kernel for 2 emission case:
  1. Evaluate all diagrams including self energies
  2. Combine according to partitioning framework
  3. Include mapping and phase space

# Summary

Current status:

- New mapping which exposes soft and collinear singularities for up to  $n$ -emissions
- Comprehensive framework for organising the singularities up to 2-emission case
- Partial implementation in Mathematica

What's next?

- Implement mapping for 2-emission case
- Improve diagrammatic algorithmic procedure
- Test implementation in Herwig parton shower

# Backup slides

# Determination of emission kernels

- Need to take collinear and soft limits which allow factorisation<sup>3</sup>
- Where  $q_k$  is momentum of final state parton  $k$ :
  - a) Soft limit  $q_k = \lambda q$ ,  $\lambda \rightarrow 0$ ,  $|\mathcal{M}_{m+1,a...}|^2 \propto 1/\lambda^2$
  - b) Collinear limit  $q_k \rightarrow (1-z)q_i/z$ ,  
 $|\mathcal{M}_{m+1,a...}|^2 \propto 1/q_i \cdot q_k$

## Dipole factorisation

- Consider  $(m+1)$  partons, factorise out parton  $k$  to give  $|\mathcal{M}_{m,a...}|^2$ 
$$|\mathcal{M}_{m+1,a...}|^2 \rightarrow |\mathcal{M}_{m,a...}|^2 \otimes V_{ik,j}$$
- $V_{ik,j}$  = singular factor including parton  $k$  and its interaction with partons  $i$  and  $j$  from the  $m$  parton amplitude

<sup>3</sup>Catani and Seymour 1997.

## Mapping in two emission case





$$q_i^\mu = (1 - (\alpha_1 + \alpha_2)) \alpha \Lambda^\mu{}_\nu p_i^\nu + y(1 - (\beta_1 + \beta_2)) n^\mu \\ - \sqrt{y\alpha_1\beta_1} n_{\perp,1}^\mu - \sqrt{y\alpha_2\beta_2} n_{\perp,2}^\mu,$$

$$k_1^\mu = \alpha_1 \alpha \Lambda^\mu{}_\nu p_i^\nu + y\beta_1 n^\mu + \sqrt{y\alpha_1\beta_1} n_{\perp,1}^\mu,$$

$$k_2^\mu = \alpha_2 \alpha \Lambda^\mu{}_\nu p_i^\nu + y\beta_2 n^\mu + \sqrt{y\alpha_2\beta_2} n_{\perp,2}^\mu,$$

$$q_k^\mu = \alpha \Lambda^\mu{}_\nu p_k^\nu, \quad (k = 1, \dots, n \quad k \neq i)$$

# References I

-  Guido Altarelli and G. Parisi. “Asymptotic Freedom in Parton Language”. In: *Nucl. Phys.* B126 (1977), pp. 298–318. DOI: 10.1016/0550-3213(77)90384-4.
  -  S. Catani and M. H. Seymour. “A General algorithm for calculating jet cross-sections in NLO QCD”. In: *Nucl. Phys.* B485 (1997). [Erratum: *Nucl. Phys.* B510,503(1998)], pp. 291–419. DOI: 10.1016/S0550-3213(96)00589-5, 10.1016/S0550-3213(98)81022-5. arXiv: hep-ph/9605323 [hep-ph].
  -  Mrinal Dasgupta et al. “Logarithmic accuracy of parton showers: a fixed-order study”. In: (2018). arXiv: 1805.09327 [hep-ph].
  -  Simon Plätzer and Malin Sjödaahl. “The Sudakov Veto Algorithm Reloaded”. In: *Eur. Phys. J. Plus* 127 (2012), p. 26. DOI: 10.1140/epjp/i2012-12026-x. arXiv: 1108.6180 [hep-ph].
- QCD and Collider Physics, Ellis, Stirling and Webber. (p.170)