

A new tool for mastering the Singular Limits

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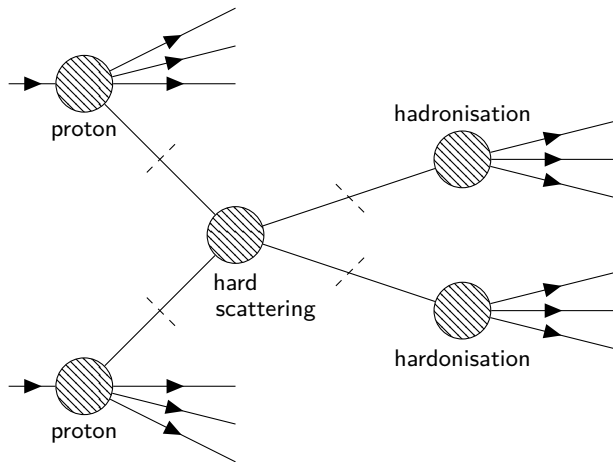
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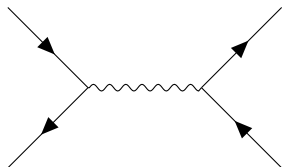
LHC Event Factorisation



$$|\mathcal{M}|^2 = \mathcal{A}_1 + \alpha_S \mathcal{A}_2 + \alpha_S^2 \mathcal{A}_3 + \dots$$

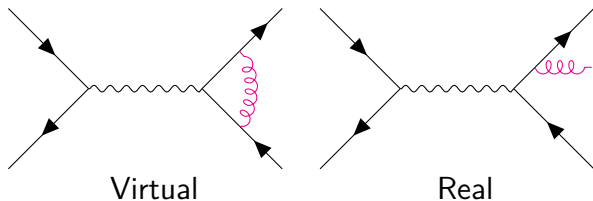
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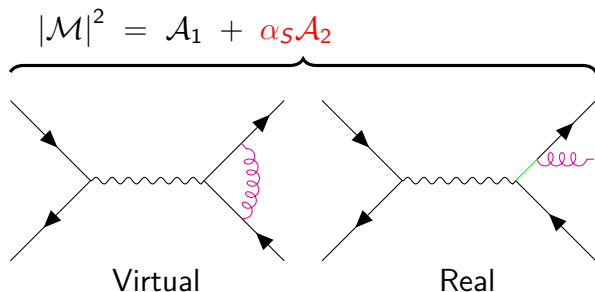


Born

$$|\mathcal{M}|^2 = \mathcal{A}_1 + \alpha_s \mathcal{A}_2$$



Problem!



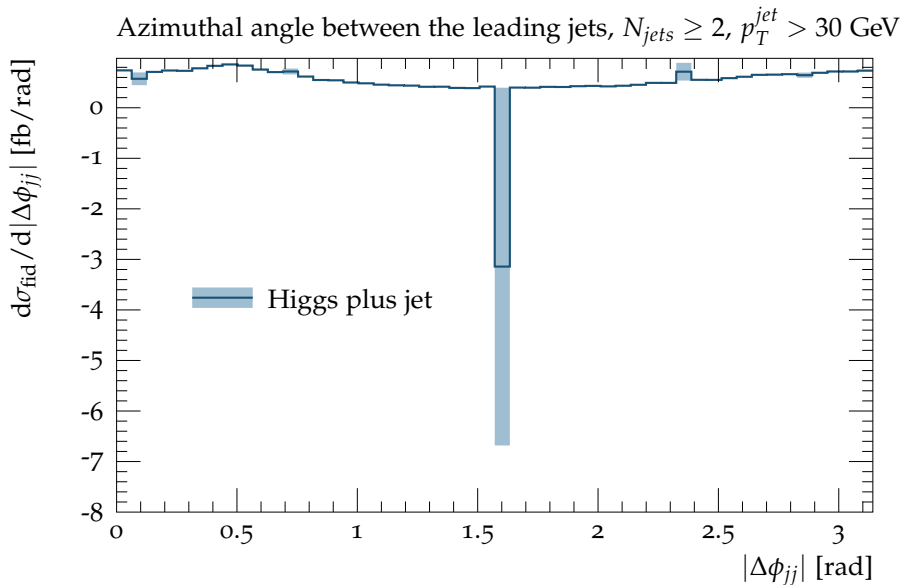
- Propagator introduces factor: $\frac{1}{(p_g + p_q)^2} = \frac{1}{2E_g E_q (1 - \cos \theta_{qg})}$
- Goes singular as gluon or quark go soft or collinear
- Resolve this using symmetry between virtual and real

Subtraction terms

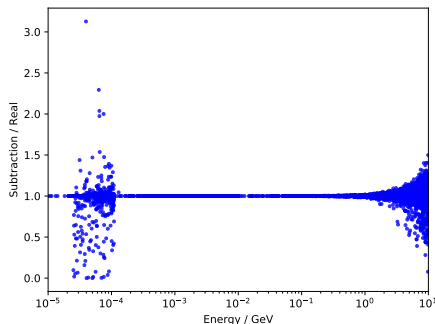
$$\begin{aligned}d\sigma^{NLO} &= d\phi_n [B + V] + d\phi_{n+1} R \\&= d\phi_n [B + V] + d\phi_{n+1} R + d\phi_{n+1} S - d\phi_{n+1} S \\&= d\phi_n \left[B + V + \int_{\phi_1} S \right] + d\phi_{n+1} [R - S]\end{aligned}$$

- The terms labelled S are called the Subtraction Dipoles
- They are defined to mimic the singular behaviour of R/V and be analytically integrable over ϕ_1
- Both integrands are now finite and can be Monte Carlo integrated 🙌

Mismatch between singularities



Energy of softest gluon in event against S/R for that event:



- $\sim 10^{-4}$ GeV is the problem area
- Difficult to work out much more though (other than problem is not numerical)
- Points are smeared because of randomised momenta

The SingularPhasespace approach

- Generate events with progressively lower gluon energy
- Only adjust other particles for the sake of momentum conservation
- Scatter plot should now trace out smooth line
- Allow user to choose other singular limits too

Conserving momentum

- 2 parton final state is easy: CoM frame, angles are constant, just vary energy
- In the >2 case though it is significantly more fiddly
- For example, $gg \rightarrow ggH_0$:

$$\begin{pmatrix} E_1 \\ E_1 \cos \phi_1 \sin \theta_1 \\ E_1 \sin \phi_1 \sin \theta_1 \\ E_1 \cos \theta_1 \end{pmatrix} + \begin{pmatrix} E_2 \\ E_2 \cos \phi_2 \sin \theta_2 \\ E_2 \sin \phi_2 \sin \theta_2 \\ E_2 \cos \theta_2 \end{pmatrix} + \begin{pmatrix} E_3 \\ \sqrt{E_3^2 - m_H^2} \cos \phi_3 \sin \theta_3 \\ \sqrt{E_3^2 - m_H^2} \sin \phi_3 \sin \theta_3 \\ \sqrt{E_3^2 - m_H^2} \cos \theta_3 \end{pmatrix} = \text{const}$$

- Thankfully though I can exploit someone else's hard work

Equation from before:

$$d\sigma^{NLO} = d\phi_n \left[B + V + \int_{\phi_1} S \right] + d\phi_{n+1} [R - S]$$

- Do not want a new S for each new matrix element
- Thankfully they can be written in the form of $d\sigma_B \otimes dA$
- But we have just generated a phase space point of ϕ_{n+1} - what do we give to $d\sigma_B(\phi_n)$?
- Hence, with the invention of Subtraction terms came a transformation $\{p_i, p_j, p_k\} \rightarrow \{\tilde{p}_{ij}, \tilde{p}_k\}$ that conserves momentum + virtuality

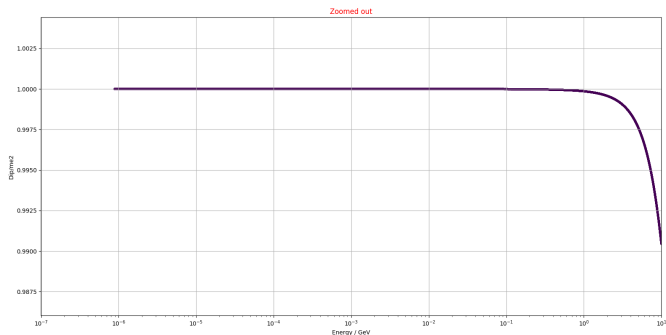
Massless case

$$\tilde{p}_k^\mu = \frac{1}{1 - y_{ij,k}} p_k^\mu$$
$$\tilde{p}_{ij}^\mu = p_i^\mu + p_j^\mu - \frac{y_{ij,k}}{1 - y_{ij,k}} p_k^\mu$$

- Soft limit corresponds to $y_{ij,k} \rightarrow 0$
- Plan: $3 \rightarrow 2 \rightarrow 2' \rightarrow 3'$
- Similar but more complex equations for non-massive case and collinear limits

Summary

- Regulating NLO singularities is an important but subtle task
- SingularPhasespace will allow one to explore these limits in a controlled manner
- Momentum conservation will be achieved with some help from the Dipole Subtraction formalism
- Cheeky proof of concept:



Thanks!