

# Hard diffraction in photoproduction

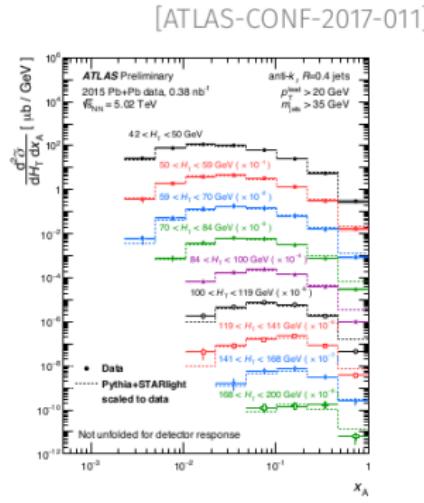
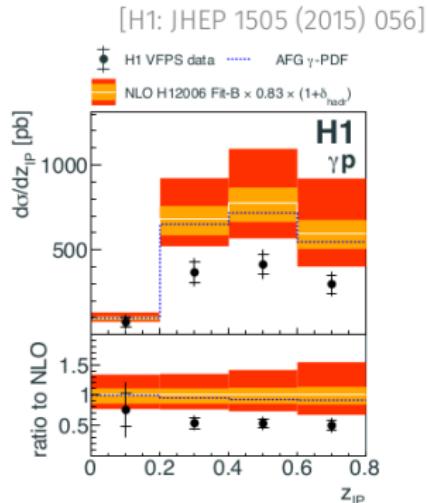
Christine O. Rasmussen

Dept. of Astronomy and Theoretical Physics

- Motivation
- Photoproduction framework in PYTHIA 8
- Hard diffraction in ep
- Hard diffraction in UPCs
- Conclusion and outlook

# Motivation

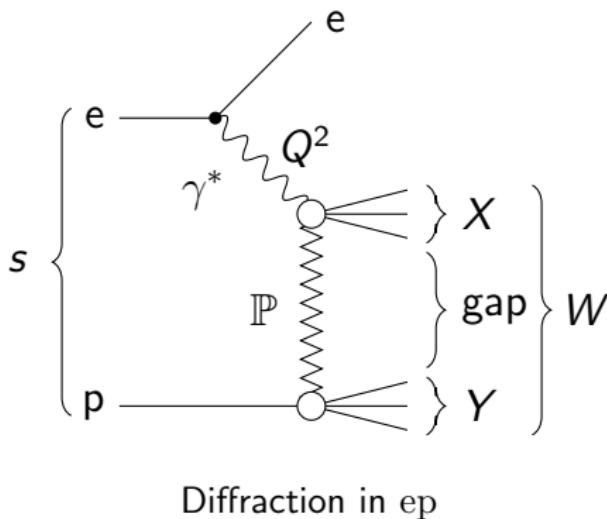
- Explanation for factorisation breaking in photoproduction regime in ep collisions.
- ATLAS feasibility study of dijets in ultraperipheral collisions (UPCs). Wish to extend to diffractive dijets.



Based on arXiv:1901.05261[hep-ph] with I. Helenius.

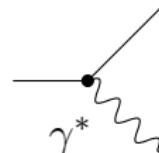
# Photoproduction framework

- Photoproduction: Low virtuality photons,  $Q^2 \lesssim 1 \text{ GeV}^2$
- Factorize photon flux,  $f_\gamma(x, Q^2)$ , from hard scattering
- Sample photon kinematics,  $x, Q^2$ , from flux
- Setup  $\gamma p$  subcollision with invariant mass of  $W_{\gamma p}$
- For single diffraction either  $Y = p$  or  $X = \rho, \phi, \omega, J/\psi$  etc.



# Photoproduction framework

Photon flux depends on beam particle:



- $Q^2$ -integrated photon flux from leptons (Weizsäcker-Williams):

$$f_{\gamma/e}(x) = \frac{\alpha_{\text{em}}}{2\pi} \frac{1 + (1 - x)^2}{x} \log \left[ \frac{Q_{\max}^2(1 - x)}{m_e^2 x^2} \right]$$

- $Q^2$ -integrated photon flux from protons (Drees-Zeppenfeld) :

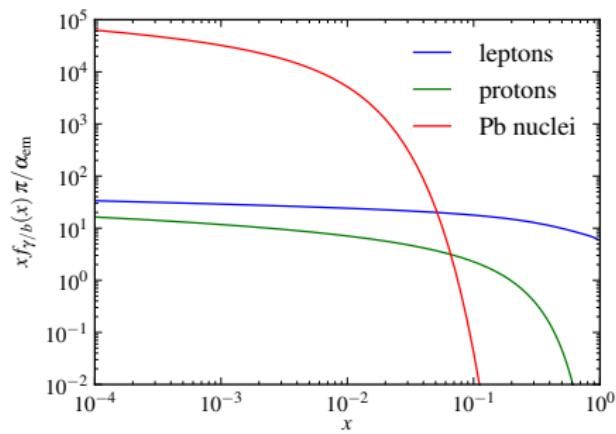
$$f_{\gamma/p}(x) = \frac{\alpha_{\text{em}}}{2\pi} \frac{1 + (1 - x)^2}{x} \left[ \log(A) - \frac{11}{6} + \frac{3}{A} - \frac{3}{2A^2} + \frac{1}{3A^3} \right],$$

$$A = 1 + Q_0^2 / Q_{\min}^2 \quad , \quad Q_0^2 = 0.71 \text{ GeV}^2$$

# Photoproduction framework

- Photon flux from nuclei found in impact-parameter space
- Reject events where nuclei overlap,  $b_{min} \sim R_{A_1} + R_{A_2}$

$$f_{\gamma/A}(x) = \frac{\alpha_{\text{em}} Z^2}{\pi x} [2\xi K_1(\xi)K_0(\xi) - \xi^2(K_1^2(\xi) - K_0^2(\xi))],$$
$$\xi = b_{\min} \times m_N$$



# Photoproduction framework

- Cross section found by convolution:

$$d\sigma(ep \rightarrow 2 \text{ jets}) = f_{\gamma/e}(x) \otimes d\sigma(\gamma p \rightarrow 2 \text{ jets})$$

- Split between **direct** and **resolved** photoproduction
- In **direct** the photon initiates the hard process:

$$d\sigma_{\text{dir}} = f_{j/p}(x_j, Q^2) \otimes d\sigma(\gamma j \rightarrow 2 \text{ jets})$$

- In **resolved** photoproduction the photon fluctuates into hadronic state before hard process: either meson state (**VMD**) or **anomalous** state described by  $\gamma \rightarrow q\bar{q}$  splitting kernel:

$$d\sigma_{\text{res}} = f_{i/\gamma}(x_\gamma, Q^2) \otimes f_{j/p}(x_j, Q^2) \otimes d\sigma(ij \rightarrow 2 \text{ jets})$$

- Requires PDF for parton in hadronic photon,  $f_{i/\gamma}(x_\gamma, Q^2)$

# Photoproduction framework

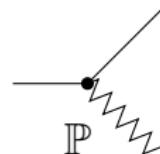
In PYTHIA 8, separation into VMD/anomalous not explicitly done in resolved photoproduction, but based on  $\gamma \rightarrow q\bar{q}$  splitting kernel present in DGLAP evolution:

$$\begin{aligned}\frac{\partial f_{i/\gamma}(x_i, Q^2)}{\partial \log(Q^2)} &= \frac{\alpha_{\text{em}}(Q^2)}{2\pi} e_i^2 P_{i\gamma}(x_i) \\ &\quad + \frac{\alpha_s(Q^2)}{2\pi} \sum_j \int_{x_i}^1 \frac{dz}{z} P_{ij}(z) f_{j/\gamma}\left(\frac{x_i}{z}, Q^2\right)\end{aligned}$$

If resolved parton  $i$  is traced back to the original photon in ISR evolution, no further ISR, MPIs allowed below the scale  $Q_0^2$  used in this splitting.

# Photoproduction framework

For diffractive dijets:



- Factorize out Pomeron flux from proton,  $f_{\mathbb{P}/p}(x_{\mathbb{P}})$ :

$$d\sigma(ep \rightarrow ep + 2 \text{ jets}) = f_{\mathbb{P}/p}(x_{\mathbb{P}}) \otimes f_{\gamma/e}(x) \otimes d\sigma(\gamma \mathbb{P} \rightarrow 2 \text{ jets})$$

- Factorize  $\gamma \mathbb{P}$ -system into hard process and Pomeron PDFs:

$$d\sigma(\gamma \mathbb{P} \rightarrow 2 \text{ jets}) = f_{j/\mathbb{P}}(x_{j/\mathbb{P}}, Q^2) \otimes d\sigma(\gamma j \rightarrow 2 \text{ jets})$$

$$+ f_{i/\gamma}(x_{\gamma}, Q^2) \otimes f_{j/\mathbb{P}}(x_{j/\mathbb{P}}, Q^2) \otimes d\sigma(ij \rightarrow 2 \text{ jets})$$

# Hard diffraction in pp with PYTHIA 8

Assume a regular PDF can be split into a diffractive (D) and non-diffractive (ND) part:

$$f_{i/p}(x_i, Q^2) = f_{i/p}^{\text{ND}}(x_i, Q^2) + f_{i/p}^{\text{D}}(x_i, Q^2) ,$$

$$f_{i/p}^{\text{D}}(x_i, Q^2) = \int_{x_i}^1 \frac{dx_{\mathbb{P}}}{x_{\mathbb{P}}} f_{\mathbb{P}/p}(x_{\mathbb{P}}) f_{i/\mathbb{P}}\left(\frac{x_i}{x_{\mathbb{P}}}, Q^2\right) ,$$

Define tentative probability for diffraction (“PDF” probability):

$$P_A^{\text{D}} = \frac{f_{i/B}^{\text{D}}(x_i, Q^2)}{f_{i/B}(x_i, Q^2)} \quad , \quad P_B^{\text{D}} = \frac{f_{i/A}^{\text{D}}(x_i, Q^2)}{f_{i/A}(x_i, Q^2)} ,$$

Gap survival introduced with MPI framework (“MPI” probability)

# MPI parameters in PYTHIA 8

- Probability for MPIs from  $2 \rightarrow 2$  QCD cross sections

$$\frac{dP_{\text{MPI}}}{dp_{\perp}^2} = \frac{1}{\sigma_{\text{ND}}(\sqrt{s})} \frac{d\sigma_{2 \rightarrow 2}}{dp_{\perp}^2}$$

- Divergent for  $p_{\perp} \rightarrow 0$
- Regularized by screening parameter  $p_{\perp 0}$

$$\frac{d\sigma_{2 \rightarrow 2}}{dp_{\perp}^2} \propto \frac{\alpha_S(p_{\perp}^2)}{p_{\perp}^4} \rightarrow \frac{\alpha_S(p_{\perp 0}^2 + p_{\perp}^2)}{(p_{\perp 0}^2 + p_{\perp}^2)^2}$$

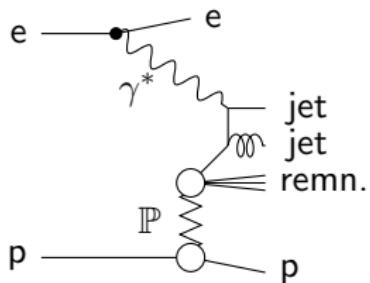
- Energy dependence through

$$p_{\perp 0}(\sqrt{s}) = p_{\perp 0}^{\text{ref}} \left( \frac{\sqrt{s}}{\sqrt{s_{\text{ref}}}} \right)^p$$

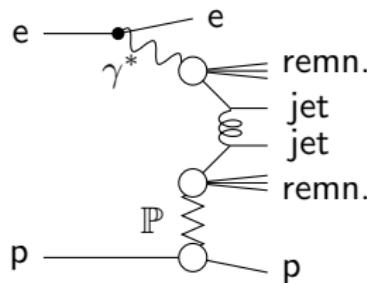
- For pp  $p_{\perp 0}^{\text{ref}} = 2.28$  GeV, while  $\gamma p$  has  $p_{\perp 0}^{\text{ref}} = 3.00$  GeV

# Hard diffraction in ep

- Tentative probability equals full probability for diffraction in **direct** events
- Gap survival introduced in **resolved** events, by requiring no additional MPIs in the  $\gamma p$  system, as these would destroy the rapidity gap



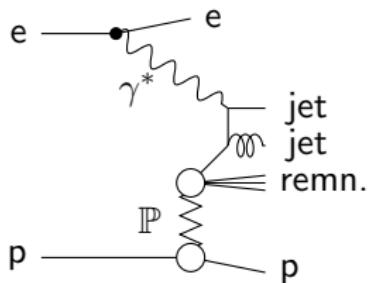
Diffractive dijets in direct photoproduction



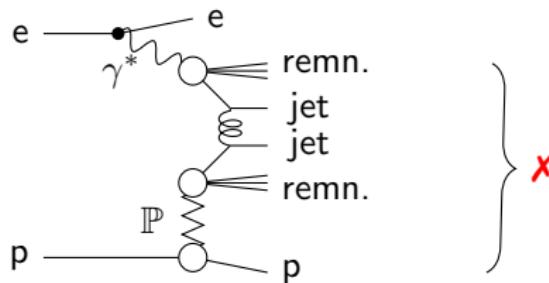
Diffractive dijets in resolved photoproduction

# Hard diffraction in ep

- Tentative probability equals full probability for diffraction in **direct** events
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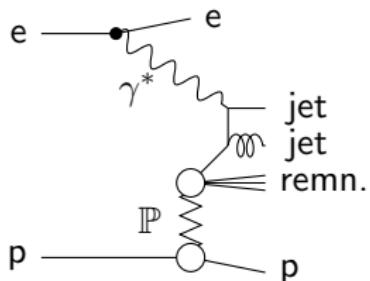
Diffractive dijets in direct photoproduction



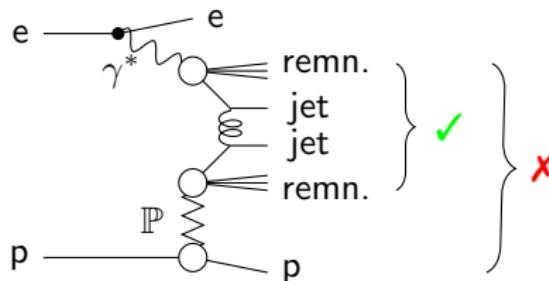
Diffractive dijets in resolved photoproduction

# Hard diffraction in ep

- Tentative probability equals full probability for diffraction in **direct** events
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Diffractive dijets in direct photoproduction



Diffractive dijets in resolved photoproduction

# Hard diffraction in ep

H1 2007 [EPJC 51 (2007) 549]

- $Q^2 < 0.01 \text{ GeV}^2$
- $E_{\perp}^{*\text{jet}\,1} > 5.0 \text{ GeV}$
- $E_{\perp}^{*\text{jet}\,2} > 4.0 \text{ GeV}$
- $-1 < \eta^{\text{jet}\,1,2} < 2.0$
- $x_{\mathbb{P}} < 0.03$

ZEUS 2008 [EPJC 55 (2008) 177]

- $Q^2 < 1 \text{ GeV}^2$
- $E_{\perp}^{\text{jet}\,1} > 7.5 \text{ GeV}$
- $E_{\perp}^{\text{jet}\,2} > 6.5 \text{ GeV}$
- $-1.5 < \eta^{\text{jet}\,1,2} < 1.5$
- $x_{\mathbb{P}} < 0.025$

Baseline setup

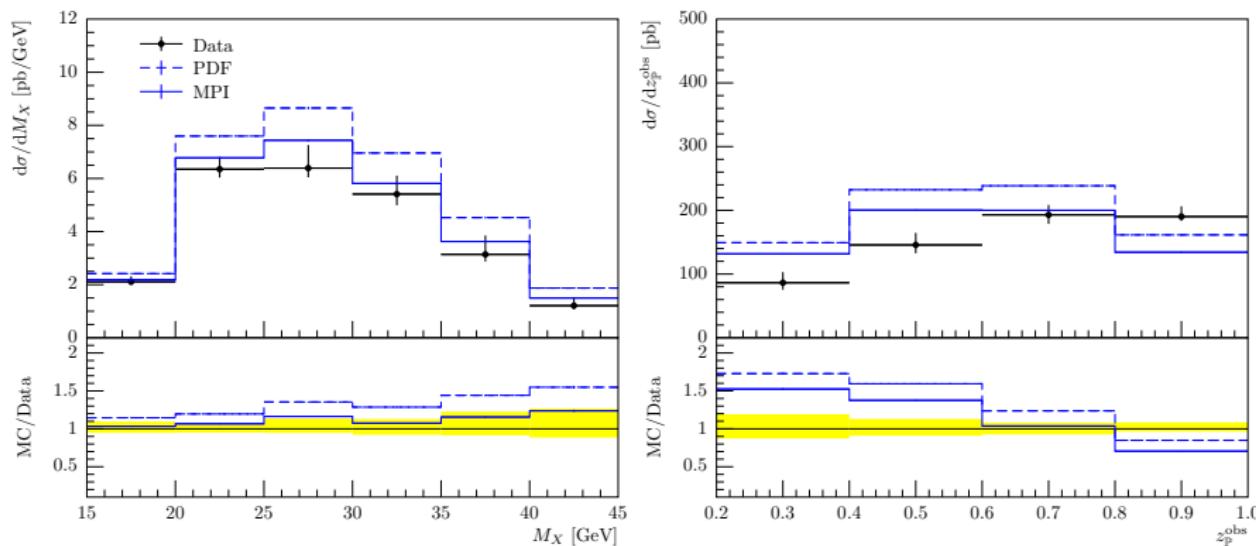
- dPDFs from H1 Fit B LO
- $\gamma$ PDFs from CJKL
- $p_{\perp 0}^{\text{ref}} = 3.00 \text{ GeV}$

Observables:

- $W$  (H1)
- $M_X$  (ZEUS)
- $z_{\mathbb{P}} = \frac{\sum_{\text{jet}=1,2}(E^{\text{jet}} + p_z^{\text{jet}})}{\sum_{i \in X}(E^i + p_z^i)}$

# Hard diffraction in ep

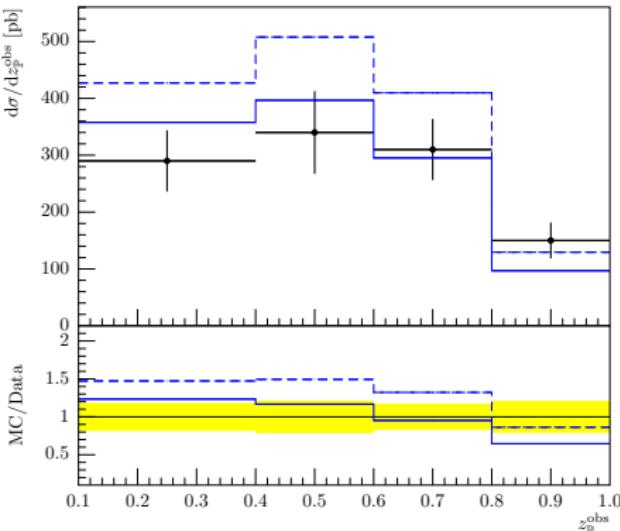
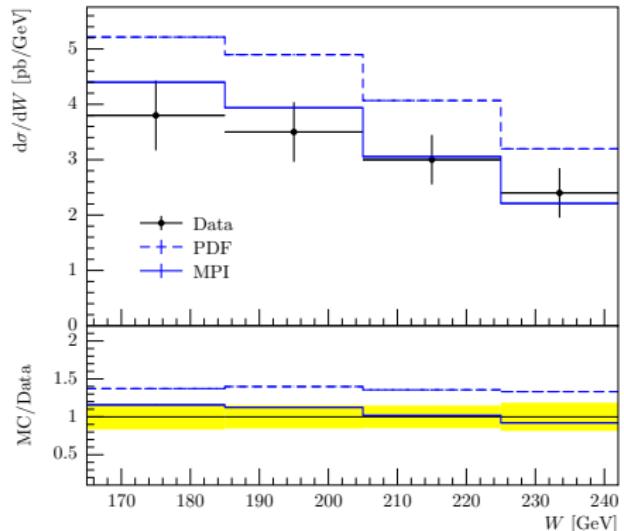
ZEUS 2008:



- PDF sample overshoots data
- Better agreement with additional gap suppression
- Some distributions not well described (eg.  $x_\gamma, z_P$ )

# Hard diffraction in ep

H1 2007:



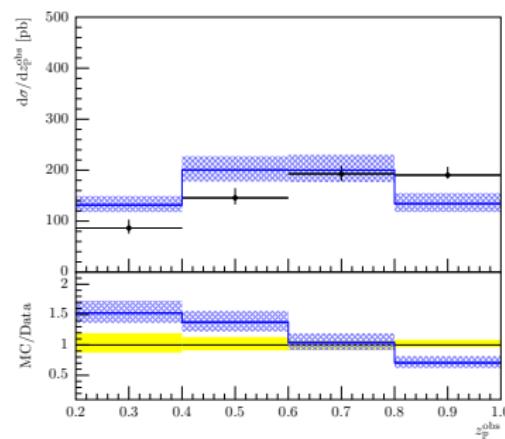
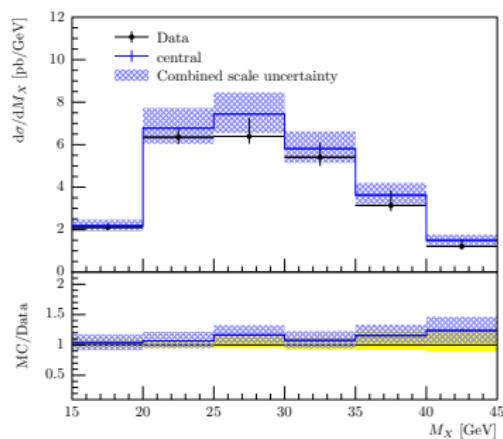
- PDF sample overshoots data
- Better agreement with additional gap suppression
- Some distributions not well described (eg.  $x_\gamma, z_{\mathbb{P}}$ )

# Hard diffraction in ep

Uncertainties arise from:

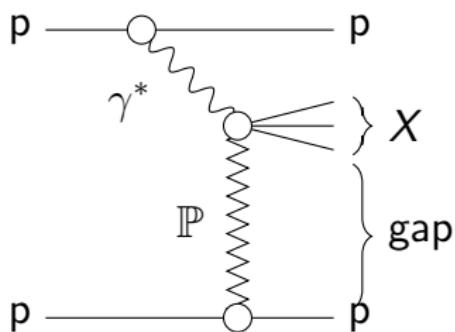
- LO ME
- dPDFs
- $\gamma$ PDFs
- MPI parameters

ZEUS 2008:

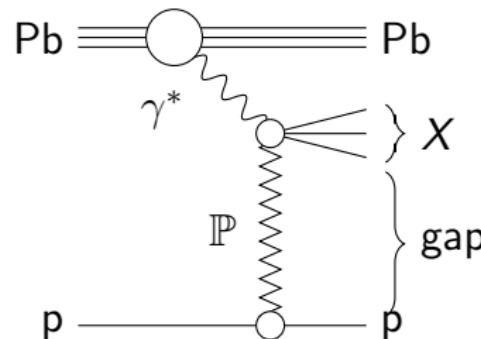


# Hard diffraction in UPCs

- Predictions for pp and pPb.
- Photon flux from either side in pp, flux from lead dominates in pPb
- Currently PbPb not possible as additional interactions between resolved photon and other nucleons should be taken into account in gap suppression



Diffraction in UPC pp



Diffraction in UPC pPb

# Hard diffraction in UPCs

Cuts:

- $\sqrt{s_{\text{NN}}} = 5.0 \text{ TeV (pp)}$ ,  
 $\sqrt{s_{\text{NN}}} = 13.0 \text{ TeV (pPb)}$
- $E_{\perp,\min}^1 = 8.0 \text{ GeV}$
- $E_{\perp,\min}^2 = 6.0 \text{ GeV}$
- $M_{\text{jets,min}} = 14.0 \text{ GeV}$
- $x_{\mathbb{P}}^{\max} = 0.025$
- $|\eta^{\max}| = 4.4$

Baseline setup

- dPDFs from H1 Fit B LO
- $\gamma$ PDFs from CJKL

Observables:

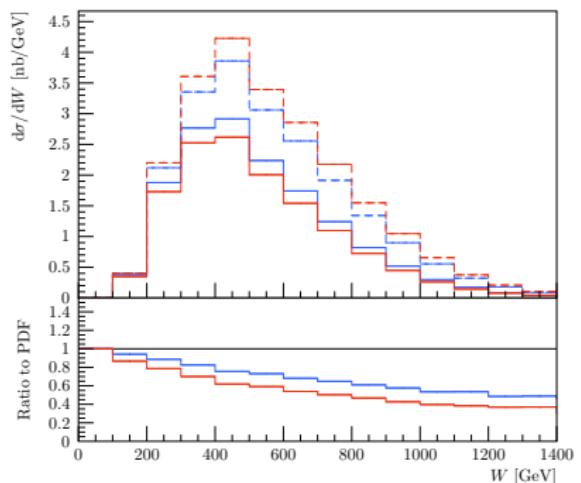
- $W$

- Two values of MPI screening parameter used:  $p_{\perp 0}^{\text{ref}} = 2.28 \text{ GeV (red lines)}$  and  $p_{\perp 0}^{\text{ref}} = 3.00 \text{ GeV (blue lines)}$
- Dashed lines are “PDF” samples, solid lines “MPI” samples

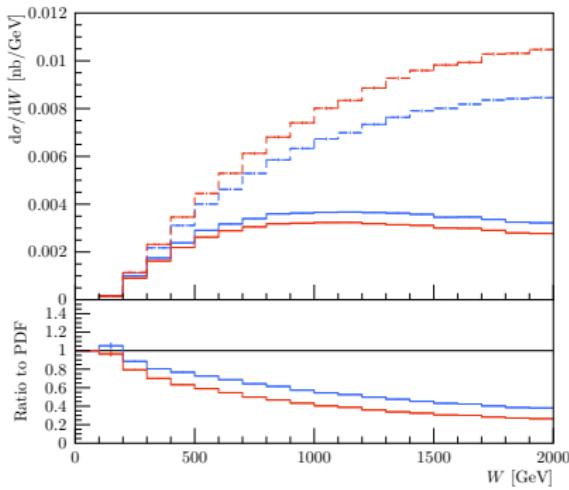
# Hard diffraction in UPCs

$$p_{\perp 0}^{\text{ref}} = 2.28 \text{ GeV} \text{ and } p_{\perp 0}^{\text{ref}} = 3.00 \text{ GeV}$$

pPb:



pp:

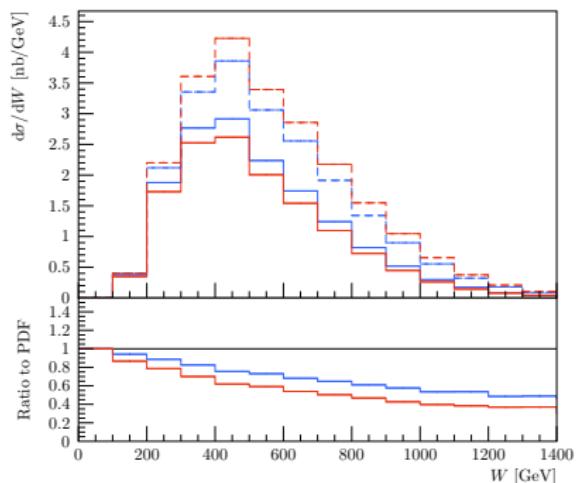


- Rejecting additional MPIs reduces cross section drastically as compared to HERA
- Suppression stronger in pp as harder flux leads to larger  $W_{\gamma p}$ , allowing for more MPIs

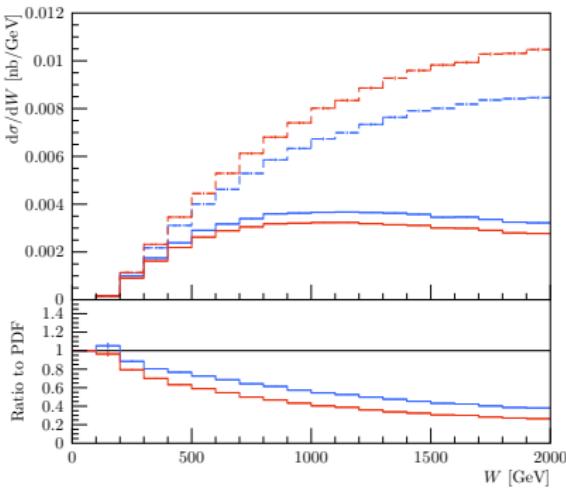
# Hard diffraction in UPCs

$$p_{\perp 0}^{\text{ref}} = 2.28 \text{ GeV} \text{ and } p_{\perp 0}^{\text{ref}} = 3.00 \text{ GeV}$$

pPb:



pp:



- Lower MPI-screening value allows for more MPIs, thus leads to larger suppression
- UPCs provide ideal place to test gap-survival model, as effects are more pronounced here w.r.t. HERA

# Conclusion and outlook

- Hard diffraction for photoproduction available, without any new parameters.
- Reasonable description of HERA data with gap suppression.
- Several theoretical uncertainties, e.g. dPDF,  $\gamma$ PDF, MPI parameters.
- Ratio DIS/photoproduction could be used to reduce theoretical uncertainties (DIS description not ready yet).
- Predictions for UPCs at LHC presented.
- Gap suppression larger here due to higher  $W_{\gamma p}$ .
- Excellent place to test model (awaiting data).
- Extension to eA or UPC PbPb with Angantyr expected.