

Limitations for heavy-ion performance in the LHC

R. Bruce, J.M. Jowett, M. Schaumann



Outline

- Introduction: Where are we now?
 - 2018 Pb run
- Encountered limitations on luminosity
 - Collisional processes: Burnoff and beam losses
- Encountered limitations on beam parameters so far
 - Collimation
 - Beam from injectors
 - Evolution in the LHC
- Future performance and plans



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- LHC designed to collide both protons and nuclei
- 8 straight sections (4 with experiments) and 8 arcs





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- So far: Pb-Pb runs at 3.5 Z TeV and 6.37 Z TeV, p-Pb runs at 4 Z TeV and 6.37 TeV







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Parameter for Pb beam	LHC design value	Achieved 2018
Beam energy	2.76 TeV / nucleon 7 Z TeV	2.51 TeV / nucleon 6.37 Z TeV
Bunch intensity	7E7	2.3E8
Number of bunches	592	733
Stored beam energy	3.8 MJ	13.9 MJ
Bunch spacing	50 ns	75 ns
Normalized emittance	1.5 µm	2.0 µm
β*	0.5 m	0.5 m
Number of collision points	1	4
Peak luminosity (Pb-Pb)	1E27 cm ⁻² s ⁻¹	6.4E27 cm ⁻² s ⁻¹
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• In addition: p-Pb : new mode of operation not foreseen in LHC design



Typical LHC Pb-Pb fill in 2018

- Leveling in ATLAS and CMS gradually increased to 6E₂₇ cm⁻²s⁻¹
- ALICE leveled at design luminosity 1E27 cm⁻²s⁻¹







Pb-Pb luminosity production so far

- Achieved in total 2.5 nb⁻¹ in ATLAS/CMS, in 1.5 nb⁻¹ ALICE and 0.26 nb⁻¹ in LHCb over all Pb-Pb runs
 - Surpassed initial ALICE design goal of 1 nb⁻¹



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- 2018: best year so far









 $\mathcal{L} = \frac{N_1 N_2 f_{\text{rev}} k_B}{4\pi \beta^* \epsilon_{xy}} F$



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intensity

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- Additional limits on integrated luminosity: availability and turnaround time







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Cross sections for Pb-Pb collision	ns at 2.76 TeV / nucleon

Process	Cross section (b)
Bound-free pair production	281
Electromagnetic dissociation	226
Hadronic nuclear inelastic	8
Total	515

• Dominating UEI processes

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Meier et al. Phys. Rev. A, 63, 032713 (2001)

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1-neutron Electromagnetic dissociation (EMD1, 96 barn)

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Ultra-peripheral electromagnetic interactions

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• All these processes create unwanted beam losses





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- Puts upper limit on luminosity
 - Limit found experimentally at 2.3E27 cm⁻²s⁻¹ in IR5.
 - Could be different at different magnets / IRs





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CERN

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Bound-free pair production secondary beams from IPs

Losses from collimation inefficiency, nuclear processes in primary collimators



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Careful setup of bumps in beginning of the run to achieve desired loss displacement.

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- Established technique in 2015: Orbit bumps are used to move the secondary beam losses to empty cryostat in order to reduce risk of quench.
- With bumps, achieved ~6E27 cm⁻²s⁻¹ in ATLAS / CMS



R. Bruce, 2018.12.05





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- Use orbit bump to distribute losses between different magnets
 - Upper luminosity limit not investigated experimentally
- ALICE anyway leveled at 1E27 cm⁻²s⁻¹
 - Too high event rate for detector otherwise
 - ALICE upgrade foreseen in LS₂

















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- TCLDs should allow luminosity increase for upgraded ALICE to run at 50 kHz

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Collimation with ion beams

- LHC collimation much less efficient with nuclear beams than with protons
 - Very high probability of nuclear breakup in primary collimator
 - Fragments very often miss downstream collimation stages
 - Different charge-to-mass ratio => fragments bent wrongly and lost in the first few dipoles



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Measured loss patterns



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Limit from collimation



- Efficiency of collimation system introduces upper limit on acceptable beam losses (Pb/s)
 - In case of large beam losses, a proportionally large leakage to cold magnets occurs
 - Beams should be dumped by beam loss monitors before quench or damage occur
 - If losses are frequently too high, frequent dumps make operation less efficient or even impossible



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 - If losses are frequently too high, frequent dumps make operation less efficient or even impossible
- Gives effectively an upper limit on the total beam intensity





Future alleviation: extra collimators

- No hard limit reached yet, although each ion run suffered from a couple of unforeseen beam dumps due to losses
 - Risk for more serious limitation expected in the future when beam intensity is increased



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 - Make space by replacing a standard dipoles by two shorter 11T dipole, with TCLD in between





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- Present baseline: install 1 IR7 TCLD per beam in LS2





LHC injectors

• Ion beam passes through a chain of injectors before reaching the LHC



LHC Large Hadron Collider SPS Super Proton Synchrotron PS Proton Synchrotron AD Antiproton Decelerator CTF3 Clic Test Facility

AWAKE Advanced WAKefield Experiment ISOLDE Isotope Separator OnLine REX/HIE Radioactive EXperiment/High Intensity and Energy ISOLDE

LEIR Low Energy Ion Ring LINAC LINear ACcelerator n-ToF Neutrons Time Of Flight HiRadMat High-Radiation to Materials



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H. Bartosik, G. Rumolo et al.



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- Initial LHC design performance has already been largely surpassed by injectors after optimization efforts
- 2016: LEIR performance reached target value, further work needed on margin and shot-to-shot stability

- 2018:

- LEIR performance confirmed, higher accumulated intensity reached
- Better reproducibility from Linac3/LEIR thanks to improved diagnostics (e.g. BPMs in injection line, Schottky monitor for energy matching)



H. Bartosik, G. Rumolo et al.





• Steady increase in injected intensity over the years





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- Bunch spacing decreased to
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 - Hope for 50 ns in the future with new hardware – see next talk
- Not evident how to further increase injected intensity in LHC with present hardware



Beam evolution in LHC



- Once injected in the LHC, beams suffer from further blowup and losses
 - Beams blow up at flat bottom due to intrabeam scattering
 - At top energy, radiation damping dominates over intrabeam scattering => blowup not a big issue at top energy
 - Need to minimize injection time and start ramp as soon as possible



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 - At top energy, radiation damping dominates over intrabeam scattering => blowup not a big issue at top energy
 - Need to minimize injection time and start ramp as soon as possible
- Some losses throughout the cycle, but transmission generally good
 - Usually a few percent of beam lost between start of ramp and start of collisions

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×m
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- Peak luminosity could be increased by focusing the beams to smaller beam size at the collision point (smaller β*)
 - Causes beam size increase in triplet => increase limited by collimation system
- We are not yet at a hard limit for Pb-Pb, however, as long as ALICE is leveled, no major gain expected $(3\sigma_x,3\sigma_y,5\sigma_t) envelope for \epsilon_x=5.52358 \times 10^{-10}m, \epsilon_y=5.52358 \times 10^{-10}m, \sigma_p=0.0001137$



×m



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- Heavy-ion runs up to now highly successful
 - Design luminosity performance surpassed by factor 6
 - Achieved goal of 1 nb^{-1} in ALICE
 - 4 experiments taking data instead of 1
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- Heavy-ion runs up to now highly successful
 - Design luminosity performance surpassed by factor 6
 - Achieved goal of 1 nb^{-1} in ALICE
 - 4 experiments taking data instead of 1
 - New operation mode: p-Pb
- Still several limitations for further performance increase
 - Collisional losses
 - High burn-off cross section
 - Secondary beams, in particular bound-free pair production
 - Collimation efficiency
 - Not easy to increase injected intensity from injector chain





Luminosity production so far

- Achieved in total in ATLAS/CMS, in ALICE and in LHCb
- Include 2018









EM interactions create ions with altered magnetic rigidity:

$$\delta = \frac{Z_0}{A_0} \frac{A}{Z} (1 + \delta_{\rm kin}) - 1$$

- These ions follow locally generated dispersion function d_{x} from IP
- Lost in localized spot where aperture A_x , d_x and δ satisfy

 $\delta d_r = A_r$

Apart from significant luminosity decay, induced heating risks to quench superconducting magnets - BFPP gives 25W beam

S. Klein, NIM A 459 (2001) 51

Beam





- IR2 has different quadrupole polarity and dispersion from IR1/IR5
- Primary BFPP loss location is further upstream from connection cryostat
- Solution is to modify connection cryostat to include a collimator to absorb the BFPP beam to be ready for LS2 installation
- With levelled luminosity in ALICE, quenches were not seen in 20145tt, Town Meeting:



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Also during LS2, further TCLD collimators will be installed between 11 T magnets in IR7 to improve Pb collimation (first application of Nb₃Sn superconductors in an operating accelerator).



Timeline of heavy-ion runs







• Fewer collisions when bunches are not fully overlapping



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• Decrease bunch length and crossing angle to minimize effect



• Fewer collisions when bunches are not fully overlapping



- Decrease bunch length and crossing angle to minimize effect
- Crossing angle limited by beam-beam separation and aperture



• Standard dump: extraction kickers fire when no beam passes



• Asynchronous dump: kicker(s) fire when beam passes – kicked beam damage could TCTs/triplets. TCDQ should protect



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What can happen if a TCT is hit?

- Impacts studied in HiRadMat
- Significant damage observed





Proton acceleration to LHC





Superconducting magnets

- "Quench" = loss of superconductivity
- Happens if working point is outside of surface in magnetic field (B), current (I) and temperature (T)





LHC lattice



Ion fragment distribution after collimator



Fig. 4. Energetic fractions of the individual isotopes emerging from the IR7 horizontal TCP from an impacting ²⁰⁸Pb⁸²⁺ beam at 3.5 *Z* TeV, simulated with FLUKA as shown in Fig. 3 with an impact parameter of $b = 3 \mu m$. The energy fraction is computed by multiplying the isotope abundance N(A, Z) with the nucleon number *A* and the momentum per nucleon p(A, Z) and normalizing by the total ion energy coming out of the collimator $\sum_{A,Z} A \times N(A, Z) \times p(A, Z)$. It is assumed that the momentum per nucleon is approximately identical for all ions $p(A, Z) \approx p$.