SMEFT@NLO in QCD

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LHCTheory ERC meeting
Université Catholique de Louvain

22nd March 2017

Outline

- Overview of our recent progress in implementations of MC event generation of SMEFT at NLO in QCD matched to parton shower ➝ towards complete public implementation
- Brief outline of SMEFT
- Top/Higgs/Electroweak SMEFT @ the LHC
  - Selection of codes & results
- Results from the implementations of operators affecting Higgs couplings to gauge bosons
  - Effect on SM inputs & EW parameters
  - Current constraints from global fits & resulting benchmark choices
  - EW production: WH, ZH and VBF

Going NLO

- The LHC is now in the precision era
  - No clear evidence for new physics as we approach the limits of the 'energy frontier'
  - Fully complementary approach to search for deviations in SM processes
  - Many channels are becoming systematics dominated & require high precision theory input from higher order corrections (FO & PS)
- EFT: theoretically consistent, model independent approach to deviations of interactions between SM fields
  - Active area of research that is moving towards NLO predictions
  - NLO important for capturing potentially large QCD K-factors in total rates → greater sensitivity
  - Verify stability of differential information beyond leading order
  - Consistent scale uncertainty estimates

Going NLO

- State-of-the-art in MC event generation is well beyond LO
  - Software like FeynRules+NLOCT+MG5_aMC@NLO provides fully automated event generation at NLO in QCD from Lagrangian
  - Other matching/merging schemes exist on a process-by-process basis, e.g., POWHEG-BOX,…
  - Individual codes exist for specific processes, up to NNLO QCD + NLO EW
- Public implementations largely restricted to SM predictions although some codes permit the inclusion of anomalous couplings
  - See, e.g., Higgs Characterisation *[Demartin, Maltoni, Mawatari, Page & Zaro; EPJC 74 (2014) 9, 3065]*
  - Several others: HAWK, VBFNLO, POs…
  - Full SM-EFT descriptions are naturally well motivated and will provide a valuable addition to the existing toolbox

SMEFT

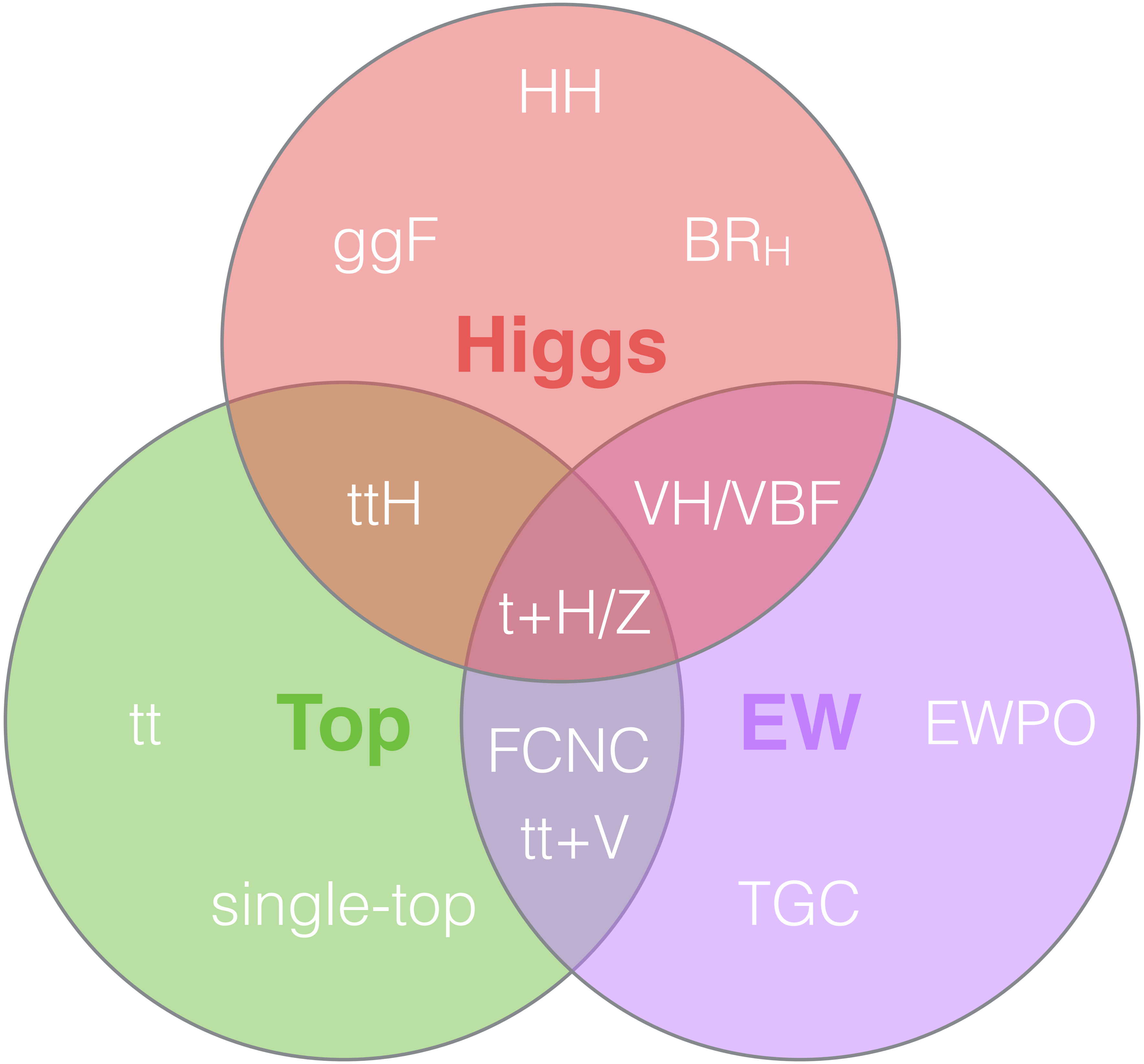
- Basic introduction to EFT is unnecessary at this meeting
  - Expansion in the cutoff scale, Λ, using only SM fields
  - Truncated at canonical dimension 6
  - Introduces higher-derivative operators to which we are sensitive through large momentum flows through vertices (i.e. tails of energy distributions)
- Operator expansion: $$\mathcal{L}_{\text{eff}} = \sum_i \frac{c_i \mathcal{O}_i^D}{\Lambda^{D-4}}$$ more: fields derivatives
- In the SM: 59 (76 real) - 2499 operators depending on assumptions regarding CP/flavour structure etc.

[Buchmuller & Wyler; Nucl.Phys. B268 (1986) 621]
[Grzadkowski et al.; JHEP 1010 (2010) 085]

- Dimension 8 now known ~ 895 (36971) operators!

[Lehman et al.; PRD 91 (2015) 105014]
[Henning et al.; Commun.Math.Phys. 347 (2016) no.2, 363-388 & arXiv:1512.03433]

SMEFT @ the LHC



+CPV, flavor,…

SM(Higgs)-EFT

- Nicolas' talk later on ggF in SMEFT @ NLO in QCD
- Public codes include partial SMEFT contributions @ NLO
  - HiGlu *[Spira; arXiv:hep-ph/9510347]*
  - SusHi (aMC-SusHi) *[Harlander, Liebler & Mantler; arXiv:1605.03190]*
- Double Higgs
  - HPAIR *[Dawson, Dittmaier & Spira; Phys. Rev. D58:115012]*
  - HiggsPair (HERWIG++) *[Goertz et al.; JHEP 1504 (2015) 167]*
- eHDECAY for Higgs branching fractions in SMEFT
  [R. Contino et al.; Comp. Phys. Comm. 185 (2014) 3412-3423]
- Full 1-loop H→$\gamma\gamma$
  [Hartmann & Trott; PRL115 (2015) 191801]
- and H→bb
  [Gauld, Scott & Pecjak; PRD94 (2016) 074045]

SM(Top+X)-EFT

- Most developed sector of SMEFT@NLO in QCD
  - Major overlap with the LHCTheory ERC collaboration
- Coloured sector, strongly coupled to the Higgs
  - Large corrections to inclusive rates (~1 K-factors)
  - Non-trivial shape corrections at differential level
  - Non-trivial renormalisation/operator mixing
- ttH *[Maltoni, Vryonidou & Zhang; JHEP 1610 (2016) 123]*
- tt+Z/γ *[Bylund et al.; JHEP 1605 (2016) 052]*
- single-top *[Zhang; PRL 116 (2016) 162002]*
- top FCNC *[Degrande et al.; PRD 91 (2015) 034024]* *[Durieux, Maltoni & Zhang; PRD 91 (2015) 074017]*

[Maltoni, Vryonidou & Zhang; JHEP 1610 (2016) 123]

ttH in SMEFT

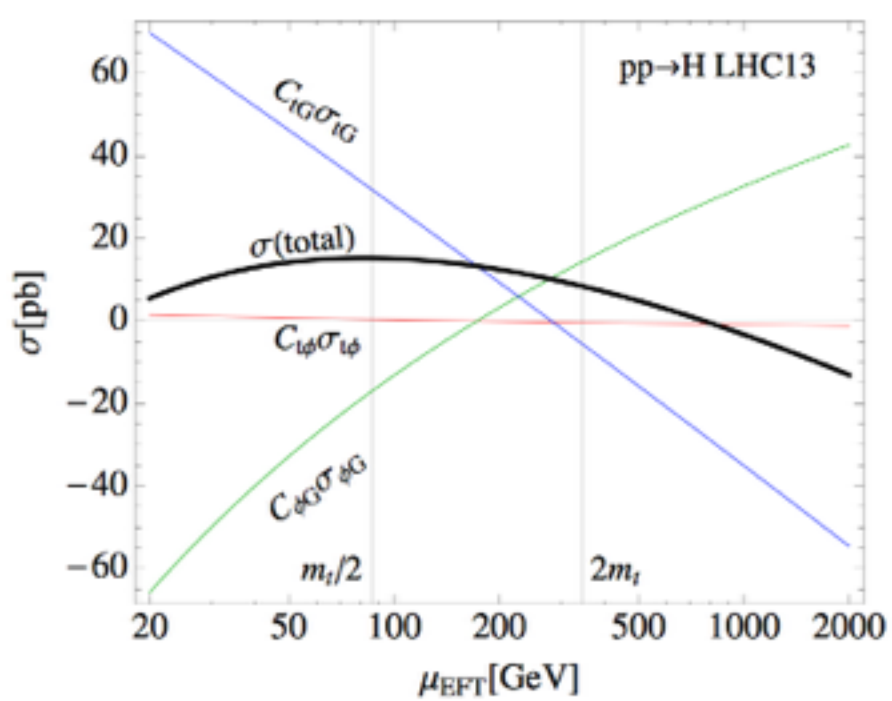
$$\mathcal{O}_{t\varphi} = (\varphi^\dagger \varphi)(\bar{Q}_L\, \tilde{\varphi}\, t_R)$$
$$\mathcal{O}_{\varphi G} = (\varphi^\dagger \varphi)\, G^A_{\mu\nu} G^{\mu\nu}_A$$
$$\mathcal{O}_{tG} = (\bar{Q}_L \sigma_{\mu\nu}\, T^A\; t_R)\, \tilde{\varphi}\, G^{\mu\nu}_A$$

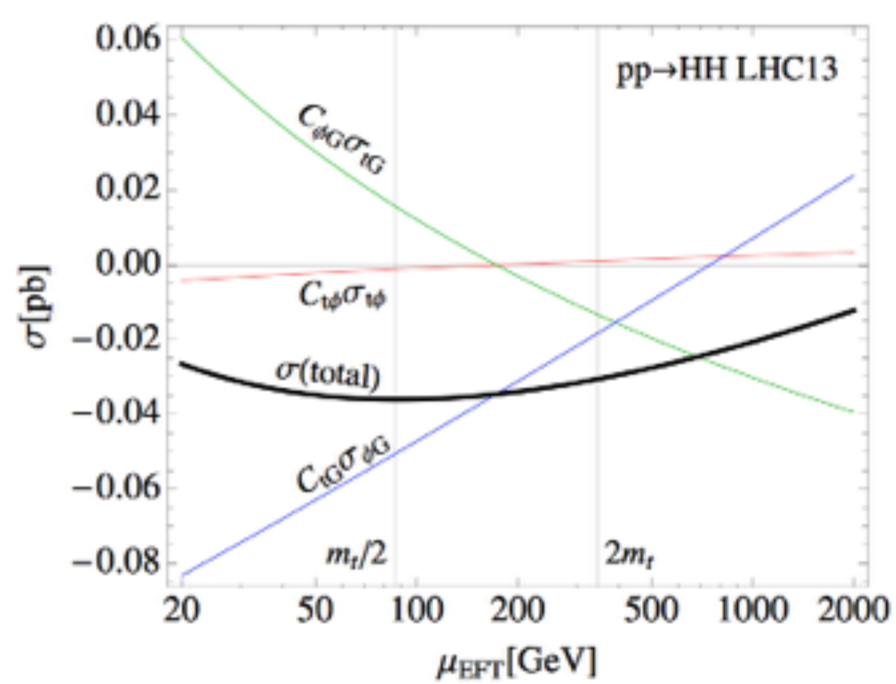
- Operators involving the top/Higgs/gluon
  - gg→H & tt production partly constrain the Wilson coefficient space
  - ttH is the only direct probe of the Top-Higgs interaction
  - In principle 3-gluon O_G and 4 fermion operators also contribute but turn out to be better constrained by tt and multi-jet measurements (next talk)
- Predictions for ttH and HH production presented
  - First inclusion of chromomagnetic dipole operator for HH
  - Demonstration of non-trivial K-factors at differential level
  - Comparison of RG improved vs. full NLO (finite terms important)

[Maltoni, Vryonidou & Zhang; JHEP 1610 (2016) 123]

New EFT scale uncertainty

- NLO calculations use scale uncertainty to approximate missing higher orders in perturbative expansion
  - EFT description contains an additional source of scale dependence from the running/mixing of Wilson coefficients
- Proposal for a new scale uncertainty component
  - Take c_i defined at scales $2\mu_0$ & $\mu_0/2$ and run back to the central scale





Does not cancel in e.g. cross section ratios for which traditional scale uncertainty drops out

SM(Higgs+EW)-EFT

- "Canonical" sector for Higgs studies in SMEFT
- Easy to implement NLOQCD compared to e.g. ttH
- Implementations build upon previous work done at LO employing the 'SILH' basis of operators
  - SMEFT Basis dependence of HEP tools can be reduced by linear redefinitions using e.g. Rosetta
- Also made use of existing implementations of SM process at NLO

MCFM/ POWHEG-BOX WH & ZH (incl. gg → ZH)	FeynRules/NLOCT UFO model via Madgraph5_aMC@NLO WH & VBF

(in backup)

EW Higgs production

- EFT effects in EW production mechanisms for the Higgs: VH & VBF
  - A small number of relevant & uncoloured operators at D=6 in SM-EFT
  - LHC can provide complementary information to existing fits to lower energy data, i.e. LEP
  - Higgs comes with some additional objects from which we can construct kinematic quantities probing the high energy regime
  - VH: Higgs p_T, M_{VH}, leading lepton p_T,…
  - VBF: Higgs p_T, $\Delta\eta_{jj}$, total H_T,…
- Investigate validity of EFT expansion given current constraints from global fits
- Considered future reach of HL-LHC

SILH operators

- SMEFT: Higgs-EW gauge boson operators in SILH basis

$$\mathcal{L}_{D6} = \frac{1}{\Lambda^2}\Big[\frac{g'^2}{4}\bar{c}_{BB}\Phi^\dagger\Phi B^{\mu\nu}B_{\mu\nu} + \frac{ig}{2}\bar{c}_W\big[\Phi^\dagger T_{2k}\overleftrightarrow{D}^\mu\Phi\big]D^\nu W^k_{\mu\nu} + \frac{ig'}{2}\bar{c}_B\big[\Phi^\dagger\overleftrightarrow{D}^\mu\Phi\big]\partial^\nu B_{\mu\nu}$$
$$+ ig\,\bar{c}_{HW}\big[D^\mu\Phi^\dagger T_{2k}D^\nu\Phi\big]W^k_{\mu\nu} + ig'\,\bar{c}_{HB}\big[D^\mu\Phi^\dagger D^\nu\Phi\big]B_{\mu\nu}$$
$$+ \frac{g'^2}{4}\tilde{c}_{BB}\Phi^\dagger\Phi B^{\mu\nu}\tilde{B}_{\mu\nu} + ig\,\tilde{c}_{HW}\big[D^\mu\Phi^\dagger T_{2k}D^\nu\Phi\big]\widetilde{W}^k_{\mu\nu} + ig'\,\tilde{c}_{HB}\big[D^\mu\Phi^\dagger D^\nu\Phi\big]\tilde{B}_{\mu\nu}\Big]$$

$$\Phi^\dagger\overleftrightarrow{D}^\mu\Phi \equiv \left(D^\mu\Phi^\dagger\right)\Phi - \Phi^\dagger\left(D^\mu\Phi\right)$$

- Anomalous couplings: new Lorentz structures (1) & (2):

$$\mathcal{L}_{HAC} = -\frac{1}{4}g^{(1)}_{hzz}Z_{\mu\nu}Z^{\mu\nu}h - g^{(2)}_{hzz}Z_\nu\partial_\mu Z^{\mu\nu}h + \frac{1}{2}g^{(3)}_{hzz}Z_\mu Z^\mu h - \frac{1}{4}\tilde{g}_{hzz}Z_{\mu\nu}\tilde{Z}^{\mu\nu}h$$
$$-\frac{1}{2}g^{(1)}_{hww}W^{\mu\nu}W^\dagger_{\mu\nu}h - \left[g^{(2)}_{hww}W^\nu\partial^\mu W^\dagger_{\mu\nu}h + \text{h.c.}\right] + g^{(3)}_{hww}W_\mu W^{\dagger\mu}h - \frac{1}{2}\tilde{g}_{hww}W^{\mu\nu}\tilde{W}^\dagger_{\mu\nu}h$$
$$-\frac{1}{2}g^{(1)}_{haz}Z_{\mu\nu}F^{\mu\nu}h - g^{(2)}_{haz}Z_\nu\partial_\mu F^{\mu\nu}h - \frac{1}{2}\tilde{g}_{haz}Z_{\mu\nu}\tilde{F}^{\mu\nu}h$$

(+ Higgs-fermion current operators not included here)

Mapping to AC/(i.e. HC)

Coupling	HEL@NLO
$g^{(1)}_{hzz}$	$\frac{e^2 v}{2\hat{c}^2_W \hat{s}^2_W}\frac{1}{\Lambda^2}\left[\hat{c}^2_W \bar{c}_{HW} + 2\hat{s}^2_W \bar{c}_{HB} - 2\hat{s}^4_W \bar{c}_{BB}\right]$
$g^{(2)}_{hzz}$	$\frac{e^2 v}{4\hat{s}^2_W \hat{c}^2_W \Lambda^2}\left[\hat{c}^2_W(\bar{c}_{HW} + \bar{c}_W) + 2\hat{s}^2_W(\bar{c}_B + \bar{c}_{HB})\right]$
$g^{(3)}_{hzz}$	$\frac{g^2 v}{2\hat{c}^2_W} + \frac{e^4 v^3}{8\hat{c}^4_W \hat{s}^2_W \Lambda^2}\left[\hat{c}^2_W \bar{c}_W + 2\bar{c}_B\right]$
$g^{(1)}_{haz}$	$\frac{e^2 v}{4\hat{s}_W \hat{c}_W \Lambda^2}\left[\bar{c}_{HW} - 2\bar{c}_{HB} + 4\hat{s}^2_W \bar{c}_{BB}\right]$
$g^{(2)}_{haz}$	$\frac{e^2 v}{4\hat{s}_W \hat{c}_W \Lambda^2}\left[\bar{c}_{HW} + \bar{c}_W - 2(\bar{c}_B + \bar{c}_{BB})\right]$
$g^{(1)}_{hww}$	$\frac{e^2 v}{2\hat{s}^2_W \Lambda^2}\bar{c}_{HW}$
$g^{(2)}_{hww}$	$\frac{v e^2}{4\Lambda^2 \hat{s}^2_W}\left[\bar{c}_W + \bar{c}_{HW}\right]$
$g^{(3)}_{hww}$	$\frac{g^2 v}{2}$

Anomalous couplings (AC) equivalent to
Higgs Characterisation (HC)

[Trott & Passarino; LHCHXSWG-DRAFT-2016-005]
[Falkowski et.al; LHCHXSWG-INT-2015-001]
[Williams, KM & Sanz; JHEP 1608 (2016) 039]
[Ge, He & Xiao; JHEP 1610 (2016) 007]
[Degrande, Fuks, Mawatari, KM, Sanz; 1609.04833]

SM inputs

$$\begin{aligned}\mathcal{O}_H &= \frac{\bar{c}_H}{2}\partial_\mu\left(\Phi^\dagger\Phi\right)\partial^\mu\left(\Phi^\dagger\Phi\right)\\ &= \frac{\bar{c}_H}{\Lambda^2}\frac{v^2}{2}\partial_\mu h\partial^\mu h + \mathcal{O}(h^3,h^2)\end{aligned}$$

$$h \to h(1+\delta h), \quad \delta h = -\frac{\bar{c}_H}{\Lambda^2}\frac{v^2}{4}$$

$$\begin{aligned}\mathcal{O}_W|_{\Phi=\langle\Phi\rangle} &= \frac{ig}{2}\bar{c}_W\left[\Phi^\dagger T_{2k}\overleftrightarrow{D}^\mu\Phi\right]D^\nu W^k_{\mu\nu}|_{\Phi=\langle\Phi\rangle}\\ &= \frac{gv^2}{16}\bar{c}_W\left[2gW^{\mu\nu}_+W^-_{\mu\nu} + g(W^{\mu\nu}_3 - g'B^{\mu\nu})W^3_{\mu\nu}\right] + \mathrm{aGC}\end{aligned}$$

$$\begin{aligned}W^\mu_\pm &\to W^\mu_\pm\left[1+\delta W\right]\\ B^\mu &\to B^\mu\left[1+\delta B\right] + yW^\mu_3\\ W^\mu_3 &\to W^\mu_3\left[1+\delta W\right] + zB^\mu\end{aligned}$$

- After EWSB, canonical mass eigenbasis, different from SM
  - Perform field redefinitions to fix their normalisation
  - Gauge coupling redefinitions can absorb part of the resulting modifications
  - Modifications of gauge bosons masses, interactions, e.g., Z→ff
  - Modifications to the SM parameters as a function of EW inputs
  - Can also affect backgrounds
- Not all tools take these into account
  - Various choices can be made that are all equivalent up to dimension-6

Limits from global fits

- A number of global fits to data deriving constraints on EFT Wilson coefficients have been performed
  - LHC, LEP & other low-energy experiments
- Marginalised constraints from EWPO + LHC Run 1 data on coefficients of interest

[Sanz et al.; JHEP 1503 (2015) 157]

Operator	Coefficient	Constraints
$\mathcal{O}_W = \frac{ig}{2}\left(H^\dagger T_{2k}\overleftrightarrow{D}^\mu H\right) D^\nu W^k_{\mu\nu}$	$\frac{m_W^2}{\Lambda^2}(\frac{\bar{c}_W}{2} - \bar{c}_B)$	(-0.035,0.005)
$\mathcal{O}_B = \frac{ig'}{2}\left(H^\dagger \overleftrightarrow{D}^\mu H\right) \partial^\nu B_{\mu\nu}$	$\frac{m_W^2}{\Lambda^2}(\frac{\bar{c}_W}{2} + \bar{c}_B)$	$(-0.0033, 0.0018)$
$\mathcal{O}_{HW} = ig(D^\mu H)^\dagger T_{2k}(D^\nu H)W^k_{\mu\nu}$	$\frac{m_W^2}{\Lambda^2}\bar{c}_{HW}$	$(-0.07, 0.03)$
$\mathcal{O}_{HB} = ig'(D^\mu H)^\dagger (D^\nu H)B_{\mu\nu}$	$\frac{m_W^2}{\Lambda^2}\bar{c}_{HB}$	$(-0.045, 0.075)$

stronger & weaker directions

See also: [Falkowksi & Riva; JHEP 1502 (2015) 039], [Berthier & Trott; JHEP 1505 (2015) 024], [Corbett et al.; JHEP 1508 (2015) 156], [Englert et al.; EPJC 76 (2016) 7, 393]

EFT Benchmarks

- To showcase the usage of both implementations, we select points in c_W, c_{HW} parameter space that:
  - Approximately saturate these limits
  - Select particular Lorentz structures in the new vertices
  - Are also motivated from a BSM point of view
- Tightly constrained direction in (c_B, c_W) forces $c_B \sim -c_W/2$

$$\mathcal{L}_{\rm new} = -\frac{1}{4} g^{(1)}_{hvv} V_{\mu\nu} V^{\mu\nu} h - g^{(2)}_{hvv} V_\nu \partial_\mu V^{\mu\nu} h$$

- We pick benchmark points that single out:

I) $V_\nu \partial_\mu V^{\mu\nu} h \, : \, g^{(1)}_{hvv} = 0, \, g^{(2)}_{hvv} \neq 0 \rightarrow \bar{c}_{HW} = 0, \, \bar{c}_W \neq 0$

II) $V_{\mu\nu} V^{\mu\nu} h \, : \, g^{(2)}_{hvv} = 0, \, g^{(1)}_{hvv} \neq 0 \rightarrow \bar{c}_W = -\bar{c}_{HW}$

EFT Benchmarks

- Pattern II) is a feature of matching conditions that arise in a large class of UV completions, e.g. 2HDM
[Gorbahn, No & Sanz; JHEP 1510 (2015) 036]
- Constraints then become tighter:

$$c_{HW} = -\bar{c}_W = (0.0008, 0.04)$$

- Summary of benchmarks used, roughly compatible with current limits

POWHEG/MCFM	$\bar{c}_{HW}$	$\bar{c}_W$	$\bar{c}_B$	$g^{(1)}_{hvv}$	$g^{(2)}_{hvv}$
I	0	0.008	0	X	✓
II	0.008	-0.008	0	✓	X
MG5_aMC					
A	0.03	0	0	✓	✓
B	0.03	-0.03	0.015	✓	X

Selection of results

- WH, VBF in FR+NLOCT/Madgraph5_aMC@NLO
- Used PYTHIA8 for Higgs decay, PS and Hadronisation
  - Rescaled rates by eHDECAY BRs to capture EFT contributions
- Events were reconstructed using Fastjet thanks to MadAnalysis5 “reco” mode and analysed according to some realistic event selection procedure also in MA5
- Theoretical uncertainties due to scale variation were quantified but not PDF uncertainties
  - Envelope of 9 combinations of (1/2, 2) x μ_0
- See backups for ZH in POWHEG-BOX/MCFM, including SM $gg \to ZH$

[Degrande, Fuks, Mawatari, KM, Sanz; arXiv:1609.04833]

HELatNLO

http://feynrules.irmp.ucl.ac.be/wiki/HELatNLO

- SMEFT implementation in FeynRules + NLOCT framework
  - Generate NLO ready UFO file
  - Simulation performed with MadGraph5_aMC@NLO ~ any process!
  - First results for VBF in SMEFT @ NLO in QCD
- Includes 5 operators affecting Higgs couplings to W/Z/γ
  - First step for EW Higgs production
- Builds upon previous LO implementation of full SILH basis

  [Alloul, Fuks & Sanz; JHEP 1404 (2014) 110]
- Modification of EW parameters taken into account in the (m_Z, α_S, G_F) input scheme

Simulation

```
generate p p > h ve e+ [QCD]
```

- For WH we specified that the Higgs decay to bb while for VBF, decay Higgs to photons
  - PYTHIA8 for Higgs decay (scaled by eHDECAY BR prediction), PS, Hadronisation
- Used NNPDF23_nlo_as_0118_qed PDFs
- Made use of recent MG5 feature to select only interference terms for comparison (LO only)
  - Specify coupling order squared , e.g., “NP^2<=2” to get interference
  - Naive measure of “validity” of EFT interpretation
- Validated results against POWHEG-BOX implementation
  - Found reasonable agreement

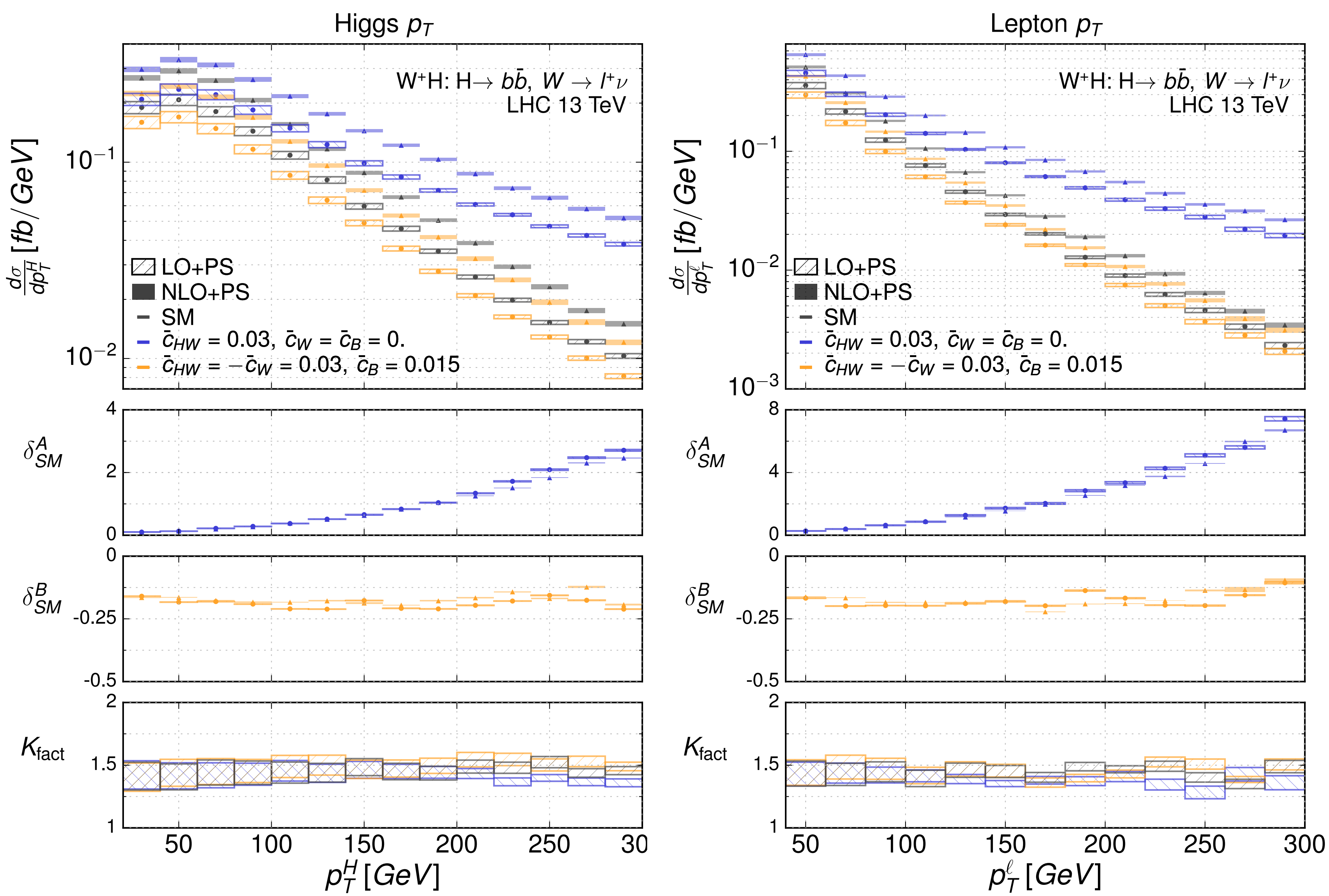
[KM, Sanz, Williams; JHEP 1608 (2016) 039]

pp → W⁺ H → l⁺ ν bb

Benchmarks correspond to ‘large’ values of Wilson coefficients as described in previous analysis since they saturate current limits

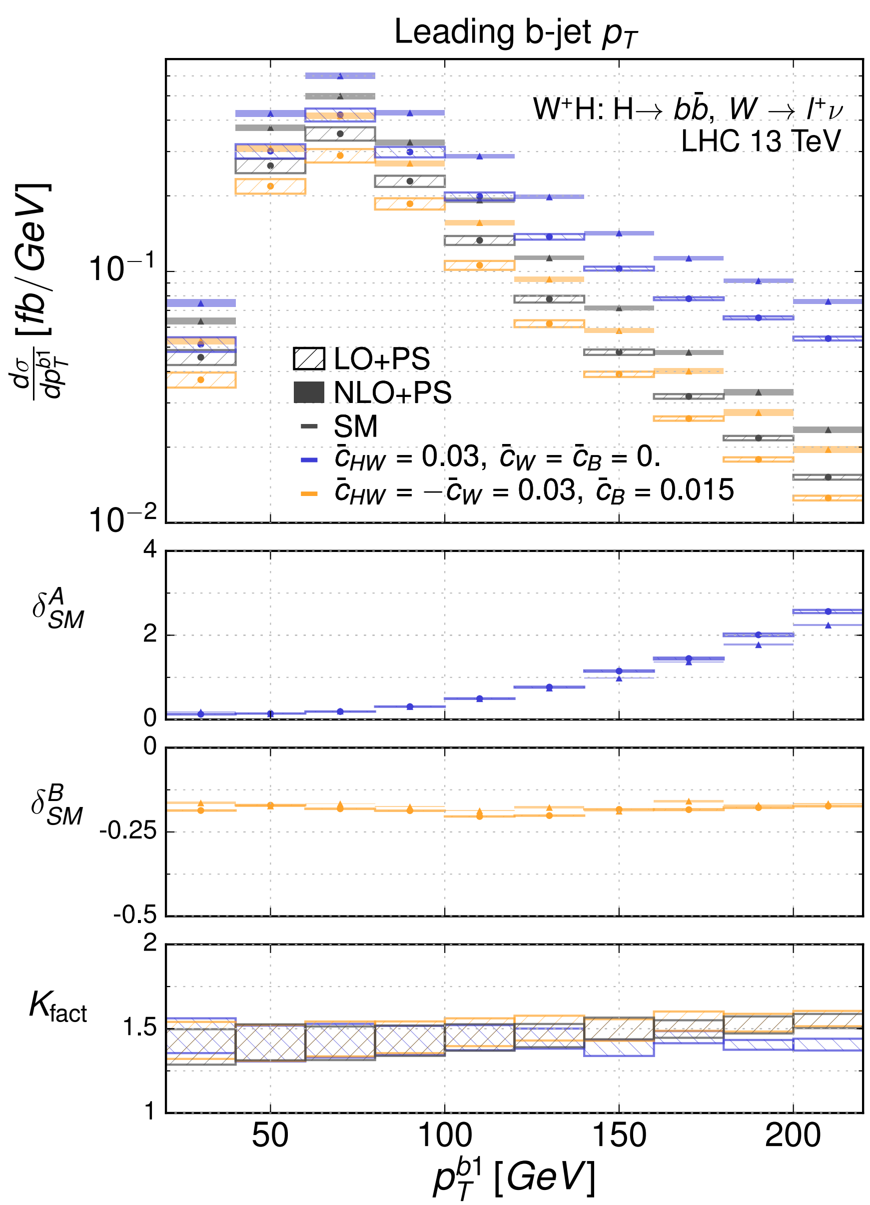
δ = TOT/SM-1

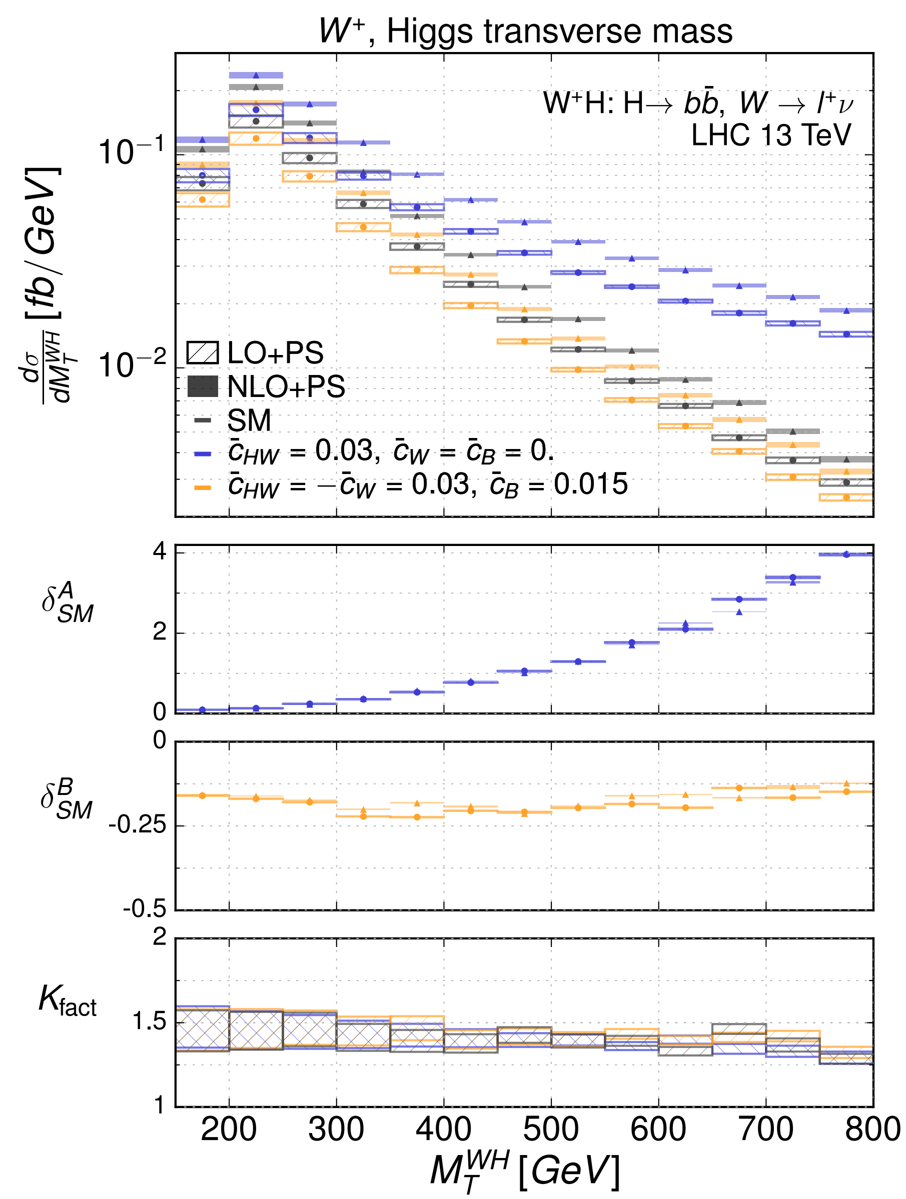
NLO/LO



pp → W⁺ H → l⁺ ν bb

Benchmark **B)** does not exhibit strong "EFT" features
The $g_{hvv}^{(2)}$ Lorentz structure is responsible for these



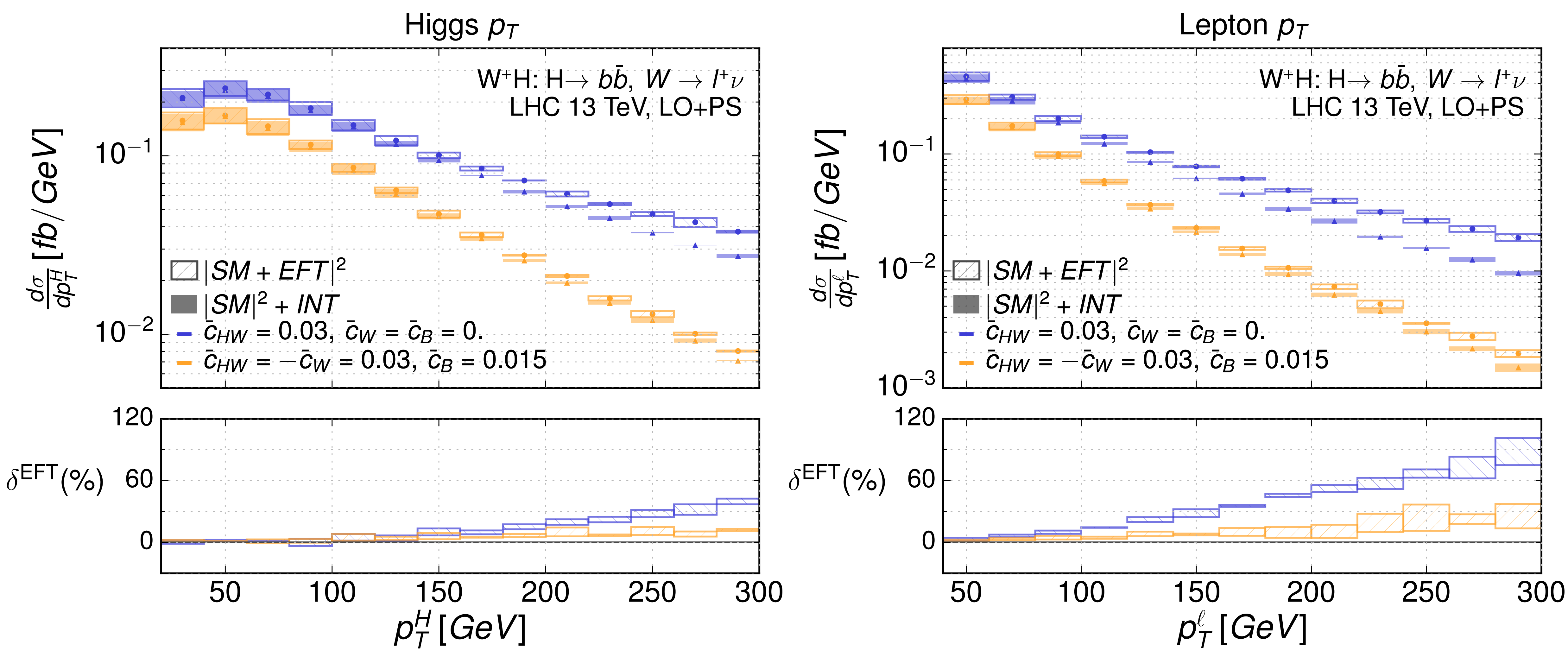


Interference only (LO)

40-80% difference for benchmarks saturating current limits

A possible way to define an additional theory uncertainty?

LHC8 + EWPO not perfect for EFT interpretation of c_W, c_{HW}



pp → H j j→ɣ ɣ j j

```
generate p p > h j j $$ w+ w- z a QCD=0 [QCD]
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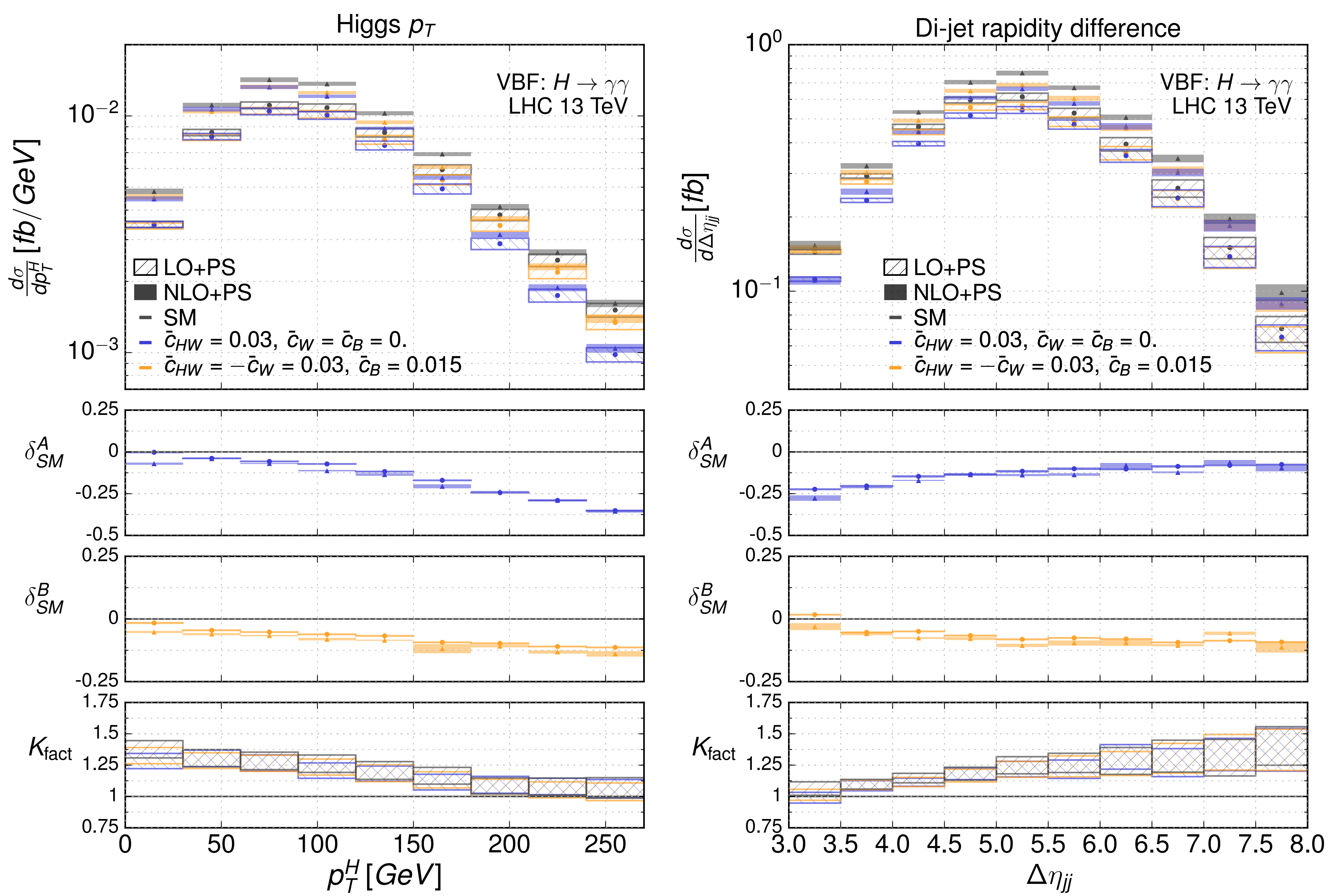
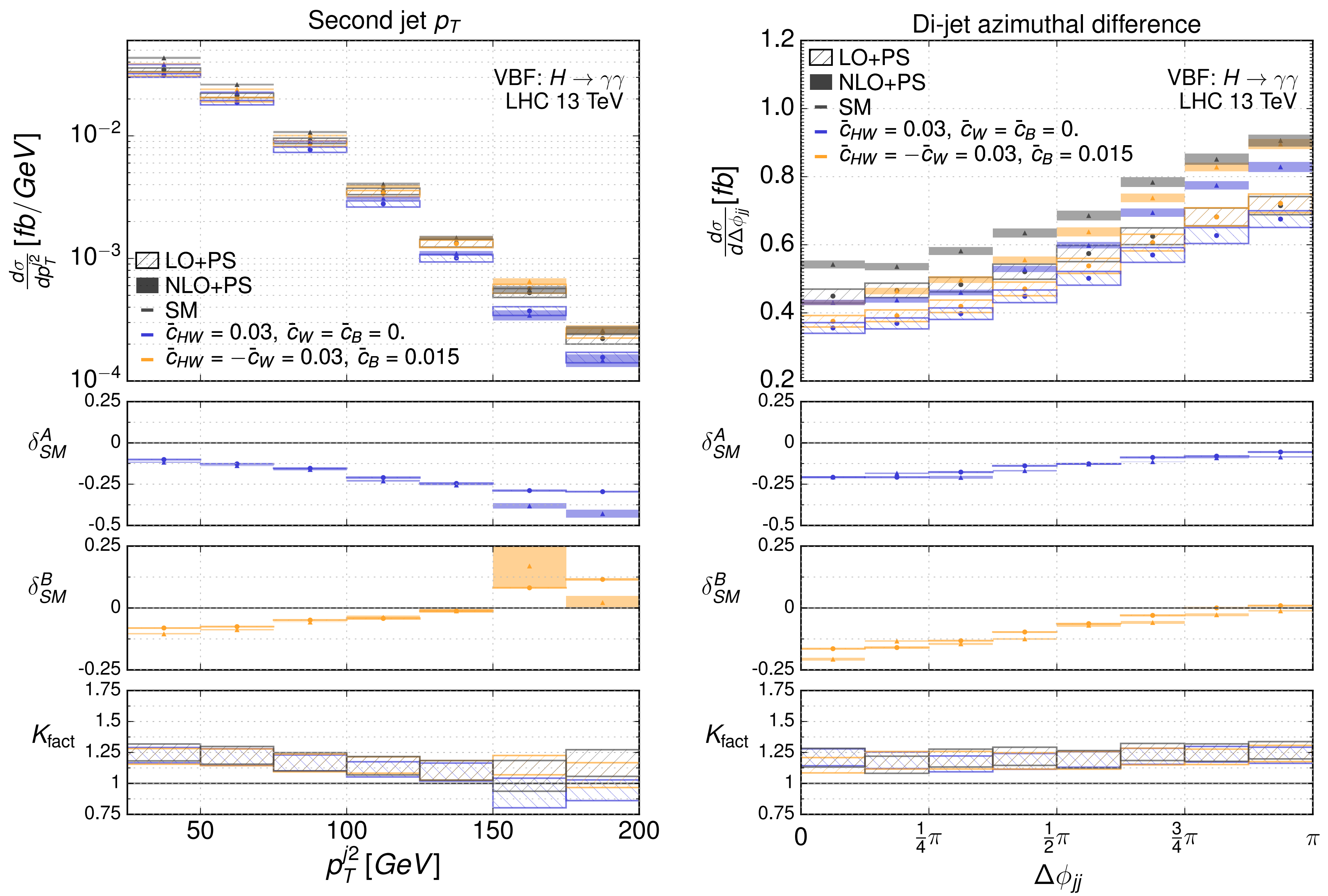
Used a fixed scale of m_W as suggested by literature



K-factors not as flat

$pp \to H\,j\,j \to \gamma\,\gamma\,j\,j$

Generally smaller effects of order 25-50% present, sensitivity to benchmark B. Correlating VH & VBF may help disentangle this coupling structure.



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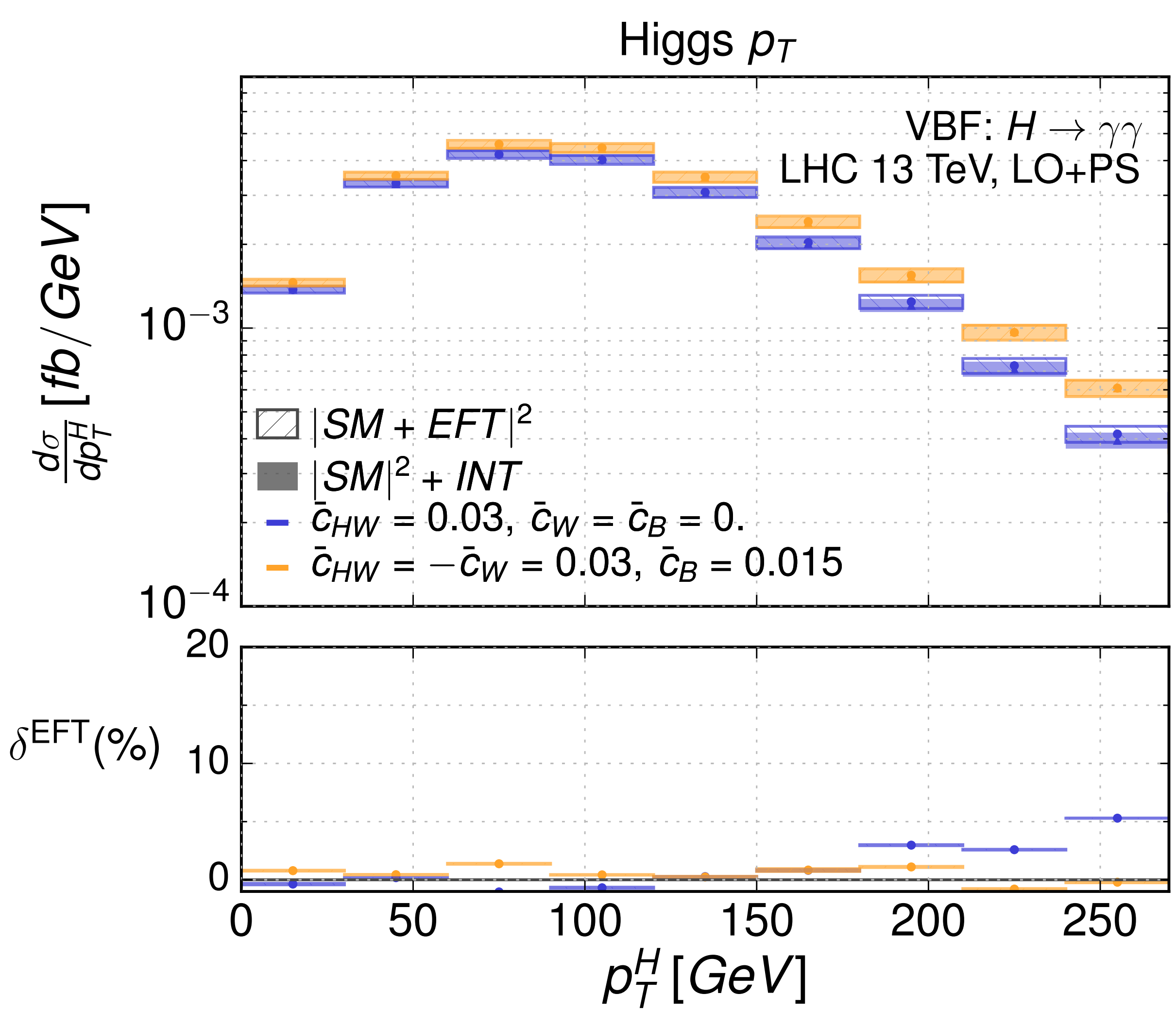
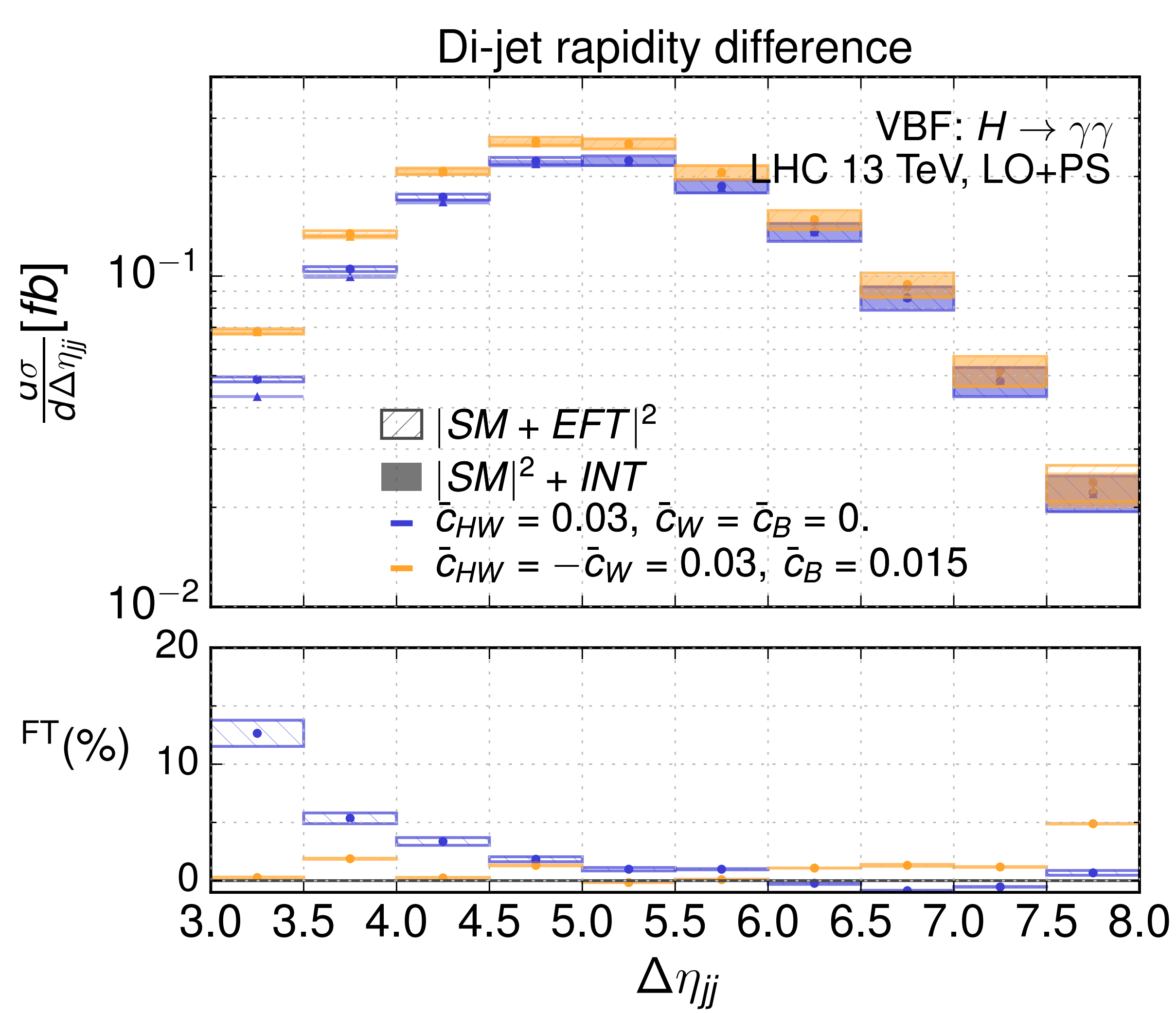
Interference only (LO)

Interference vs. square much more under control.

~10% difference





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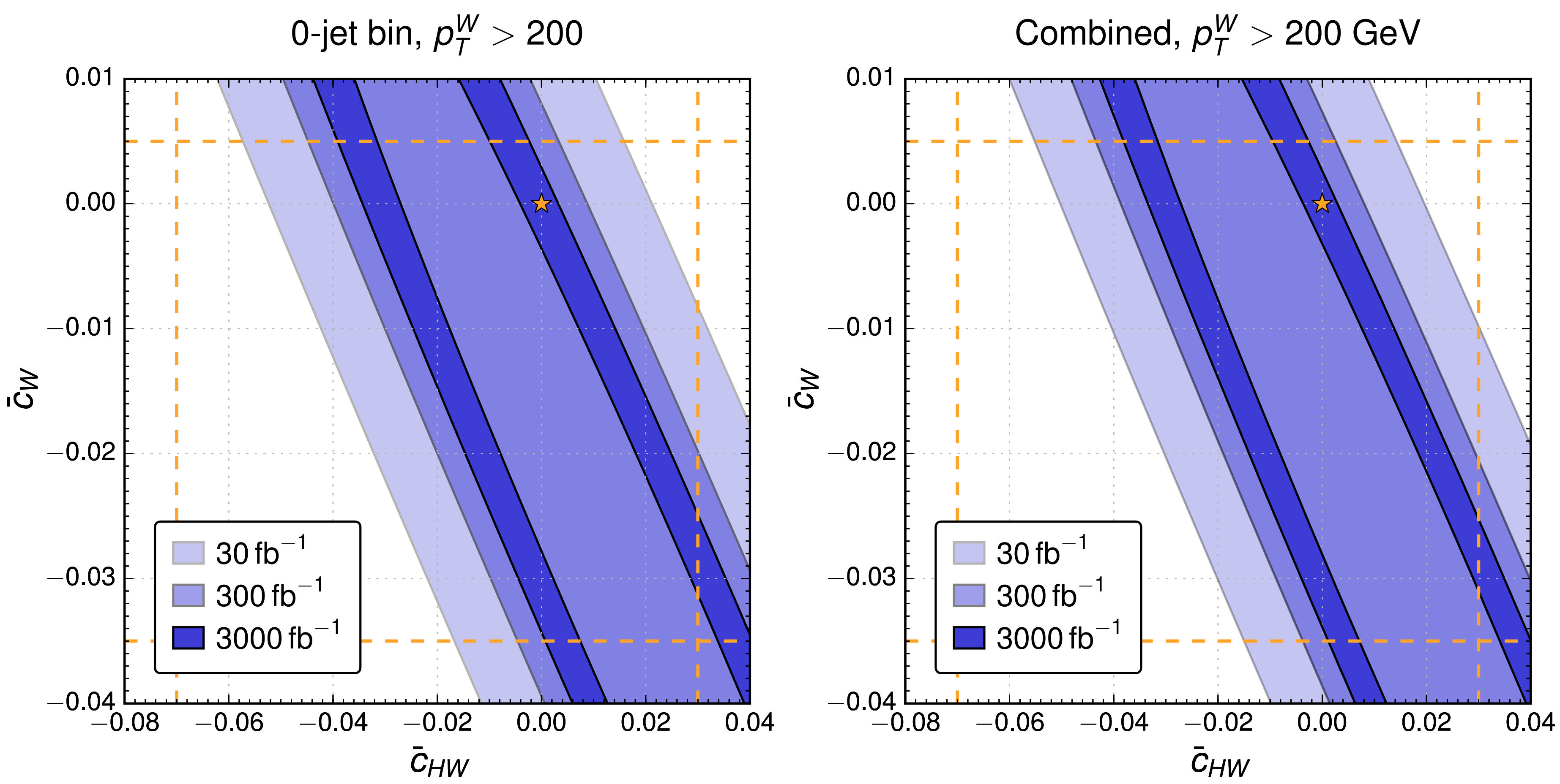
HL-LHC prospects in VH

- 8 & 13 TeV analyses searching for VH → llbb
    - 13 TeV uses multivariate methods = difficult to recast without further info
    - Naive projection of PTV overflow bin of cut-based 8 TeV analysis
- ~per mille sensitivity to c_{HW}, c_W with 3 ab^{-1}



Future

- Several separate implementations of SMEFT operators in different sectors now exist
- Working on a “merge” of these to obtain a complete SMEFT model at NLO in QCD
  - Full set of operators contributing to EW Higgs production processes
  - Validation of anomalous dimension matrix calculation
- Basis independent predictions will be accessible via Rosetta translation tool

  http://rosetta.hepforge.org
- Ultimate goal is to incorporate NLO QCD corrections in a global fit to LHC + low energy data

BACKUP

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LHCTheory ERC meeting
Université Catholique de Louvain

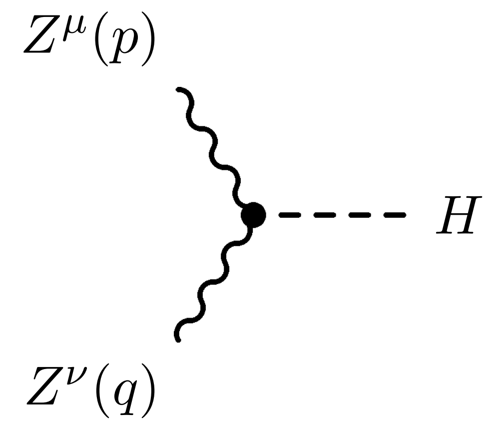
22nd March 2017

Going NLO

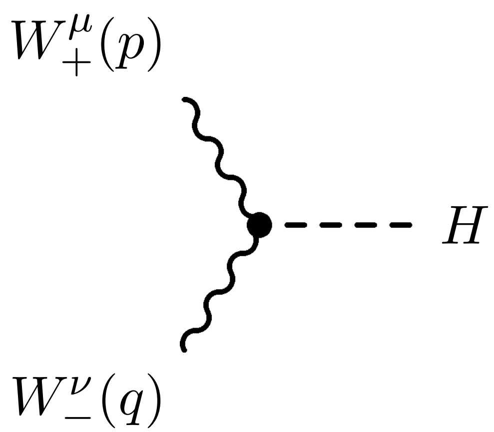
- NLO+PS accurate predictions in QCD are a necessary step for precision EFT analysis at LHC run 2 & beyond
- Other important avenues…
- NLO EW corrections
  - Potentially important but much harder
  - Available for specific processes
  - Automation on the way via Madgraph5_aMC@NLO
- RG-improved predictions thanks to recent anomalous dimension matrix calculation
  - Very helpful for cross checking NLO implementations

[Alonso, Jenkins, Manohar & Trott; JHEP 1310 (2013) 087, JHEP 1401 (2014) 035 & JHEP 1404 (2014) 159*]*

Feynman Rules

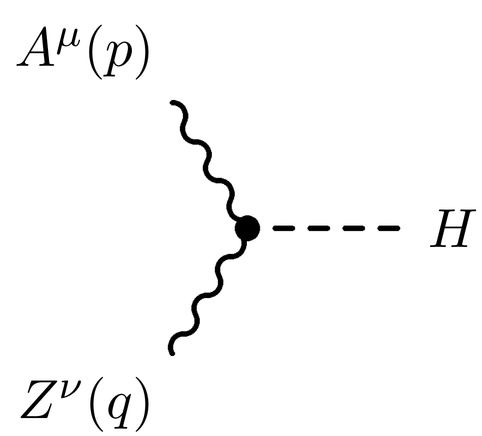


$$:\quad i\Big[\eta^{\mu\nu}\Big(\tfrac{g}{\cos\theta_W}M_Z + g^{(1)}_{hzz}\,p\cdot q + g^{(2)}_{hzz}\,(p^2+q^2)\Big) - g^{(1)}_{hzz}\,q^\mu p^\nu - \tilde{g}_{hzz}\epsilon^{\mu\nu\rho\sigma}\,q_\rho p_\sigma - g^{(2)}_{hzz}\,(p^\mu p^\nu + q^\mu q^\nu)\Big]$$



$$:\quad i\Big[\eta^{\mu\nu}\Big(gM_W + g^{(1)}_{hww}\,p\cdot q + g^{(2)}_{hww}\,(p^2+q^2)\Big) - g^{(1)}_{hww}\,q^\mu p^\nu - \tilde{g}_{hww}\epsilon^{\mu\nu\rho\sigma}\,q_\rho p_\sigma - g^{(2)}_{hww}\,(p^\mu p^\nu + q^\mu q^\nu)\Big]$$

BSM→



$$:\quad i\Big[\eta^{\mu\nu}\Big(g^{(1)}_{haz}\,p\cdot q + g^{(2)}_{haz}\,p^2\Big) - g^{(1)}_{haz}\,q^\mu p^\nu - \tilde{g}_{haz}\epsilon^{\mu\nu\rho\sigma}\,q_\rho p_\sigma - g^{(2)}_{haz}\,p^\mu p^\nu\Big]$$

*Inflowing momenta

POWHEG-BOX/MCFM

- Higgs associated production with a leptonically decaying W or Z at NLO in QCD matched to parton shower
  - Include EFT effects via a mapping to AC/HC (also CP violating)
- At NLO, the initial state current factorises from the final state, even when the Higgs decays to b's
  - Drell-Yan-like NLO corrections which are well known
- Builds upon previous work in the SM matched to parton shower in the same framework as well as fixed order predictions including anomalous couplings
- Matrix elements based on MCFM code interfaced with POWHEG-BOX for which the SM process was already implemented

Simulation (POWHEG)

- For definiteness we specified that the Higgs decay to bb, allowing PYTHIA to perform the decay but scaling the rates by the BR predicted by eHDECAY
- Used CTEQ10 PDFs for NLO predictions and CTEQ6L1 PDFs for LO comparisons
- Modification of EW parameters taken into account in the (m_Z, m_W, G_F) input scheme
- Scale uncertainty determined by varying μ_R, μ_F together around a central scale of $\mu_0 = m_{VH}$
  - Envelope of $\mu_0/2$ and $2\ \mu_0$

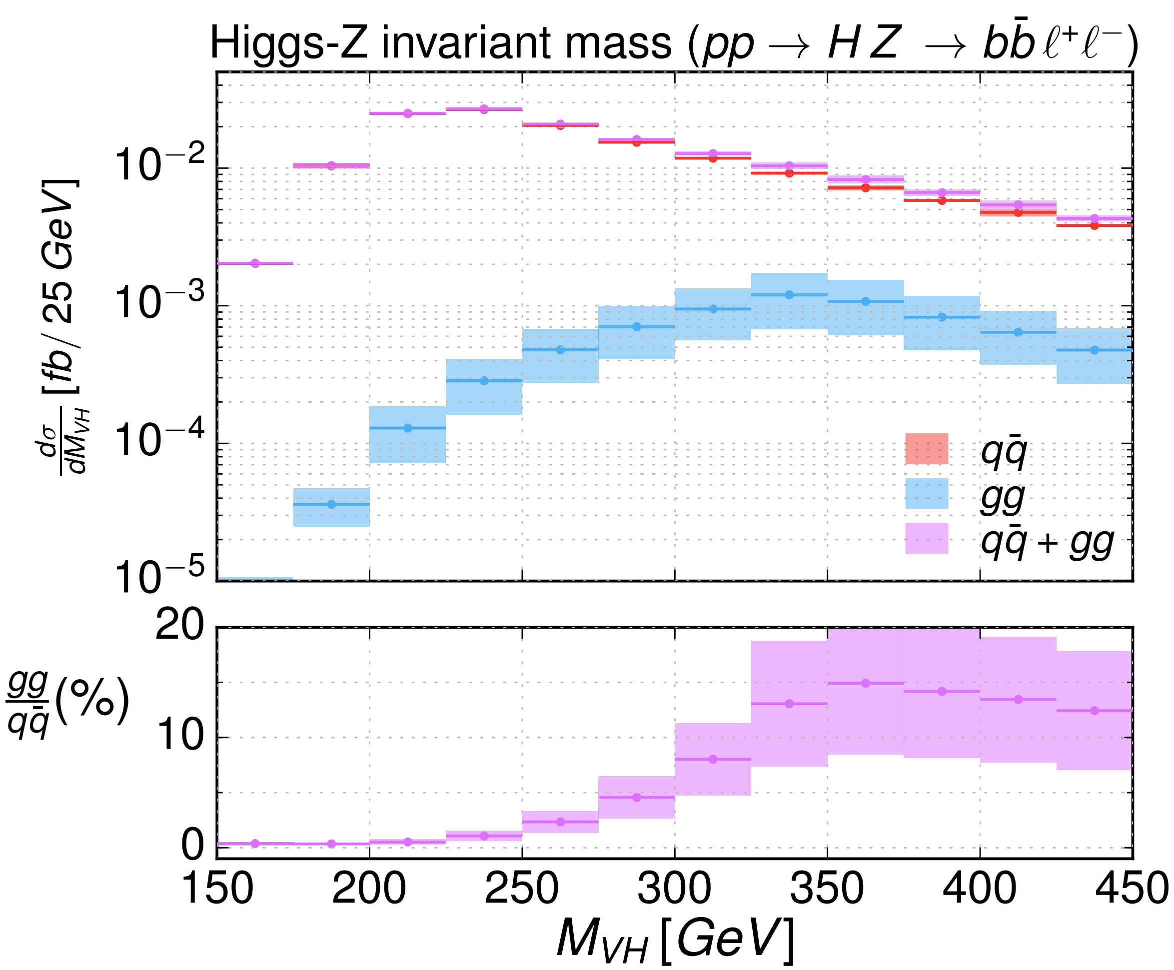
Selection

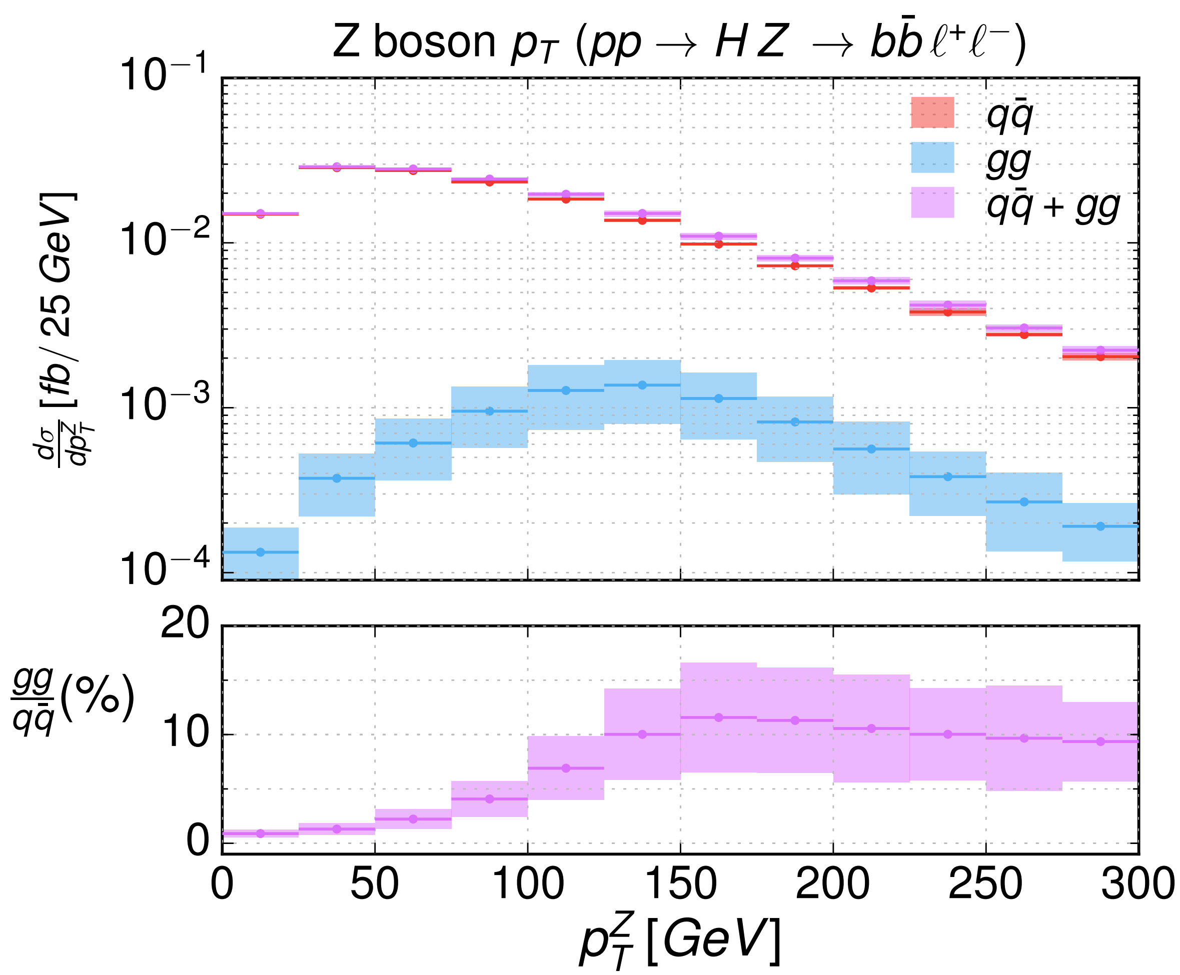
Process	
$H\,Z \to b\bar{b}\,\ell^+\ell^-$	$H\,W \to b\bar{b}\,\ell\nu$
Jets	
k_T algorithm: ΔR=0.4, $p_T > 25$ GeV & $\eta_b < 2.5$	
Cuts	
2 b-jets, $p_T > 25$ GeV, $\eta_b < 2.5$	
1 lepton, $\ell^\pm$ (e or μ)	2 leptons, ℓ^+, ℓ^- (e or μ)
$p_T^\ell < 25$ GeV, $\lvert\eta_\ell\rvert < 2.5$	

MA5 performs b-jet identification based on truth level jet information (presence of b-hadrons in jet)

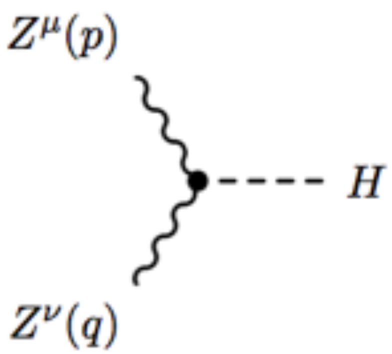
$gg \to Z\,H \to l^+l^-\, bb$

- gg initiated process (formally NNLO)
  - Gluon PDF plus kinematics of EFT searches warrant its inclusion
  - Well known to ‘mimic’ EFT effects if not properly taken into account





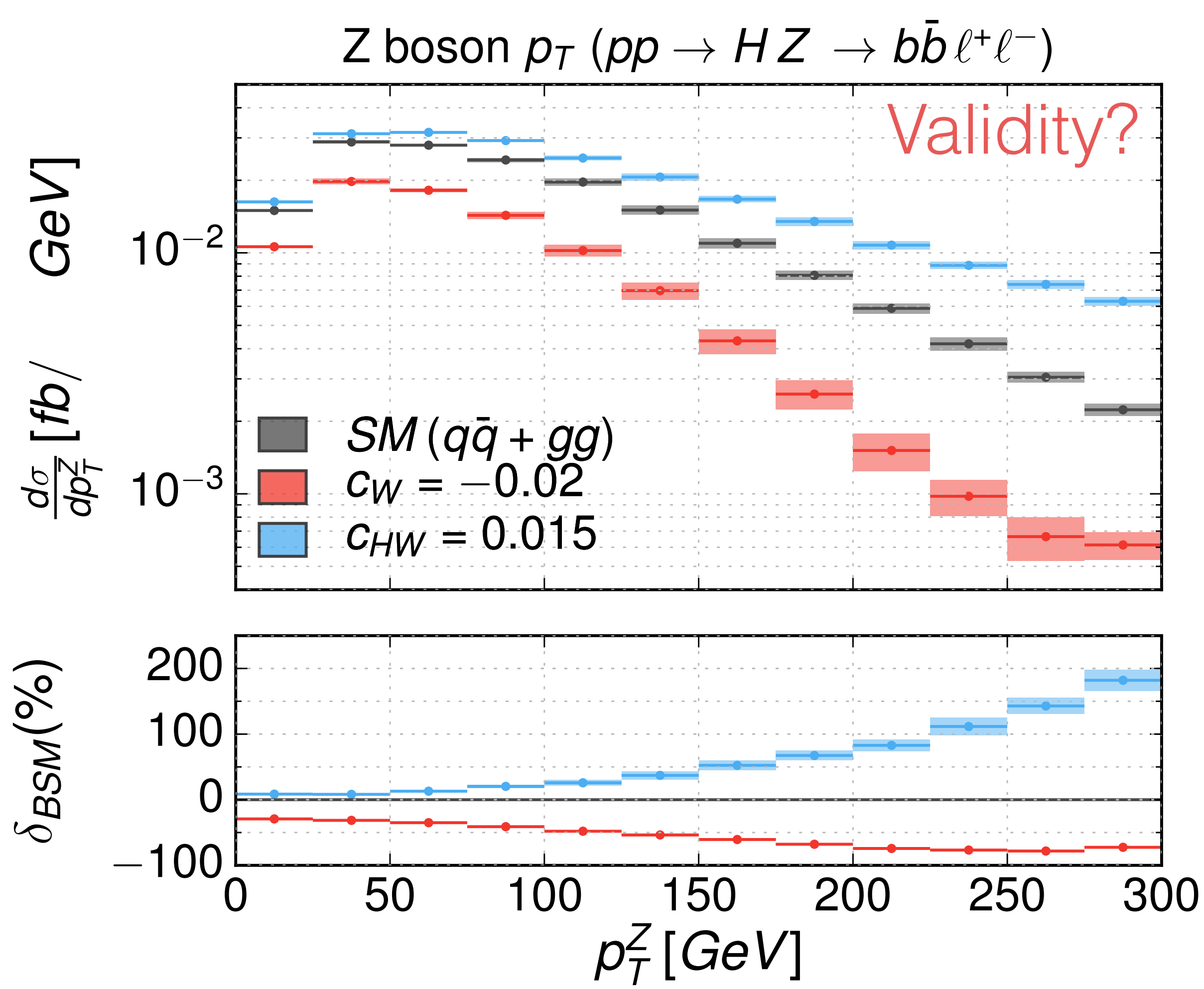
pp → Z H → l+l- bb

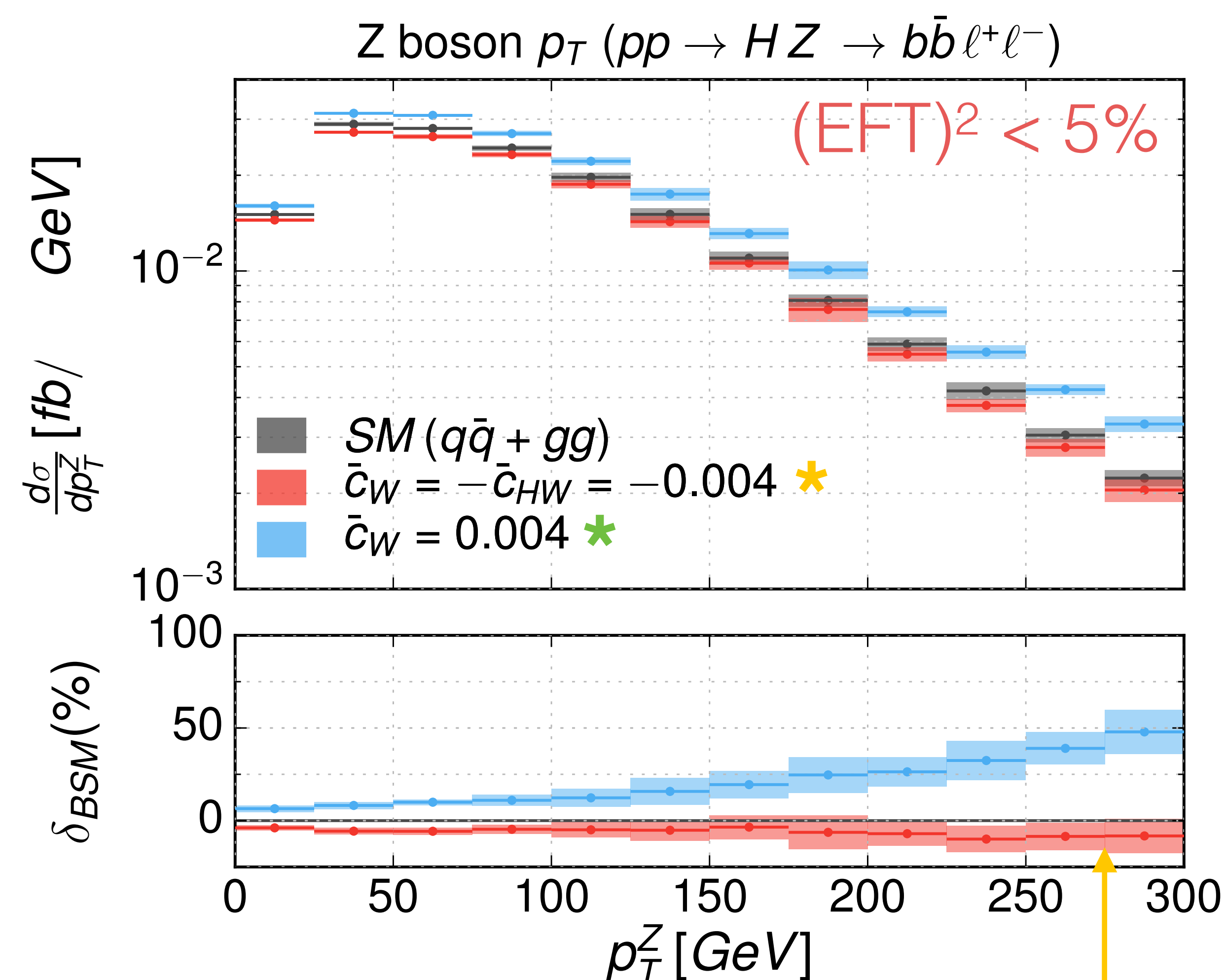


$$g^{(1)}_{hzz} \propto \bar{c}_{HW}, \quad g^{(2)}_{hzz} \propto (\bar{c}_{HW} + \bar{c}_W)$$

$$i\left[\frac{g}{\cos\theta_W} M_Z + g^{(1)}_{hzz}\,(\eta^{\mu\nu} p\cdot q - q^\mu p^\nu) + g^{(2)}_{hzz}\,\left((p^2+q^2)\eta^{\mu\nu} - p^\mu p^\nu + q^\mu q^\nu\right)\right]$$

"BM II" "BM I"

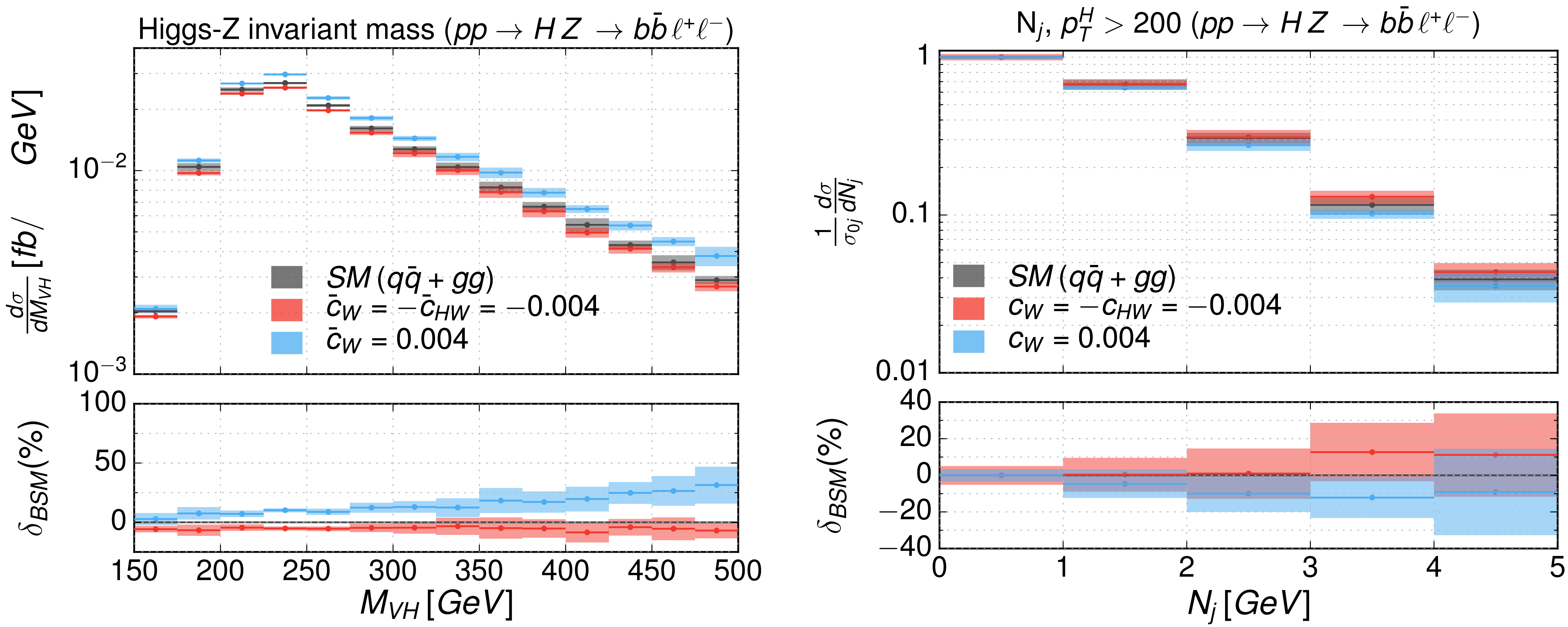




* Benchmark II does not show "EFT-like" features

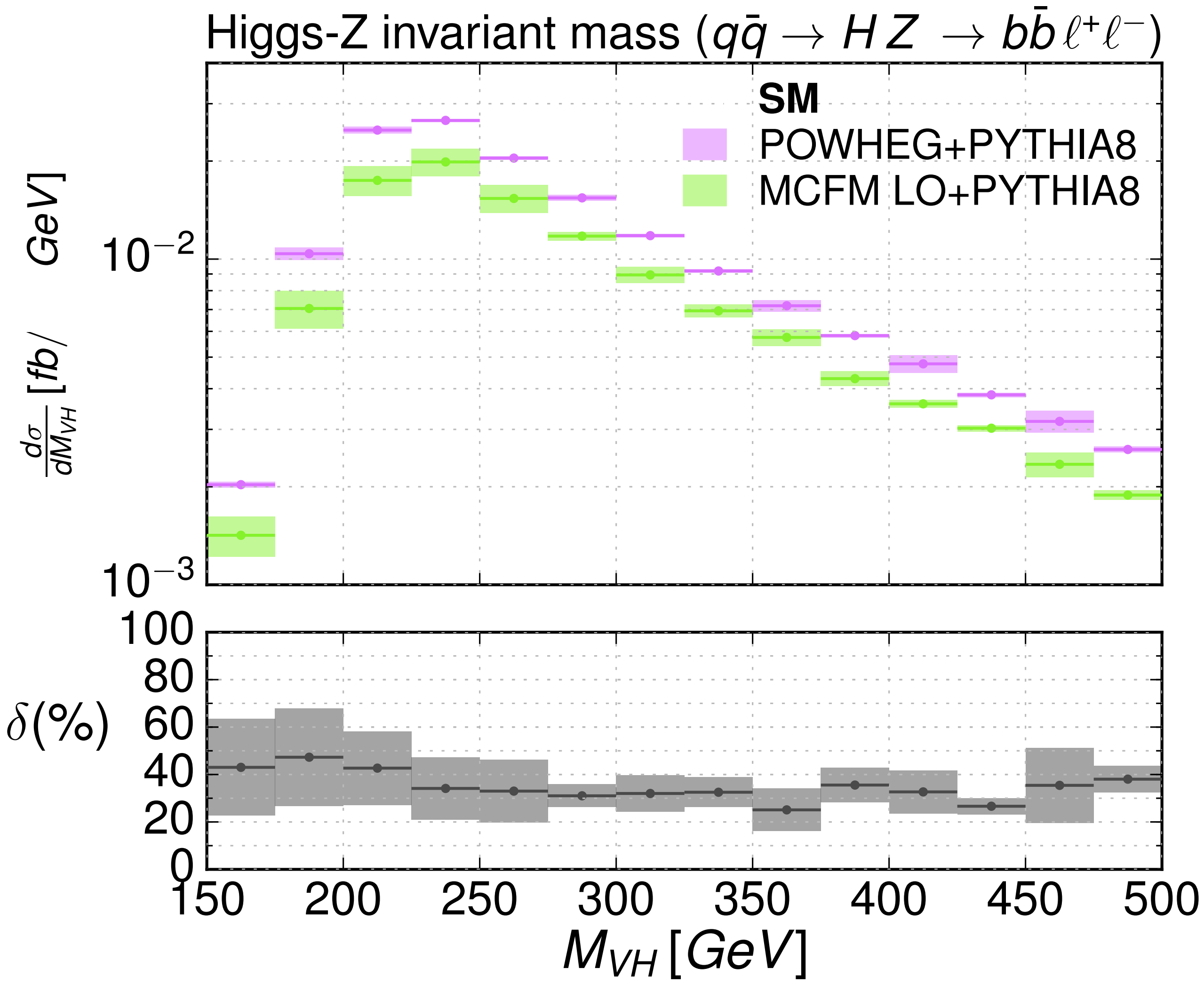
pp → Z H → l+l- bb

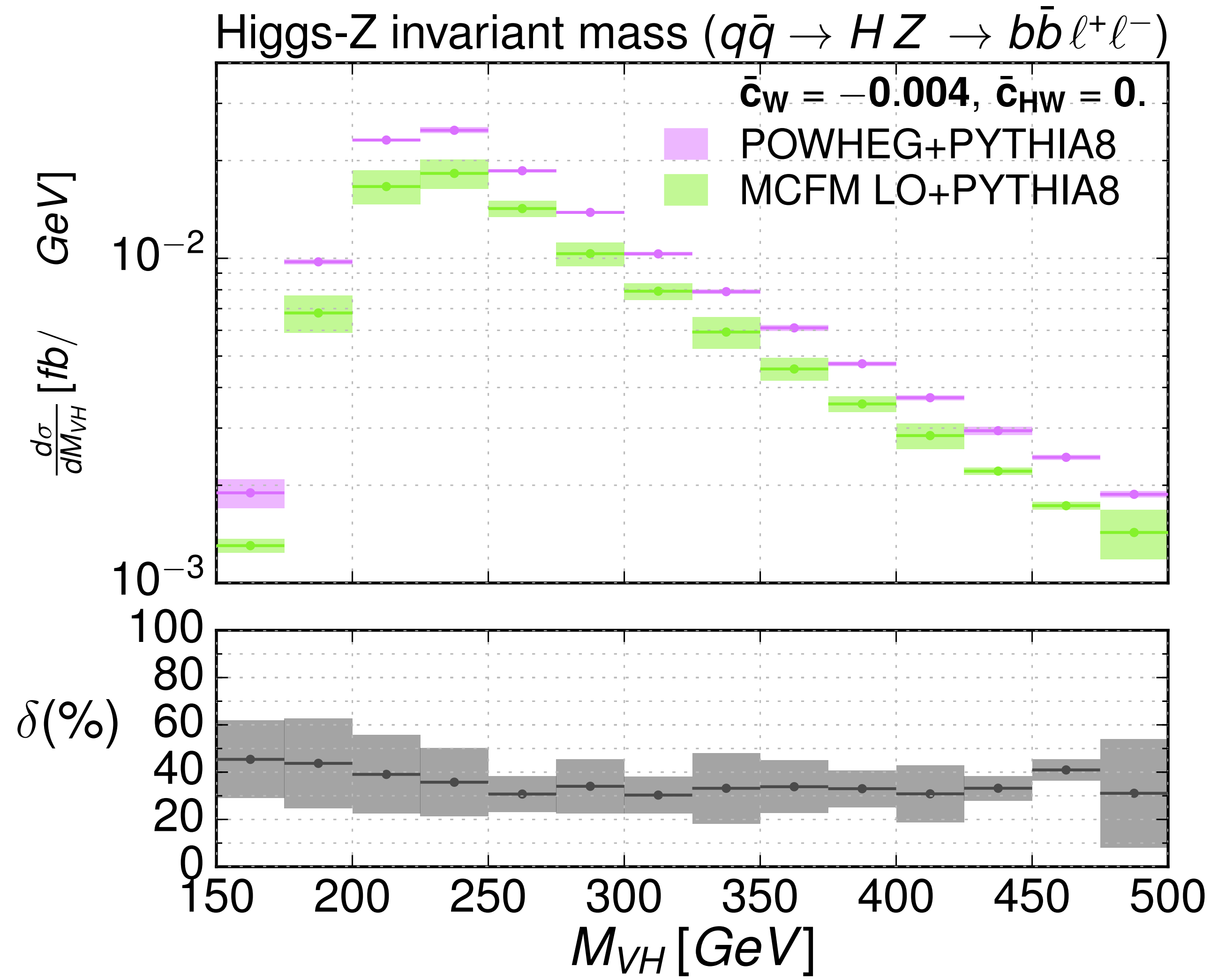
Nj exhibits some difference but stats too low to distinguish



K-factors

No significant difference between SM & EFT
Relatively flat





[R. Contino et al.; Comp. Phys. Comm. 185 (2014) 3412-3423]

eHDECAY

https://www.itp.kit.edu/~maggie/eHDECAY/

- Extension of HDECAY
  [A. Djouadi, J. Kalinowski, M. Spira, Comp. Phys. Comm. 108 (1998) 56]
- QCD corrections
  - qq: Interpolation between massive NLO corrections near threshold & massless $O(a_S^4)$ far above threshold
  - gg: N^3LO in heavy quark limit neglecting higher order terms in $(m_H/\Lambda)^2$
  - $\gamma\gamma$: NLO
  - WW, ZZ, $Z\gamma$: LO
- EFT contribution truncated at $(1/\Lambda)^2$
  - Anomalous coupling (“non-linear”) Lagrangian
  - Alternative ‘SILH’ basis input maps to the anomalous couplings
  - SILH input also includes some NLO EW corrections
  - Unclear if modifications to EW parameters form SM inputs included

Conclusion

- We are near the beginning of a long and fruitful programme of Higgs characterisation via EFT
- Current precision of global fit allow for much room for large EFT deviations and precision is now improved to NLO in QCD+PS
- For VH, k-factors are relatively flat, nothing crazy going on, as expected
  - Also validated implementations between one another for WH
- Seemingly, benchmarks which saturate current limits are slightly beyond the naive “validity” of the EFT for WH while VBF is under control