

Top EFT at lepton colliders

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Introduction

The NP EFT

provides a systematic parametrization of the theory space
in direct vicinity of the SM

- ▶ in a low-energy limit
- ▶ through a proper QFT
- ▶ consistent when global

EFT analysis recipe:

1. Go global!
2. Combine observables!
3. Offer yourself NLO!
 - FCNCs
 - top pair production
 - single top production
 - $t\bar{t}Z$, $t\bar{t}\gamma$
 - $t\bar{t}h$
 - four-fermion operators

[Degrande et al, 14']

[Franzosi et al. 15']

[Zhang 16']

[Bylund et al. 16']

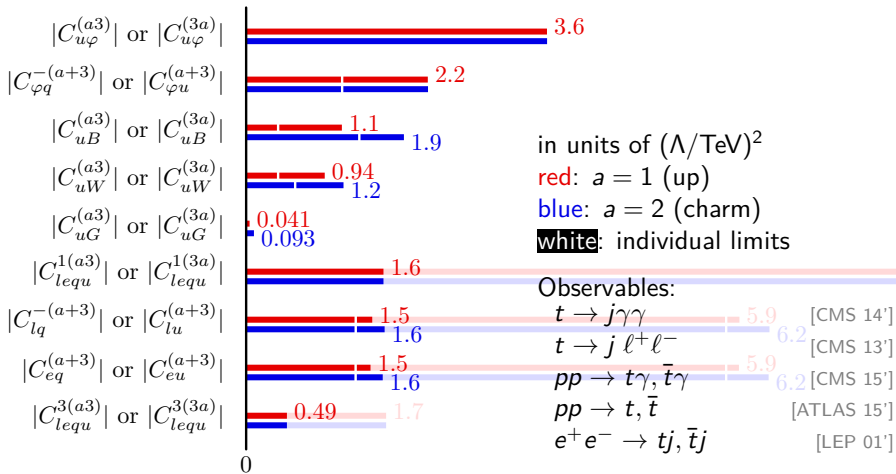
[Maltoni et al. 16']

[$\bar{l}l\bar{q}q$ OK, $\bar{q}q\bar{q}q$ ongoing]

Global top FCNC constraints at NLO in QCD

Showcase example:

[GD et al. 14']

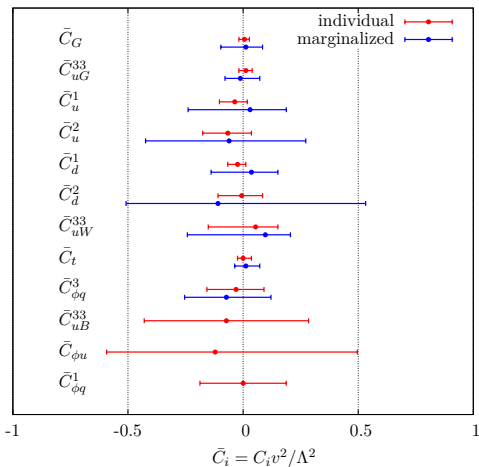


Global top EFT constraints from LHC

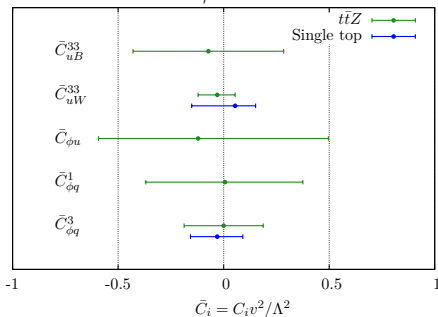
- 195 observables (174 from differential distributions), 12 operators
- mainly from $t\bar{t}$, then single top, charge asymmetries, associated production, W helicity fraction in decay
- standard-model only (N)NLO k-factors in each bin

[Buckley *et al.* Jun. 15]

[Buckley *et al.* Dec. 15]



individual 95% CL limits
from $t\bar{t}Z$ and $t\bar{t}\gamma$:

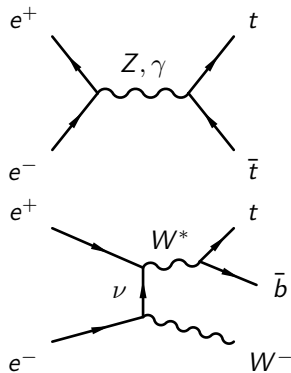
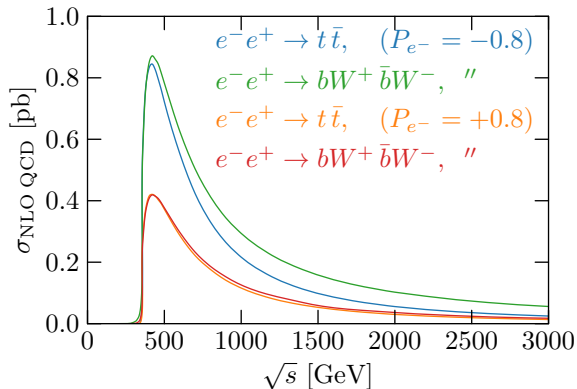


Top EFT at lepton colliders

Top physics at lepton colliders

Handles:

- ▶ $t\bar{t}$ threshold scan \rightarrow mass, width, Yukawa, etc.
- ▶ $t\bar{t}$ continuum:
 - total rate
 - production distributions
 - decay distributions
- ▶ Polarized beams and \sqrt{s} dependence (at linear colliders)
- ▶ + $\sigma(t\bar{t}h)$ above $\sqrt{s} \simeq 500$ GeV (+ $t\bar{t}Z$, $t\bar{t}\gamma$, etc.)



Up-sector neutral current operators

[Grzadkowski et al 10']

Two-quark operators:

Scalar: $O_{u\varphi} \equiv \bar{q} u \tilde{\varphi} \varphi^\dagger \varphi,$

Vector: $O_{\varphi q}^1 \equiv \bar{q} \gamma^\mu q \varphi^\dagger \overleftrightarrow{D}_\mu \varphi \equiv [O_{\varphi q}^+ + O_{\varphi q}^-]/2,$

$O_{\varphi q}^3 \equiv \bar{q} \gamma^\mu \tau^I q \varphi^\dagger \overleftrightarrow{D}_\mu^I \varphi, \equiv [O_{\varphi q}^+ - O_{\varphi q}^-]/2$ (CC also)

$O_{\varphi u} \equiv \bar{u} \gamma^\mu u \varphi^\dagger \overleftrightarrow{D}_\mu \varphi,$

$O_{\varphi ud} \equiv \bar{u} \gamma^\mu d \tilde{\varphi}^\dagger \overleftrightarrow{D}_\mu \varphi,$ (CC only, m_b int.)

Tensor: $O_{uB} \equiv \bar{q} \sigma^{\mu\nu} u \tilde{\varphi} B_{\mu\nu},$

$O_{uW} \equiv \bar{q} \sigma^{\mu\nu} \tau^I u \tilde{\varphi} W_{\mu\nu}^I,$ (CC also)

$O_{dW} \equiv \bar{q} \sigma^{\mu\nu} \tau^I d \tilde{\varphi} W_{\mu\nu}^I,$ (CC only, m_b int.)

$O_{uG} \equiv \bar{q} \sigma^{\mu\nu} T^A u \tilde{\varphi} G_{\mu\nu}^A.$

Two-quark–two-lepton operators:

Scalar: $O_{1e q u}^S \equiv \bar{l} e \varepsilon \bar{q} u,$ (CC also, m_e int.)

$O_{1e d q} \equiv \bar{l} e \bar{d} q,$ (CC only, m_e int.)

Vector: $O_{1q}^1 \equiv \bar{l} \gamma_\mu l \bar{q} \gamma^\mu q \equiv [O_{1q}^+ + O_{1q}^-]/2,$

$O_{1q}^3 \equiv \bar{l} \gamma_\mu \tau^I l \bar{q} \gamma^\mu \tau^I q \equiv [O_{1q}^+ - O_{1q}^-]/2,$ (CC also)

$O_{1u} \equiv \bar{l} \gamma_\mu l \bar{u} \gamma^\mu u,$

$O_{eq} \equiv \bar{e} \gamma^\mu e \bar{q} \gamma_\mu q,$

$O_{eu} \equiv \bar{e} \gamma_\mu e \bar{u} \gamma^\mu u,$

Tensor: $O_{1e q u}^T \equiv \bar{l} \sigma_{\mu\nu} e \varepsilon \bar{q} \sigma^{\mu\nu} u.$ (CC also, m_e int.)

Independent coefficients

Two-quark operators: 10 real degrees of freedom

Scalar: $C_{u\varphi}^{(33)}$,

Vector: $C_{\varphi q}^{+(33)} = C_{\varphi q}^{+(33)*}$, (down-Z, tbW)
 $C_{\varphi q}^{V(33)} = C_{\varphi q}^{V(33)*} \equiv C_{\varphi u}^{(33)} + C_{\varphi q}^{-(33)}$, (up-Z, tbW)
 $C_{\varphi q}^{A(33)} = C_{\varphi q}^{A(33)*} \equiv C_{\varphi u}^{(33)} - C_{\varphi q}^{-(33)}$, (up-Z, tbW)
 $C_{\varphi ud}^{(33)}$

Tensor: $C_{uA}^{(33)} \equiv C_{uW}^{(33)} + C_{uB}^{(33)}$,
 $C_{uZ}^{(33)} \equiv \cotan \theta_W C_{uW}^{(33)} - \tan \theta_W C_{uB}^{(33)}$,
 $C_{uG}^{(33)}$

Two-quark–two-lepton operators: 9×3^2 real degrees of freedom

Scalar: $C_{lequ}^{S(33)}$,

Vector: $C_{1q}^{+(33)} = C_{1q}^{+(33)*}$, (up- ν , down- ℓ)
 $C_{1q}^{V(33)} = C_{1q}^{V(33)*} \equiv C_{1u}^{(33)} + C_{1q}^{-(33)}$, (up- ℓ)
 $C_{1q}^{A(33)} = C_{1q}^{A(33)*} \equiv C_{1u}^{(33)} - C_{1q}^{-(33)}$, (up- ℓ)
 $C_{eq}^{V(33)} = C_{eq}^{V(33)*} \equiv C_{eu}^{(33)} + C_{eq}^{(33)}$, (up- ℓ , down- ℓ)
 $C_{eq}^{A(33)} = C_{eq}^{A(33)*} \equiv C_{eu}^{(33)} - C_{eq}^{(33)}$,

Tensor: $C_{lequ}^{T(33)}$.

Anomalous vertices

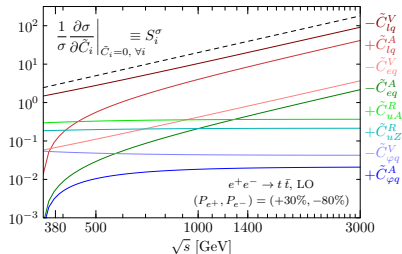
$$\begin{aligned}
 t\bar{t}\gamma: & \quad \gamma^\mu \overbrace{(F_{1V}^\gamma + \gamma_5 F_{1A}^\gamma)}^{\sim \phi} + \frac{\sigma_{\mu\nu} i q^\nu}{2m_t} \overbrace{(F_{2V}^\gamma + i\gamma_5 F_{2A}^\gamma)}^{\sim \text{Re, Im}\{C_{uA}\}} \\
 t\bar{t}Z: & \quad \gamma^\mu \overbrace{(F_{1V}^Z + \gamma_5 F_{1A}^Z)}^{\sim C_{\varphi u}^V, C_{\varphi u}^A} + \frac{\sigma_{\mu\nu} i q^\nu}{2m_t} \overbrace{(F_{2V}^Z + i\gamma_5 F_{2A}^Z)}^{\sim \text{Re, Im}\{C_{uZ}\}} \\
 t\bar{t}W: & \quad \gamma^\mu \overbrace{(F_{1V}^W + \gamma_5 F_{1A}^W)}^{\sim C_{\varphi ud}, C_{\varphi q}^+ - \frac{1}{2}(C_{\varphi q}^V - C_{\varphi q}^A)} + \frac{\sigma_{\mu\nu} i q^\nu}{2m_t} \overbrace{(F_{2V}^W + i\gamma_5 F_{2A}^W)}^{\sim \text{Re, Im}\{s_W^2 C_{uA} + s_W c_W C_{uZ}\}}
 \end{aligned}$$

Insufficiencies:

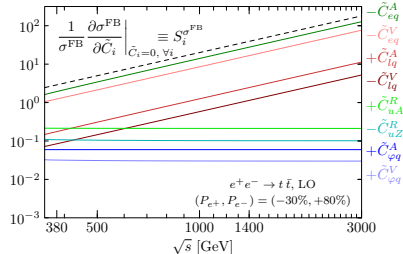
- Miss four-fermion operators
- Conflict with gauge invariance
 - Do not allow for radiative corrections to be computed
- Complex couplings where the tree-level EFT prescribes real ones
- Hide correlations induced by gauge invariance
 - Preclude the combination of measurements in various sectors

Operator sensitivities in $e^+e^- \rightarrow t\bar{t}$

Total cross section (left pol.):

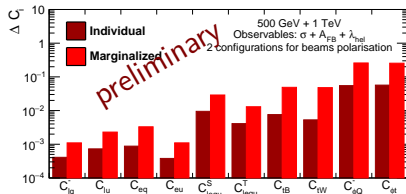


FB-integrated cross section (right pol.):



Few features:

- quadratic energy growth for four-fermion operators
- no growth for two-fermion operators
Some azimuthal $bW^+ \bar{b}W^-$ observables have linearly growing sensitivity to C_{uZ}, C_{uA} .
- p -wave $\beta = \sqrt{1 - m_t^2/s}$ suppression of axial vectors at threshold
- enhanced sensitivity of axial vector operators in σ^{FB}
- sensitivity sign flip for $C_{\varphi q}^V$ and C_{uZ}^R when polarization is reversed
- etc.



Helicity amplitude decomposition in $bW^+\bar{b}W^-$

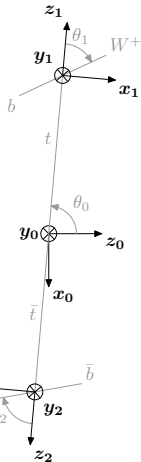
[Jacob,Wick 59']

Production amplitudes: $A_1(++) \sim \frac{2m_t}{\sqrt{s}}(V+A) + \sqrt{s}(T - \beta\tilde{T})$
 $A_2(--)\sim \frac{2m_t}{\sqrt{s}}(V+A) + \sqrt{s}(T + \beta\tilde{T})$
 $A_3(+-)\sim (V + \beta A) + 2m_t(T + \tilde{T})$
 $A_4(-+)\sim (V - \beta A) + 2m_t(T + \tilde{T})$

[Schmidt 95']

In terms of $\Omega = \{\theta_0, \theta_1, \phi_1, \theta_2, \phi_2\}$ helicity angles:

	+3/4	$(A_3 ^2 + A_4 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$(1 + \cos^2 \theta_0)$			
	+3/4	$(A_3 ^2 - A_4 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$(1 + \cos^2 \theta_0)$		$\cos \theta_2$	
	+3/4	$(A_3 ^2 - A_4 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$(1 + \cos^2 \theta_0)$	$\cos \theta_1$		
	+3/4	$(A_3 ^2 + A_4 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$(1 + \cos^2 \theta_0)$	$\cos \theta_1$		$\cos \theta_2$
	-3/2	$(A_3 ^2 - A_4 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$\cos \theta_0$			
	-3/2	$(A_3 ^2 + A_4 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\cos \theta_0$			$\cos \theta_2$
	-3/2	$(A_3 ^2 + A_4 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$\cos \theta_0$	$\cos \theta_1$		
	-3/2	$(A_3 ^2 - A_4 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\cos \theta_0$	$\cos \theta_1$		$\cos \theta_2$
	+3/2	$(A_1 ^2 + A_2 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$\sin^2 \theta_0$			
	-3/2	$(A_1 ^2 - A_2 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin^2 \theta_0$			$\cos \theta_2$
	+3/2	$(A_1 ^2 - A_2 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$\sin^2 \theta_0$	$\cos \theta_1$		
	-3/2	$(A_1 ^2 + A_2 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin^2 \theta_0$	$\cos \theta_1$		$\cos \theta_2$
	+3/2	$\sqrt{2} \operatorname{Re}(A_1^* A_4)$	$ a_2 ^2 - a_4 ^2$	$ b_3 ^2$	$\sin \theta_0(1 + \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\cos \phi_1$
	+3/2	$\sqrt{2} \operatorname{Re}(A_1^* A_4)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2$	$\sin \theta_0(1 + \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\cos \phi_1$
	+3/2	$\sqrt{2} \operatorname{Re}(A_2^* A_3)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2$	$\sin \theta_0(1 - \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\cos \phi_1$
	+3/2	$\sqrt{2} \operatorname{Re}(A_2^* A_3)$	$ a_2 ^2 - a_4 ^2$	$ b_3 ^2$	$\sin \theta_0(1 - \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\cos \phi_1$
	-3/2	$\sqrt{2} \operatorname{Re}(A_2^* A_4)$	$ a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0(1 + \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\cos \phi_2$
	-3/2	$\sqrt{2} \operatorname{Re}(A_2^* A_4)$	$ a_2 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0(1 + \cos \theta_0)$	$(1 - \cos \theta_1)$	$\sin \theta_2$	$\cos \phi_2$
	-3/2	$\sqrt{2} \operatorname{Re}(A_1^* A_3)$	$ a_2 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0(1 - \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\cos \phi_2$
	-3/2	$\sqrt{2} \operatorname{Re}(A_1^* A_3)$	$ a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0(1 - \cos \theta_0)$	$(1 - \cos \theta_1)$	$\sin \theta_2$	$\cos \phi_2$
	-3	$\operatorname{Re}(A_1^* A_2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin^2 \theta_0$	$\sin \theta_1$	$\sin \theta_2$	$\cos(\phi_1 + \phi_2)$
	-3/2	$\operatorname{Re}(A_3^* A_4)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin^2 \theta_0$	$\sin \theta_1$	$\sin \theta_2$	$\cos(\phi_1 - \phi_2)$
	+3/2	$\sqrt{2} \operatorname{Im}(A_1^* A_4)$	$ a_2 ^2 - a_4 ^2$	$ b_3 ^2$	$\sin \theta_0(1 + \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\sin \phi_1$
	+3/2	$\sqrt{2} \operatorname{Im}(A_1^* A_4)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2$	$\sin \theta_0(1 + \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\sin \phi_1$
	+3/2	$\sqrt{2} \operatorname{Im}(A_2^* A_3)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2$	$\sin \theta_0(1 - \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\sin \phi_1$
	-3/2	$\sqrt{2} \operatorname{Im}(A_2^* A_3)$	$ a_2 ^2 - a_4 ^2$	$ b_3 ^2$	$\sin \theta_0(1 - \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\sin \phi_1$
	+3/2	$\sqrt{2} \operatorname{Im}(A_2^* A_4)$	$ a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0(1 + \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$
	+3/2	$\sqrt{2} \operatorname{Im}(A_2^* A_4)$	$ a_2 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0(1 + \cos \theta_0)$	$(1 - \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$
	-3/2	$\sqrt{2} \operatorname{Im}(A_1^* A_3)$	$ a_2 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0(1 - \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$
	-3/2	$\sqrt{2} \operatorname{Im}(A_1^* A_3)$	$ a_2 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0(1 - \cos \theta_0)$	$(1 - \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$



NLO in QCD for $e^+e^- \rightarrow bW^+\bar{b}W^-$

For various beam polarizations and center-of-mass energies:

pol	\sqrt{s} [GeV]	σ_{SM} [fb]	$ A_{lq}^+ $	$ A_{lq}^- $	$ A_{lq}^0 $	$ V_{lq}^+ $	$ V_{lq}^- $	σ_{γ}^b [fb]	$ R_{W^+}^b $	$ R_{W^-}^b $	$ R_{W^+}^b $	$ R_{W^-}^b $	$ R_{W^+}^b $	$ R_{W^-}^b $
00	300	$2.92^{+1\%}_{-1.15\%}$	$0.353^{+1\%}_{-1.15\%}$	$-0.0855^{+2\%}_{-1.27\%}$	$0.14^{+2\%}_{-1.2\%}$	$-0.621^{+2\%}_{-1.34\%}$	$-0.303^{+2\%}_{-1.31\%}$	$-0.136^{+2\%}_{-1.21\%}$	$0.349^{+2\%}_{-1.33\%}$	$0.32^{+3\%}_{-1.33\%}$	$-0.000225^{+9\%}_{-9\%}$	$-0.000125^{+9\%}_{-9\%}$	$0.000214^{+6\%}_{-6\%}$	
00	380	$825^{+2\%}_{-1.18\%}$	$77.1^{+3\%}_{-1.34\%}$	$-55.6^{+2\%}_{-1.32\%}$	$53.1^{+1\%}_{-1.2\%}$	$-99.3^{+2\%}_{-1.19\%}$	$-63.5^{+1\%}_{-1.16\%}$	$-62.9^{+1\%}_{-1.19\%}$	$111^{+2\%}_{-1.14\%}$	$323^{+2\%}_{-1.18\%}$	$0.107^{+2\%}_{-1.36\%}$	$-0.434^{+10\%}_{-10\%}$	$0.25^{+10\%}_{-9\%}$	
00	500	$669^{+0.4\%}_{-0.95\%}$	$250^{+0.5\%}_{-0.95\%}$	$-233^{+0.5\%}_{-0.95\%}$	$49.2^{+0.6\%}_{-0.96\%}$	$-1230^{+0.8\%}_{-0.92\%}$	$-750^{+0.8\%}_{-0.92\%}$	$-1.10^{+0.8\%}_{-0.92\%}$	$102^{+0.5\%}_{-0.97\%}$	$263^{+0.7\%}_{-0.82\%}$	$-2.08^{+5\%}_{-5\%}$	$1.78^{+8\%}_{-8\%}$	$0.715^{+3\%}_{-3\%}$	
00	1000	$221^{+0.5\%}_{-0.89\%}$	$756^{+0.9\%}_{-1.21\%}$	$-475^{+2.0\%}_{-0.83\%}$	$15.6^{+4\%}_{-0.84\%}$	$-1070^{+2.0\%}_{-0.784\%}$	$-940^{+6\%}_{-0.95\%}$	$-15.5^{+3\%}_{-0.944\%}$	$36.1^{+5\%}_{-0.995\%}$	$87.3^{+3\%}_{-0.89\%}$	$0.392^{+2\%}_{-2.00\%}$	$-8.74^{+4\%}_{-4\%}$	$0.907^{+5\%}_{-5\%}$	
00	1400	$132^{+0.6\%}_{-0.936\%}$	$391^{+7\%}_{-0.555\%}$	$-412^{+3\%}_{-0.903\%}$	$8.29^{+2\%}_{-0.803\%}$	$-1460^{+0.9\%}_{-0.926\%}$	$-816^{+3\%}_{-0.794\%}$	$-8.9^{+2\%}_{-0.859\%}$	$21.6^{+18\%}_{-1.03\%}$	$59.8^{+18\%}_{-1.08\%}$	$1.8^{+8\%}_{-0.8\%}$	$0.257^{+2\%}_{-0.908\%}$	$0.414^{+10\%}_{-10\%}$	
00	3000	$40.2^{+1.0\%}_{-1.1\%}$	$1080^{+1\%}_{-1.1\%}$	$-70.4^{+2\%}_{-0.128\%}$	$0.323^{+10\%}_{-0.60\%}$	$-1270^{+2\%}_{-0.981\%}$	$-689^{+10\%}_{-0.688\%}$	$-0.717^{+3\%}_{-0.400\%}$	$10.3^{+3\%}_{-0.53\%}$	$14.5^{+2\%}_{-0.45\%}$	$1.85^{+10\%}_{-0.44\%}$	$2.15^{+10\%}_{-0.394\%}$	$-0.261^{+2\%}_{-0.6\%}$	
+-	300	$2.73^{+0.2\%}_{-1.14\%}$	$0.351^{+0.2\%}_{-1.14\%}$	—	$0.126^{+2\%}_{-1.19\%}$	$-0.62^{+2\%}_{-1.34\%}$	—	$-0.14^{+2\%}_{-1.22\%}$	$0.376^{+2\%}_{-1.23\%}$	$0.241^{+2\%}_{-1.28\%}$	$6.25e-06^{+200\%}_{-50\%}$	$2.7e-06^{+200\%}_{-30\%}$	$0.000197^{+4\%}_{-9\%}$	
+-	380	$579^{+2\%}_{-1.19\%}$	$73^{+1\%}_{-1.35\%}$	—	$36.1^{+0.5\%}_{-1.18\%}$	$-968^{+2\%}_{-1.16\%}$	—	$-0.5^{+1\%}_{-0.9\%}$	$165^{+0.4\%}_{-1.18\%}$	$190^{+0.6\%}_{-1.18\%}$	$0.44^{+7\%}_{-9\%}$	$-0.324^{+10\%}_{-10\%}$	$0.185^{+10\%}_{-9\%}$	
+-	500	$469^{+0.4\%}_{-0.96\%}$	$287^{+0.5\%}_{-0.93\%}$	—	$31.7^{+0.4\%}_{-0.95\%}$	$-1270^{+2\%}_{-0.972\%}$	—	$-1.1^{+0.8\%}_{-0.982\%}$	$130^{+1\%}_{-0.918\%}$	$164^{+0.5\%}_{-0.94\%}$	$1.35^{+5\%}_{-5\%}$	$-0.442^{+20\%}_{-8\%}$	$0.554^{+10\%}_{-8\%}$	
+-	1000	$160^{+0.9\%}_{-0.902\%}$	$470^{+1.0\%}_{-0.92\%}$	—	$11.8^{+0.6\%}_{-0.931\%}$	$-1450^{+0.8\%}_{-0.920\%}$	—	$-0.5^{+0.8\%}_{-0.918\%}$	$44.9^{+0.8\%}_{-0.918\%}$	$52^{+1\%}_{-0.883\%}$	$-0.663^{+7\%}_{-7\%}$	$5.09^{+10\%}_{-8\%}$	$0.587^{+10\%}_{-8\%}$	
+-	1400	$84.9^{+2\%}_{-0.81\%}$	$507^{+3\%}_{-0.75\%}$	—	$7.57^{+0.6\%}_{-1.07\%}$	$-1230^{+0.6\%}_{-0.7\%}$	—	$-0.7^{+2\%}_{-0.835\%}$	$22.2^{+2\%}_{-0.908\%}$	$29.9^{+0.9\%}_{-0.908\%}$	$-1.22^{+8\%}_{-10\%}$	$-2.38^{+10\%}_{-10\%}$	$0.281^{+10\%}_{-10\%}$	
+-	3000	$23.8^{+1\%}_{-0.787\%}$	$1060^{+1\%}_{-0.414\%}$	—	$0.574^{+10\%}_{-0.338\%}$	$-1270^{+2\%}_{-0.88\%}$	—	$-1.08^{+2\%}_{-0.543\%}$	$6.28^{+3\%}_{-1.25\%}$	$1.93^{+6\%}_{-10\%}$	$8.36^{+3\%}_{-5\%}$	$0.197^{+7\%}_{-10\%}$		
+-	300	$0.218^{+0.4\%}_{-1.37\%}$	—	$-0.0855^{+2\%}_{-1.27\%}$	$0.0147^{+0.1\%}_{-1.42\%}$	$-0.302^{+3\%}_{-1.51\%}$	$0.00343^{+0.1\%}_{-1.59\%}$	$-0.0259^{+4\%}_{-1.5\%}$	$0.0799^{+4\%}_{-1.5\%}$	$3.38e-06^{+300\%}_{-5\%}$	$-7.78e-06^{+300\%}_{-0.07\%}$	$1.84e-05^{+4\%}_{-9\%}$		
+-	380	$249^{+0.4\%}_{-1.19\%}$	—	$-51.6^{+1\%}_{-1.18\%}$	$16.2^{+2\%}_{-1.18\%}$	$-649^{+2\%}_{-1.18\%}$	$3.31^{+1\%}_{-1.15\%}$	$-41.8^{+1\%}_{-1.15\%}$	$124^{+1\%}_{-1.19\%}$	$0.0946^{+2\%}_{-1\%}$	$-0.0633^{+2\%}_{-10\%}$	$0.000205^{+200\%}_{-0.0614\%}$		
+-	500	$203^{+0.5\%}_{-0.948\%}$	—	$-213^{+0.5\%}_{-0.958\%}$	$15.8^{+0.9\%}_{-0.975\%}$	$-810^{+0.8\%}_{-0.933\%}$	$0.767^{+4\%}_{-0.783\%}$	$-34.2^{+0.9\%}_{-0.923\%}$	$99.7^{+0.9\%}_{-0.909\%}$	$0.316^{+4\%}_{-4\%}$	$0.187^{+20\%}_{-20\%}$	$0.255^{+10\%}_{-9\%}$		
+-	1000	$63.4^{+0.9\%}_{-0.911\%}$	—	$-327^{+0.9\%}_{-0.820\%}$	$4.88^{+3\%}_{-0.820\%}$	$-810^{+3\%}_{-0.907\%}$	$-0.24^{+10\%}_{-0.907\%}$	$-10.4^{+6\%}_{-0.812\%}$	$34.1^{+9\%}_{-0.90\%}$	$-0.832^{+7\%}_{-10\%}$	$0.255^{+30\%}_{-1.28\%}$	$0.39^{+10\%}_{-8\%}$		
+-	1400	$33.8^{+0.8\%}_{-0.917\%}$	—	$-493^{+0.8\%}_{-0.93\%}$	$2.86^{+1\%}_{-0.897\%}$	$-850^{+2\%}_{-0.842\%}$	$-0.208^{+0.4\%}_{-1.06\%}$	$-4.75^{+2\%}_{-0.666\%}$	$10.2^{+0.9\%}_{-0.504\%}$	$0.448^{+3\%}_{-1.23\%}$	$0.475^{+20\%}_{-0.962\%}$	$0.258^{+3\%}_{-9\%}$		
+-	3000	—	—	$-424^{+6\%}_{-1.09\%}$	$0.226^{+4\%}_{-0.2\%}$	—	$0.146^{+0\%}_{-0\%}$	$-1.5^{+7\%}_{-0.651\%}$	$-497^{+0\%}_{-0\%}$	$2.52^{+4\%}_{-4\%}$	$2110^{+0\%}_{-0\%}$	$-0.0453^{+2\%}_{-0.04\%}$		

(MG, complex mass scheme, $m_b \rightarrow 0$, total width EFT dependence not included)

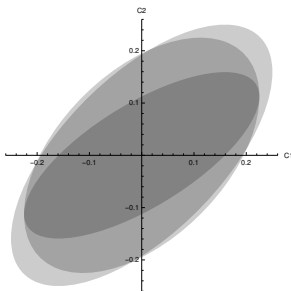
A couple of additional ideas

Statistically optimal observables

minimize the one-sigma ellipsoid in EFT parameter space.

(joint efficient set of estimator, saturating the Rao-Cramér-Fréchet bound: $V^{-1} = I$)

For small C_i , with a phase-space distribution $\sigma(\phi) = \sigma_0(\phi) + \sum_i C_i \sigma_i(\phi)$,
the statistically optimal set of observables is: $O_i(\phi) = \sigma_i(\phi)/\sigma_0(\phi)$.



e.g. $\sigma(\phi) = 1 + \cos(\phi) + C_1 \sin(\phi) + C_2 \sin(2\phi)$

1. asymmetries: $O_i \sim \text{sign}\{\sin(i\phi)\}$

2. moments: $O_i \sim \sin(i\phi)$

3. statistically optimal: $O_i \sim \frac{\sin(i\phi)}{1 + \cos\phi}$

\Rightarrow area ratios 1.9 : 1.7 : 1
(total rate not used)

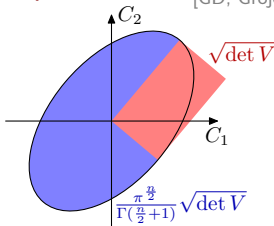
Global determinant parameter (GDP)

[GD, Grojean, Gu, Wang]

In a n -dimensional Gaussian fit

$$\text{GDP} \equiv \sqrt[2n]{\det V}$$

provides a geometric average of the constraints strengths.

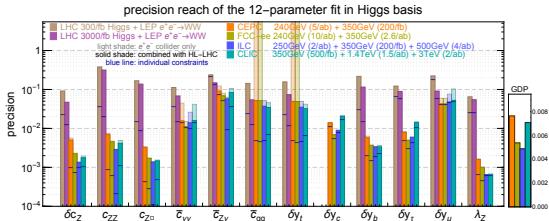


Interestingly, GDP ratios are operator-basis independent!

- as the volume scales linearly with coupling normalization
- as the volume is invariant under rotations

⇒ conveniently assess constraint strengthening.

e.g. Higgs fit at lepton colliders:



Summary

The EFT parametrizes systematically the parameter space in direct vicinity of the standard model.

A global analysis of future-lepton-collider constraints on the top EFT is ongoing.

Statistically optimal observables are surprisingly unexploited.

Global determinant parameter ratios assess the strengthening of global constraints, basis independently.