

BERNHARD MISTLBERGER



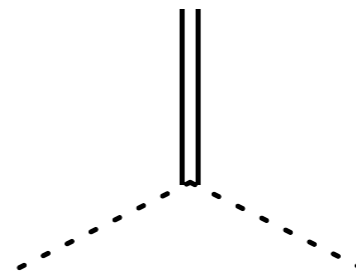
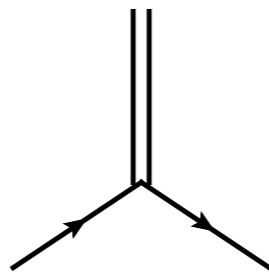
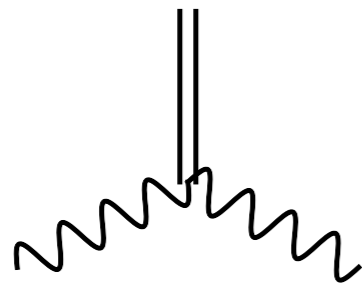
HIGGS-DIFFERENTIAL CROSS SECTIONS

STEPS TO DIFFERENTIAL N3LO

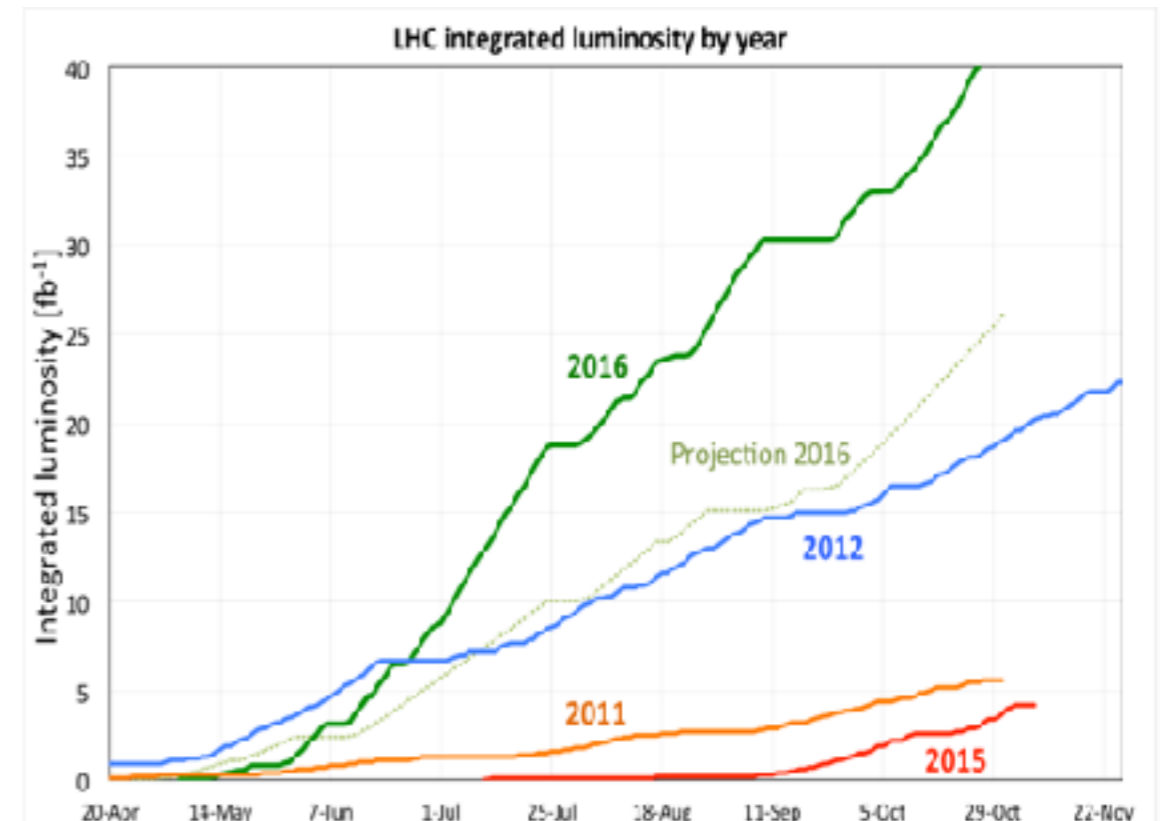
with Babis Anastasiou, Falko Dulat, Simone Lionetti, Andrea Pelloni, Caterina Specchia
[arXiv:1704.08220]

HIGGS BOSON

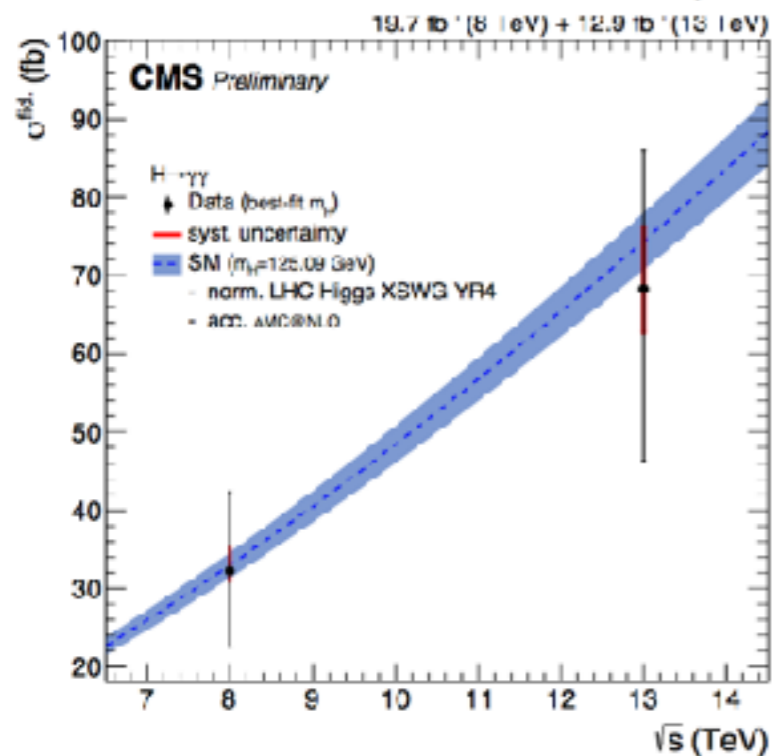
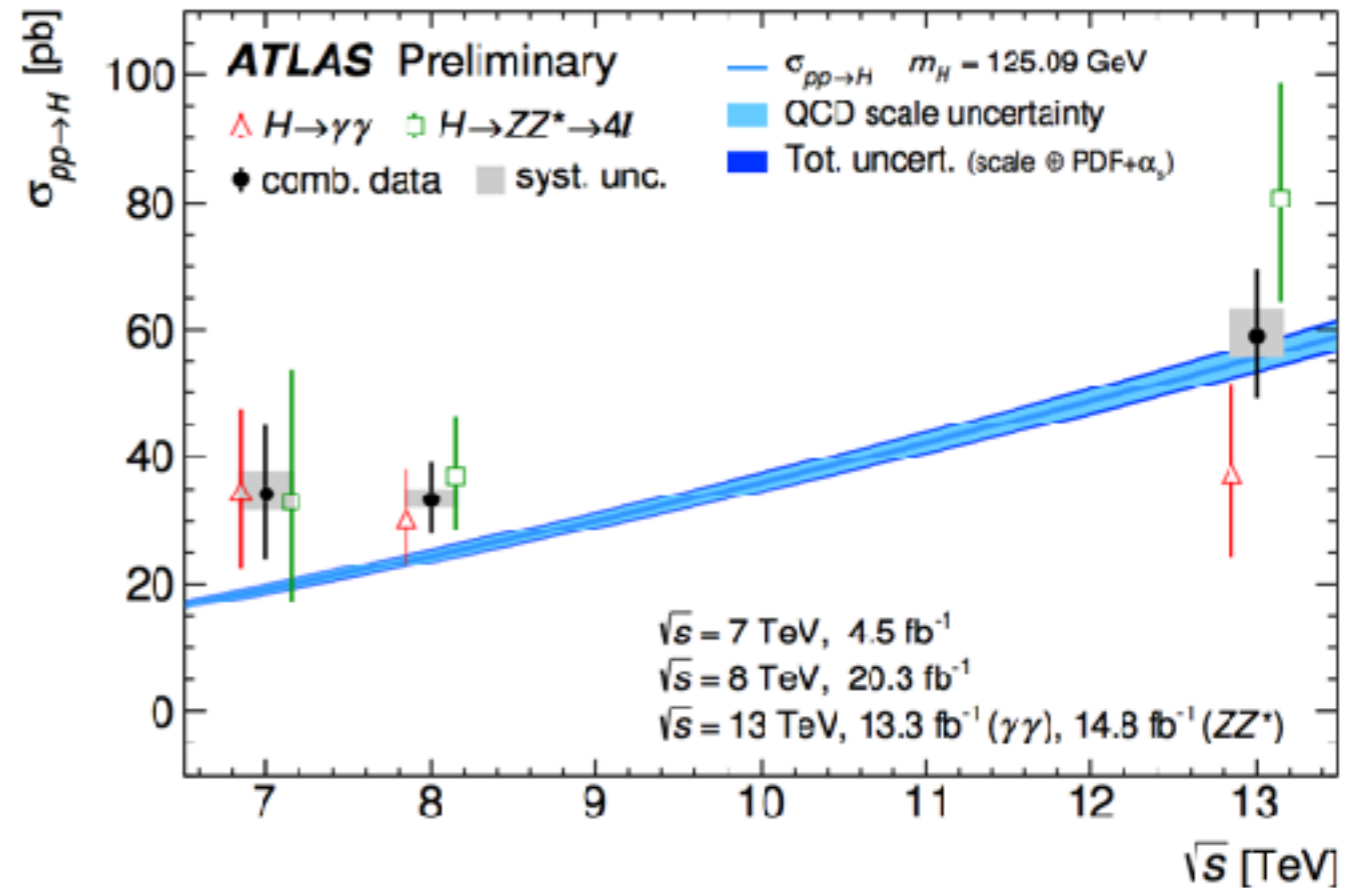
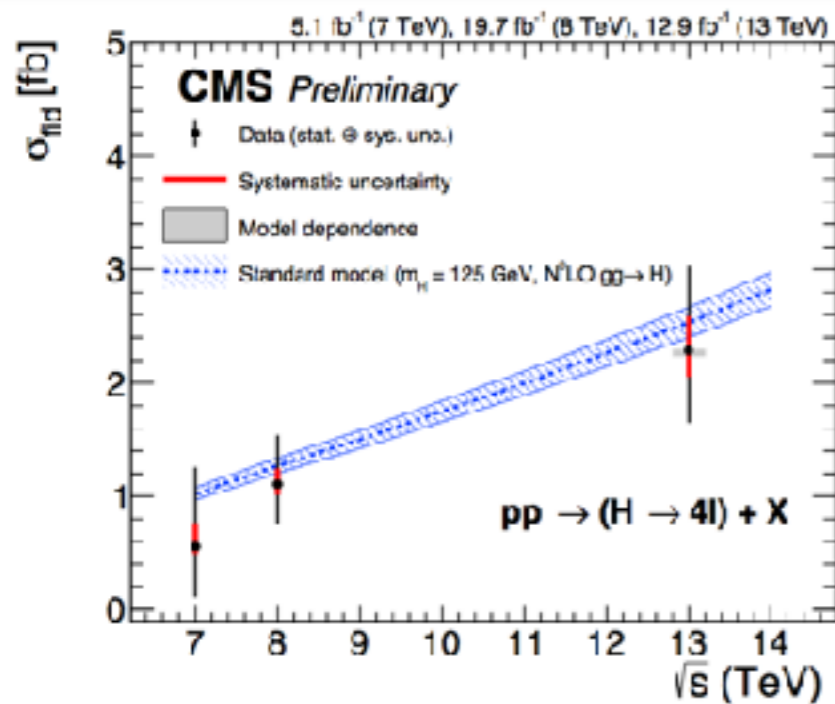
- ▶ **4th of July 2012:** The begin of the precision physics age of Higgs boson phenomenology
- ▶ Immediately after the discovery of the Higgs boson we started to ask questions about it's nature:
Couplings, spin, parity, mass, cross sections ...



- ▶ The basis for testing our understanding of nature is on the one side precise measurements that are sensitive to the Higgs boson properties.
- ▶ LHC provides the input!
Run 2: Data, data, data



ENTER THE AGE OF PRECISION HIGGS PHYSICS




- ▶ Incredible agreement of data and theory
- ▶ Triumph of SM predictions
- ▶ Higgs production
~10 sigma observed

DEMAND FOR PRECISION ON THEORY SIDE

- ▶ Testing our understanding of nature:
Compare experiment and theory!

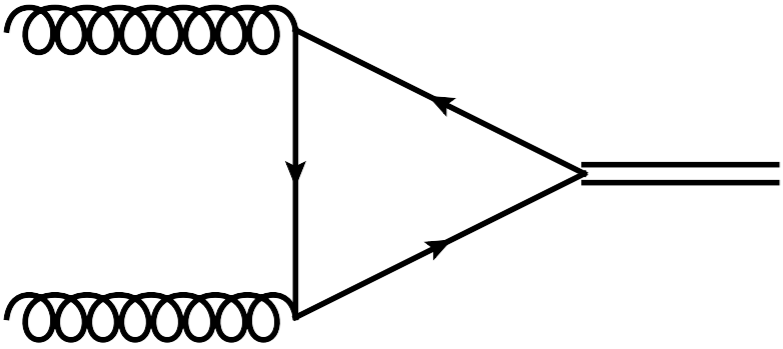
- ▶ The key to theoretical predictions at the LHC:

$$\sigma \sim \int dx dy f(x) f(y) \hat{\sigma} + \mathcal{O}\left(\frac{\Lambda}{Q}\right)$$

$$\hat{\sigma} = \hat{\sigma}^{(0)} + \alpha_S^1 \hat{\sigma}^{(1)} + \alpha_S^2 \hat{\sigma}^{(2)} + \alpha_S^3 \hat{\sigma}^{(3)} \dots$$


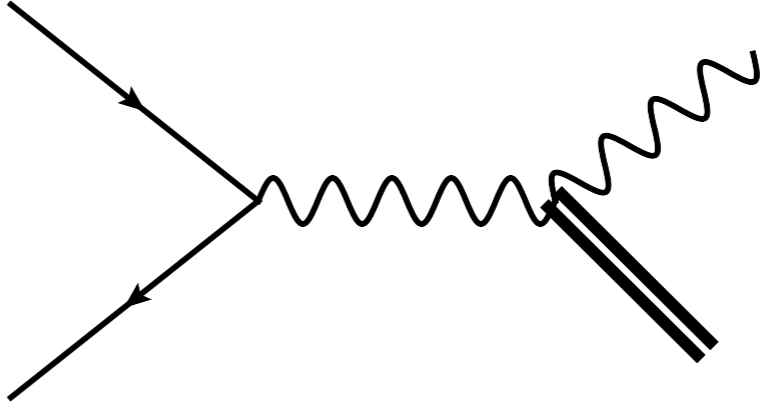
- ▶ Compute perturbative corrections from first principle
QFT: Standard Model
- ▶ Allows for % - level predictions for - experimental
precision will reach comparable levels!

4 WAYS TO PRODUCE A HIGGS



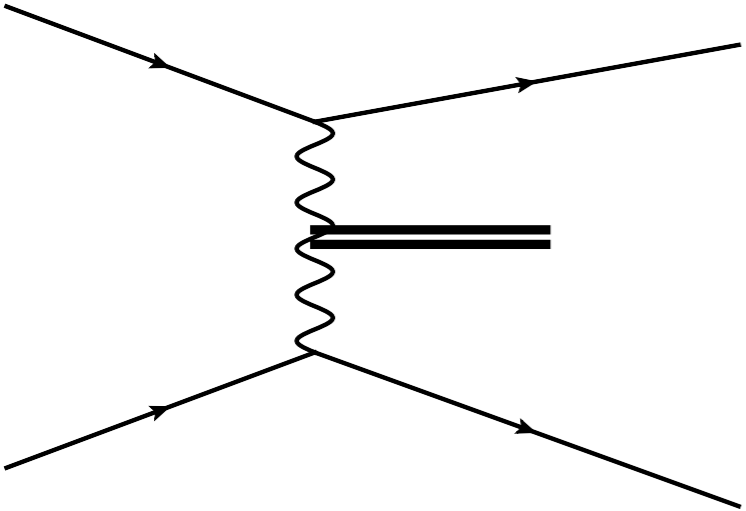
ggF

~88.2%



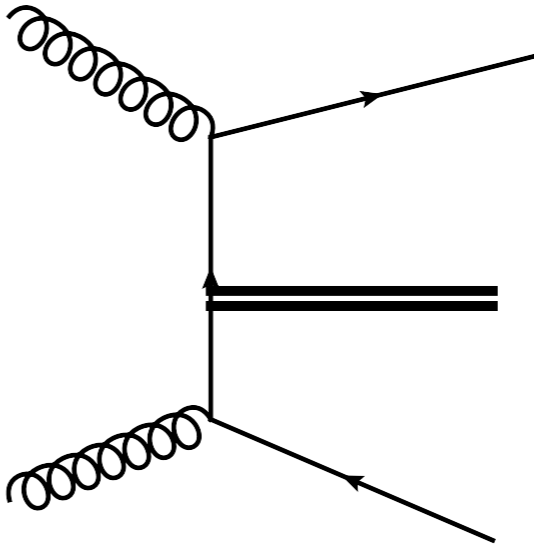
Associated

~4.1%



VBF

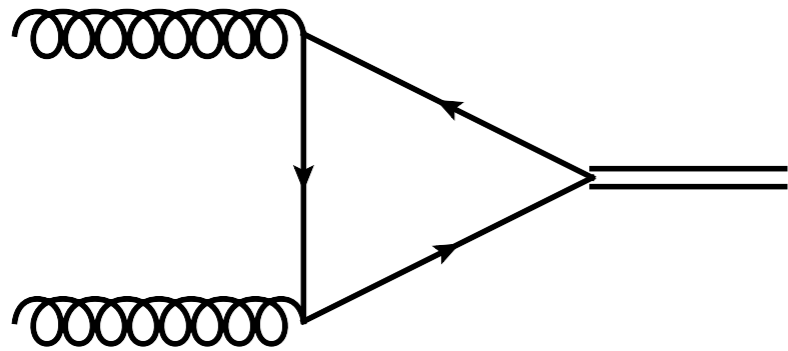
~6.8%



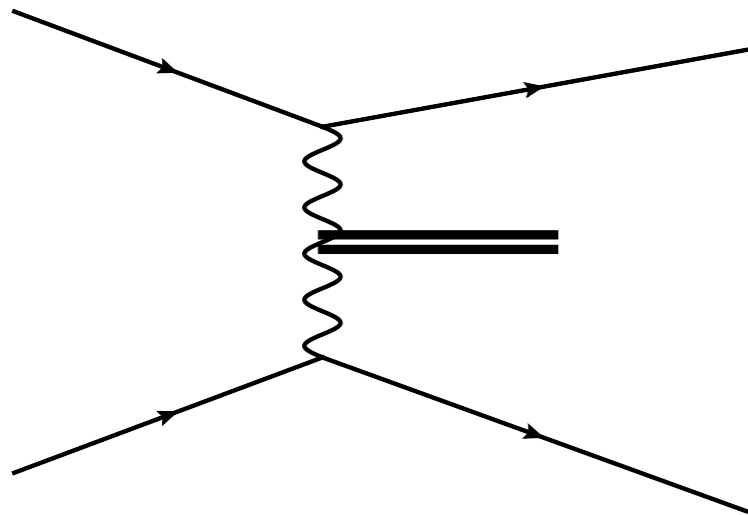
ttH

~0.9%

DOMINANT QCD CORRECTIONS

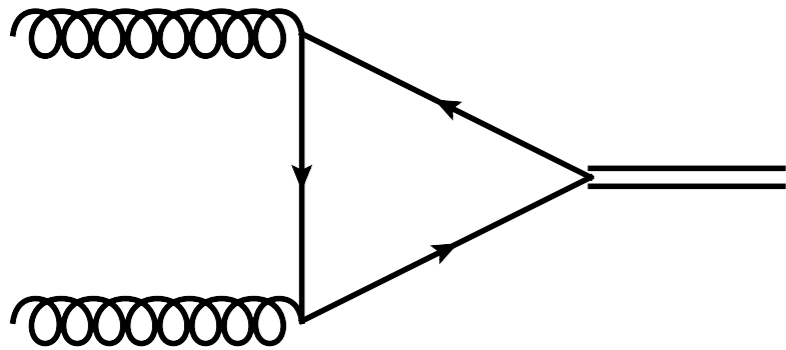


- ▶ Inclusive cross section: N3LO
- ▶ Differential cross section: NNLO
- ▶ H+J: NNLO

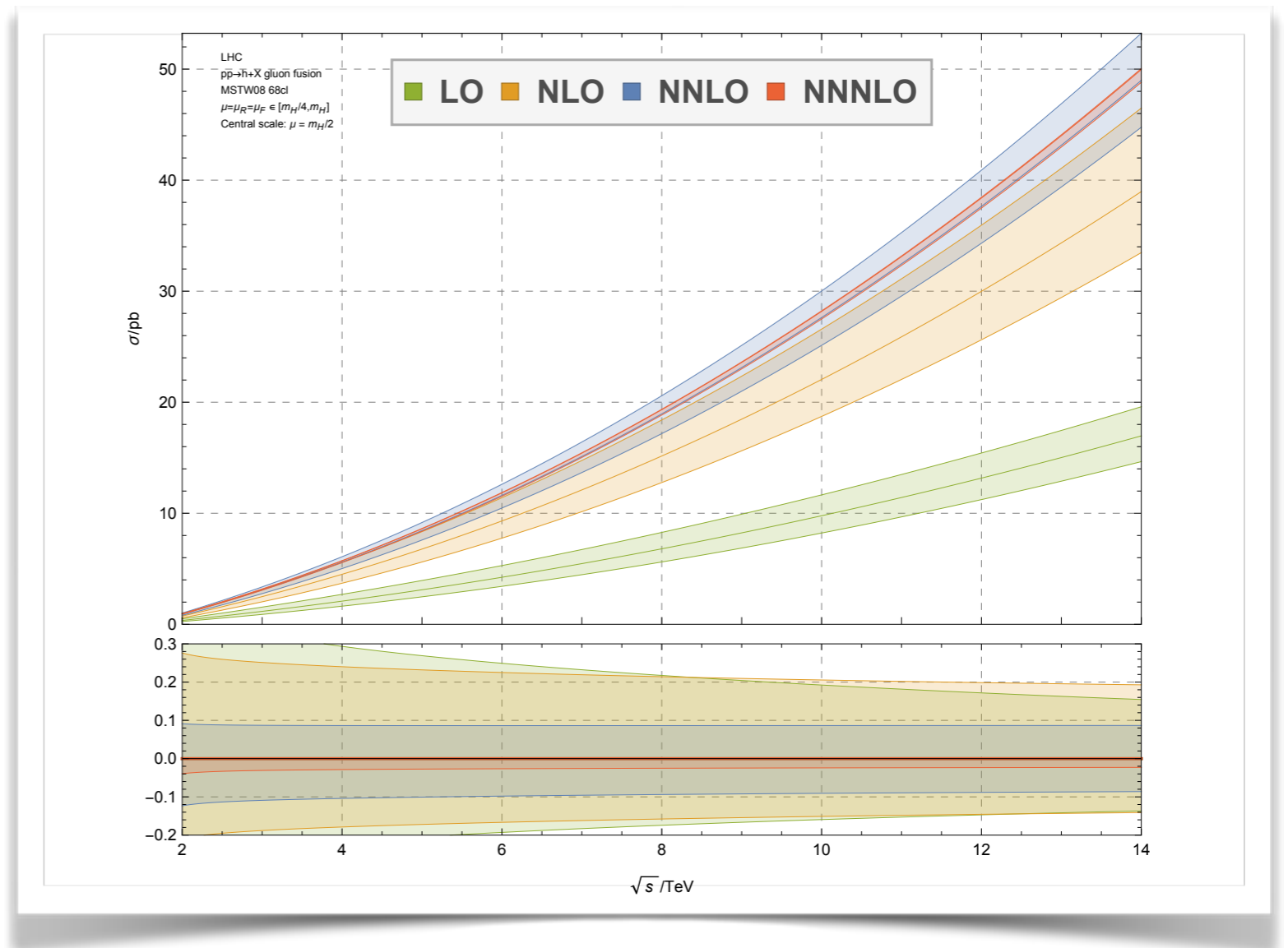


- ▶ Inclusive cross section: N3LO
- ▶ Differential cross section: NNLO

MOTIVATION FROM INCLUSIVE CROSS SECTION



- ▶ N3LO corrections stabilise perturbative expansion.
- ▶ We would like to see a similar pattern for differential cross sections.



CHALLENGES OF PERTURBATIVE PREDICTIONS

- ▶ Analytic complexity of high order perturbative computation

- ▶ Complicated mathematical structures: Elliptic / multiple polylogarithms, couple differential equations, algebraic complexity, ...

- ▶ Numerical integration over complicated and "divergent" final state configurations:

- ▶ Infrared subtraction at 2-loops and beyond.
- ▶ Main challenge of the last couple of years.
- ▶ Many methods available now.

- **Sector decomposition**

- **Non-Linear Mappings**

- **qT**

- **FKS+**

- **N-Jettiness**

H+J

- **Antenna**

- **Colourful**

- **Projection-To-Born**

VBF

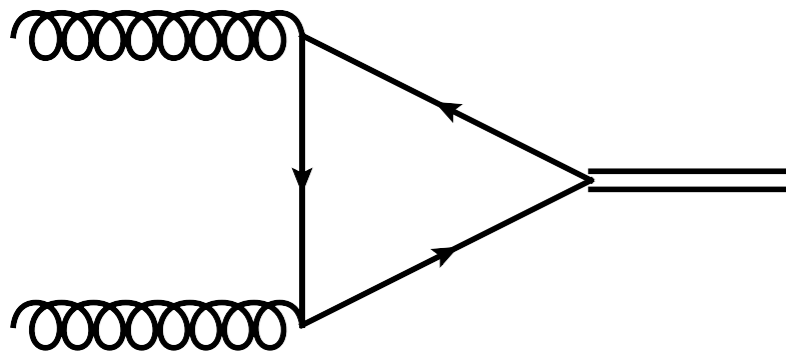
- ...

HIGGS – DIFFERENTIAL CROSS SECTIONS



- ▶ Introduce a framework that allows to compute differential cross sections at N3LO.
- ▶ Circumvent problems of NNLO infrared subtraction.
- ▶ Applicable for real life observables at the LHC.

Specifically: Differential Higgs Production in QCD



$$P P \rightarrow H + X \rightarrow \gamma\gamma + X$$

$$P P \rightarrow H + X \rightarrow 4l + X$$

- ▶ Today: Recent Progress, NNLO, Obstacles, Method

HIGGS – DIFFERENTIAL CROSS SECTIONS

- ▶ **Focus on the degrees of freedom of the Higgs boson:**

$$p_h \equiv (E, p_x, p_y, p_z) = \left(\sqrt{p_T^2 + m_h^2} \cosh Y, p_T \cos \phi, p_T \sin \phi, \sqrt{p_T^2 + m_h^2} \sinh Y \right)$$

- ▶ Entirely described in terms of p_T and Y (and a trivial azimuthal angle).

$$Y = \frac{1}{2} \log \left(\frac{E + p_z}{E - p_z} \right)$$

$$p_T = \sqrt{E^2 - p_z^2 - m_h^2}.$$

HIGGS – DIFFERENTIAL CROSS SECTIONS

- ▶ **Focus on the degrees of freedom of the Higgs boson:**

$$p_h \equiv (E, p_x, p_y, p_z) = \left(\sqrt{p_T^2 + m_h^2} \cosh Y, p_T \cos \phi, p_T \sin \phi, \sqrt{p_T^2 + m_h^2} \sinh Y \right)$$

- ▶ Entirely described in terms of p_T and Y (and a trivial azimuthal angle).

$$Y = \frac{1}{2} \log \left(\frac{E + p_z}{E - p_z} \right)$$

$$p_T = \sqrt{E^2 - p_z^2 - m_h^2}.$$

- ▶ Higgs-differential cross section:

$$\sigma_{PP \rightarrow H+X} [\mathcal{O}] = \sum_{i,j} \int_{-\infty}^{+\infty} dY \int_0^{\infty} dp_T^2 \int_0^{2\pi} \frac{d\phi}{2\pi} \int_0^1 dx_1 dx_2 f_i(x_1) f_j(x_2)$$

$$\times \frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} (S, x_1, x_2, m_h^2, Y, p_T^2) \mathcal{J}_{\mathcal{O}}(Y, p_T^2, \phi, m_h^2)$$

HIGGS – DIFFERENTIAL CROSS SECTIONS

- ▶ Focus on the degrees of freedom of the Higgs boson:

$$p_h \equiv (E, p_x, p_y, p_z) = \left(\sqrt{p_T^2 + m_h^2} \cosh Y, p_T \cos \phi, p_T \sin \phi, \sqrt{p_T^2 + m_h^2} \sinh Y \right)$$

- ▶ Entirely described in terms of pT and Y (and a trivial azimuthal angle).

$$Y = \frac{1}{2} \log \left(\frac{E + p_z}{E - p_z} \right)$$

$$p_T = \sqrt{E^2 - p_z^2 - m_h^2}$$

- ▶ Higgs-differential cross section:

$$\sigma_{PP \rightarrow H+X} [\mathcal{O}] = \sum_{i,j} \int_{-\infty}^{+\infty} dY \int_0^{\infty} dp_T^2 \int_0^{2\pi} \frac{d\phi}{2\pi} \int_0^1 dx_1 dx_2 f_i(x_1) f_j(x_2)$$

XS for observable \mathcal{O}

$$\times \frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} (S, x_1, x_2, m_h^2, Y, p_T^2) \mathcal{J}_{\mathcal{O}}(Y, p_T^2, \phi, m_h^2)$$

HIGGS - DIFFERENTIAL CROSS SECTIONS

- ▶ Focus on the degrees of freedom of the Higgs boson:

$$p_h \equiv (E, p_x, p_y, p_z) = \left(\sqrt{p_T^2 + m_h^2} \cosh Y, p_T \cos \phi, p_T \sin \phi, \sqrt{p_T^2 + m_h^2} \sinh Y \right)$$

- ▶ Entirely described in terms of pT and Y (and a trivial azimuthal angle).

$$Y = \frac{1}{2} \log \left(\frac{E + p_z}{E - p_z} \right)$$

$$p_T = \sqrt{E^2 - p_z^2 - m_h^2}$$

- ▶ Higgs-differential cross section:

$$\sigma_{PP \rightarrow H+X} [\mathcal{O}] = \sum_{i,j} \int_{-\infty}^{+\infty} dY \int_0^{\infty} dp_T^2 \int_0^{2\pi} \frac{d\phi}{2\pi} \int_0^1 dx_1 dx_2 f_i(x_1) f_j(x_2)$$

**Integrate over
Higgs degrees of freedom**

$$\times \frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} (S, x_1, x_2, m_h^2, Y, p_T^2) \mathcal{J}_{\mathcal{O}}(Y, p_T^2, \phi, m_h^2)$$

HIGGS – DIFFERENTIAL CROSS SECTIONS

- ▶ Focus on the degrees of freedom of the Higgs boson:

$$p_h \equiv (E, p_x, p_y, p_z) = \left(\sqrt{p_T^2 + m_h^2} \cosh Y, p_T \cos \phi, p_T \sin \phi, \sqrt{p_T^2 + m_h^2} \sinh Y \right)$$

- ▶ Entirely described in terms of pT and Y (and a trivial azimuthal angle).

$$Y = \frac{1}{2} \log \left(\frac{E + p_z}{E - p_z} \right)$$

$$p_T = \sqrt{E^2 - p_z^2 - m_h^2}$$

- ▶ Higgs-differential cross section:

$$\sigma_{PP \rightarrow H+X} [\mathcal{O}] = \sum_{i,j} \int_{-\infty}^{+\infty} dY \int_0^\infty dp_T^2 \int_0^{2\pi} \frac{d\phi}{2\pi} \int_0^1 dx_1 dx_2 f_i(x_1) f_j(x_2)$$

Partonic Higgs-differential cross section

$$\times \frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} (S, x_1, x_2, m_h^2, Y, p_T^2) \mathcal{J}_{\mathcal{O}}(Y, p_T^2, \phi, m_h^2)$$

HIGGS - DIFFERENTIAL CROSS SECTIONS

- ▶ Focus on the degrees of freedom of the Higgs boson:

$$p_h \equiv (E, p_x, p_y, p_z) = \left(\sqrt{p_T^2 + m_h^2} \cosh Y, p_T \cos \phi, p_T \sin \phi, \sqrt{p_T^2 + m_h^2} \sinh Y \right)$$

- ▶ Entirely described in terms of pT and Y (and a trivial azimuthal angle).

$$Y = \frac{1}{2} \log \left(\frac{E + p_z}{E - p_z} \right)$$

$$p_T = \sqrt{E^2 - p_z^2 - m_h^2}$$

- ▶ Higgs-differential cross section:

$$\sigma_{PP \rightarrow H+X} [\mathcal{O}] = \sum_{i,j} \int_{-\infty}^{+\infty} dY \int_0^\infty dp_T^2 \int_0^{2\pi} \frac{d\phi}{2\pi} \int_0^1 dx_1 dx_2 f_i(x_1) f_j(x_2)$$

Definition of observable - measurement function $\times \frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} (S, x_1, x_2, m_h^2, Y, p_T^2) \mathcal{J}_{\mathcal{O}}(Y, p_T^2, \phi, m_h^2)$

For example: $\mathcal{J}_{(Y, p_T^2, \phi, m_h^2)} = \theta(p_T^2 > 20 \text{ GeV})$

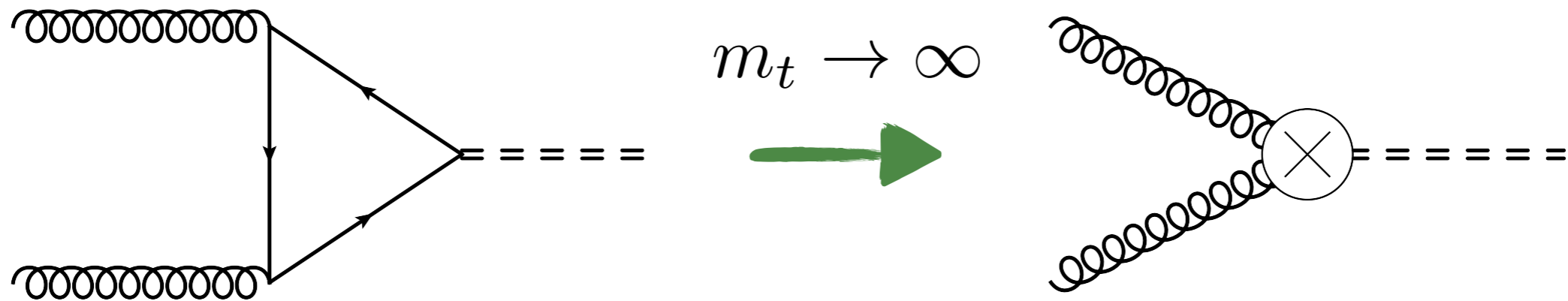
PARTONIC HIGGS – DIFFERENTIAL CROSS SECTIONS

$$\frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} \sim \sum_X \int d\Phi_n \left| \mathcal{M}_{ij \rightarrow H+X} \right|^2$$

PARTONIC HIGGS - DIFFERENTIAL CROSS SECTIONS

$$\frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} \sim \sum_X \int d\Phi_n \left| \mathcal{M}_{ij \rightarrow H+X} \right|^2$$

- ▶ Compute all required matrix elements of different final states X to a given order in perturbation theory.
- ▶ Work in effective theory: Excellent approximation!



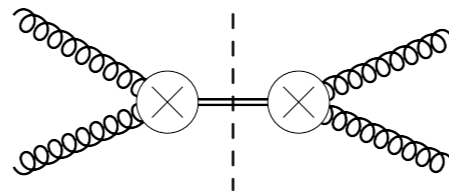
$$\mathcal{L} = \mathcal{L}_{QCD,5} - \frac{1}{4v} C_1 H G_{\mu\nu}^a G_a^{\mu\nu}$$

PARTONIC HIGGS - DIFFERENTIAL CROSS SECTIONS

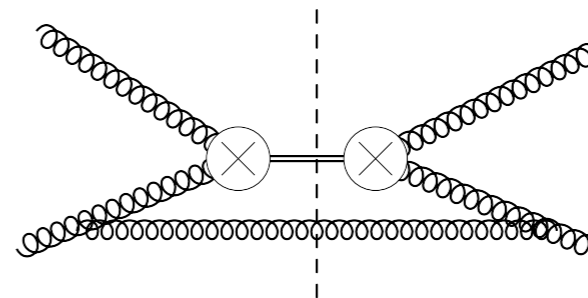
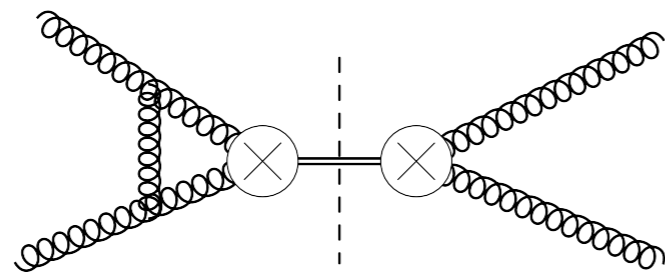
$$\frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} \sim \sum_X \int d\Phi_n \left| \mathcal{M}_{ij \rightarrow H+X} \right|^2$$

- ▶ Compute all required matrix elements of different final states X to a given order in perturbation theory.

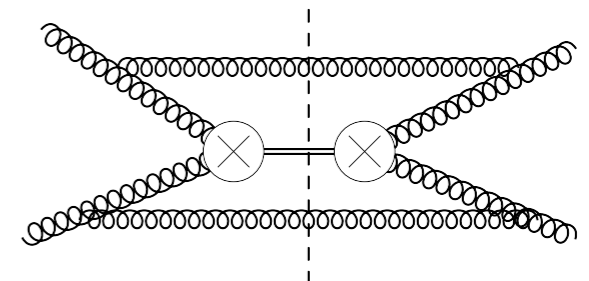
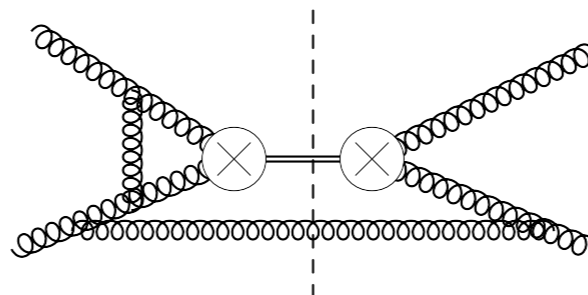
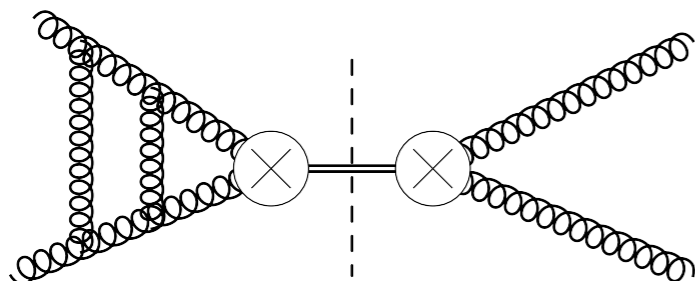
LO:



NLO:

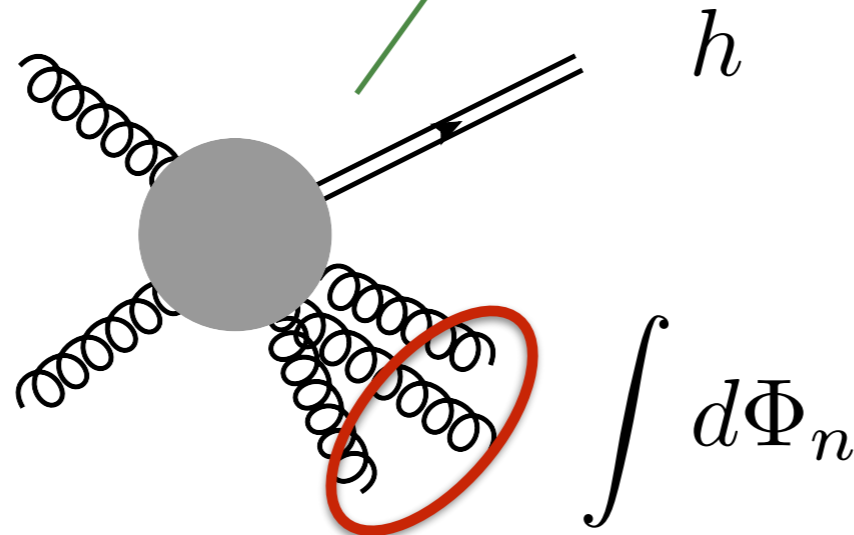


NNLO:



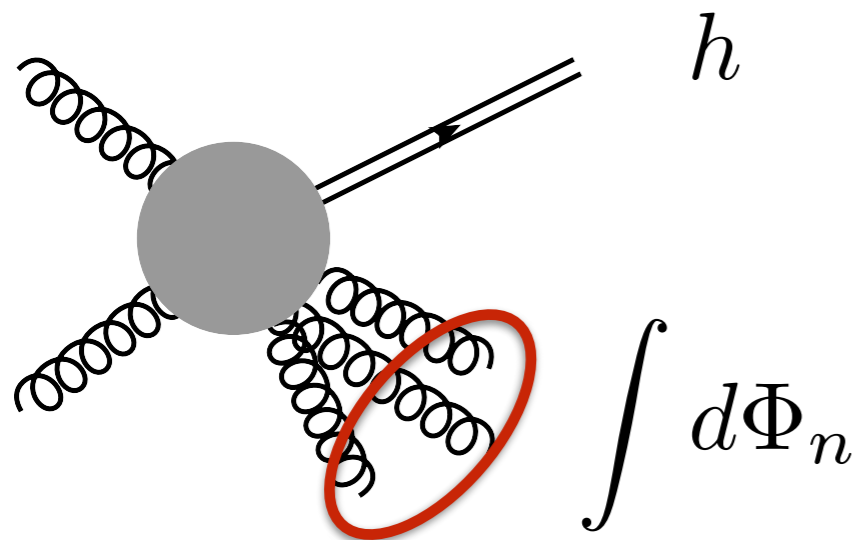
PARTONIC HIGGS – DIFFERENTIAL CROSS SECTIONS

$$\frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} \sim \sum_X \int d\Phi_n \left| \mathcal{M}_{ij \rightarrow H+X} \right|^2$$



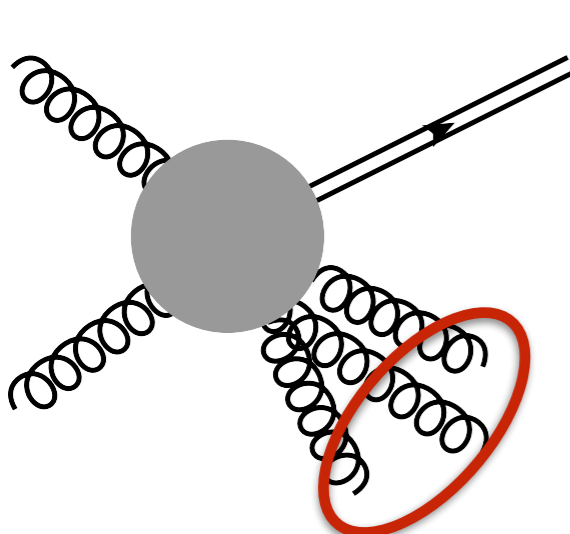
- ▶ Phase space integral over partonic final state phase space momenta
- ▶ Integrate over as many partons as there are in X .
- ▶ Integration over fixed multiplicity matrix elements is divergent! (KLN).

PARTONIC HIGGS – DIFFERENTIAL CROSS SECTIONS



- ▶ Perform integration over parton phase space analytically
- ▶ Everything else effectively covered by H+J
- ▶ Rely on tools to perform analytic computation learned from inclusive N3LO
- ▶ Make singularities of final state parton integrations manifest using dimensional regularisation. $d = 4 - 2\epsilon$

PARTONIC HIGGS – DIFFERENTIAL CROSS SECTIONS



$$\int d\Phi_n \sim \int \prod_{i=1}^n d^d p_i \delta_+(p_i^2)$$

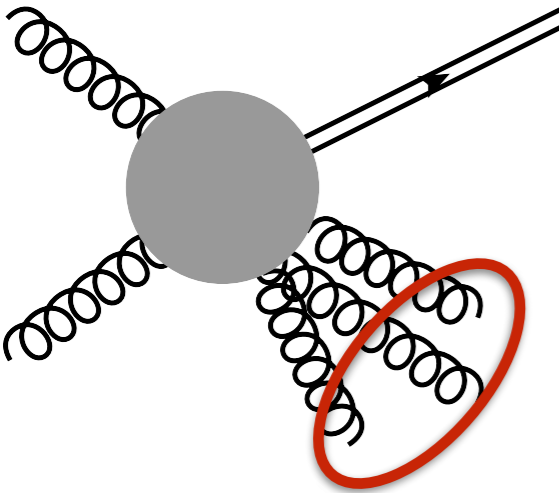
with $\delta_+(p_i^2) = \theta(E_i) \delta(p_i^2)$

REVERSE UNITARITY FRAMEWORK:

- ▶ Replace on-shell constraints with cut propagators

$$\delta_+(p_i^2) \sim \lim_{\delta \rightarrow 0} \left[\frac{1}{p_i^2 + i\delta} - \frac{1}{p_i^2 - i\delta} \right] = \left[\frac{1}{p_i^2} \right]_c$$

REVERSE UNITARITY FRAMEWORK:



A Feynman diagram showing a central grey circle vertex. Two external lines enter from the top: a double line with an arrow pointing right, labeled h , and a wavy line. Two wavy lines exit from the bottom. One of the bottom wavy lines is circled in red, indicating a cut propagator.

$$\int d\Phi_n \sim \int \prod_{i=1}^n d^d p_i \left[\frac{1}{p_i^2} \right]_c$$

- ▶ Opens the door to large variety of loop integral technology!
 - ▶ **IBPs + Differential equations**
- ▶ Key observation: Cut propagators can be differentiated similar to usual propagators.

REVERSE UNITARITY FRAMEWORK:

$$\begin{aligned} \frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} &\sim \sum_X \int d\Phi_n \left| \mathcal{M}_{ij \rightarrow H+X} \right|^2 \\ &= \sum_X \sum_i c_{X,i} F_{X,i}(S, p_T, Y, m_h^2) \end{aligned}$$

- ▶ Coefficient: Rational function of remaining kinematic variables.
- ▶ Master Integral: Integrated Feynman integrals: Polylogarithms, rational functions of remaining kinematic variables.
- ▶ Explicit Laurent series in dimensional regulator.

HIGGS - DIFFERENTIAL CROSS SECTIONS

$$\sigma_{PP \rightarrow H+X}[\mathcal{O}] = \sum_{i,j} \int_{-\infty}^{+\infty} dY \int_0^{\infty} dp_T^2 \int_0^{2\pi} \frac{d\phi}{2\pi} \int_0^1 dx_1 dx_2 f_i(x_1) f_j(x_2) \\ \times \frac{d^2 \hat{\sigma}_{ij}}{dY dp_T^2} (S, x_1, x_2, m_h^2, Y, p_T^2) \mathcal{J}_{\mathcal{O}}(Y, p_T^2, \phi, m_h^2)$$

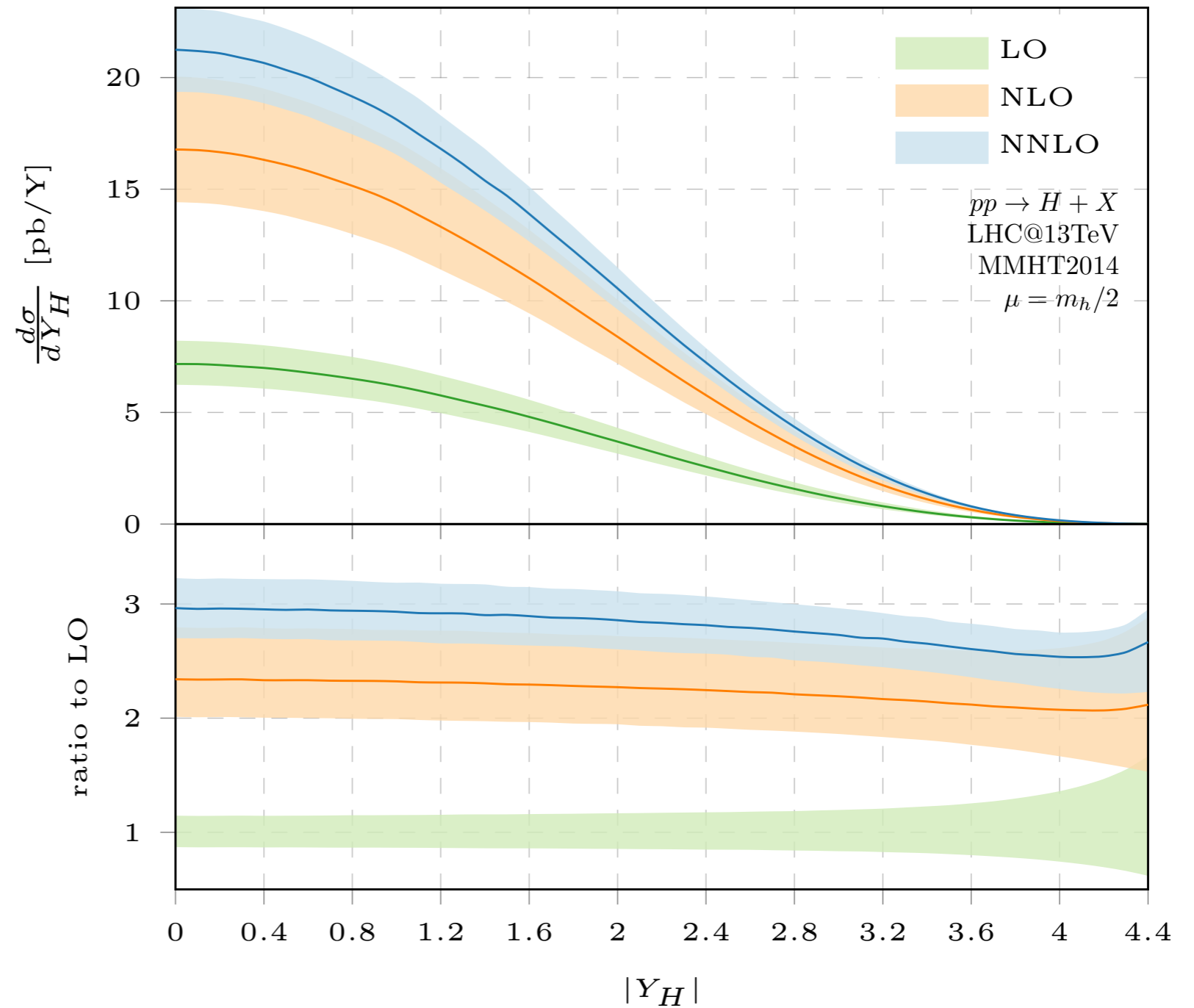
- ▶ UV - renormalisation ($\alpha_S \rightarrow \alpha_S^R Z$, etc.)
- ▶ Initial state collinear singularities: Redefine parton distributions functions:

$$f(x) \rightarrow f_R(x) \circ \Gamma$$

Currently: **Complete NNLO Higgs-differential cross section.**

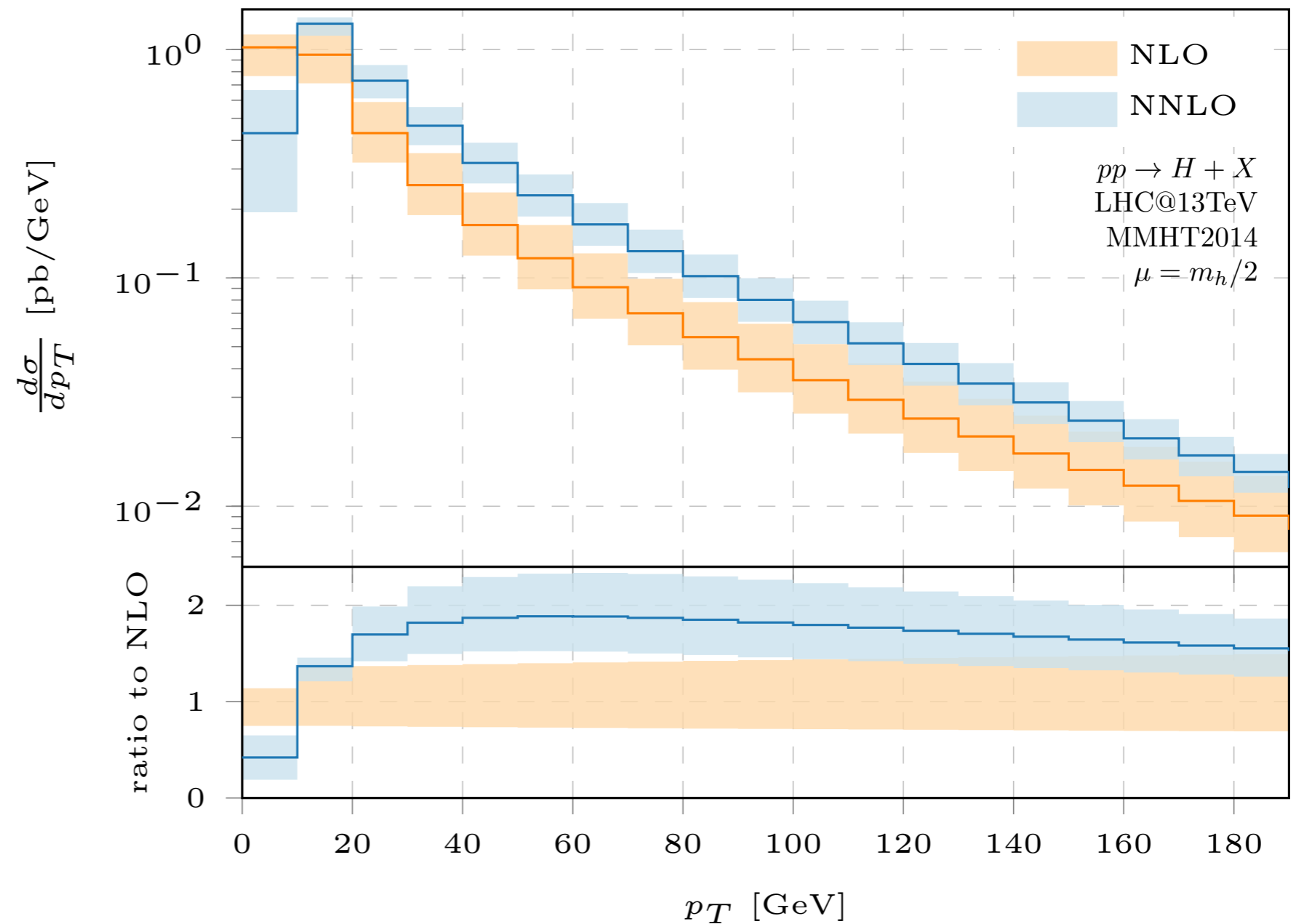
HIGGS - DIFFERENTIAL CROSS SECTIONS: RAPIDITY

- ▶ Inclusive rapidity distribution
- ▶ Large K-factors



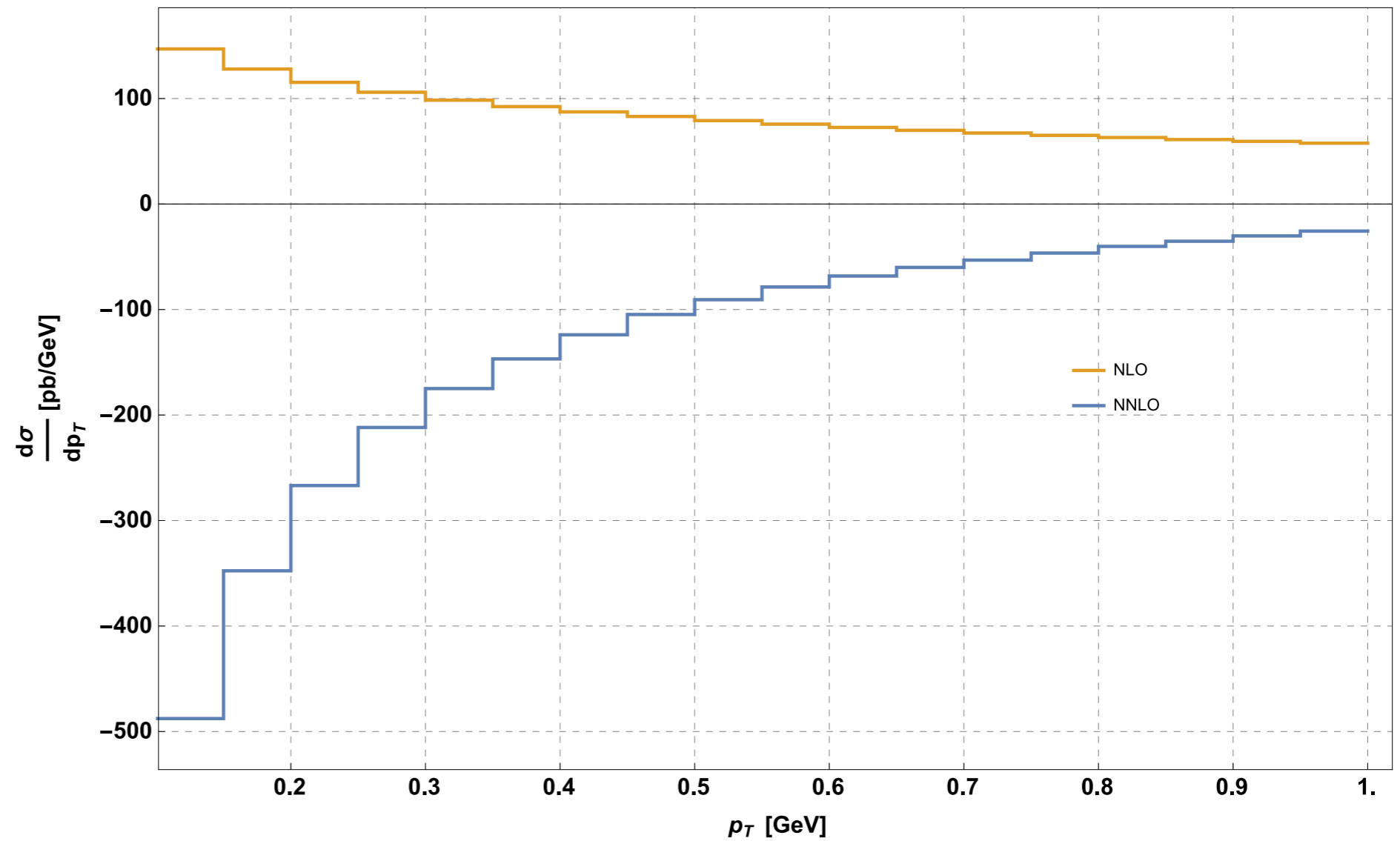
HIGGS - DIFFERENTIAL CROSS SECTIONS: p_T

► Inclusive p_T distribution



HIGGS - DIFFERENTIAL CROSS SECTIONS: PT

- ▶ Inclusive PT distribution - extreme regime



HIGGS – DIFFERENTIAL CROSS SECTIONS: FIDUCIAL XS

- ▶ Decay of the Higgs boson to two photons (Narrow width approximation)
- ▶ Employ realistic selection cuts (ATLAS)

$$\eta_{\gamma} < 2.37$$

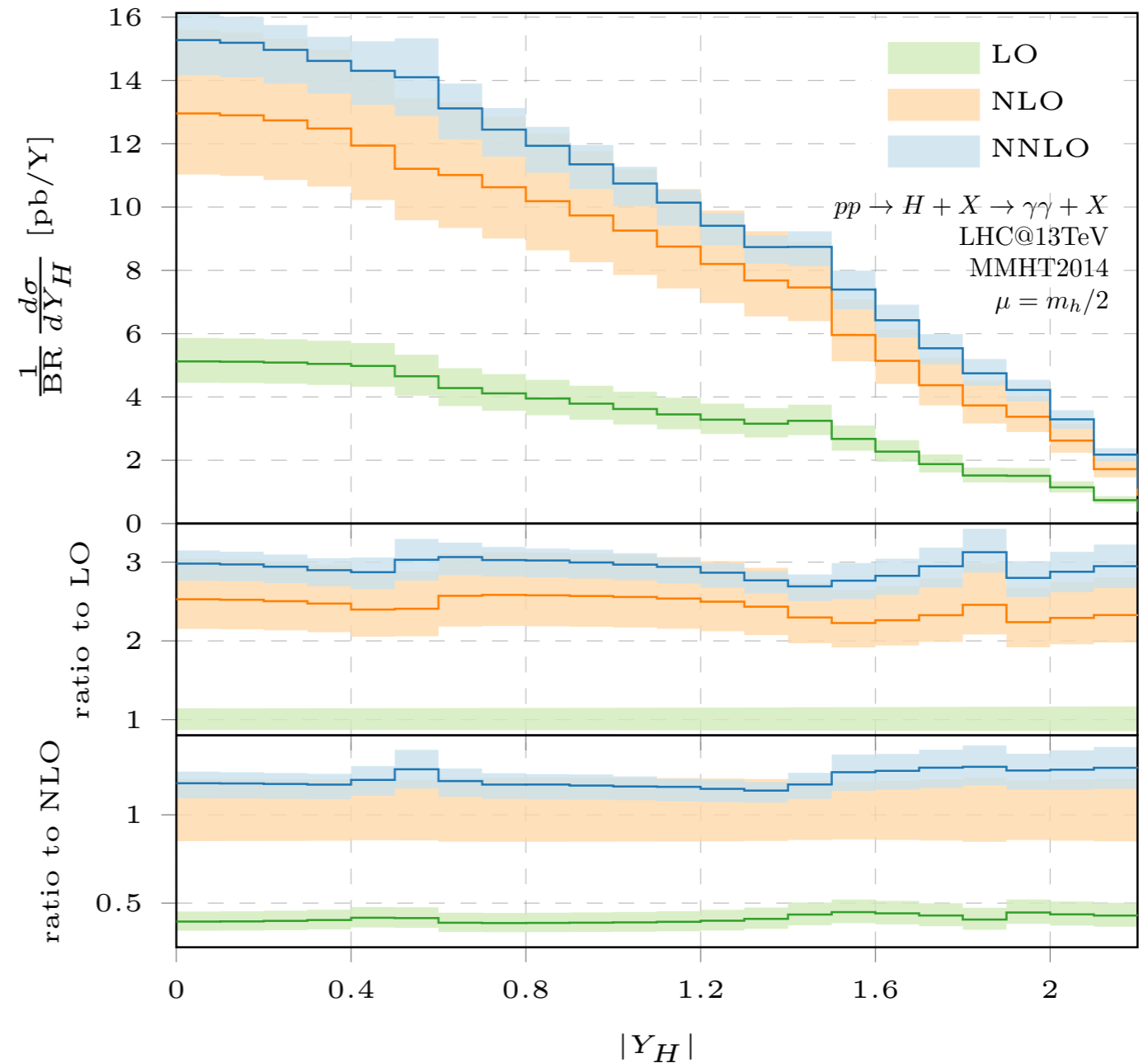
$$\eta_{\gamma} \notin [1.37, 1.52]$$

$$p_{T, \gamma_1} > 0.35 m_h$$

$$p_{T, \gamma_2} > 0.25 m_h$$

HIGGS - DIFFERENTIAL CROSS SECTIONS: FIDUCIAL XS

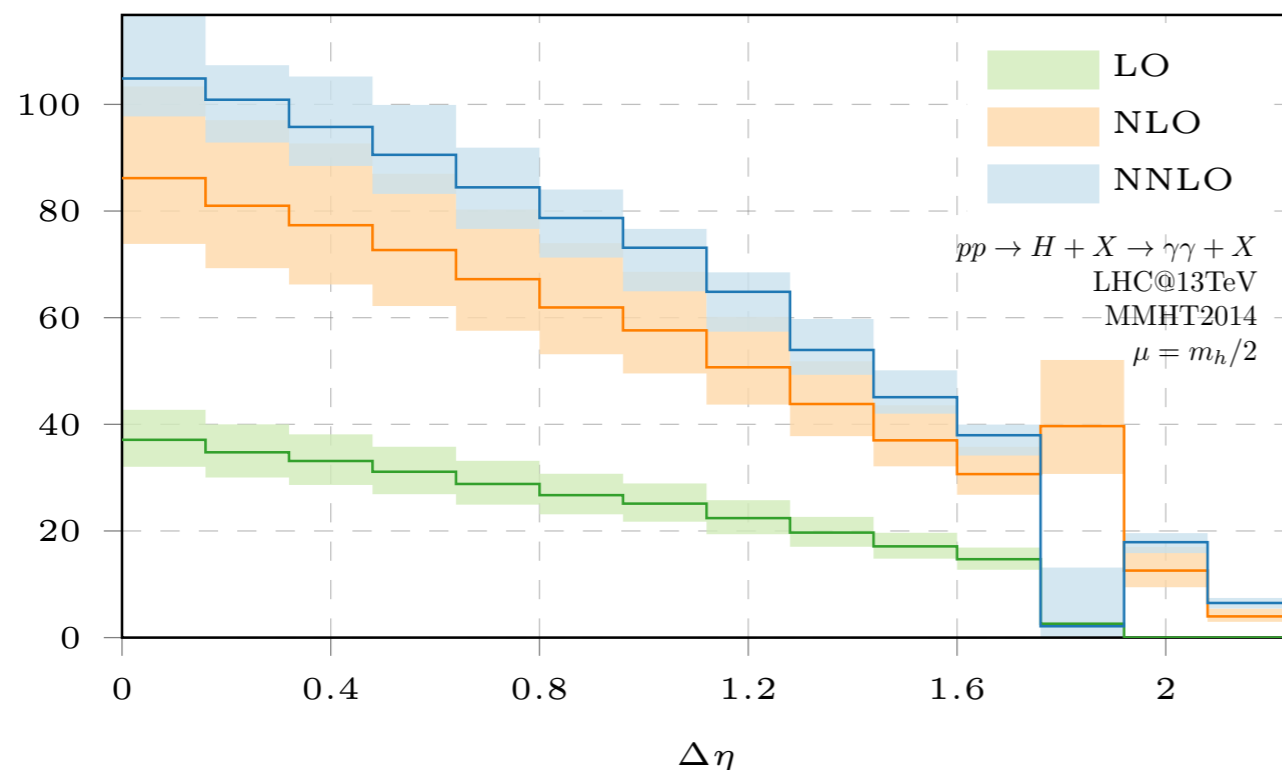
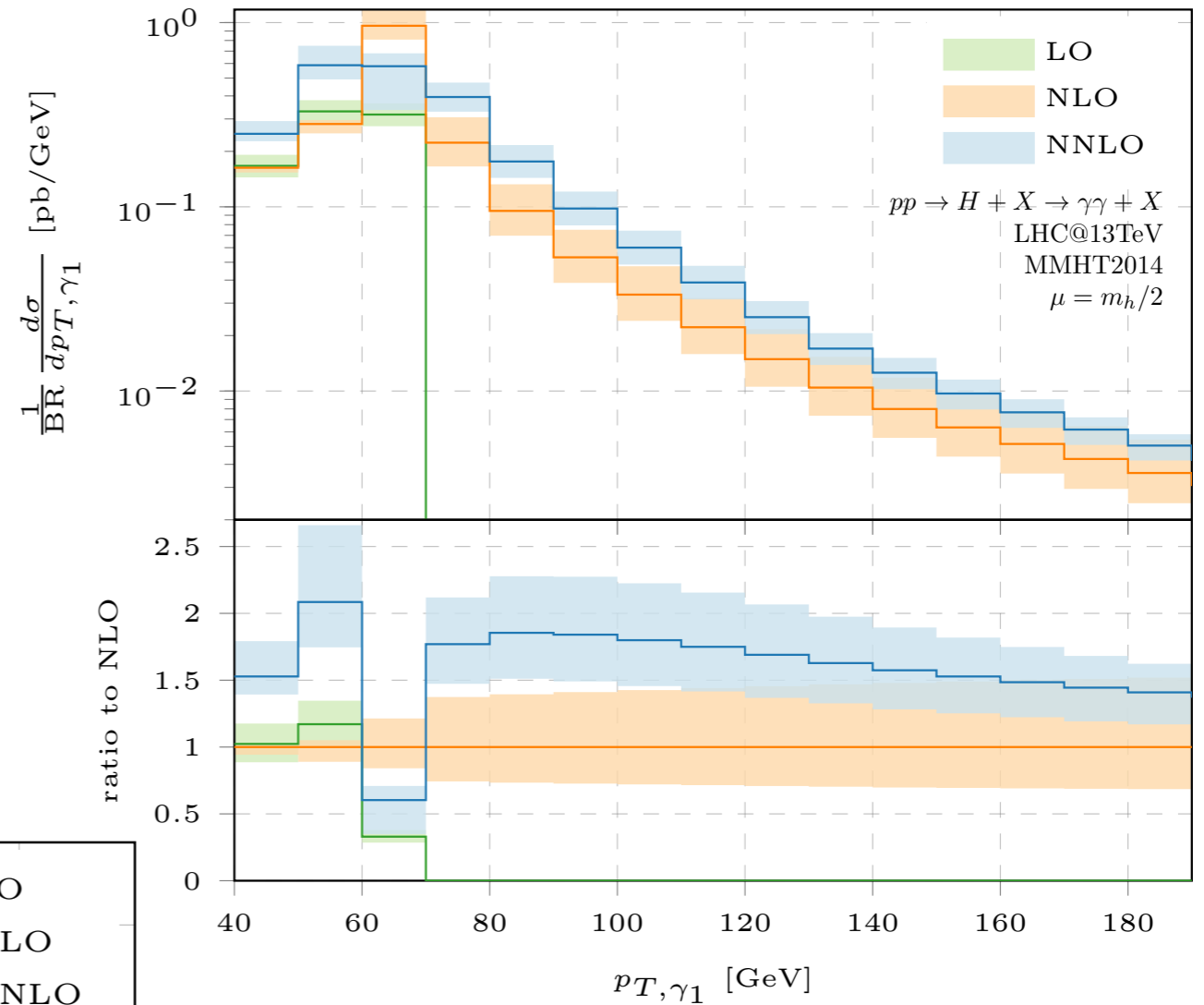
- ▶ Fiducial rapidity distribution.
- ▶ Non-trivial features due to selection criteria.
- ▶ Relatively flat K-factors
- ▶ Similar perturbative behaviour as inclusive distribution



HIGGS - DIFFERENTIAL CROSS SECTIONS: FIDUCIAL XS

▶ Distributions of the photon momenta:

- ▶ Leading Photon p_T
- ▶ Pseudo - rapidity difference



$$\Delta\eta = |\eta_{\gamma_1} - \eta_{\gamma_2}|$$

BEYOND NNNLO



WHAT DID WE LEARN FROM NNLO

- ▶ Higgs-differential cross sections: fast and stable framework for fiducial cross sections.
- ▶ Analytic computation at NNLO comparably simple.

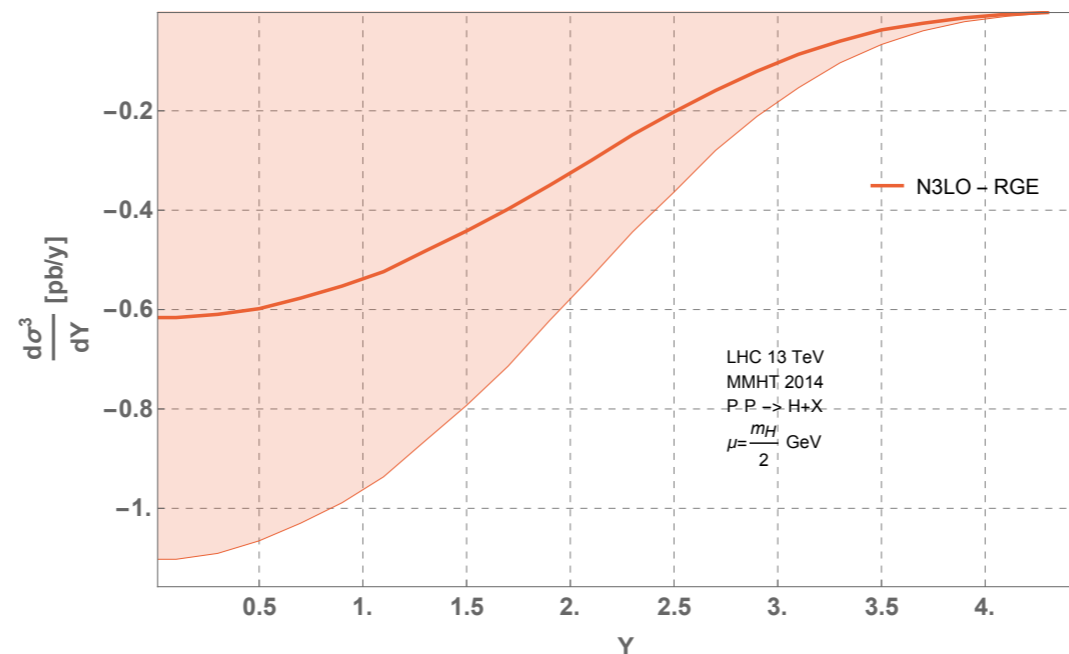
MAIN CHALLENGES FOR N3LO

- ▶ Rapid growth in analytic complexity: Many more integrals to compute, large rational expressions as a result
- ▶ Numerical stability vs. speed in evaluation of analytic coefficients.

UV RENORMALISATION AND IR FACTORISATION

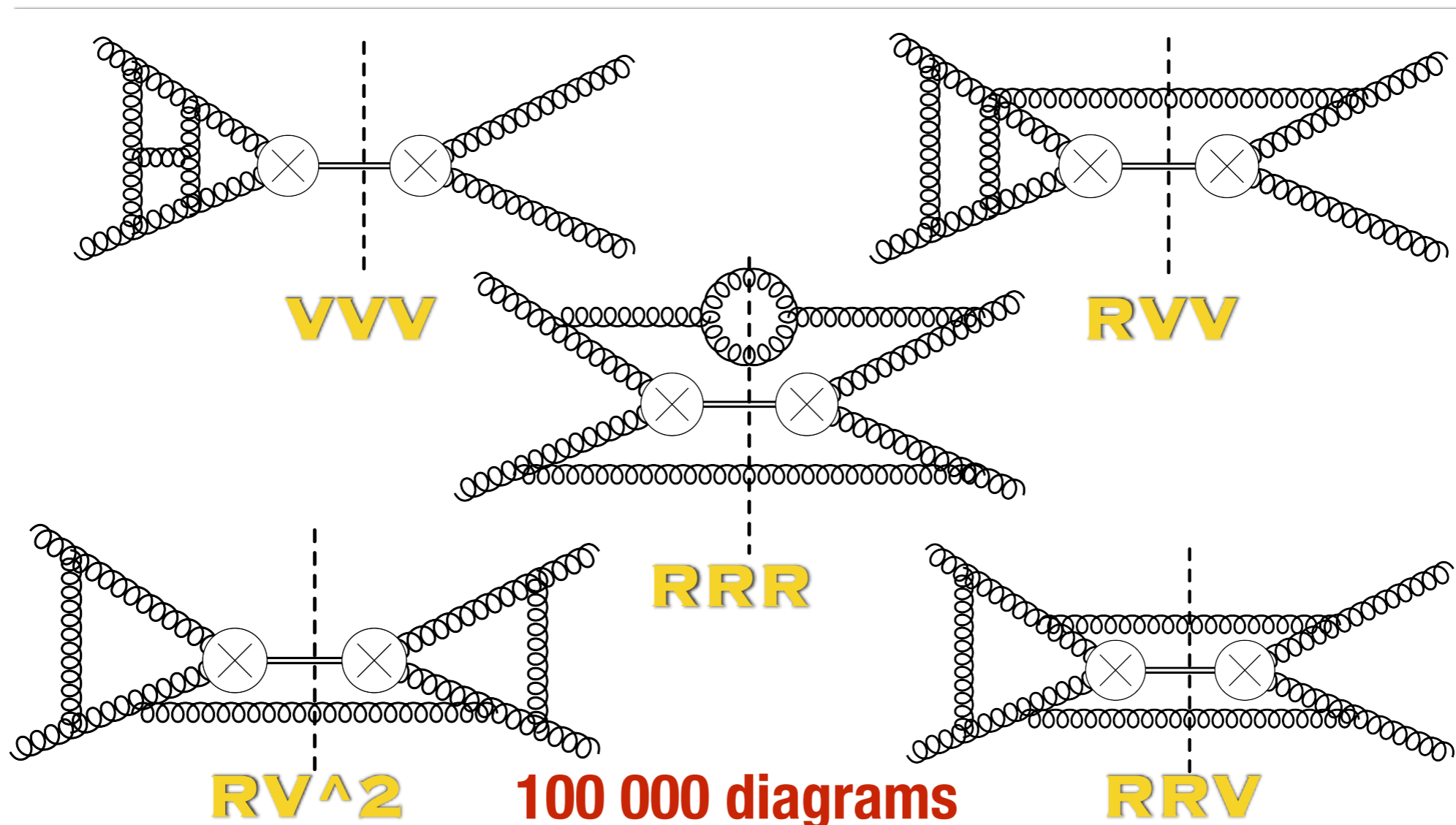
- ▶ To derive UV counter terms and IR subtraction terms we require NNLO cross sections computed beyond the finite term in ϵ
- ▶ Allow to derive complete N3LO scale variation from DGLAP

$$\hat{\sigma}^{(3)} = \hat{\sigma}_0^{(3)} + \hat{\sigma}_1^{(3)} \log\left(\frac{m_h^2}{\mu^2}\right) + \hat{\sigma}_2^{(3)} \log^2\left(\frac{m_h^2}{\mu^2}\right) + \hat{\sigma}_3^{(3)} \log^3\left(\frac{m_h^2}{\mu^2}\right)$$



FIXED ORDER MATRIX ELEMENTS

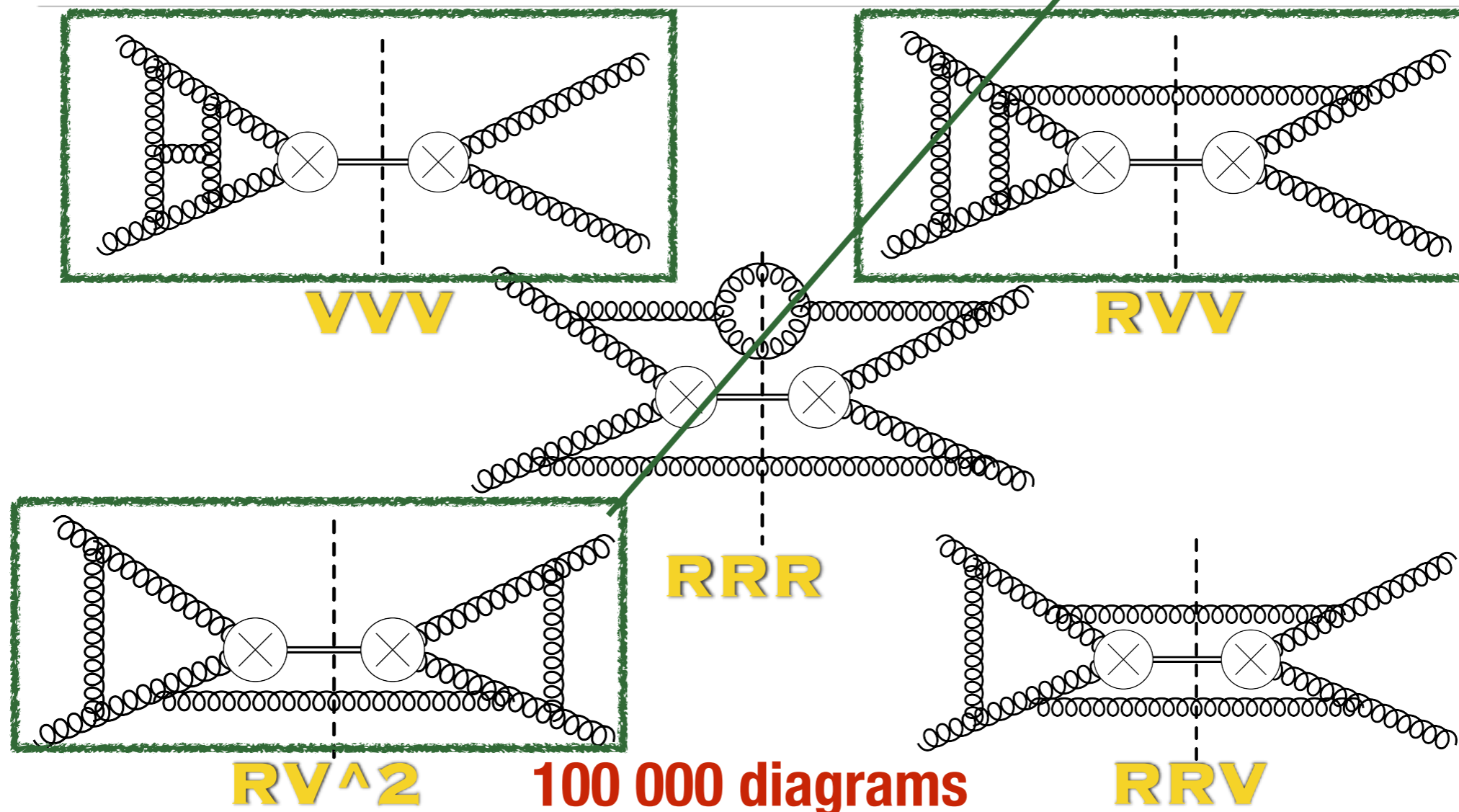
Rapid growth in complexity



FIXED ORDER MATRIX ELEMENTS

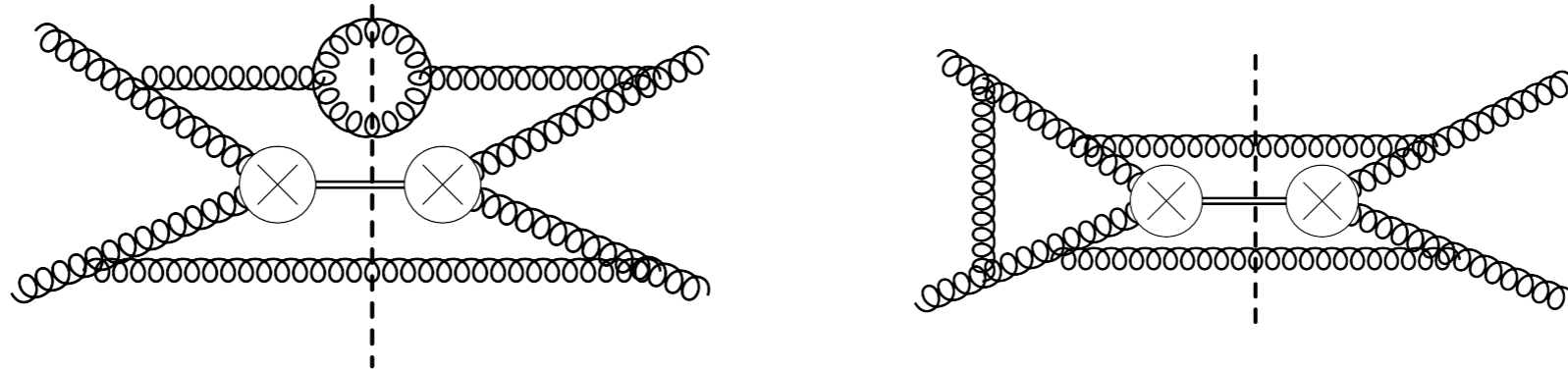
Known already! (Inclusive / H+J @NNLO)

Rapid growth in complexity



100 000 diagrams

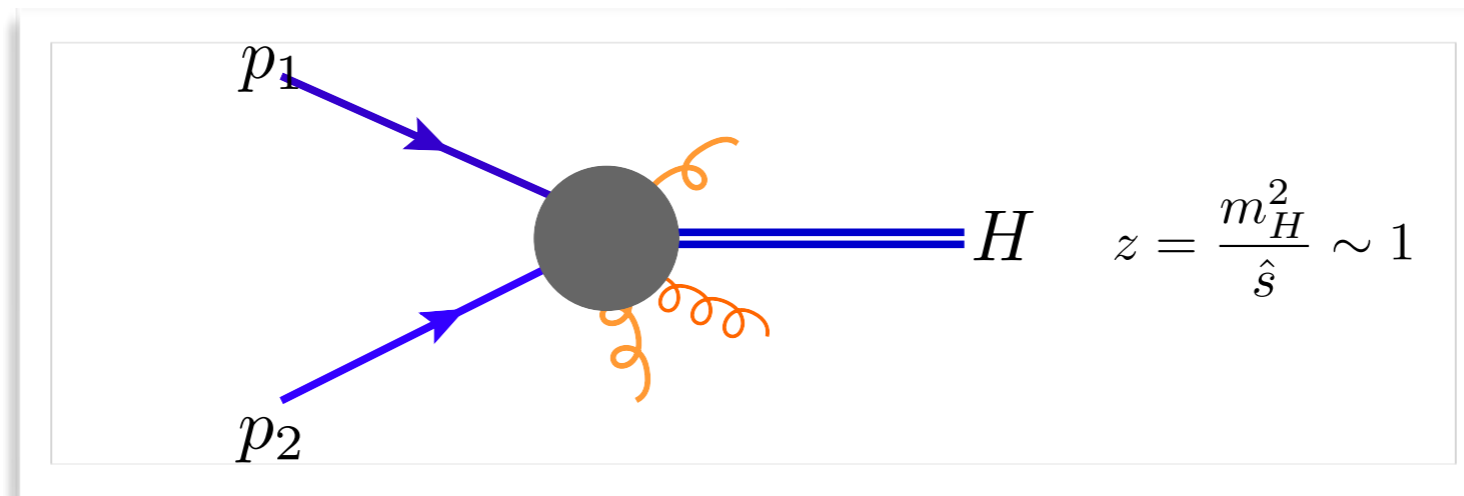
FIXED ORDER MATRIX ELEMENTS



- ▶ Missing matrix elements with 2 or 3 final state partons.
- ▶ Same strategy as for NNLO: Analytic computation using reverse unitarity, master integrals and differential equations.
- ▶ Number of master integrals required: $100 \times$ NNLO.
- ▶ Solving differential equations for master integrals:
Need boundary conditions = Master integrals evaluated at one single point.

EXPAND CROSS SECTION AROUND PRODUCTION THRESHOLD

- ▶ Inclusive Cross Section: Computed as a Threshold Expansion



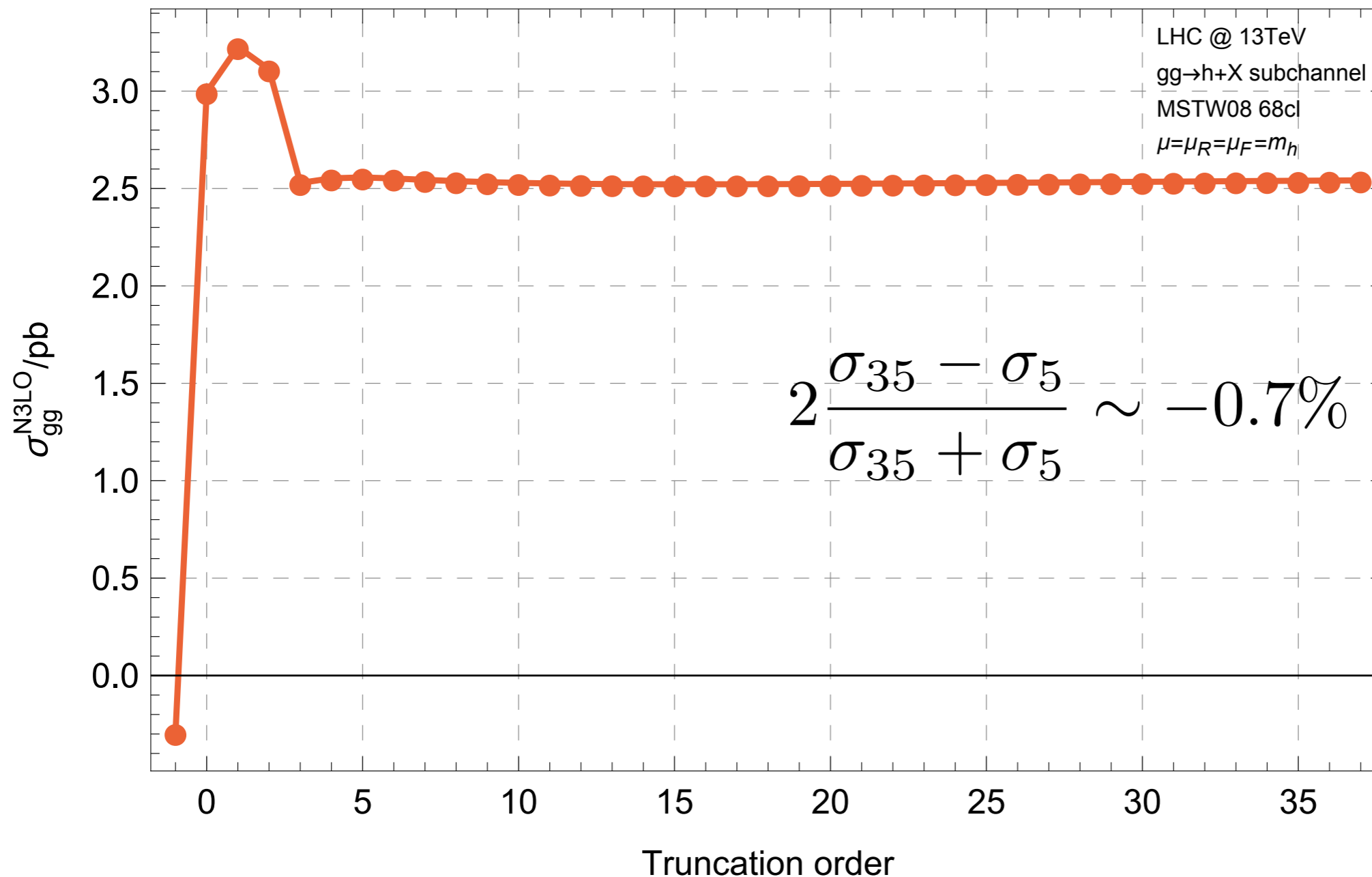
$$\bar{z} = 1 - z \quad \longrightarrow \quad \hat{\sigma}(\bar{z}) = \sigma^{SV} + \sigma^{(0)} + \bar{z}\sigma^{(1)} + \dots$$

- ▶ Served as an excellent approximation for inclusive cross section.
- ▶ Reason **Nr.1**:
Crucial analytic information a full calculation relies on - boundary conditions.
+ checks, testing ground for technology, etc.
- ▶ Reason Nr. 2: Can we use it for phenomenology?

TRUNCATION ORDER @ N3LO: INCLUSIVE

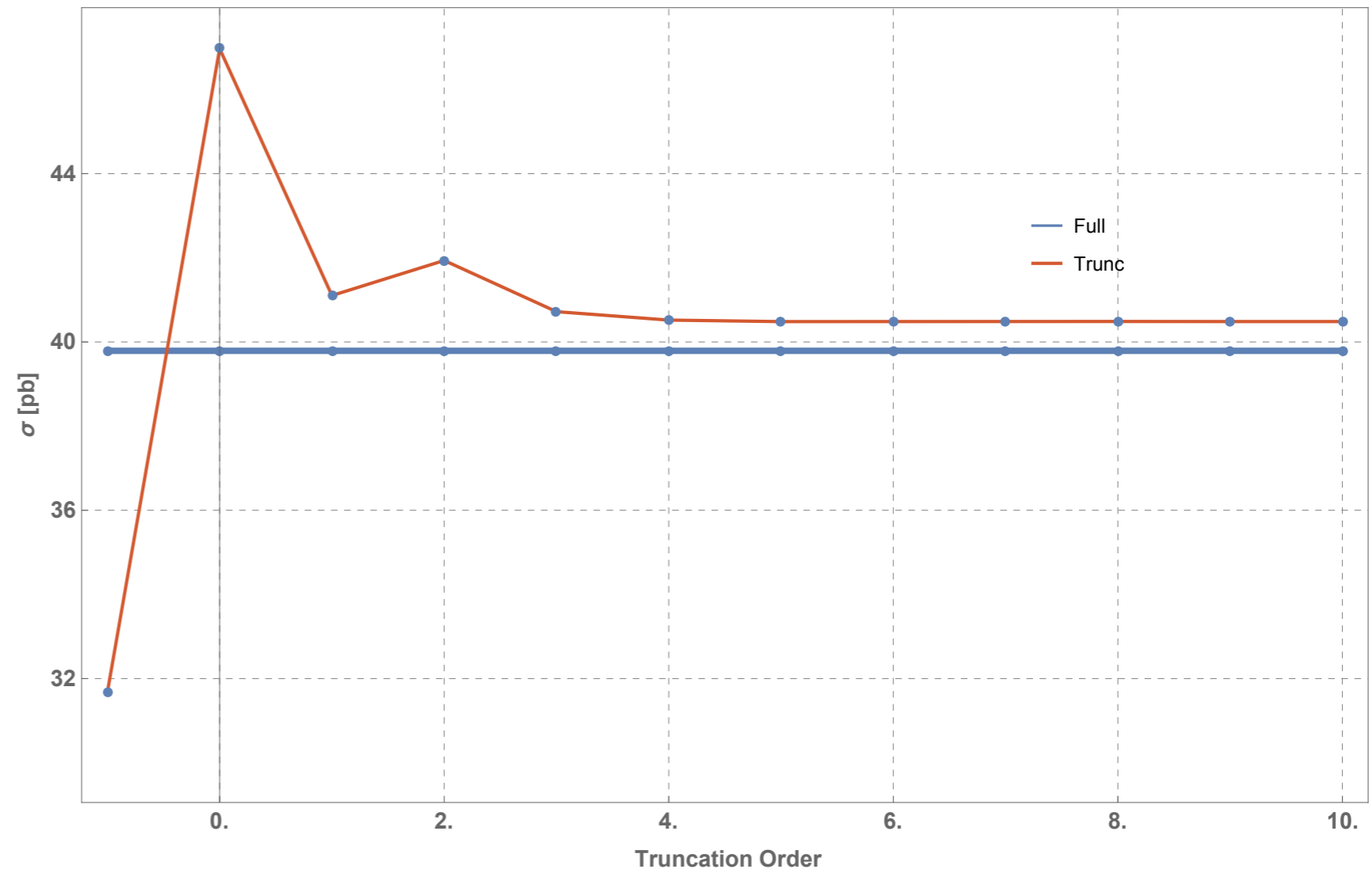
► How well did it work for N3LO inclusive?

$$\delta\sigma^{N3LO} \approx 0.5\%$$



TRUNCATION ORDER @ NNLO: INCLUSIVE

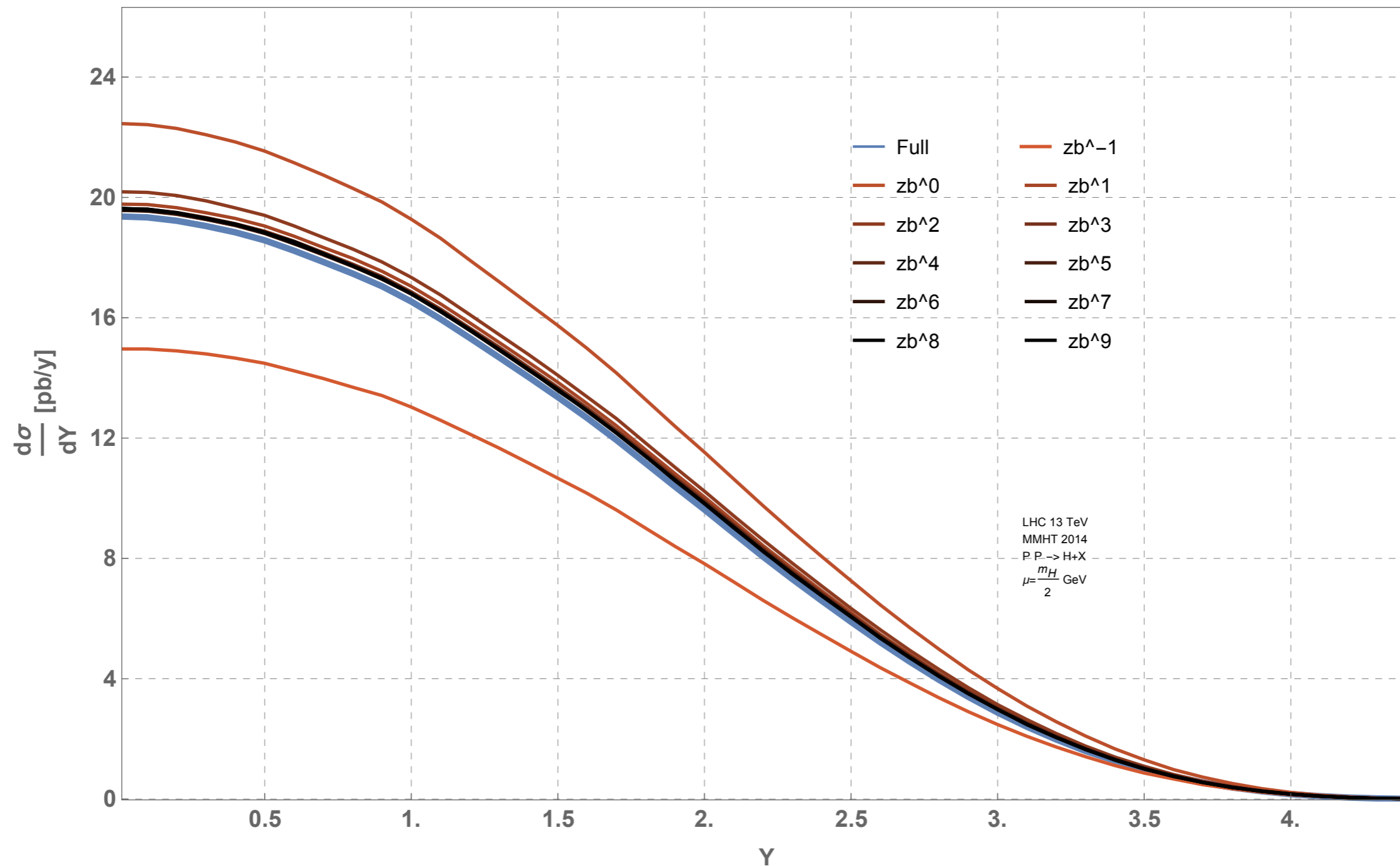
- ▶ Let's test at NNLO:
- ▶ Good approximation reached with only 5-10 terms
- ▶ ~3% off-set from full result for the NNLO correction -
> ~0.5 % on the total



- ▶ How well does it work differentially?

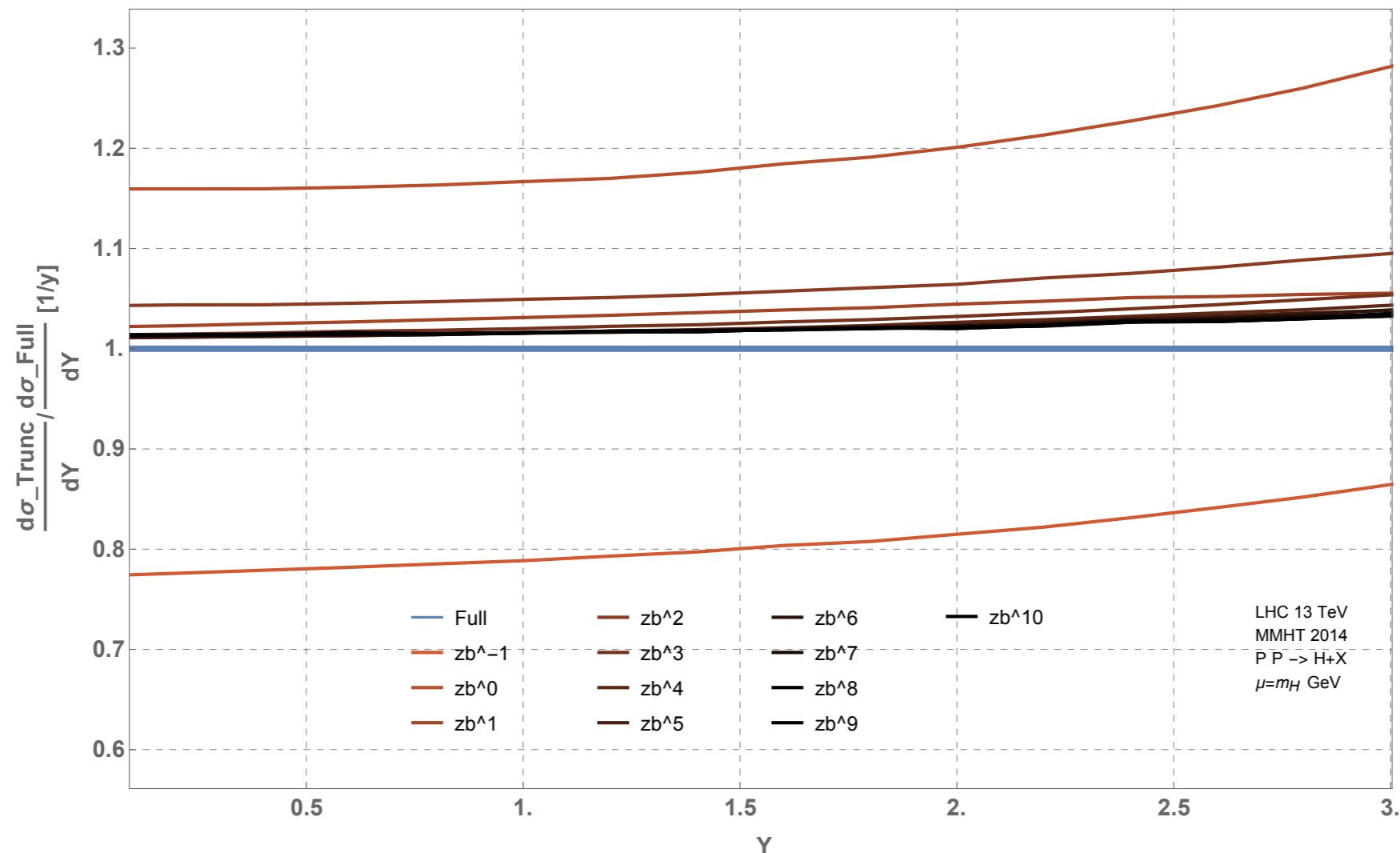
THRESHOLD EXPANSION @ NNLO DIFFERENTIAL

► Rapidity distribution



THRESHOLD EXPANSION @ NNLO DIFFERENTIAL

- ▶ Rapidity distribution normalised to true value.



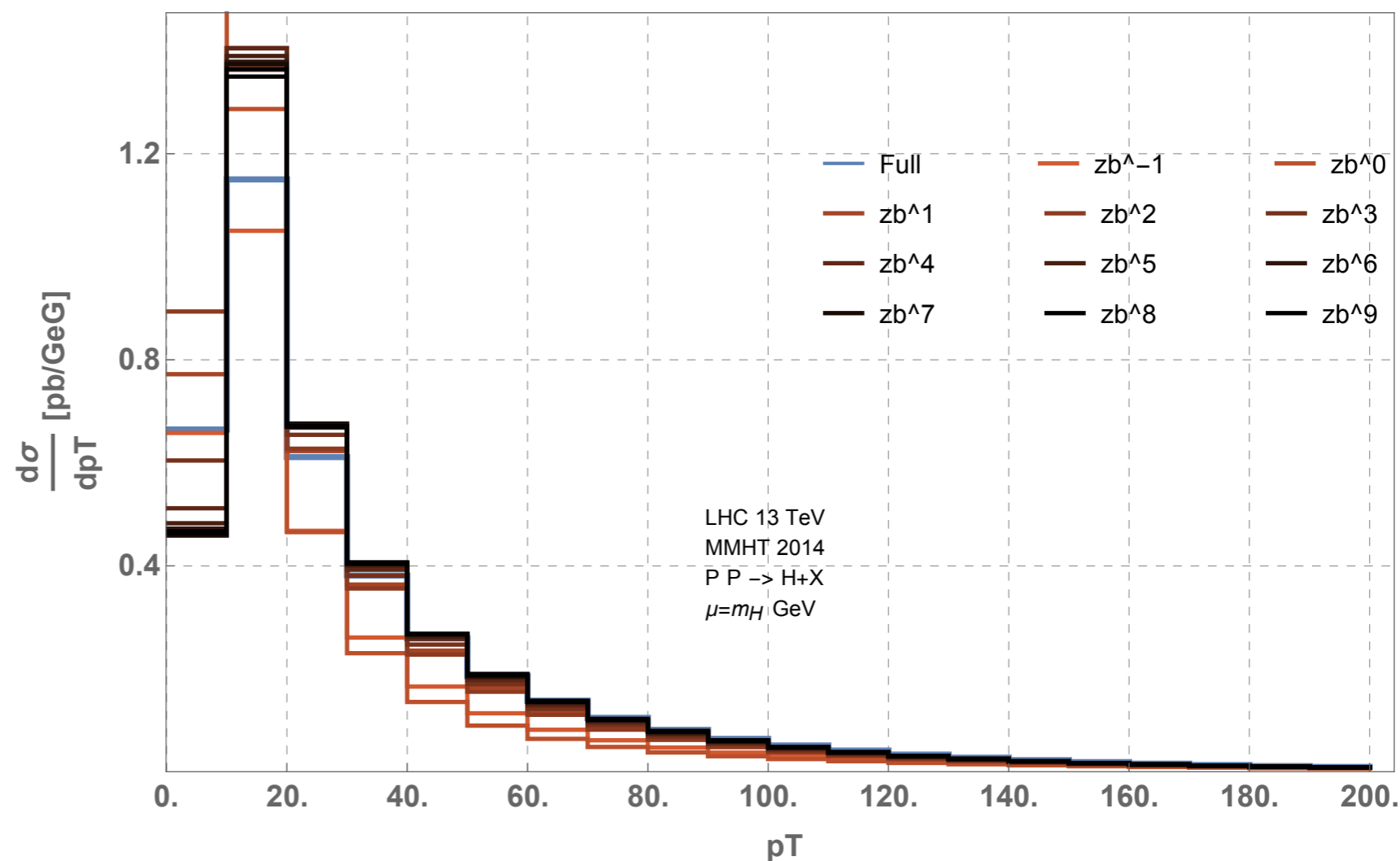
- ▶ Bulk of XS is described well with a couple of terms
- ▶ Systematic improvement possible

THRESHOLD EXPANSION @ NNLO DIFFERENTIAL

▶ PT distribution

▶ Bad convergence at low pT

▶ On-set of distribution at NLO while threshold limit is tree level.

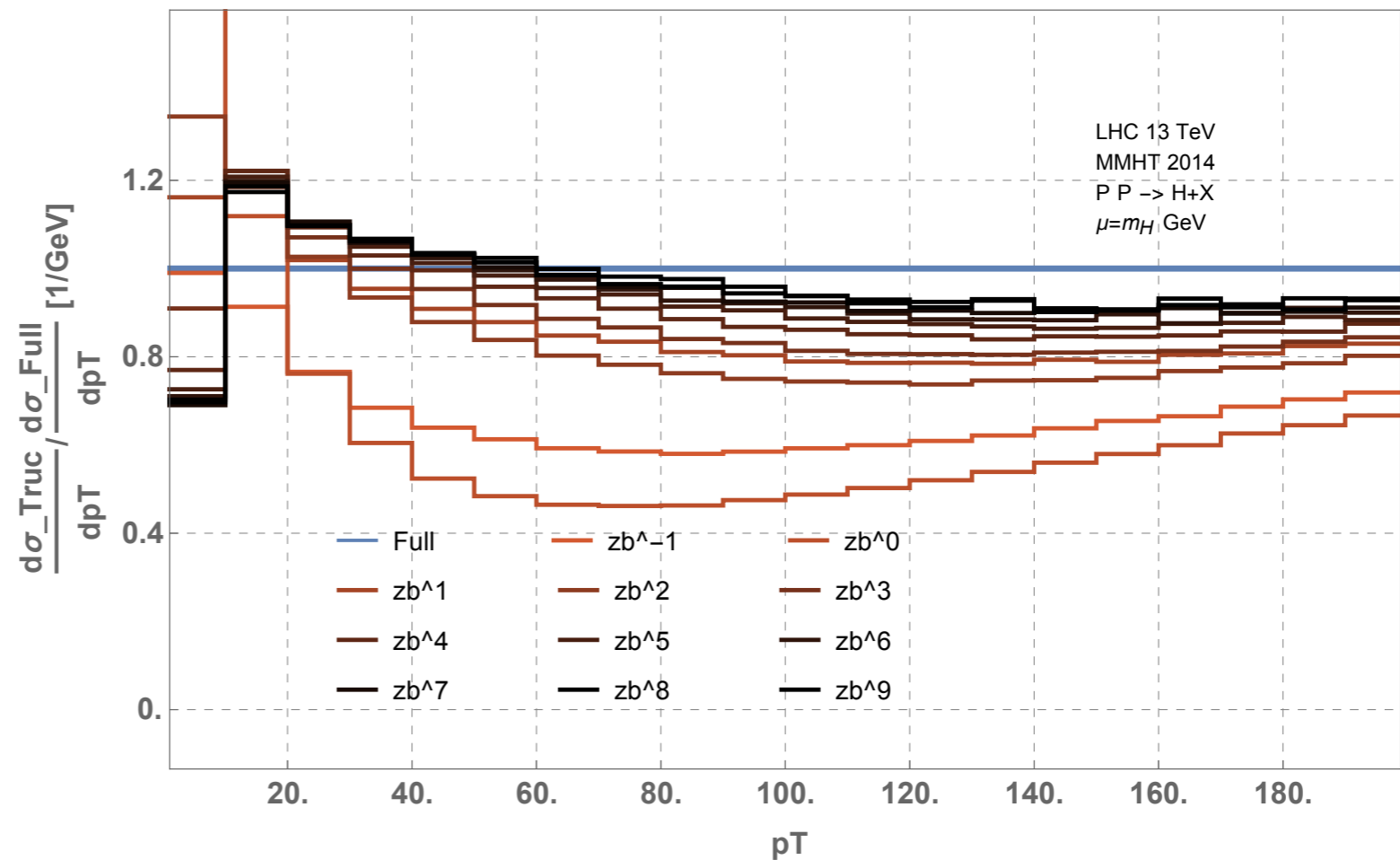


THRESHOLD EXPANSION @ NNLO DIFFERENTIAL

- ▶ PT distribution normalised to true value.

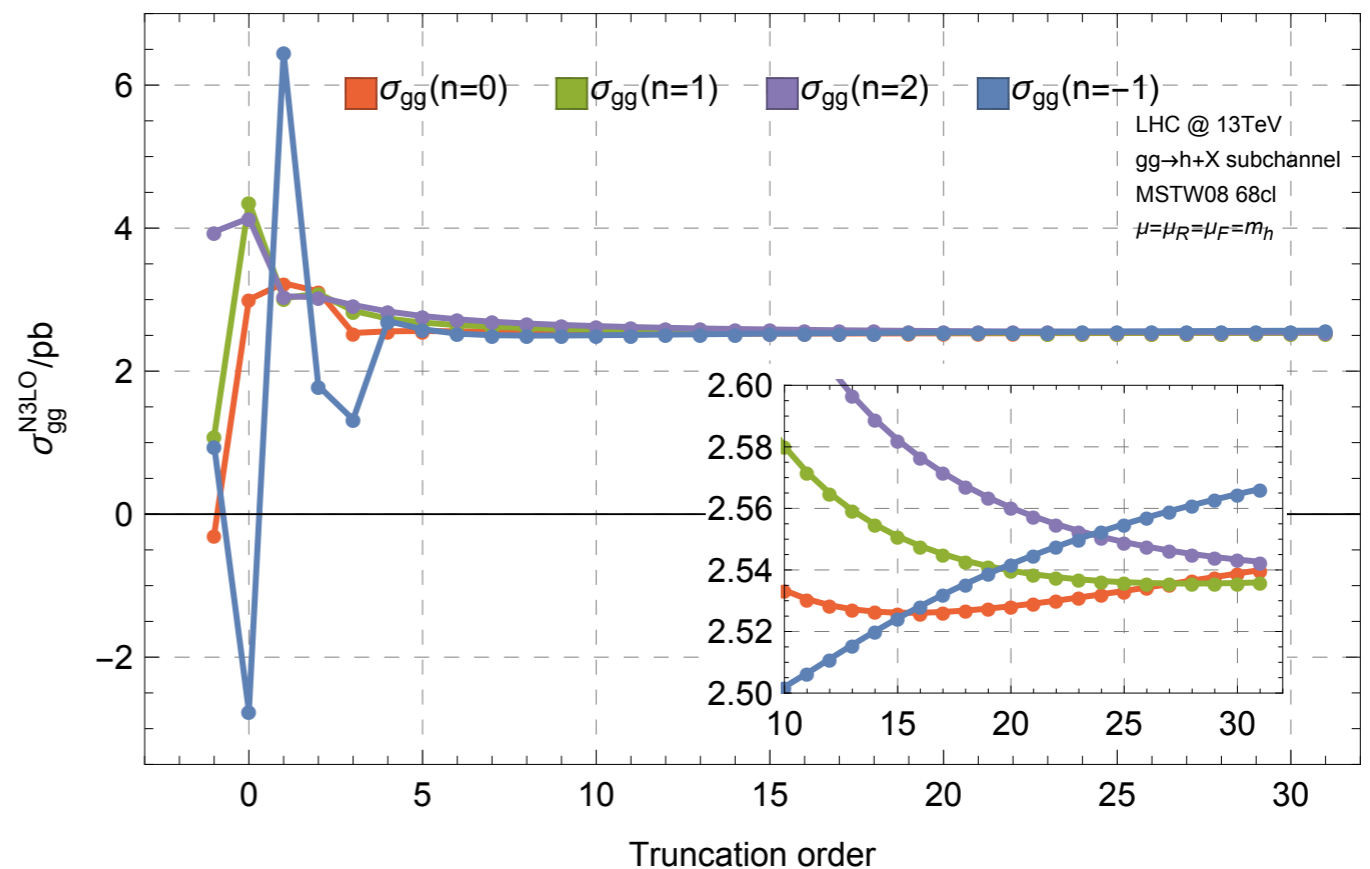
- ▶ Even with 10 terms marginally within 20 %

- ▶ Quality of expansion is subject to observable: Threshold sensitivity



THRESHOLD EXPANSION

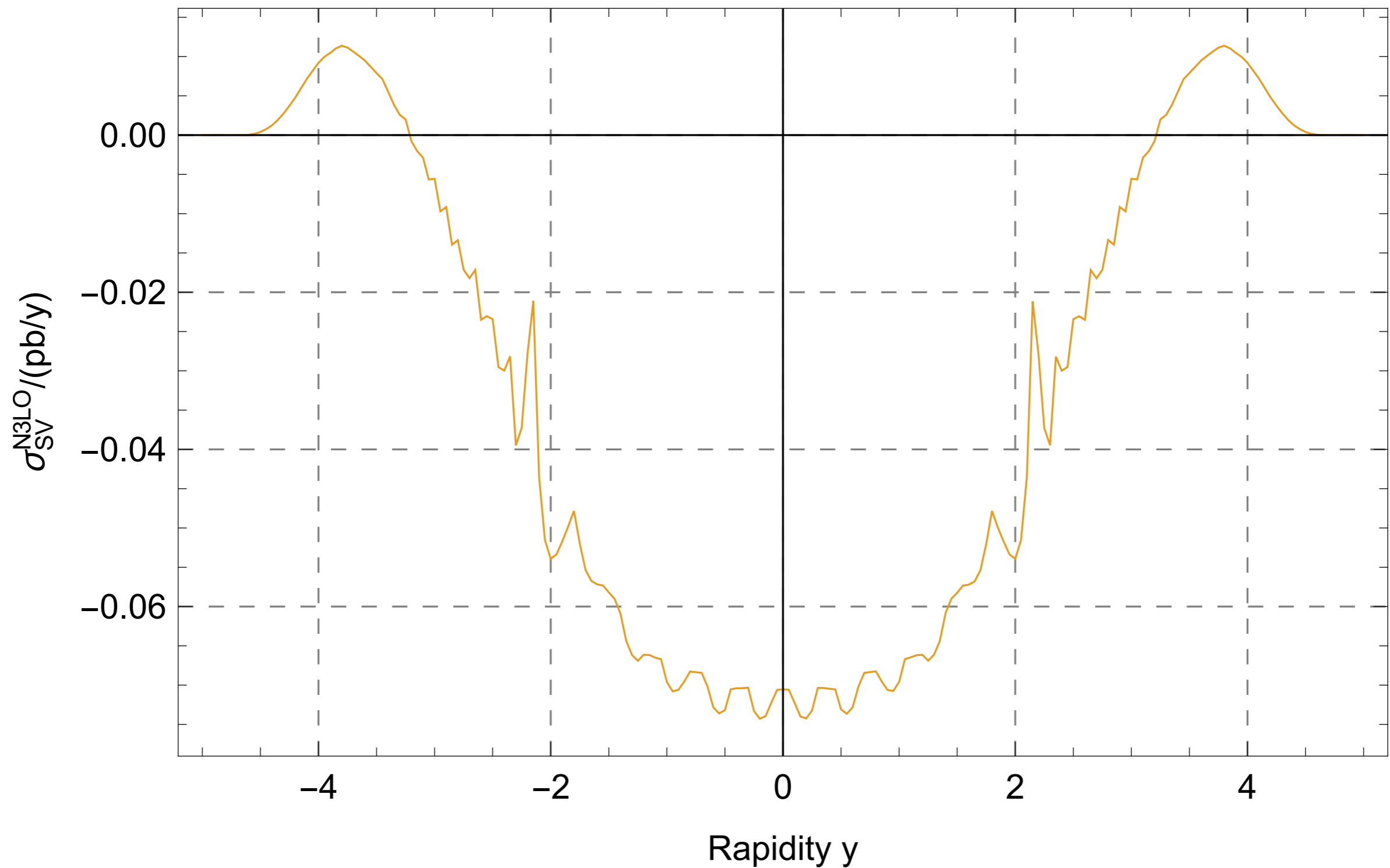
- ▶ Systematically improvable approximation.
- ▶ Soft expansion gives the opportunity to study differential distribution
- ▶ Doing phenomenology in this approximation requires careful case by case analysis to see if the approximation is valid!
- ▶ Ambiguity of higher order terms has to be controlled.
Example:
Inclusive N3LO:



THE ROAD TO N3LO

- ▶ Extend analytic techniques for automatic soft amplitude expansions.
- ▶ Apply reverse unitarity, differential equations, Multiple PolyLog, IBPs, Symbol tools,
- ▶ Compute 110 new double differential soft master integrals.
- ▶ Compute the first terms (Soft-Virtual SV) at N3LO
- ▶ Put into code and look at the N3LO corrections to the rapidity distribution and ...

N3LO CORRECTIONS TO THE RAPIDITY DISTRIBUTION



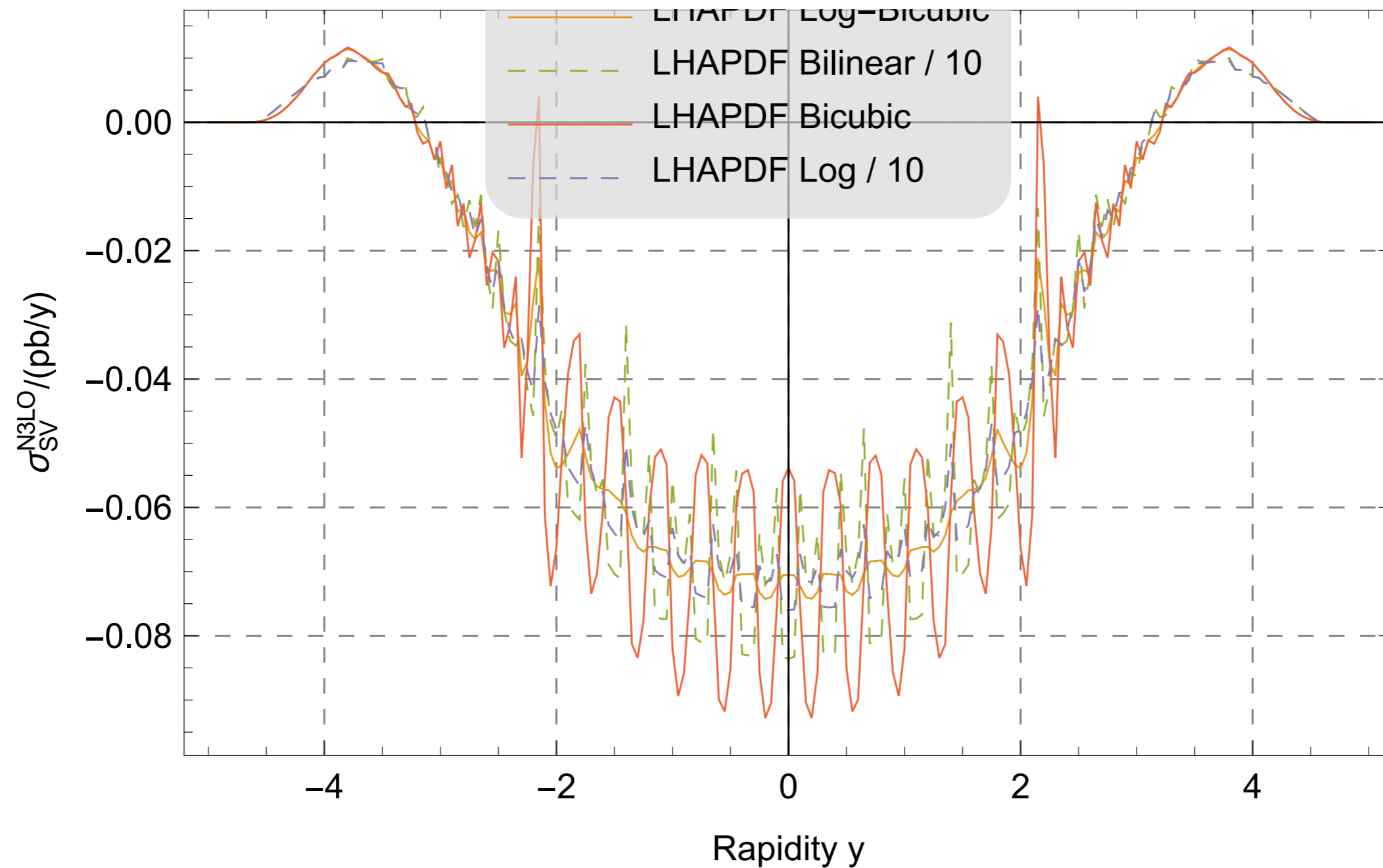
BE CAREFUL WHEN YOU DO SOMETHING NEW

$$\sigma \sim \int dz \mathcal{L}_{gg}(z) \left[\frac{\log^5(1-z)}{1-z} \right]_+$$

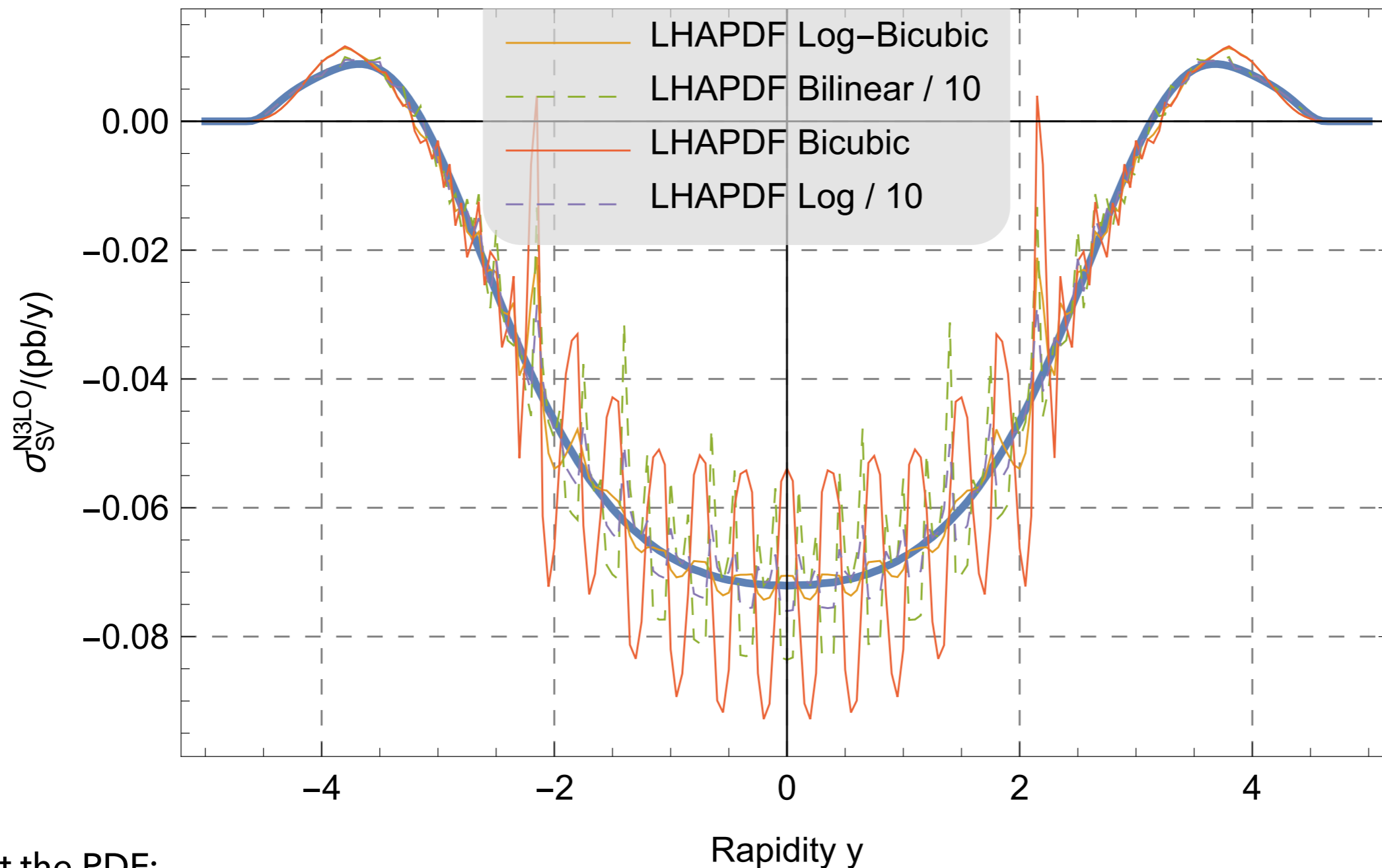
LHAPDF

- ▶ LHAPDF: Grid of points for PDFs in x and Q
- ▶ Interpolation between points with certain precision
- ▶ Not meant to be precise enough for N3LO plus distributions yet
....
- ▶ Improvements required: New interpolator, evolve from smooth PDF ?

N3LO CORRECTIONS TO THE RAPIDITY DISTRIBUTION



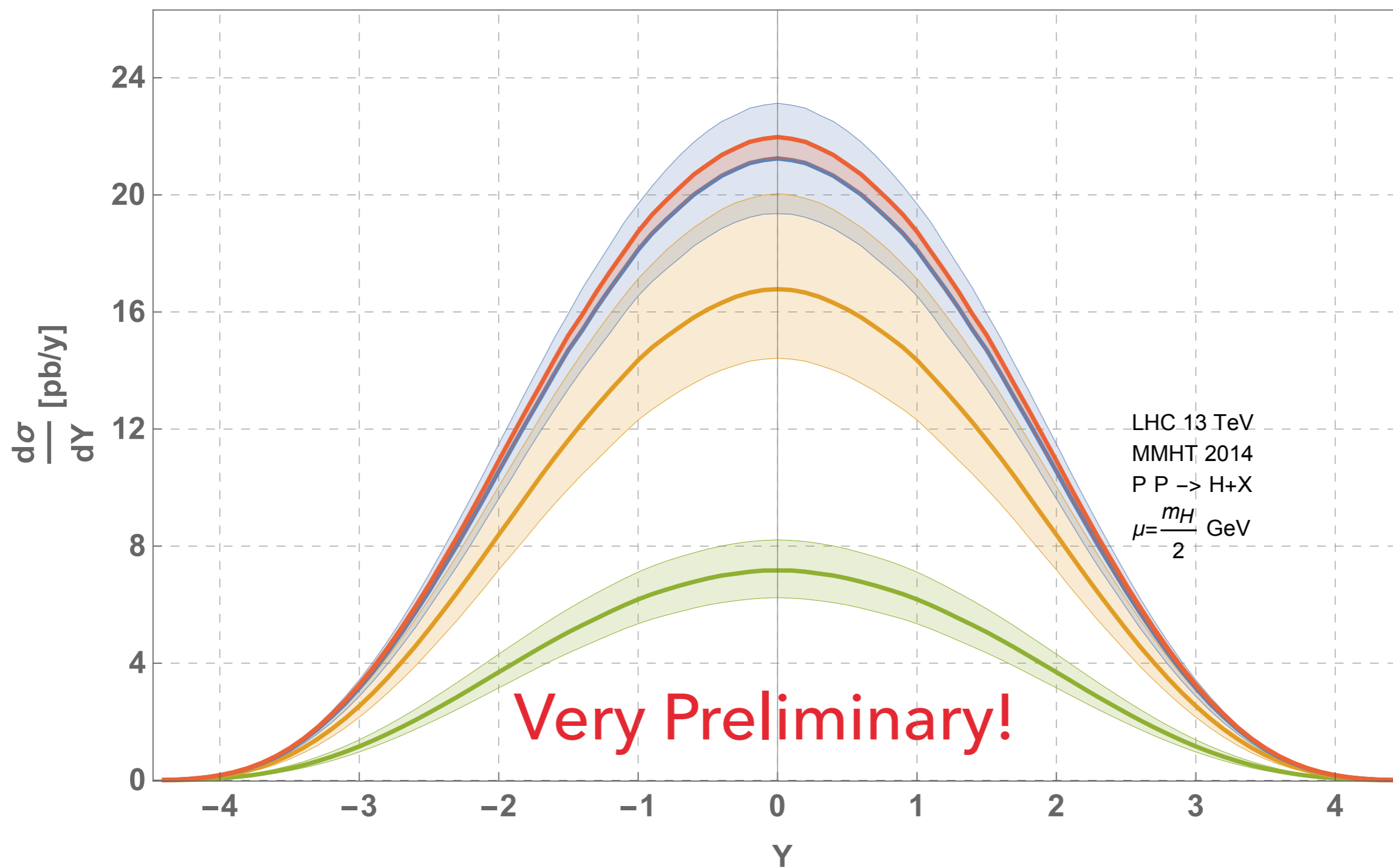
N3LO CORRECTIONS TO THE RAPIDITY DISTRIBUTION



Fit the PDF:

$$f(x, Q = 125) = c_1(1-x)^{e_1}x^{e_2} + c_2(1-x)^{e_3}(1-c_3\sqrt{x}) + c_4x + c_5\log^2(x) + c_6\log^4(x) + c_7\log^4(x)$$

INCLUDE RESCALED SV CORRECTIONS TO RAPIDITY AT N3LO



CONCLUSIONS

- ▶ Progress towards differential N3LO.
- ▶ Higgs-differential framework for realistic final states looks promising.
- ▶ Threshold expansions provide a key ingredient for analytic computation.
- ▶ Threshold expansion can be used at the differential level to approximate differential cross section predictions.
- ▶ Many interesting things to be encountered when going to higher order.

Thank you!