

**Exclusive photoproduction of a
lepton pair (TCS)
Nucleon and Nuclear Generalized
Gluon distributions**

High Energy Photon Collisions at the LHC - CERN - 24 avril 2008

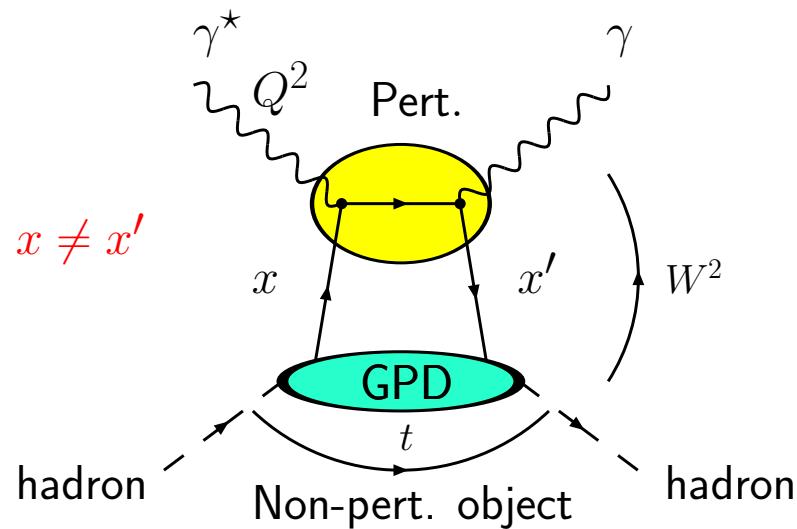
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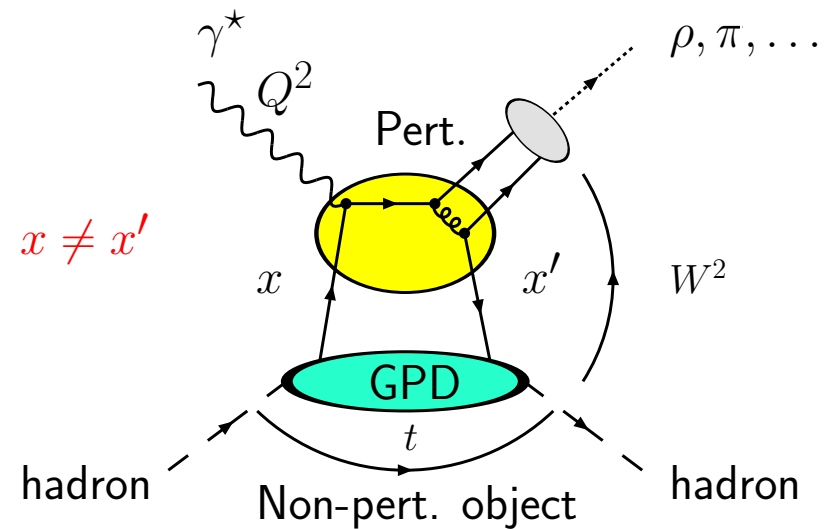
from work with L. Szymanowski, M. Diehl, J. Wagner ...

QCD factorization in Exclusive processes

DVCS



Meson Production



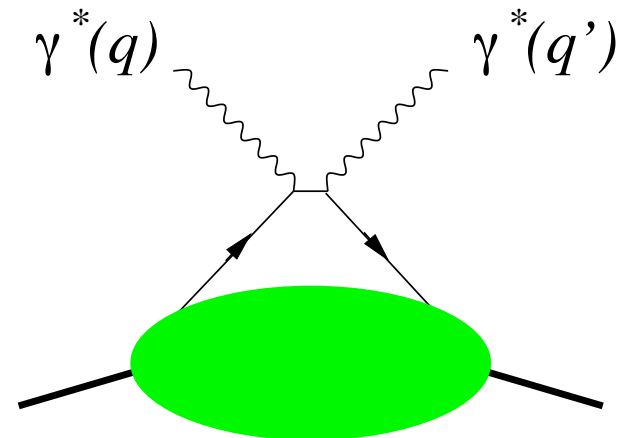
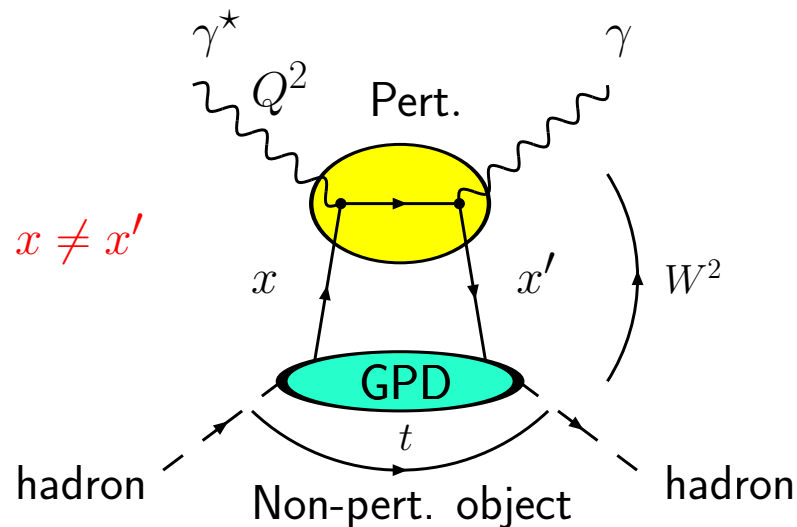
⇒ **Factorisation** between a hard part (perturbatively calculable) and a soft part (non-perturbative) *Generalized Parton Distribution*

demonstrated for $Q^2 \rightarrow \infty, x_B = \frac{Q^2}{Q^2 + W^2}$ fixed, t small

experimentally shown for $Q^2 > 2\text{GeV}^2$, at HERA and JLab

Generalised Parton Distributions

Same operators as in DIS but **non diagonal** matrix elements
 = soft part of the amplitude for exclusive reactions



$H(x, \xi, t) =$ **Fourier Transform of matrix elements**

$$\langle N(p', \lambda') | \bar{\psi}(-z/2)_\alpha [-z/2; z/2] \psi(z/2)_\beta | N(p, \lambda) \rangle \Big|_{z^+=0, z_T=0}$$

$$p' - p = \Delta \quad \Delta^2 = t \quad \Delta^+ = -\xi(p + p')^+ \quad x - x' = 2\xi$$

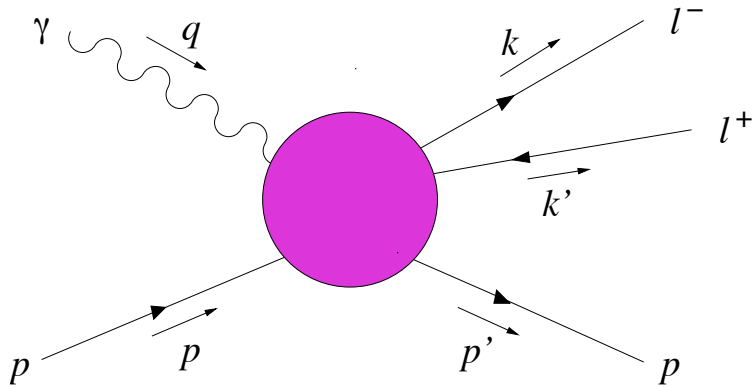
From spacelike to Timelike

Initial Photon Beam allows to study **crossed** reaction.

At lowest order, same amplitude → critical check of the
universality of GPDs.

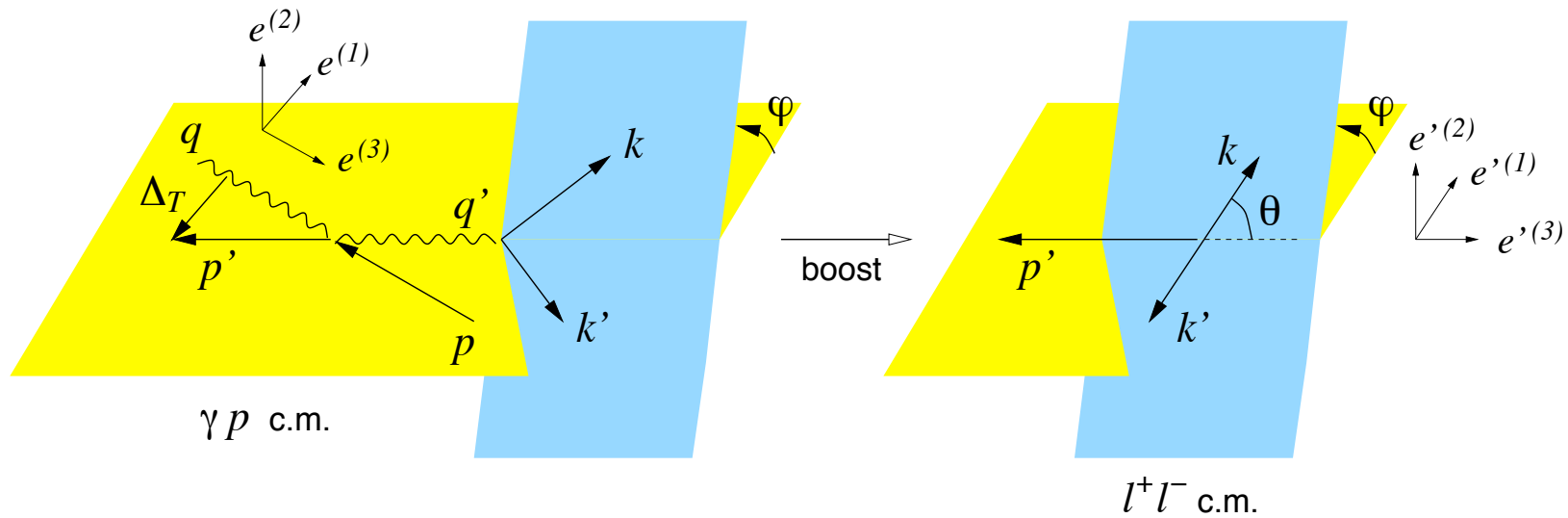
At higher orders, significant differences under control
thanks to **analyticity** properties.

Kinematics of exclusive lepton pair production

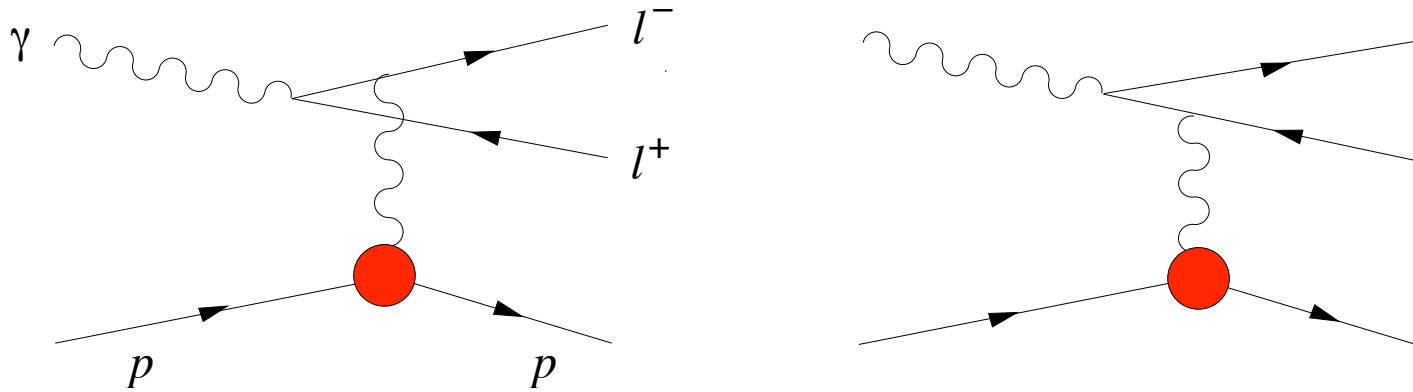


$$\tau = \frac{(k+k')^2}{s-M^2}$$

$$\eta = \frac{(p-p')^+}{(p+p')^+} = \frac{\tau}{2-\tau}$$



Bethe-Heitler process



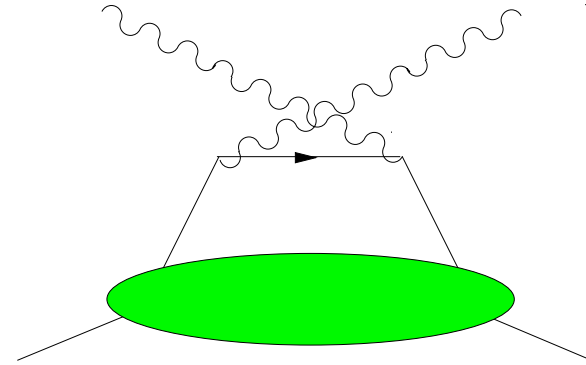
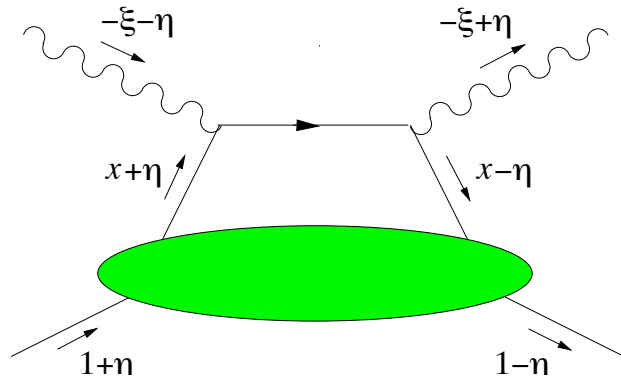
$$\frac{d\sigma_{BH}}{dQ'^2 dt d(\cos\theta) d\varphi} \approx \frac{\alpha_{em}^3}{2\pi s^2} \frac{1}{-t} \frac{1 + \cos^2\theta}{\sin^2\theta} \left[\left(F_1^2 - \frac{t}{4M^2} F_2^2 \right) \frac{2}{\tau^2} \frac{\Delta_T^2}{-t} + (F_1 + F_2)^2 \right].$$

restrict to $\sin^2\theta > 1/2$ to keep B-H far from $\frac{1}{\sin^2\theta}$ singularity

$$\frac{\Delta_T^2}{-t} \approx 1 \quad \rightarrow \quad \text{first term dominant at small } \tau$$

\rightarrow B-H Cross section almost constant in s at fixed Q^2

The TCS amplitude at lowest order



The hadronic tensor is $T^{\alpha\beta} = i \int d^4x e^{-iq \cdot x} \langle p(p') | T J_{em}^\alpha(x) J_{em}^\beta(0) | p(p) \rangle =$

$$-\frac{1}{(p+p')^+} \bar{u}(p') \left[g_T^{\alpha\beta} \left(\mathcal{H}_1 \gamma^+ + \mathcal{E}_1 \frac{i\sigma^{+\rho} \Delta_\rho}{2M} \right) + i\epsilon_T^{\alpha\beta} \left(\tilde{\mathcal{H}}_1 \gamma^+ \gamma_5 + \tilde{\mathcal{E}}_1 \frac{\Delta^+ \gamma_5}{2M} \right) \right] u(p)$$

with

$$\mathcal{H}_1(\xi, \eta, t) = \sum_q e_q^2 \int_{-1}^1 dx \left(\frac{1}{\xi - x - i\epsilon} - \frac{1}{\xi + x - i\epsilon} \right) H^q(x, \eta, t) \dots$$

$H^q(x, \eta, t)$ is the quark GPD in the target

$$\text{TCS : } -\xi = \eta = \frac{\tau}{2-\tau}$$

Resulting cross-section at LO

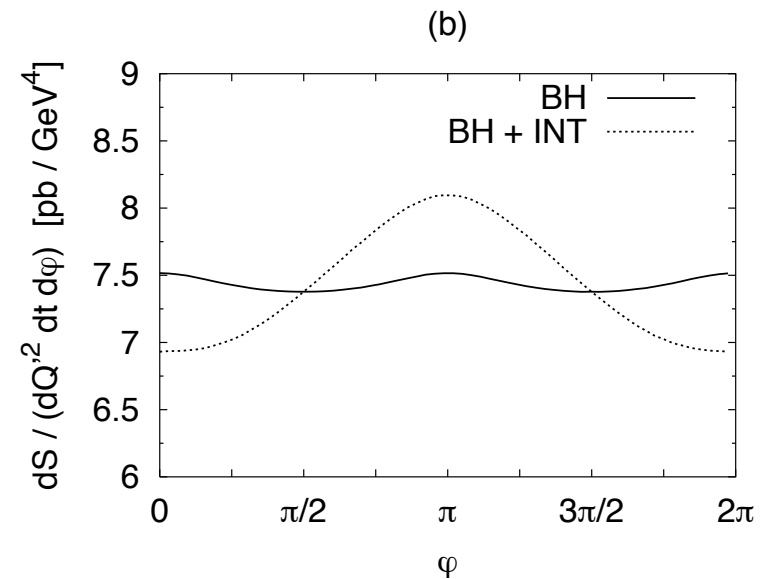
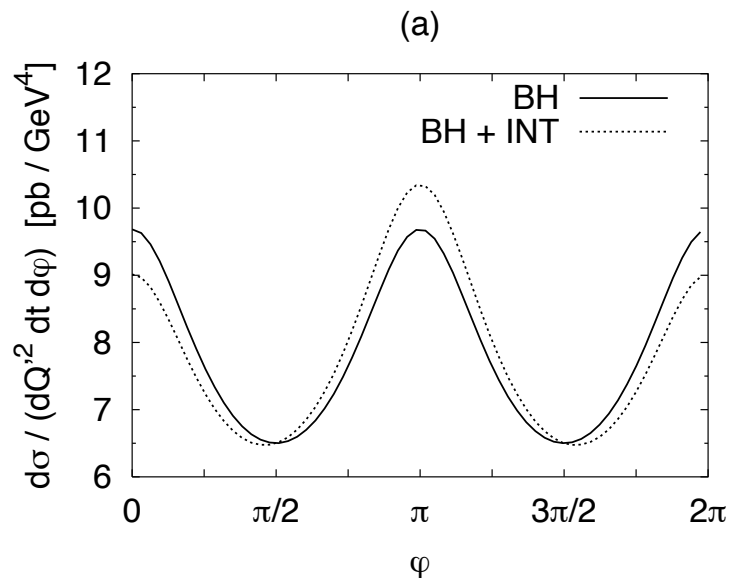
$$\frac{d\sigma_{TCS}}{dQ'^2 dt d(\cos\theta) d\varphi} \approx \frac{\alpha_{em}^3}{8\pi s^2} \frac{1}{Q'^2} \frac{1 + \cos^2\theta}{4} \sum_{\lambda, \lambda'} |M^{\lambda'-, \lambda-}|^2.$$

$$\begin{aligned} \frac{1}{2} \sum_{\lambda, \lambda'} |M^{\lambda'-, \lambda-}|^2 &= (1 - \eta^2) \left(|\mathcal{H}_1|^2 + |\tilde{\mathcal{H}}_1|^2 \right) - 2\eta^2 \mathbf{Re} \left(\mathcal{H}_1^* \mathcal{E}_1 + \tilde{\mathcal{H}}_1^* \tilde{\mathcal{E}}_1 \right) \\ &\quad - \left(\eta^2 + \frac{t}{4M^2} \right) |\mathcal{E}_1|^2 - \eta^2 \frac{t}{4M^2} |\tilde{\mathcal{E}}_1|^2, \end{aligned}$$

where \mathcal{H}_1 , $\tilde{\mathcal{H}}_1$, \mathcal{E}_1 , $\tilde{\mathcal{E}}_1$ are to be evaluated at $-\xi = \eta$.

Results at low energy

cf. E. Berger, M. Diehl, B.P., Eur. Phys. J. C 23, 675 (2002)



B-H dominant; TCS dominated by quark GPDs

Charge asymmetry \sim **interference** of B-H and TCS contributions

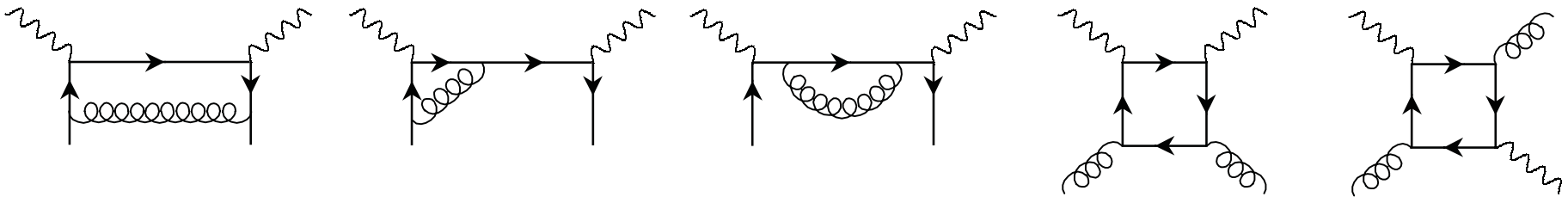
Factorization at NLO

$$\mathcal{F}_i(\xi, \eta, \Delta^2; Q^2) = \sum_{a=q,g} \int_{-1}^1 dx C_i^{a[-]}(x, \xi, \eta; Q^2/\mu^2) F^a(x, \eta, \Delta^2; \mu^2) + \mathcal{O}(Q^{-2}),$$

$F^q, F^g =$ Quark and Gluon GPDs

$(\xi = -\eta)$

$$C_i^{a[\pm]}(x, \xi, \eta; Q^2/\mu^2) = C_{i(0)}^{a[\pm]}(x, \xi) + \frac{\alpha_s}{2\pi} C_{i(1)}^{a[\pm]}(x, \xi, \eta; Q^2/\mu^2) + \mathcal{O}(\alpha_s^2).$$



Coefficient Functions

cf. Belitsky-Radyushkin, Phys. Rep. 418 (2005)

$$C_{i(1)}^{q[\pm]}(x, \xi, \eta; Q^2/\mu^2) \equiv C_F Q_q^2 \left[c_{i(1)}^{q[\pm]}(x, \xi, \eta) + \kappa_{i(1)}^{q[\pm]}(x, \xi, \eta) \ln(Q^2/\mu^2) \right] ,$$

$$C_{i(1)}^{g[\pm]}(x, \xi, \eta; Q^2/\mu^2) \equiv 2T_F \sum_q Q_q^2 \left[c_{i(1)}^{g[\pm]}(x, \xi, \eta) + \kappa_{i(1)}^{g[\pm]}(x, \xi, \eta) \ln(Q^2/\mu^2) \right] .$$

$$\begin{aligned} c_{1(1)}^{g[-]}(x, \eta, \xi) = & \frac{4\xi^2 - 4x\xi + x^2 - \eta^2}{2(x^2 - \eta^2)^2} \ln\left(1 - \frac{x}{\xi}\right) - \frac{2\xi^2 - 2x\xi + x^2 - \eta^2}{4(x^2 - \eta^2)^2} \ln^2\left(1 - \frac{x}{\xi}\right) \\ & + \frac{(\xi - \eta)(x^2 - 4\eta\xi - \eta^2)}{2\eta(x^2 - \eta^2)^2} \ln\left(1 - \frac{\eta}{\xi}\right) - \frac{(\xi - \eta)(x^2 - 2\eta\xi - \eta^2)}{4\eta(x^2 - \eta^2)^2} \ln^2\left(1 - \frac{\eta}{\xi}\right) \\ & + (\xi \leftrightarrow -\xi) . \end{aligned}$$

$$\begin{aligned} \kappa_{1(1)}^{g[-]}(x, \xi, \eta) = & - \frac{2\xi^2 - 2x\xi + x^2 - \eta^2}{2(x^2 - \eta^2)^2} \ln\left(1 - \frac{x}{\xi}\right) \\ & - \frac{(\xi - \eta)(x^2 - 2\eta\xi - \eta^2)}{2\eta(x^2 - \eta^2)^2} \ln\left(1 - \frac{\eta}{\xi}\right) + (\xi \leftrightarrow -\xi) , \end{aligned}$$

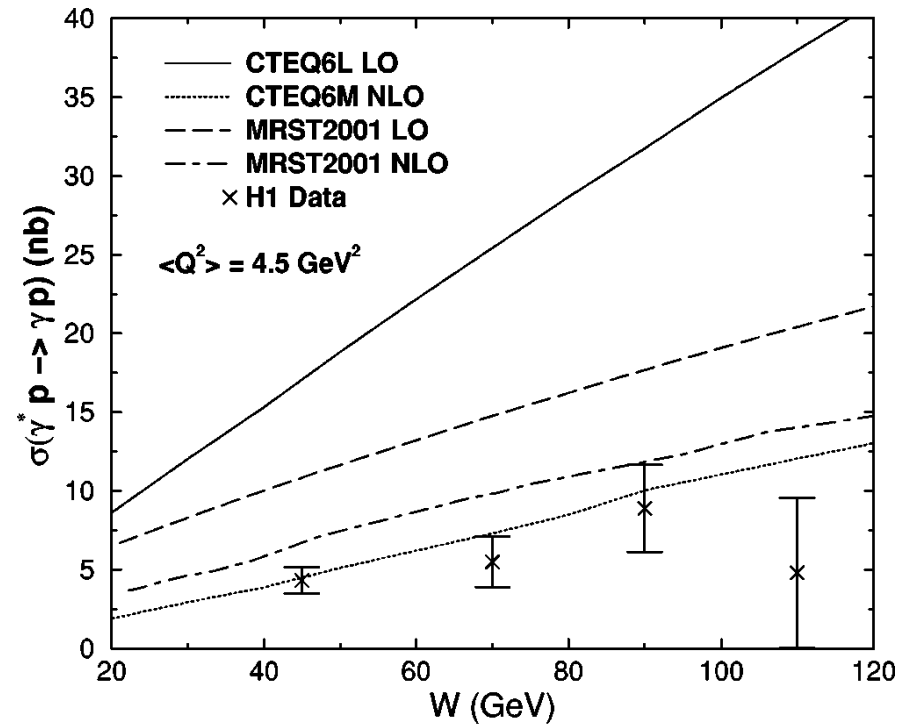
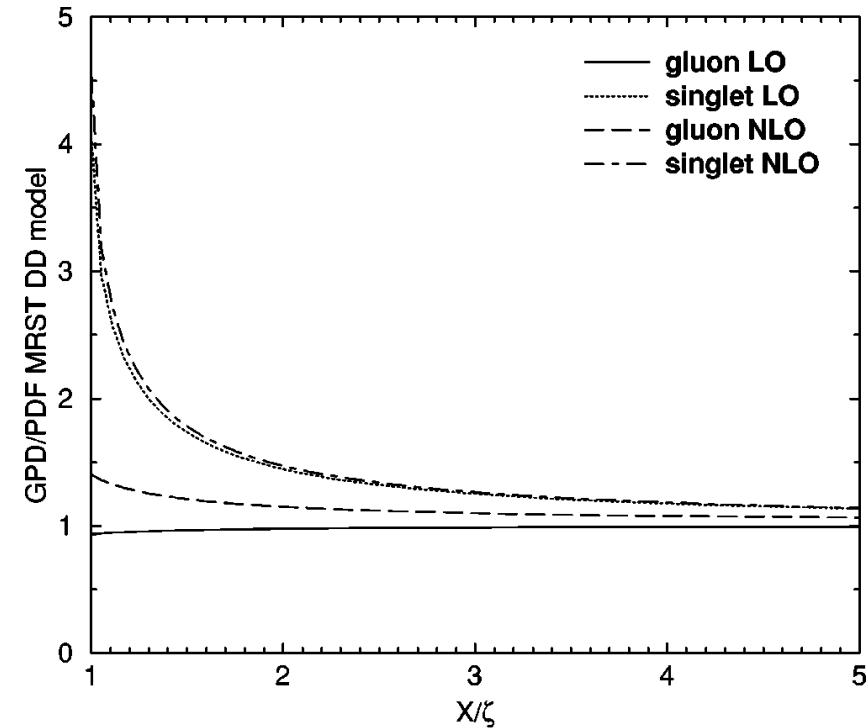
$$\kappa_{1(1)}^{g[+]}(x, \xi, \eta) = - \frac{2x\xi - x^2 - \eta^2}{2(x^2 - \eta^2)^2} \ln\left(1 - \frac{x}{\xi}\right) + \frac{x(\xi - \eta)}{(x^2 - \eta^2)^2} \ln\left(1 - \frac{\eta}{\xi}\right) - (\xi \leftrightarrow -\xi)$$

Lessons from HERA

cf. A.Freund, M.McDermott, M. Strikman

GPD/PDF at $\xi = 10^{-4}$

DVCS cross-section
($10^{-4} < \xi < 10^{-3}$)



gluon dominant ; flat $s = W^2$ -dependence

Prospects for LHC

→ Timelike Compton scattering may be measurable ;

$$Q^2 \sim 2-10 \text{ GeV}^2, \tau \sim 10^{-2} - 10^{-4}$$

→ Possibility to probe GPDs in the **small** x regime

→ Order of magnitude to estimate

→ **Nuclear** effects : Nuclear gluon GPDs cf. V.Guzey, M. Strikman

→ "differential" EMC effect