

# Perspectives of probing small $x$ dynamics in protons and nuclei in ultraperipheral AA and pA collisions at LHC

*Mark Strikman, PSU*

*CERN, April 23, 08*

Workshop on High Energy Photon Collisions at the LHC



Based on our study

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## The physics of ultraperipheral collisions at the LHC

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## *Main thrusts of the HERA small $x$ QCD physics:*

- Small  $x$  parton densities
- Inclusive hard diffractive processes
- Hard exclusive processes: vector meson production, dijets, ...

### *Main issues:*

- high gluon densities, violation of DGLAP,
- diffractive pdf's - leading twist vs higher twist;
- generalized parton densities at small  $x$

Theory - gluons are most interesting for small  $x$ :

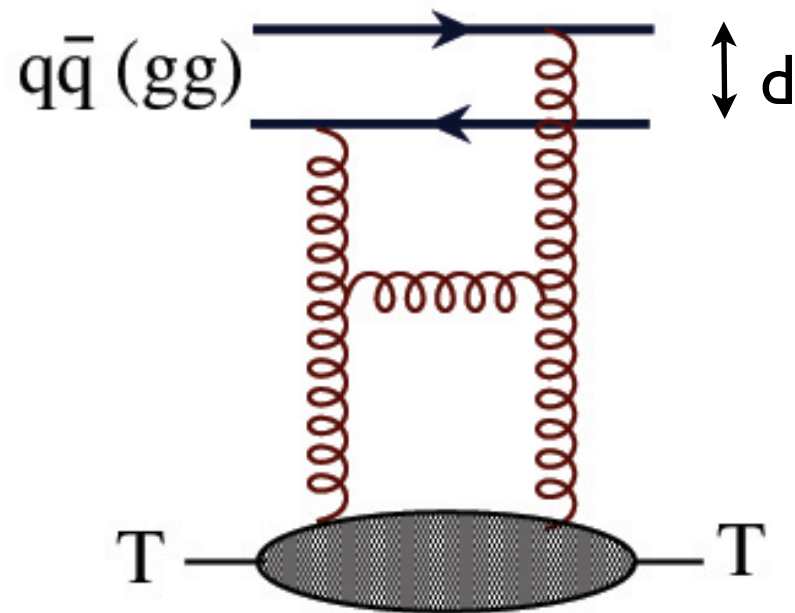
- ◆ they drive evolution and quark sea,
- ◆ interaction in the gluon sector is much stronger

Theory & HERA experience: photoproduction of dijets,  
heavy quarks, exclusive heavy meson production are  
good “*gluonometers*”



**Summary:** How strong is the interaction of small dipoles?

Consider first “small dipole - hadron” cross section



$$\sigma_{inel} = \frac{\pi^2}{3} F^2 d^2 \alpha_s (\lambda/d^2) x G_T(x, \lambda/d^2)$$

Baym et al 93

$F^2$  Casimir operator of color SU(3)

$F^2$  (quark) = 4/3

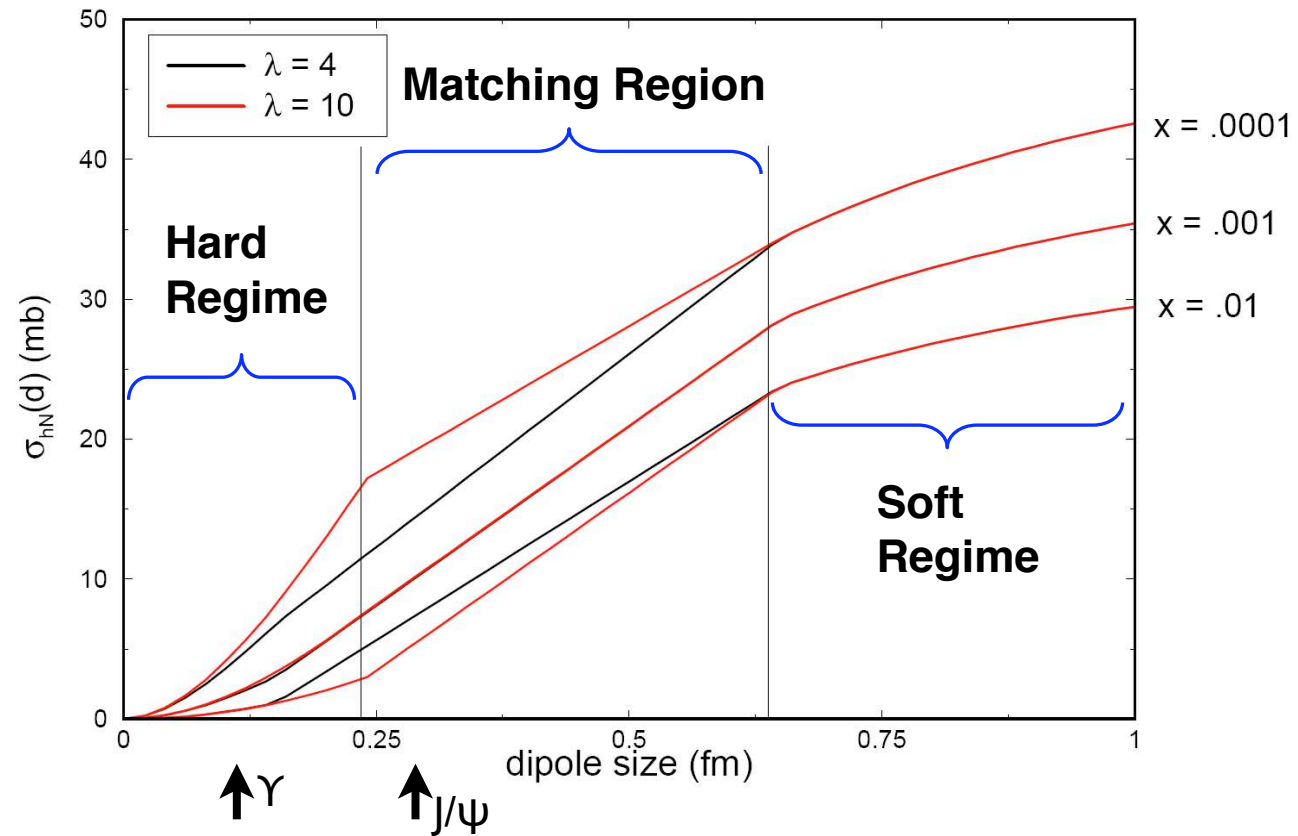
$F^2$  (gluon) = 3

Comment: This simple picture is valid only in LO. NLO would require introducing mixing of different components. Also, in more accurate expression there is an integral over x, and an extra term due to quark exchanges

New high energy QCD regime: regime of complete absorption for small  $\alpha_s$ :  
 limit - fixed  $Q$  & large energies -black disk regime (BDR)

*Evidence for proximity to BDR at HERA*

$Q^2 = 3.0 \text{ GeV}^2$



studies of the “quark-antiquark dipole” (transverse size  $d$ ) - nucleon cross section based pQCD and HERA data

Provided a reasonable prediction for  $\sigma_L$

Frankfurt et al  
2000-2001

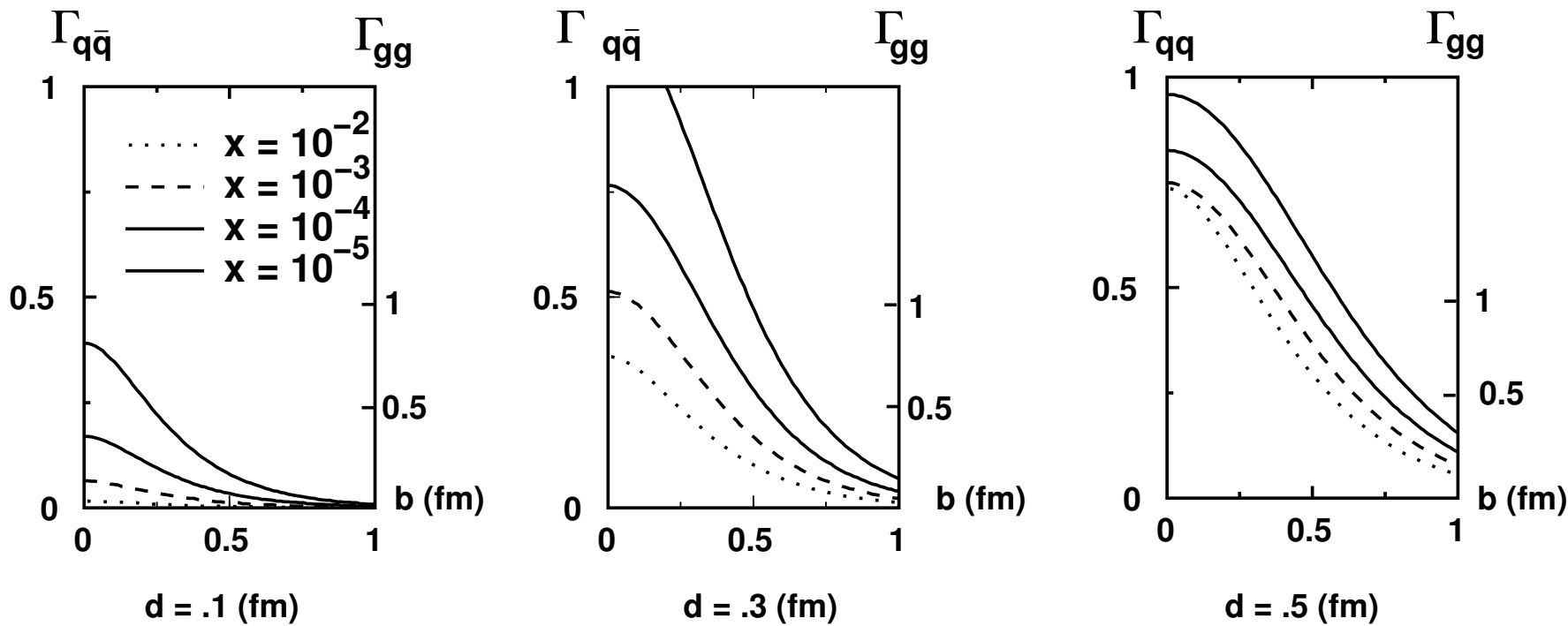


Combine with: analysis of exclusive hard processes  
(t-dependence of the dipole - nucleon scattering)

determine impact factors for elastic  $q\bar{q} - N$  scattering

$$\Gamma_h(s, b) = \frac{1}{2is} \frac{1}{(2\pi)^2} \int d^2\vec{q} e^{i\vec{q}\vec{b}} A_{hN}(s, t)$$

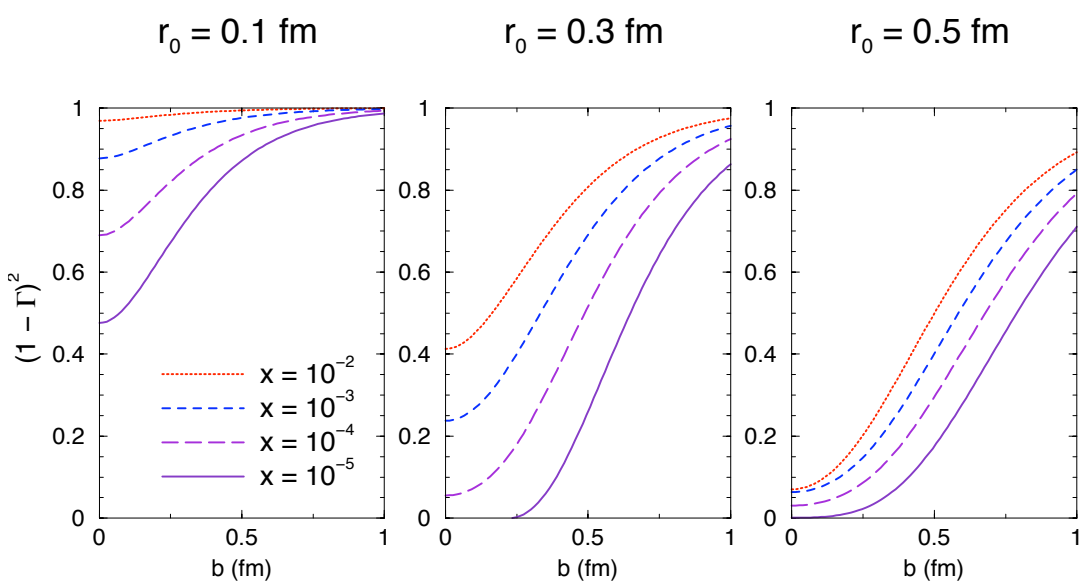
$\Gamma = 1$  corresponds to regime of complete absorption - BDR



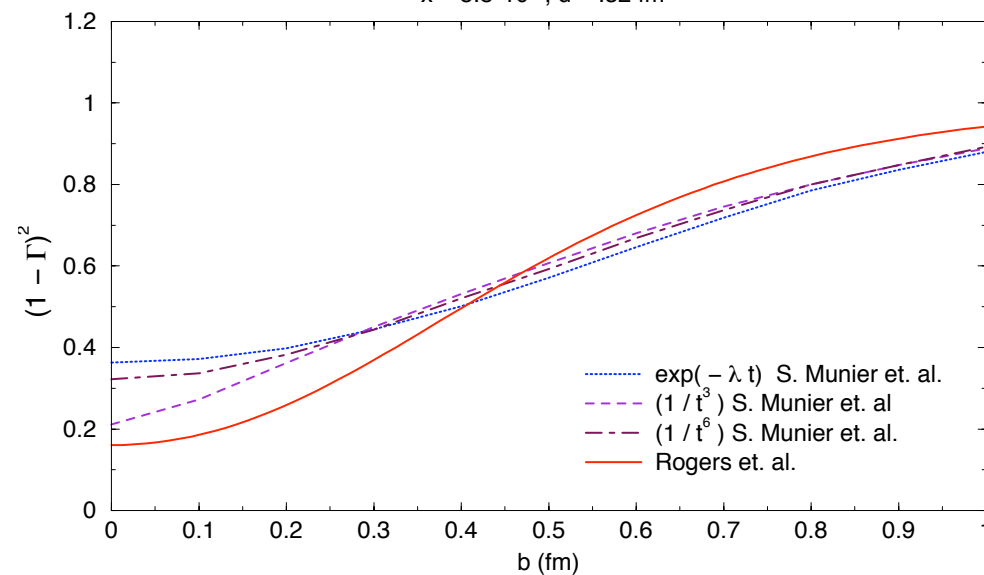
T.Rogers et al

In the case gg-N scattering we assume pQCD relation

$$\Gamma_{gg} = \frac{9}{4} \Gamma_{q\bar{q}}$$



$(1 - \Gamma)^2$  vs.  $b$   
 $x = 5.8 \cdot 10^{-4}$ ,  $d = .32 \text{ fm}$



$|1 - \Gamma(b)|^2$  -  
 probability not to interact  
 at given  $b$



*gg -N interaction seems close to BDR for  $Q^2 \sim 4 \text{ GeV}^2, x \sim 10^{-4}$*

Large gluon  
induced diffraction  
-observed - see below



*for  $Q^2 \sim 4 \text{ GeV}^2, x \sim 10^{-3}$  gg - Pb interaction at  $b=0$  is deep in BDR  $q\bar{q}$  - Pb interaction in BDR*

*for these x nuclear leading twist gluon shadowing effect is rather small*



Significant fractional energy losses and  $p_t$  broadening for partons propagating through black media (FS 01-03)

Suppression of the leading hadron production in pA scattering at large  $p_t$  comparable to the scale of Black disk regime at given energy (FS 01-06)



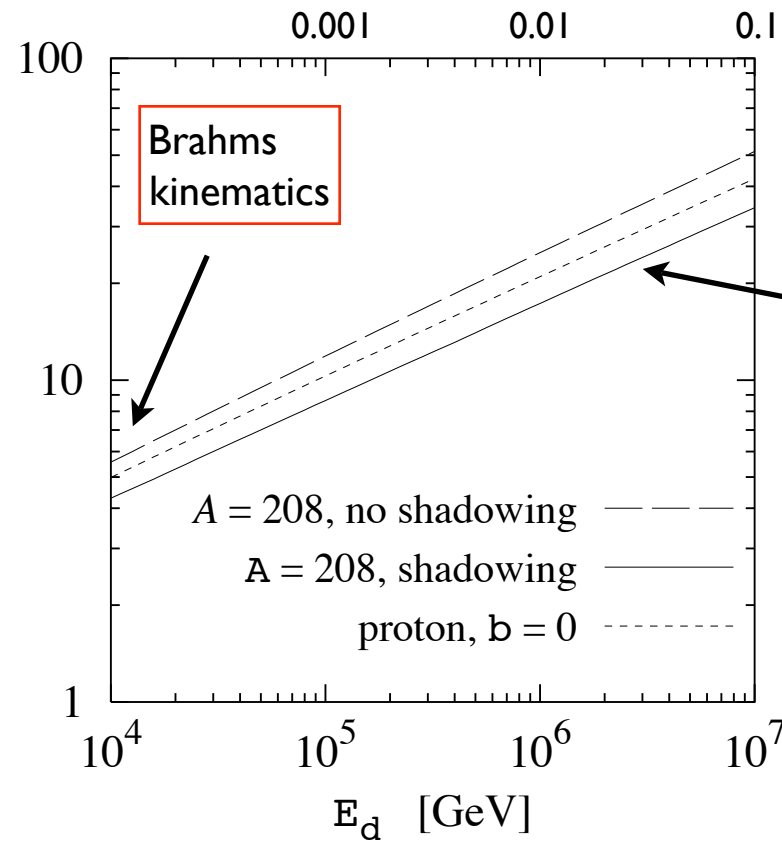
Natural explanation of the BRAHMS result at RHIC, the only one consistent with the STAR data on correlations

$$p_{t \text{ BDR}} \sim \frac{\pi}{2d}$$

where  $d$  is the minimal size of the  $gg$  ( $qq$ ) dipole for which  $\Gamma(b=0) \geq 1/2$  in LT

$$p_{t \text{ BDR}}^2(\text{gluon}) \approx 2p_{t \text{ BDR}}^2(\text{quark})$$

$x_F$  for  $pp$  at LHC



Gluon densities in nuclei and proton at  $b=0$  are very similar!!!

Difference is in the spread in  $b$



# One of fundamental questions:

How small color singlets (dipoles,...) propagate through nuclear media

*Intermediate energies* - hundred GeV (lab) - color transparency - observed at FNAL in  $\pi + A \rightarrow 2\text{jets} + A$ ,  $\gamma + A \rightarrow J/\psi + A$ ,

*High energies* -  $x_{\text{eff}} = Q^2_{\text{eff}}/s < 0.01$  - onset of color opacity regime both due to pQCD effects of LT gluon shadowing and proximity to black disk regime

⇒ strong screening of total cross section of dipole -nucleus scattering

$$\sigma_{tot}^{dA} \propto A \implies \sigma_{tot}^{dA} \propto A^{2/3}$$

*coherent photoproduction of  $J/\psi, \Upsilon$*

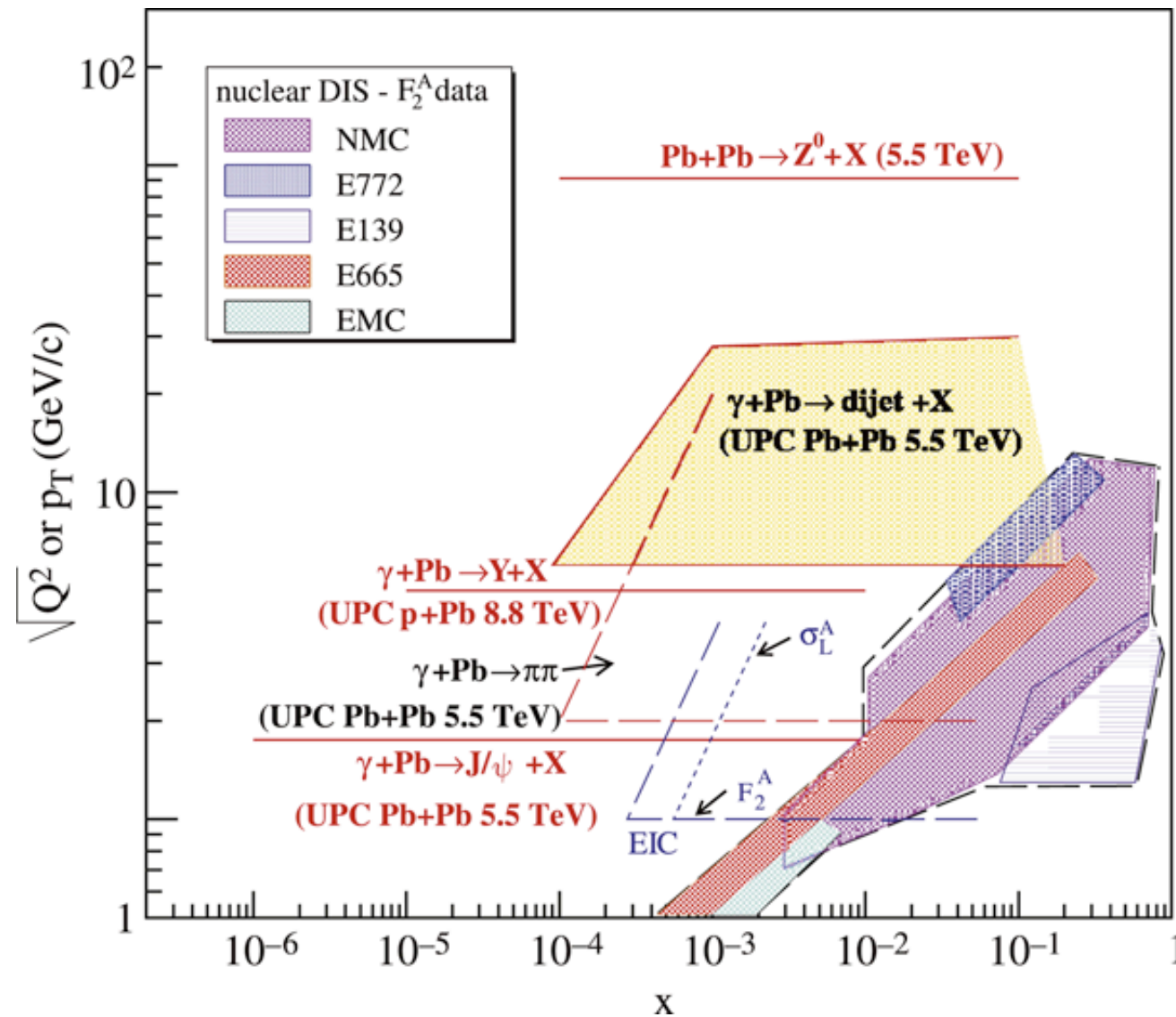
⇒ survival probability,  $P$ , for propagation through the nucleus center drops to zero only rim contributes:  $P \propto A \implies P \propto A^{1/3}$

*quasielastic photoproduction of  $J/\psi, \Upsilon$*

*large  $t$  rapidity gap photoproduction of light vector mesons*

*diffractive (rapidity gap between VM and A) photoproduction of  $J/\psi, \Upsilon$*

# Ultra-peripheral Collisions $\equiv$ UPC



The kinematic range in which UPCs at the LHC can probe gluons in protons and nuclei in quarkonium production, dijet and dihadron production. The  $Q$  value for typical gluon virtuality in exclusive quarkonium photoproduction is shown for  $J/\psi$  and  $\Upsilon$ . The transverse momentum of the jet or leading pion sets the scale for dijet and  $\pi\pi$  production respectively. For comparison, the kinematic ranges for  $J/\psi$  at RHIC,  $F_2^A$ ,  $\sigma_L^A$  and  $Z^0$  hadroproduction at the LHC are also shown.



What can be measured/discovered at LHC in UPC to follow up on HERA?

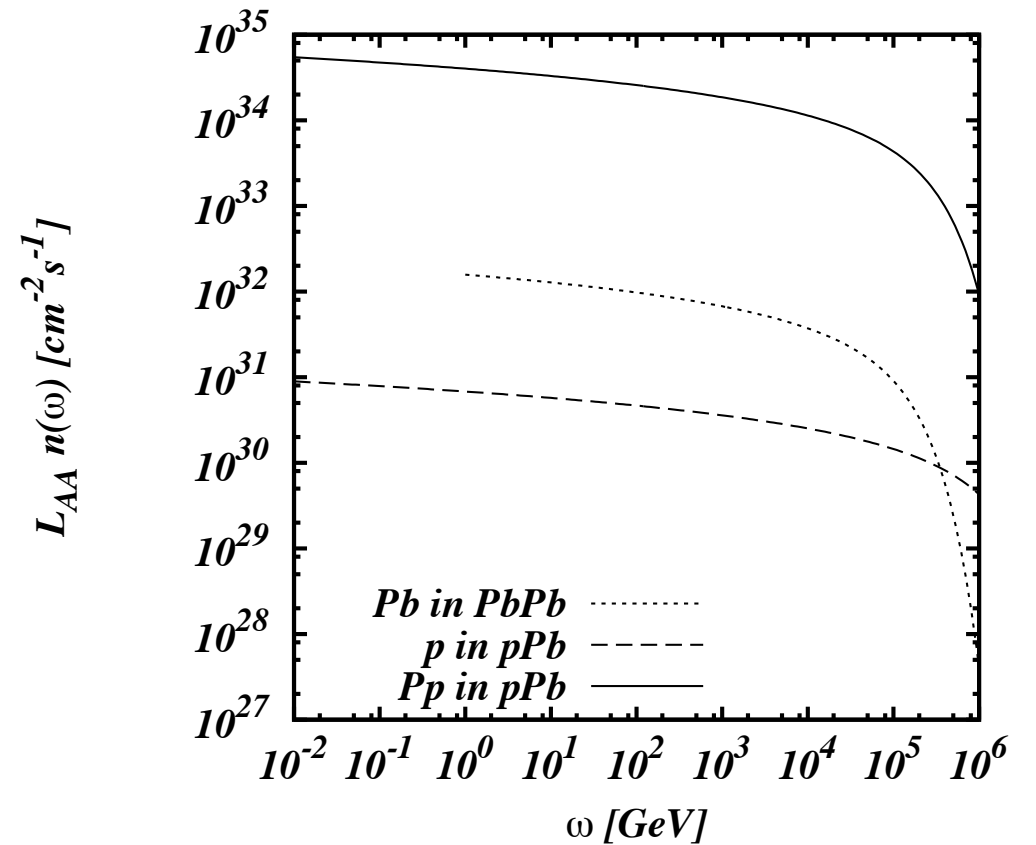
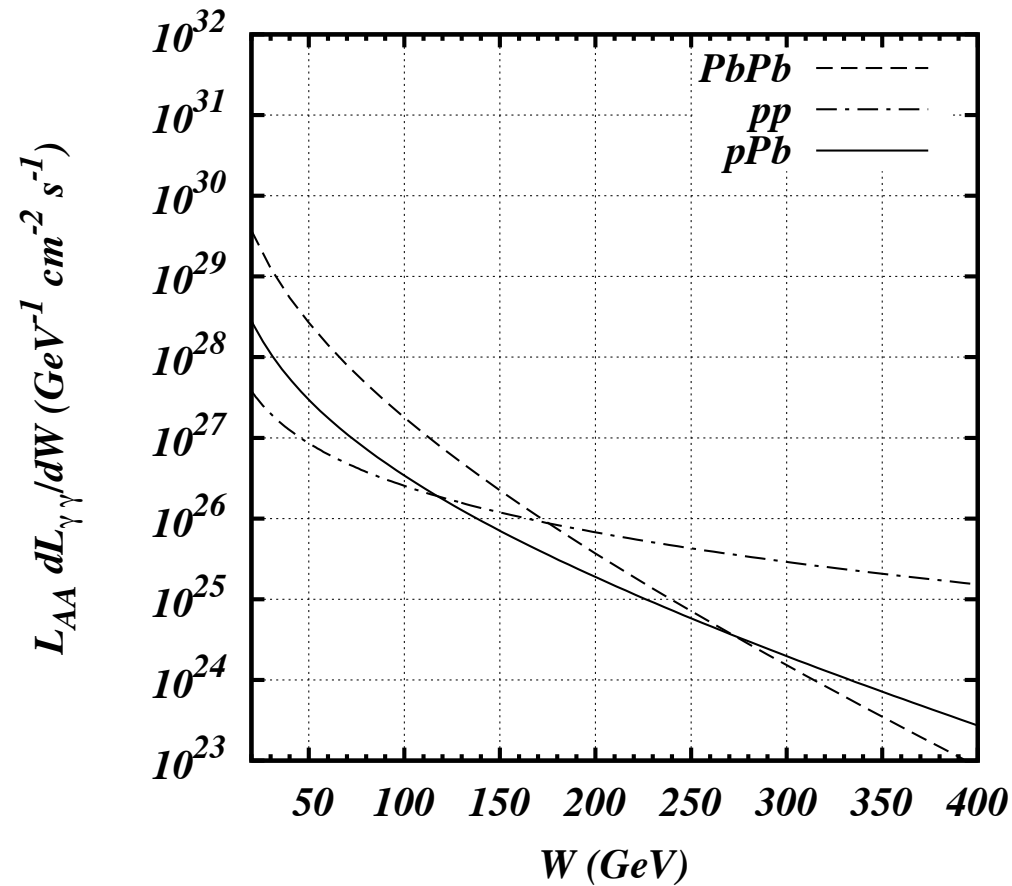


Trigger: One or both nuclei remain intact

Breakup of nuclei due to the Coulomb excitations are allowed (emission of few soft (in the nucleus rest frame) neutrons. Contribution of strong interactions due to nucleus-nucleus scattering at  $b \sim 2R_A$  is a small correction (weak  $A$ -dependence & small probability of diffraction). One can also study asymmetric UPC -  $pA, \& AA$

Counting rates are large up to

$$s_{eff}^{\gamma A}(LHC) \sim (1TeV)^2, \sim 10s_{max}, HERA(\gamma p)$$



(a) The effective  $\gamma A$  luminosity,  $L_{AB}n(\omega)$ , is shown for the cases where the photon is emitted from the proton ( $\gamma Pb$ ) and the ion ( $\gamma p$ ) as well as when the proton is emitted from the ion in a  $Pb+Pb$  collision ( $\gamma Pb@Pb+Pb$ ).

(b) The photon-photon luminosities,  $L_{AB}dL_{\gamma\gamma}/dW$ , are compared for  $pp$ ,  $pPb$  and  $Pb+Pb$  collisions at the LHC.

# Study of elastic dipole - nucleus scattering: exclusive vector meson production

$$\frac{d\sigma(AA \rightarrow VAA)}{dy} = N_\gamma(y)\sigma_{\gamma A \rightarrow VA}(y) + N_\gamma(-y)\sigma_{\gamma A \rightarrow VA}(-y).$$

$$\text{rapidity } y = \frac{1}{2} \ln \frac{E_V - p_3^V}{E_V + p_3^V} = \ln \frac{2k}{m_V}.$$

The flux of the equivalent photons  $N_\gamma(y)$  is

$$N(y) = \frac{Z^2\alpha}{\pi^2} \int d^2b \mathbf{R}_{AA}(\vec{b}) \frac{1}{b^2} X^2 \left[ K_1^2(X) + \frac{1}{\gamma} K_0^2(X) \right].$$

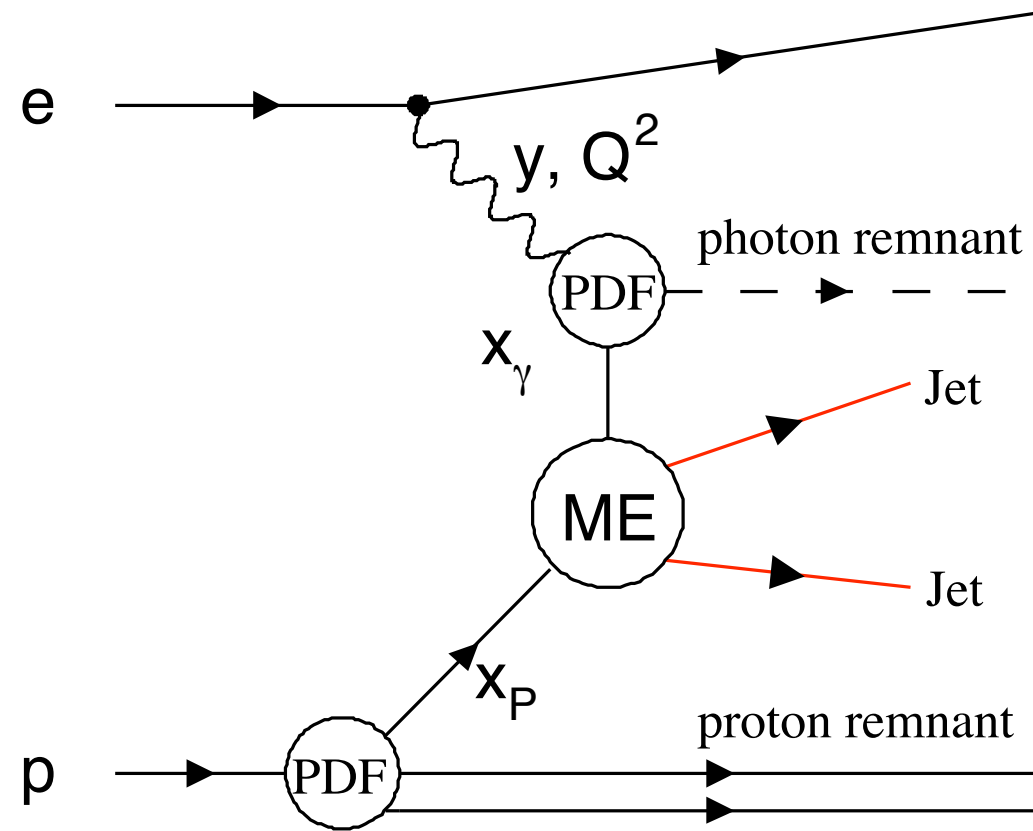
$K_0(X), K_1(X)$  – modified Bessel functions with argument  $X = \frac{bm_V e^y}{2\gamma}$ ,  $\gamma$  is Lorentz factor and  $\vec{b}$  is the impact parameter.  $R_{AA}(\vec{b}) = 1 - |1 - \Gamma(\vec{b})|^2$

$\rho, J/\psi$  photoproduction off nuclei is studied at RHIC - talks of Boris Grube & Sebastian White - reasonable agreement with the theory



# Another lesson from HERA

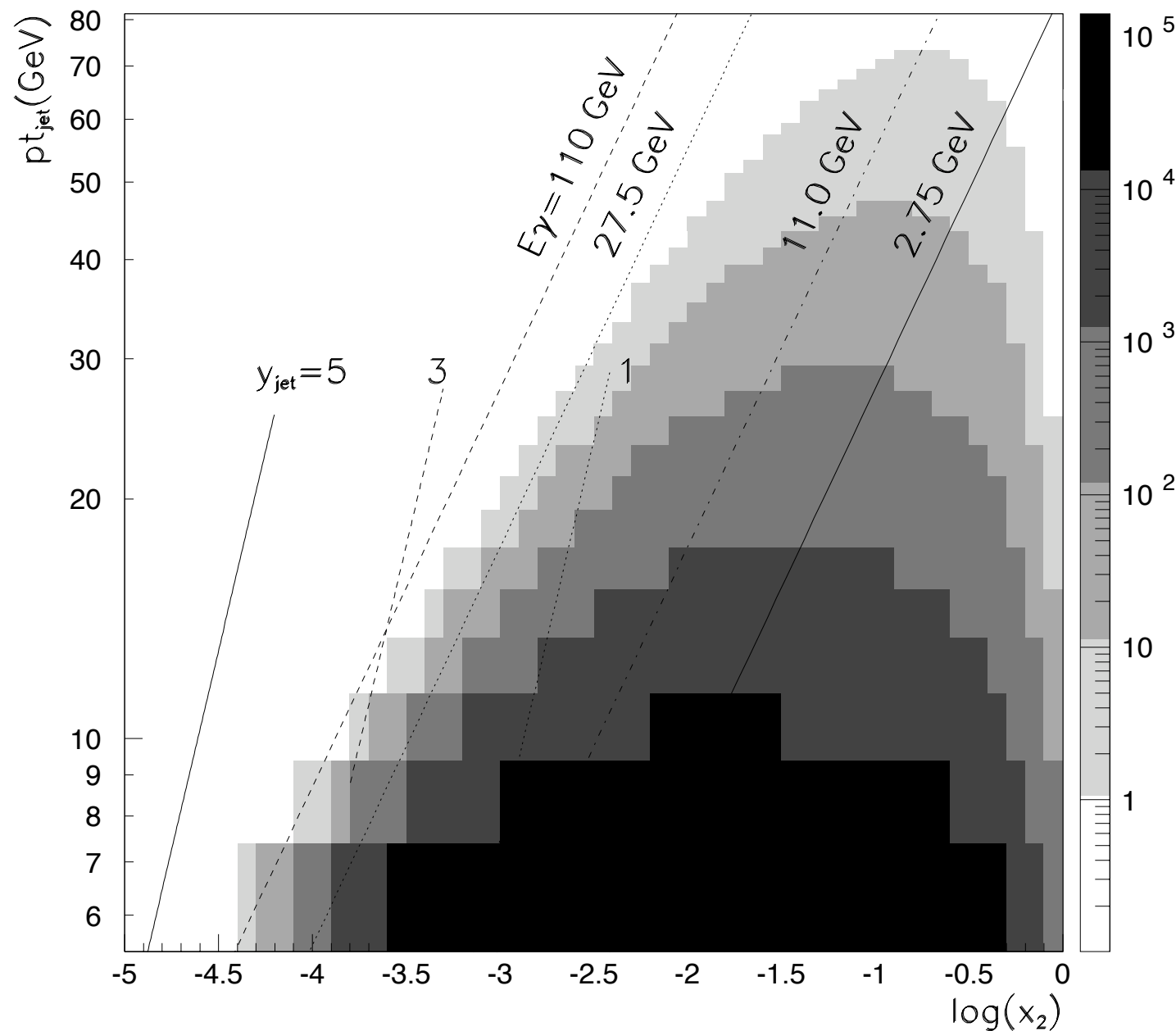
Real photon was effectively used for the QCD studies



*Schematic view of dijet production in ep scattering studied at HERA*

*$x_\gamma$  and  $x_p$  are light cone fractions of partons of photon and proton*

Still open problems - see Tuesday talks



Expected rate of dijet photoproduction for a 1 month LHC Pb+Pb run at  $0.4 \times 10^{27} \text{ cm}^{-2} \text{ s}^{-1}$ . Rates are counts per bin of  $0.25 \times x_2$  and  $2 \text{ GeV}/c$  in  $p_T$ .

Many more interesting questions to study - like inclusive leading pion A-dependence as a function of  $p_T$ , associated multiplicity at different rapidities,...

R.Vogt, S.White, MS

More details in S.White talk

# Nonlinear effects: AA UPC at LHC vs HERA and eRHIC

The parameter to compare is:

gluon density/unit area \* strength of interaction

$$\frac{C\alpha_s(Q^2)xG(x, Q^2)}{Q^2 \text{ "area"}}$$

where  $C_g \approx 9/4C_q$

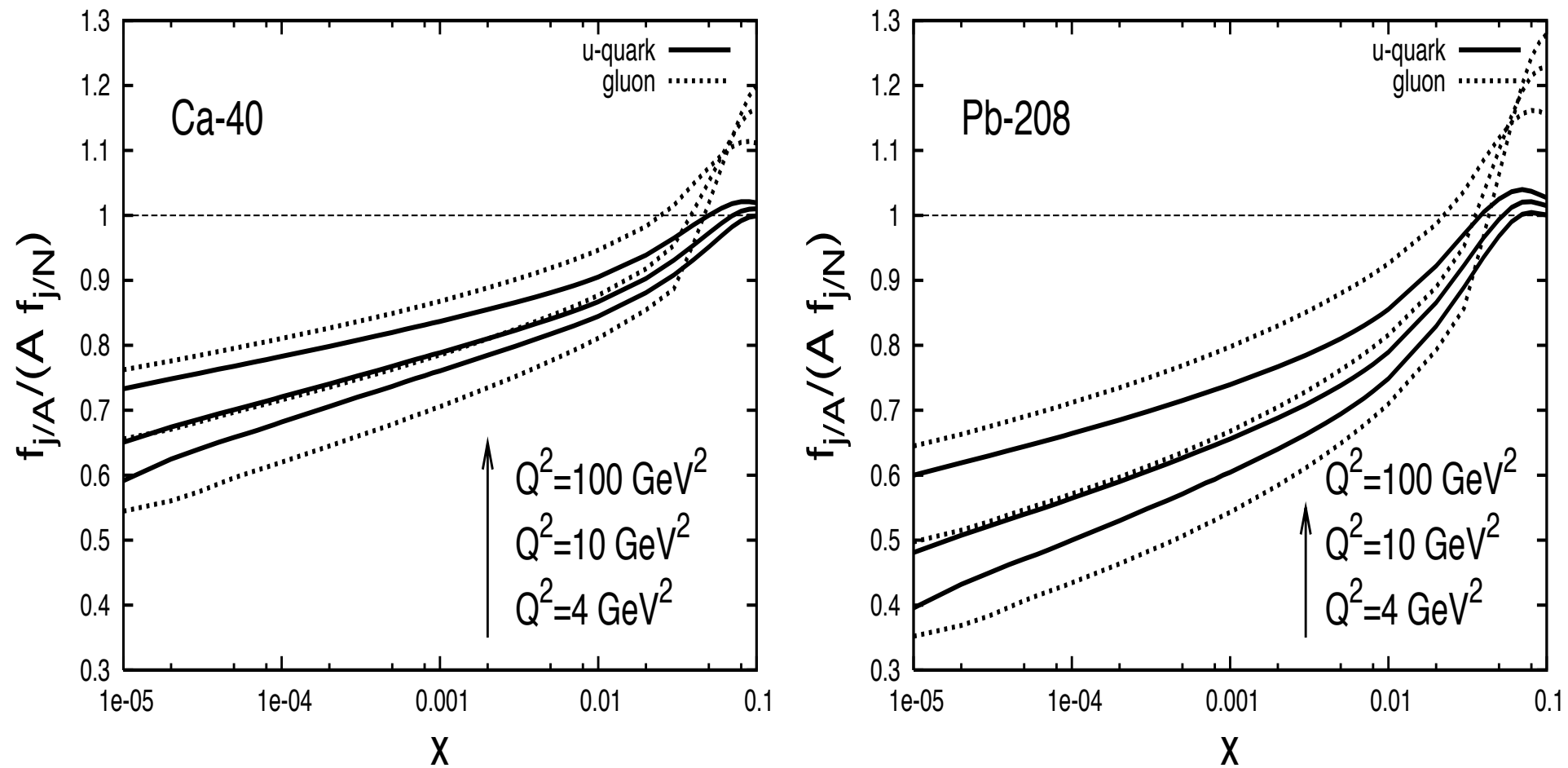
LHC vs ep HERA  $\frac{(9/4)A^{1/3}\alpha_s(p_T^2)xG_N(x \sim 5 \cdot 10^{-5}, p_T^2)/p_T^2}{\alpha_s(Q^2)xG_N(x \sim 10^{-4}, Q^2)/Q^2} \sim 3$

for central  $\gamma A$  collisions (with no centrality trigger the gain is a factor of two smaller). *A factor of 3 gain = change in x by a factor ~100.*

LHC vs eRHIC:  $eA$  at  $Q=2, x=10^{-3}$  the gain is a factor of 1.5

Will be possible to study energy dependence of the dijet cross section in the x range between  $10^{-2}$  and  $10^{-4}$  and check whether taming of the increase is happening at the smallest x.

Are significant nuclear effects expected in the UPC AA kinematics at LHC? The leading twist theory FS 98 based on AGK cutting rules and Collins factorization theorem for diffraction indicates that effects are likely to be significant (will briefly discuss later)



Shadowing for nuclear pdfs using HI 2006 diffractive pdfs (difference between in the shadowing results for HI B and jet fits is small)



Expected suppression is large enough in the UPC kinematics (a factor of two) to be measured. Can be further enhanced with centrality trigger.

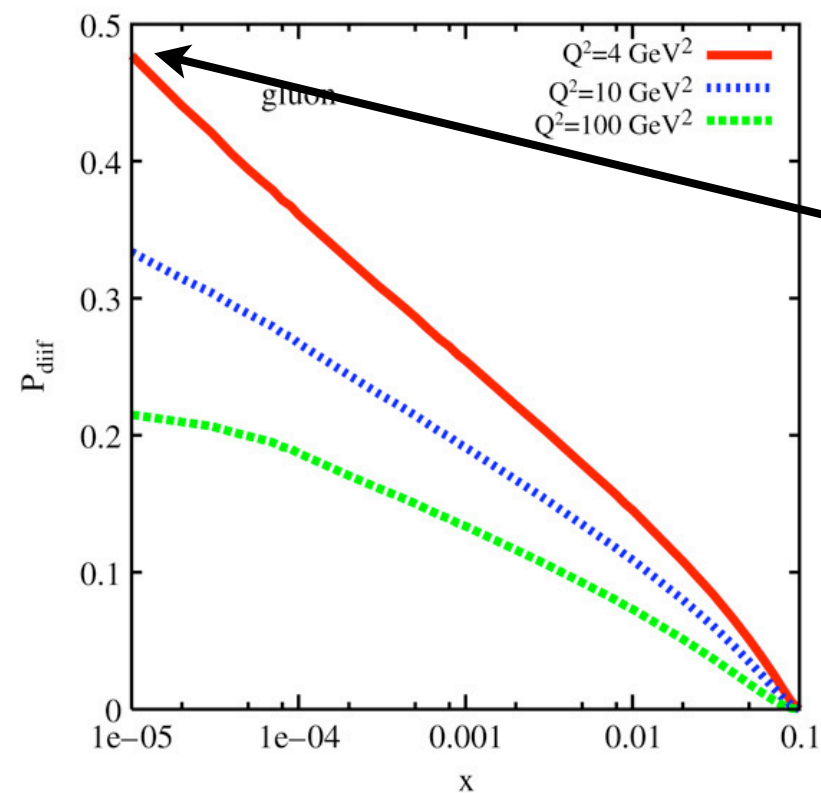
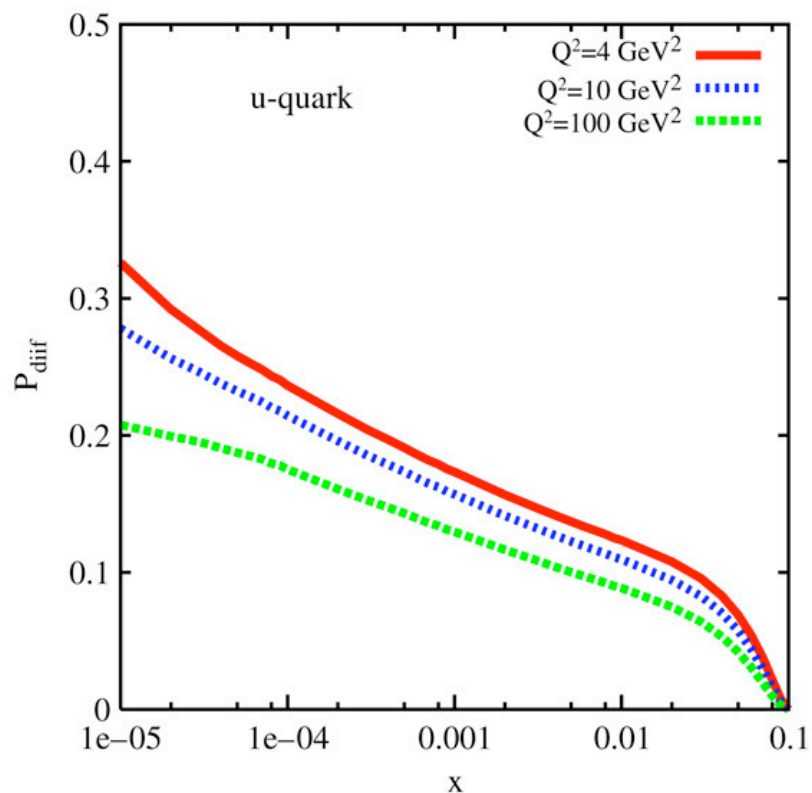
**Another critical measurement is hard diffraction:**

$\gamma A \rightarrow jet_1 + jet_2 + X + A$  for direct photon:  $\beta_\gamma \approx 1$

In the black disk regime

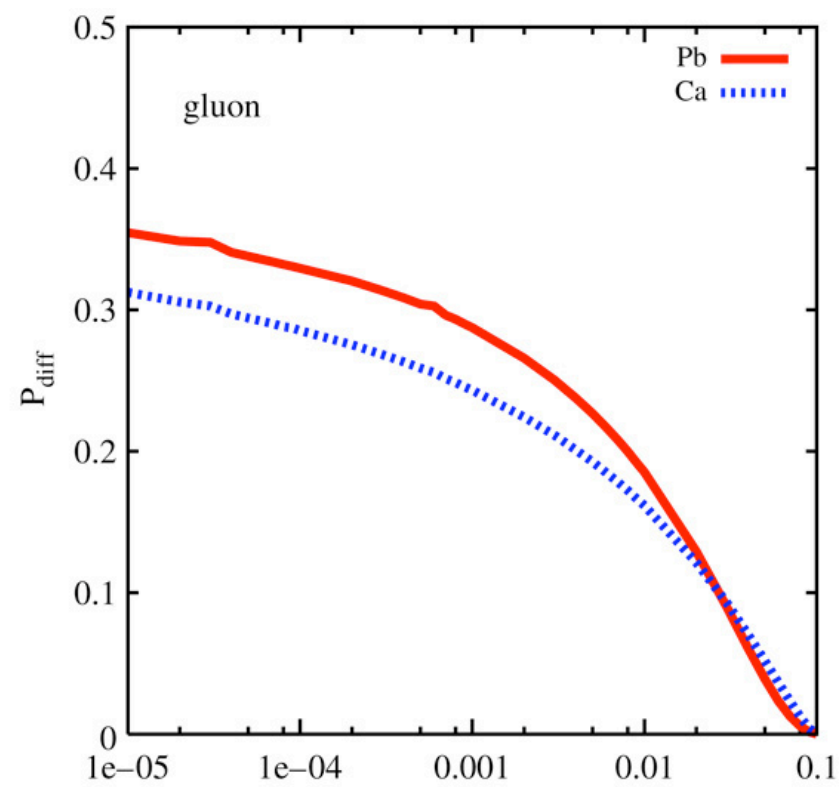
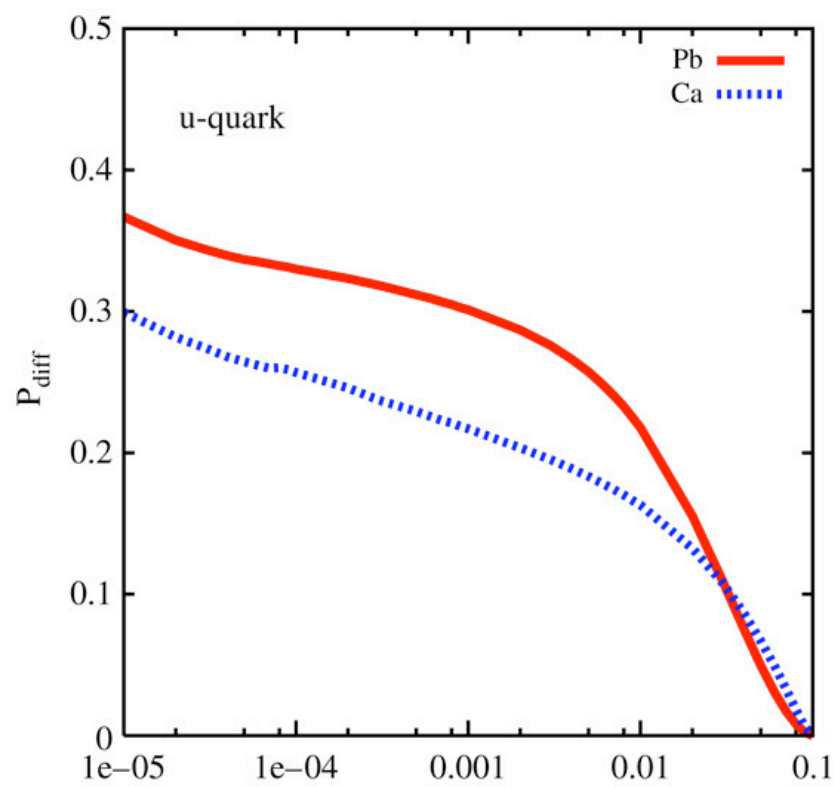
$$\frac{\sigma(\gamma A \rightarrow jet_1 + jet_2 + X + A)}{\sigma(\gamma A \rightarrow jet_1 + jet_2 + X)} \approx 0.5$$

Nuclear diffractive pdfs were calculated by Guzey et al 03 in the same approximations as LT nuclear pdf's



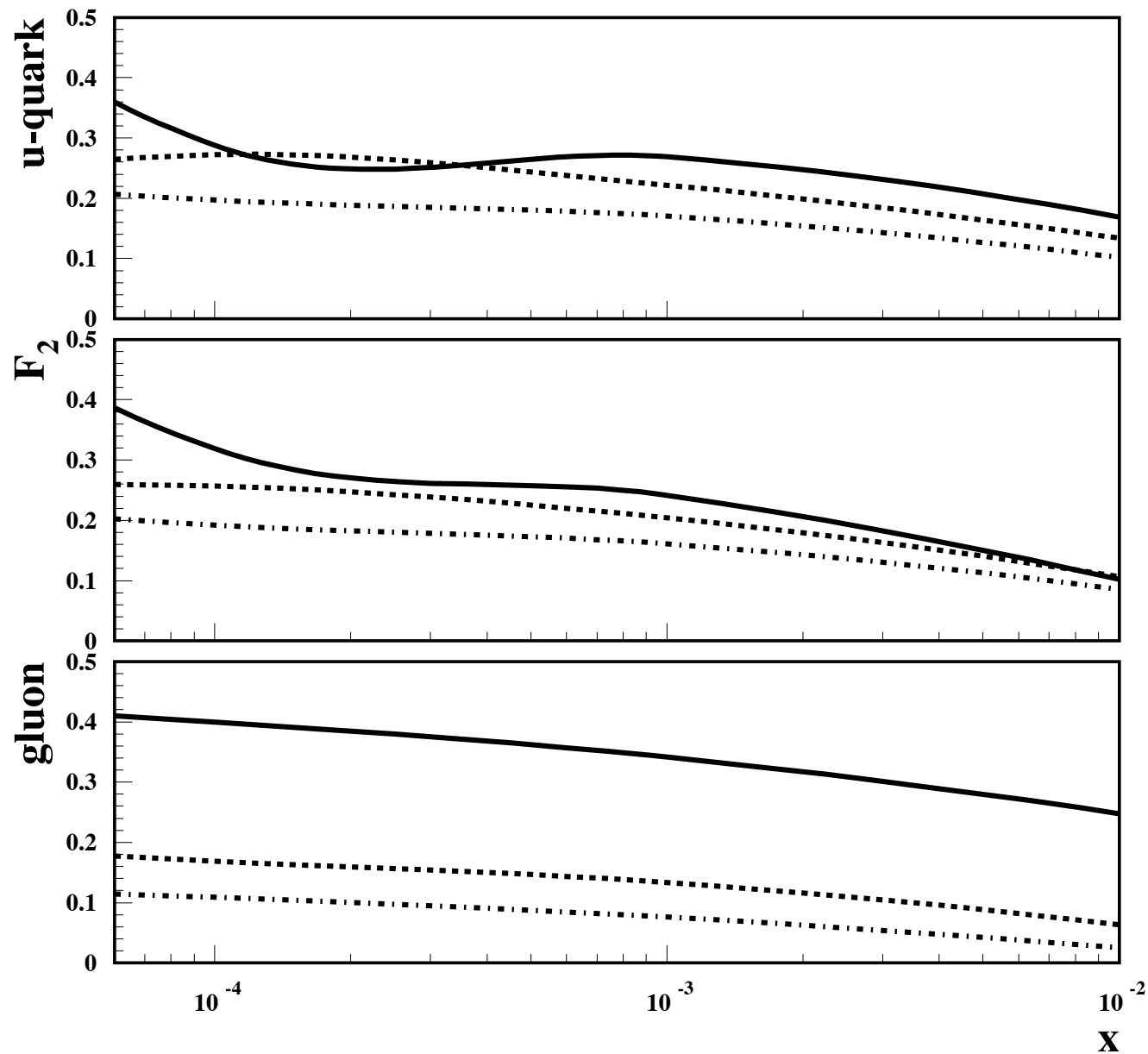
**Black limit**

The probability of hard diffraction on the nucleon,  $P_{j \text{ diff}}$  as a function of  $x$  and  $Q^2$  for u quarks (left) and gluons (right).



The probability of hard diffraction,  $P_{j \text{ diff}}$  as a function of  $x$  and  $Q^2$  for u quarks (left) and gluons (right) for  $Q^2=4 \text{ GeV}^2$

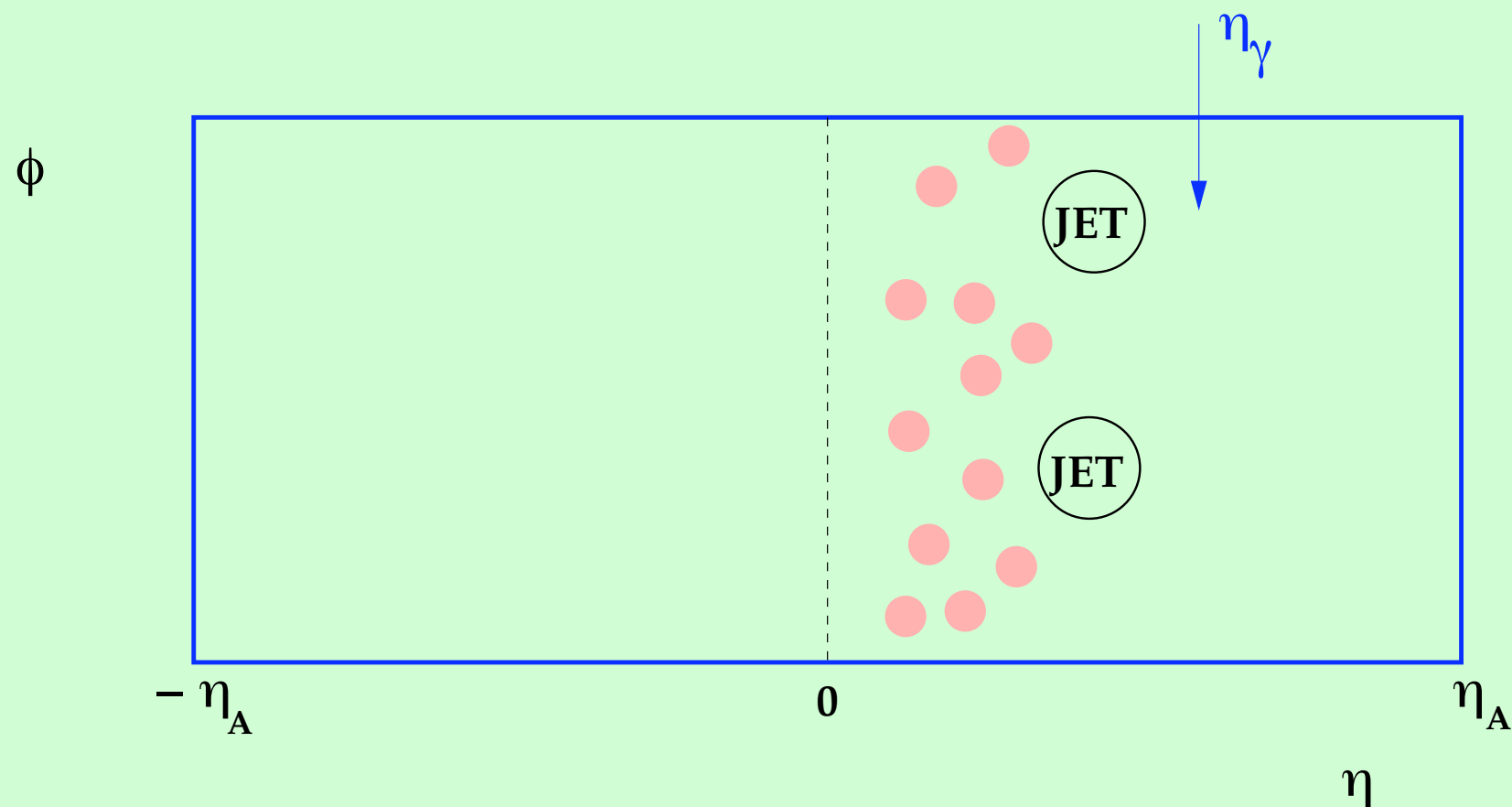
# Probability of diffraction remains large up to very large Q



*Proximity of the hard interactions with nuclei to BBL leads also to a large probability of diffractive events in nuclei (larger than in the proton). Results of the calculation within the leading twist model (Guzey, et al, 03) are shown the ratios  $f_{j/A}^{D(2)} / f_{j/A}$  for the  $u$ -quarks and gluons and NLO  $F_{2A}^{D(2)} / F_{2A}$  for  $^{208}\text{Pb}$  at  $Q=2, 10, 100 \text{ GeV}$ .*



In AA scattering it will be possible to measure gluon nuclear diffractive pdfs (or at least rapidity gap probabilities) in most of the small  $x$  kinematic range where measurements of nuclear gluon pdfs will be feasible. The key element is the possibility to use the direct photon mechanism to determine which of the nuclei has emitted the photon



UPC induced direct photon hard diffraction:  $AA \rightarrow AA + 2\text{jets} + X$



# Studies of exclusive photoproduction processes:

## *Hard physics:*



Onium production



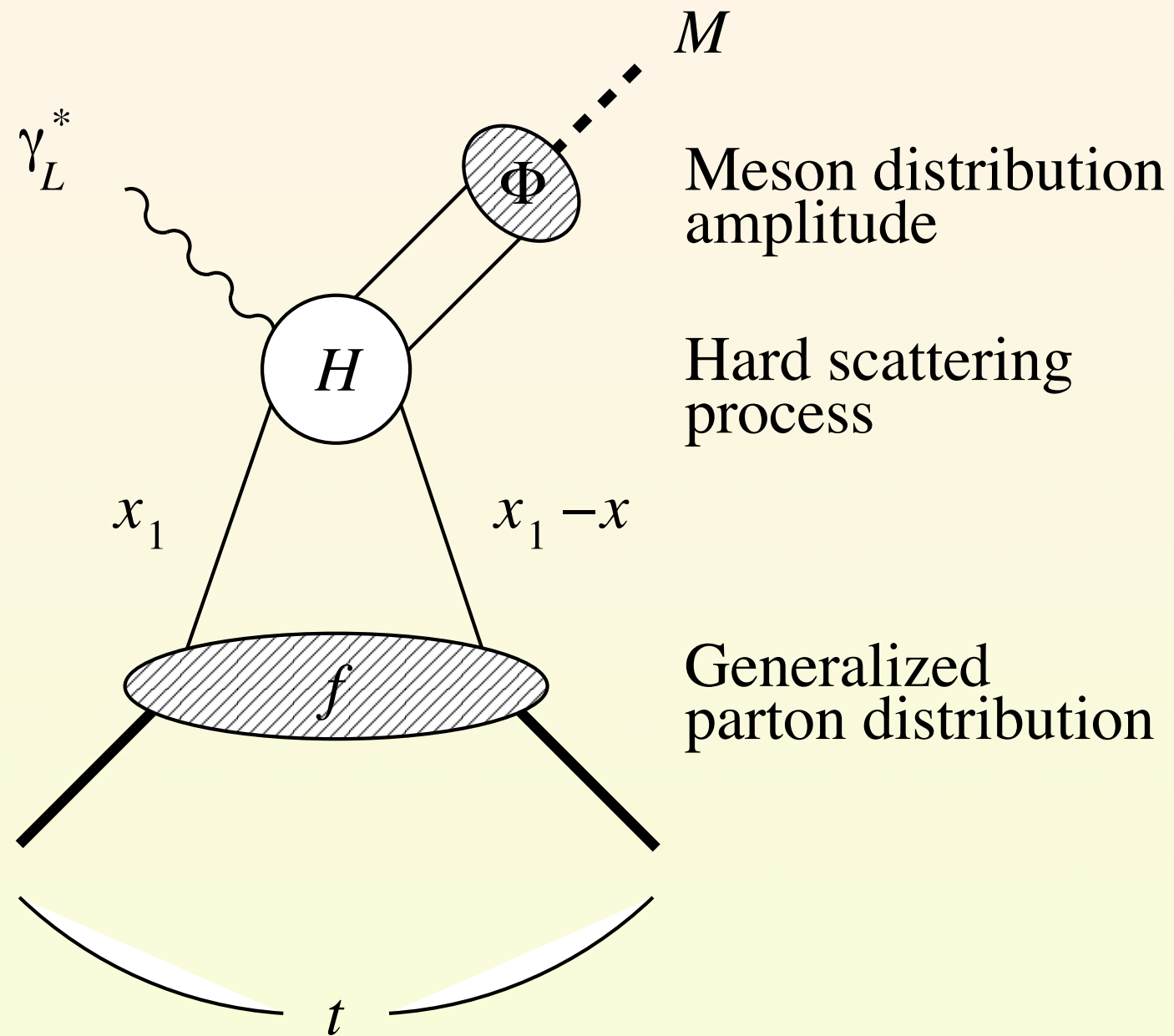
Diffraction into two, three jets

## *Soft (Pomeron) physics:*



Energy dependence of production of  $\rho, \varphi$ -mesons

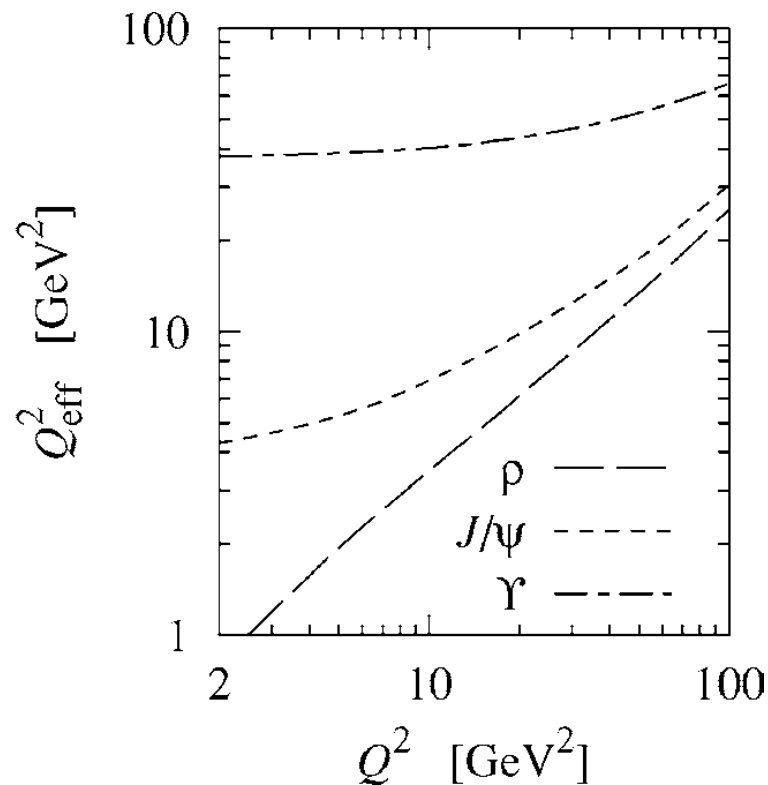
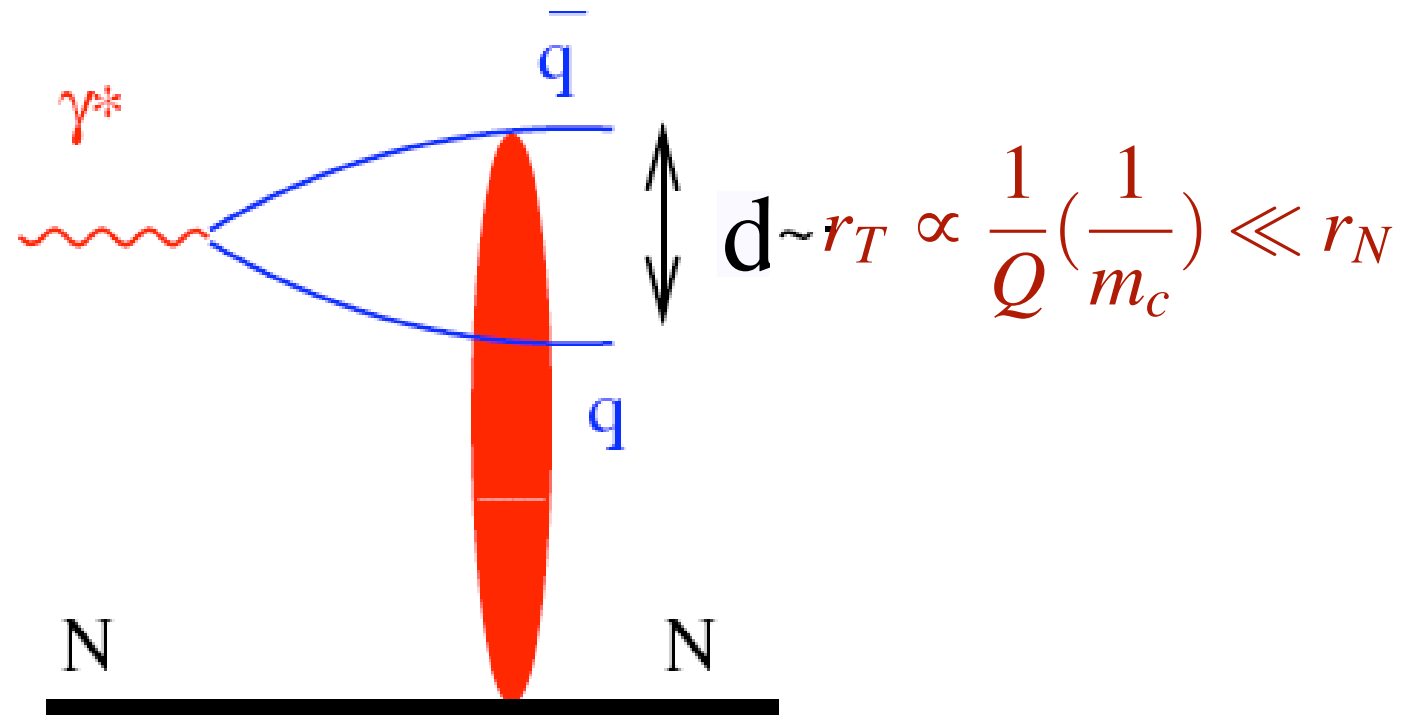
# QCD factorization theorem for DIS exclusive processes (Brodsky, Frankfurt, Gunion, Mueller, MS 94 - vector mesons, small $x$ ; general case Collins, Frankfurt, MS 97)



# Vector meson production

## Models:

- dipole model - effectively LO with choice of scale based on transverse size of the dipole and large HT effects due to comparable transverse sizes of  $\gamma^*_{L}$  and VM  $q\bar{q}$  dipoles

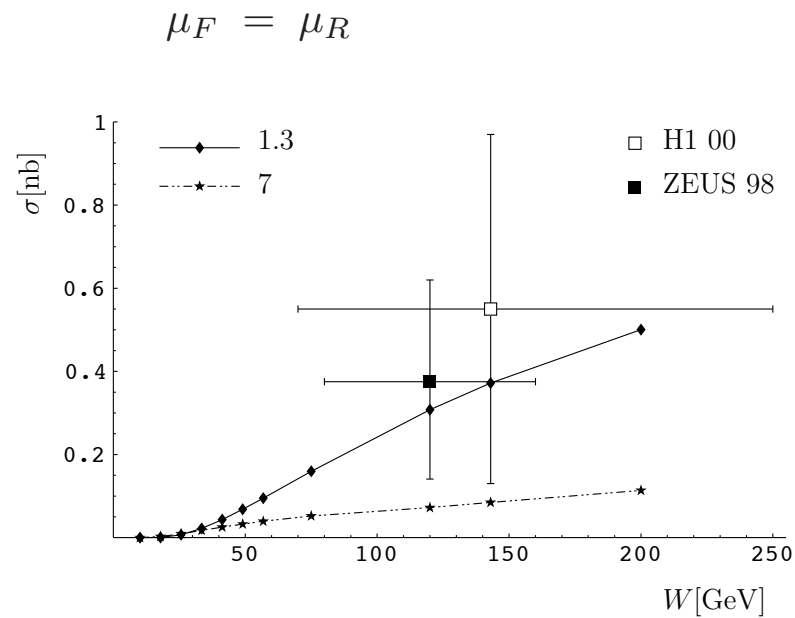


Frankfurt, Koepf, MS 95-97

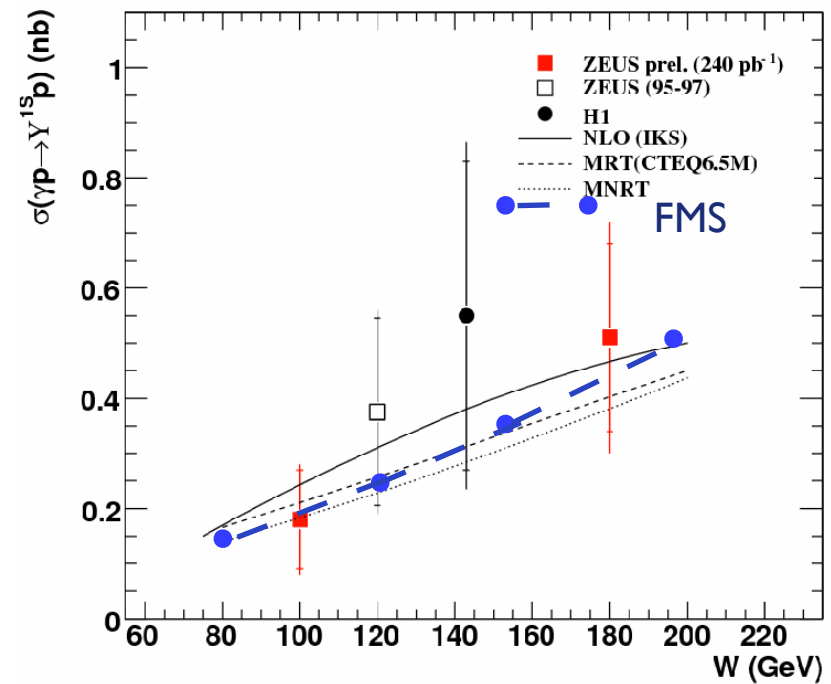
- NLO approach Ivanov, Szymanowski et al 04-06  
open questions - energy conservation and related issues with gauge invariance. treatment of the meson wave function

# Upsilon production

Ivanov et al (IKS)



Strong dependence of  
NLO result on  $\mu_R$



New ZEUS 240 pb<sup>-1</sup> two data points

NLO calculations done by Ivanov, Krasnikov,  
Szymanowski (IKS)[hep-ph/0412235]

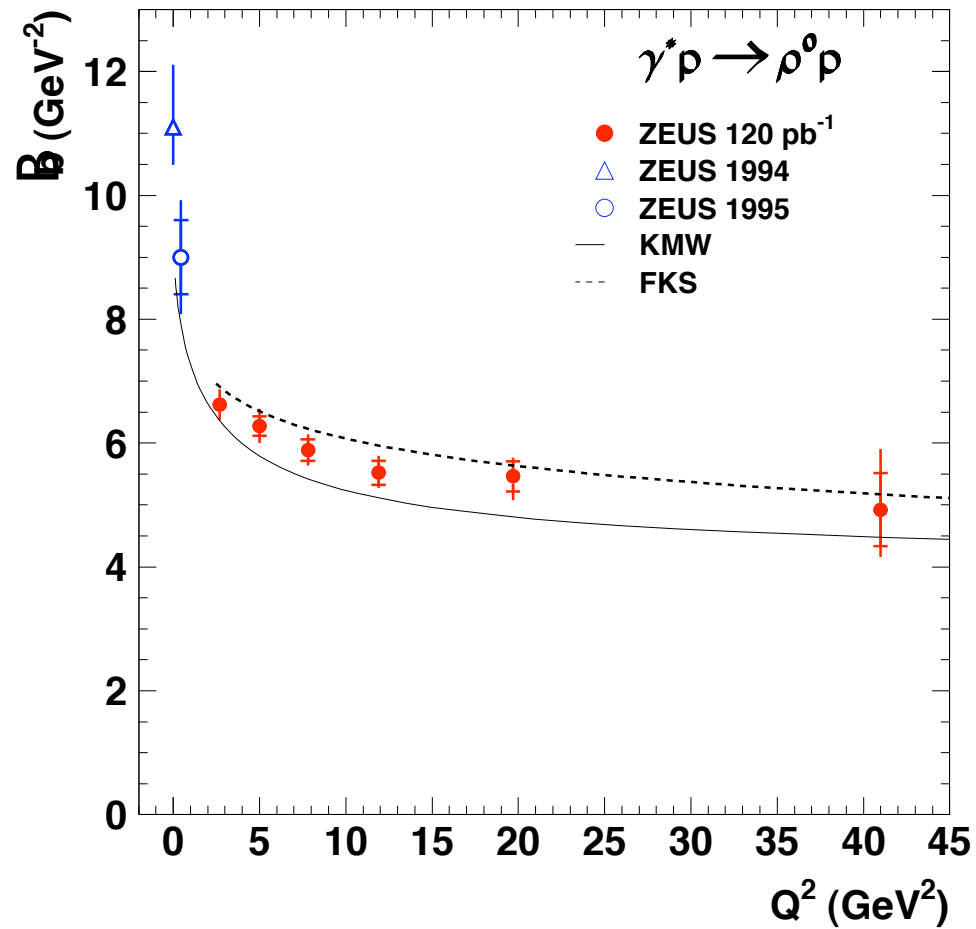
MRT – Martin, Ryskin, Teubner, based on  
CTEQ6.5M gluon.

MNRT – Martin, Nockles, Ryskin, Teubner, based  
on diffractive J/Ψ data alone.

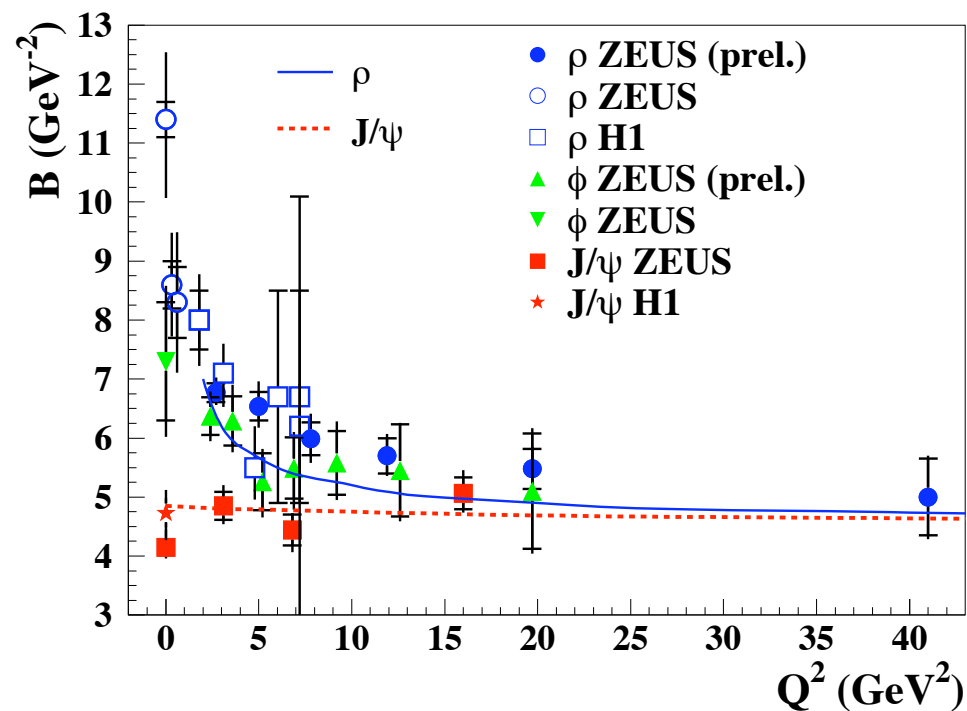
FMS - Frankfurt, McDermott, Strikman 98



# ZEUS



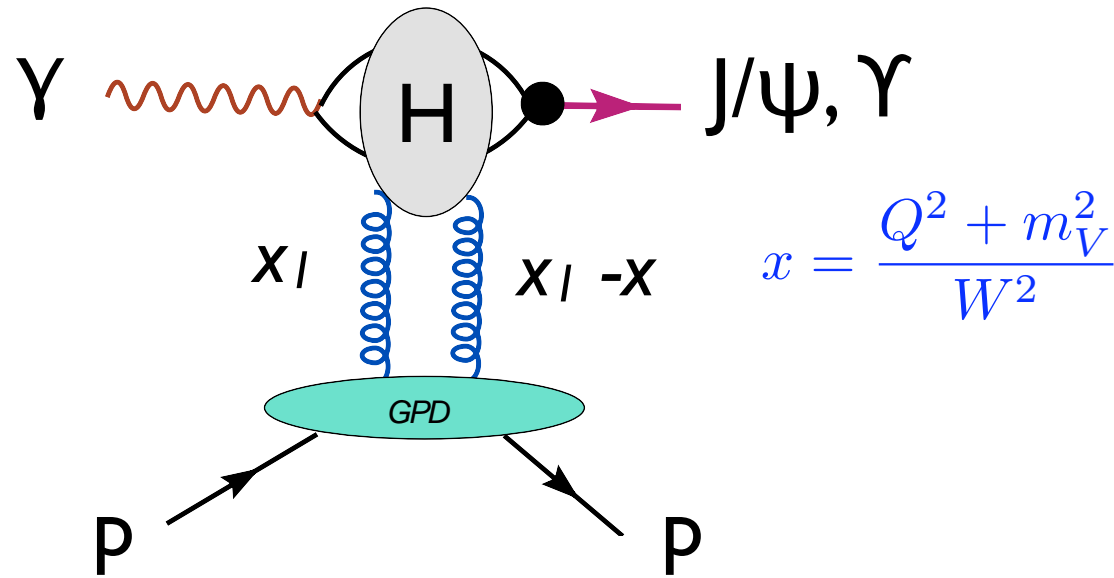
Drop of  $B$  is well reproduced by dipole models (in case of FKS actually a prediction)



Convergence of the t-slopes,  $B$  of  $\rho$ -meson electroproduction to the slope of  $J/\psi$  photo(electro) production.

⇒ Transverse distribution of gluons can be extracted from  $\gamma + p \rightarrow J/\psi + N$

- ⇒ Presence of small size  $q\bar{q}$  Fock components in light mesons is unambiguously established
- ⇒ At transverse separations  $d \leq 0.3$  fm pQCD reasonably describes “small  $q\bar{q}$  - dipole”- nucleon interaction for  $10^{-4} < x < 10^{-2}$
- ⇒ Color transparency is established for the interaction of small dipoles with nucleons and with nuclei (for  $x \sim 10^{-2}$  )



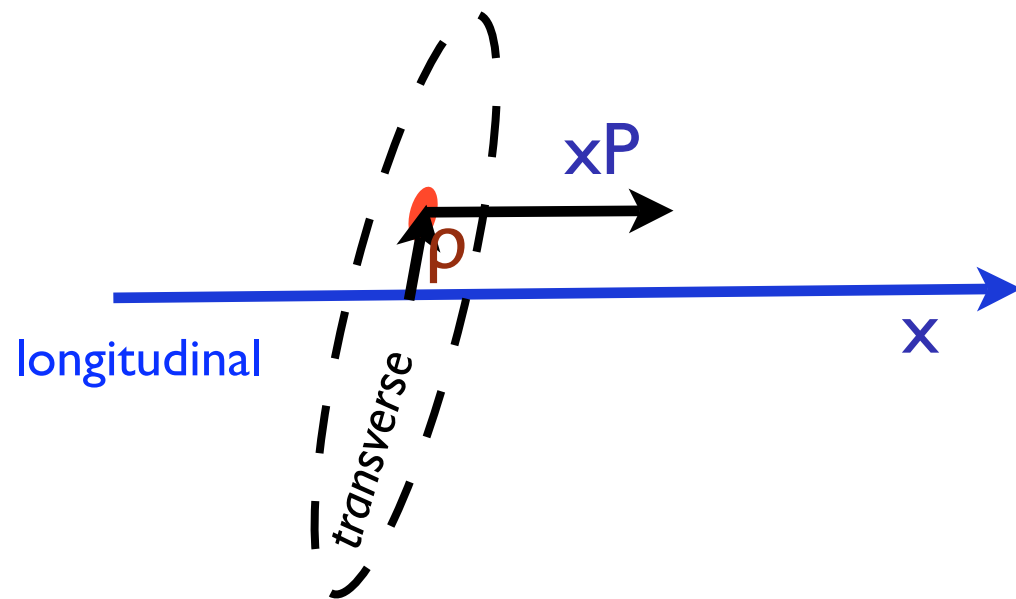
In LT limit  $x_1 - x \ll x_1$

however due to DGLAP evolution skewed GPD kinematics for large  $Q$  probes diagonal GPD at  $Q_0$  scale

$$A(\gamma^* + p \rightarrow \text{"Onium"} + p) \propto G(x_1, x_1 - x, t)$$

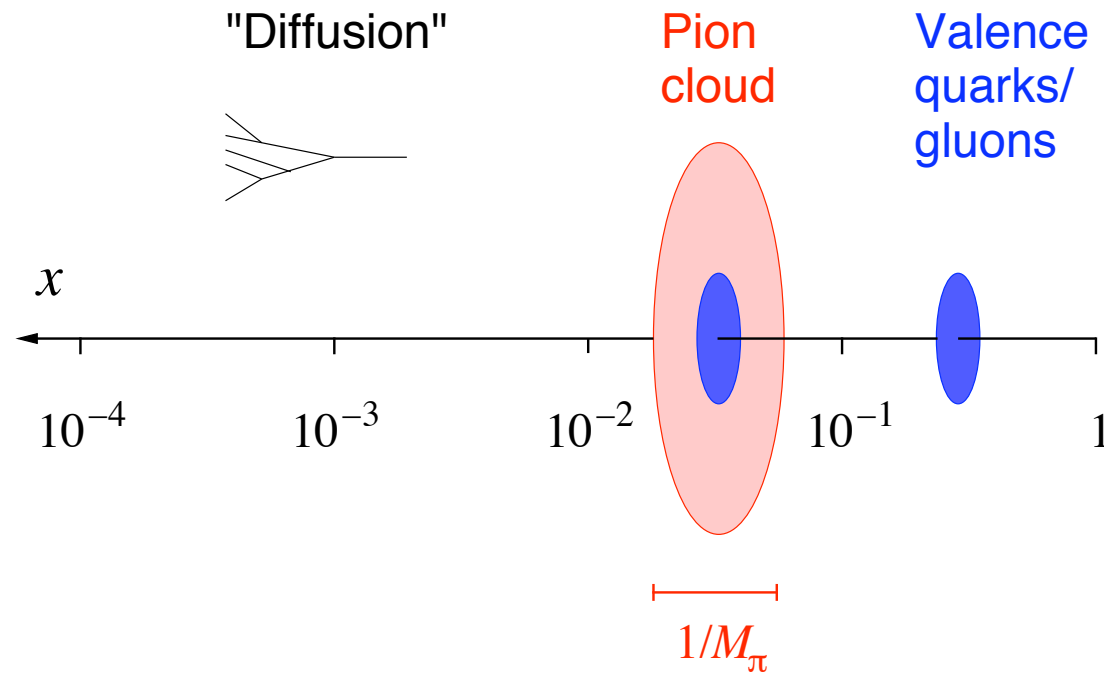
$$G(x, x, t) \equiv G(x, t) = \int d^2 \rho e^{-i \vec{\Delta}_\perp \rho} G(x, \rho)$$

← — transverse spatial distribution of gluons



$$\int d^2 \rho G(x, \rho) = G(x) \quad \text{total gluon density}$$

● Gluonic transverse size - x dependence



Gluon transverse size decreases with increase of x

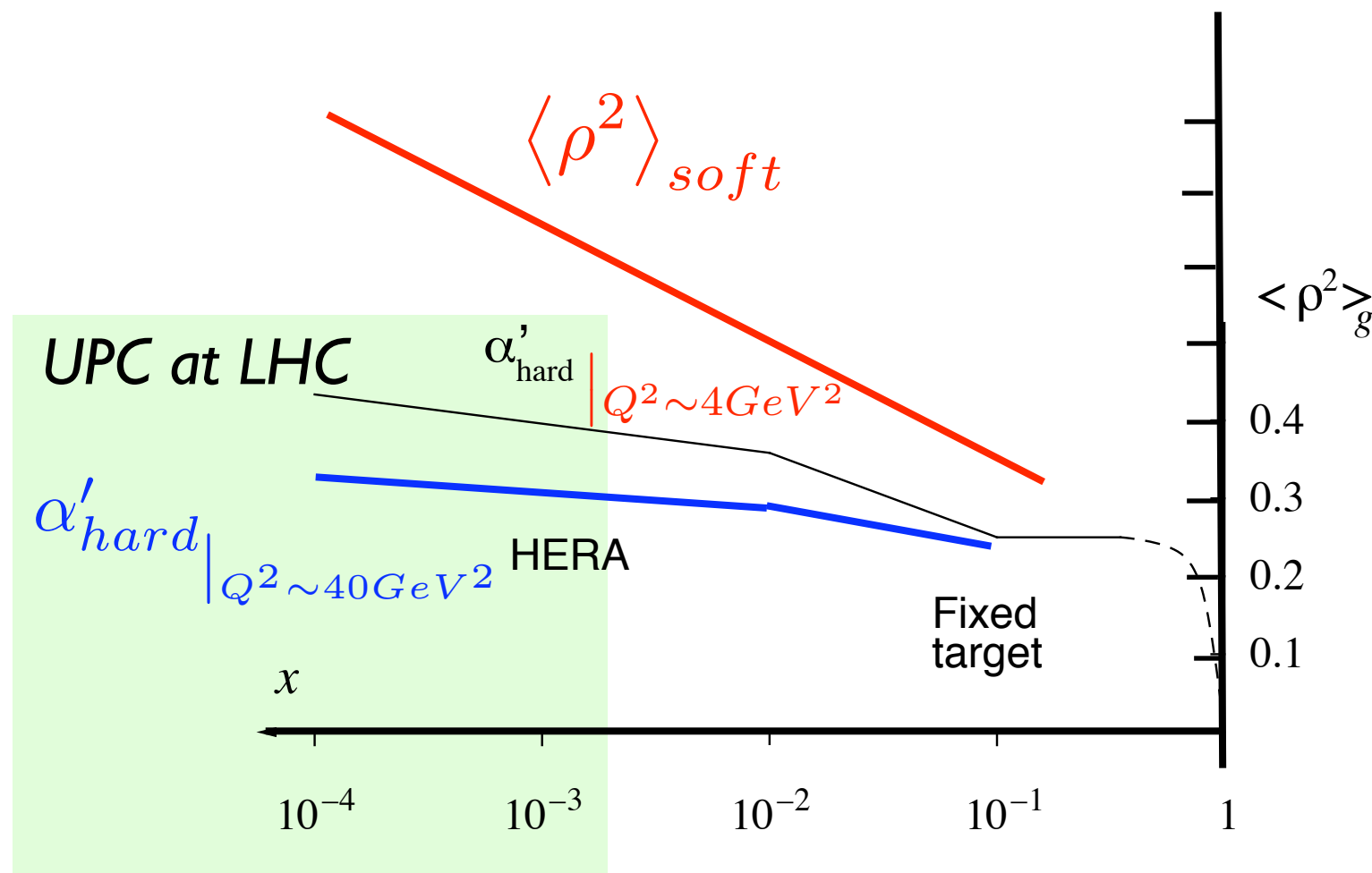
Pion cloud contributes for  $x < M_\pi/M_N$  [MS & C.Weiss 03]

Transverse size of large x partons is much smaller than the transverse range of soft strong interactions

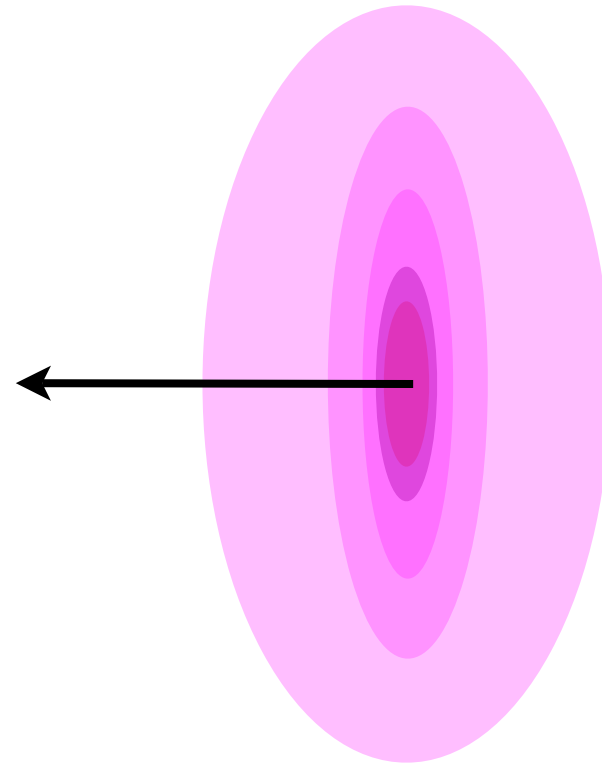
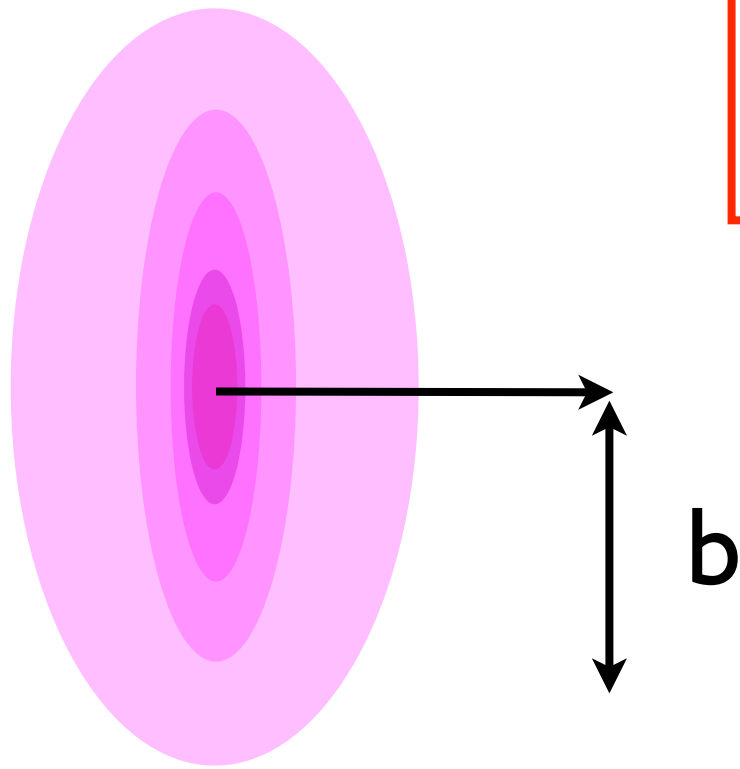
$$\langle \rho^2(x > 10^{-2}) \rangle \ll R_{soft}^2$$



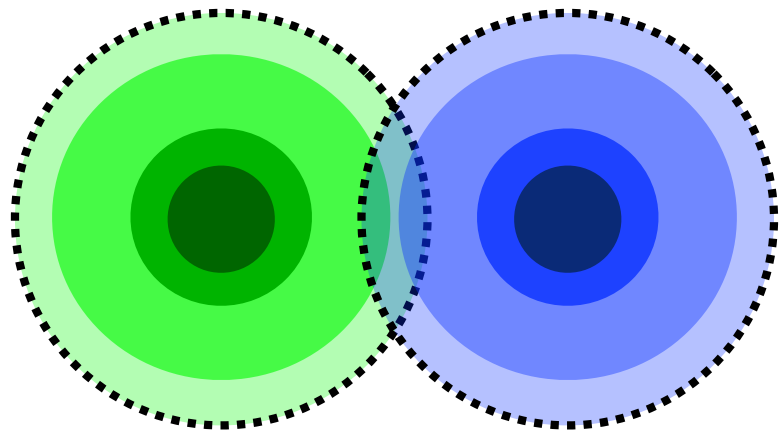
Two scale picture



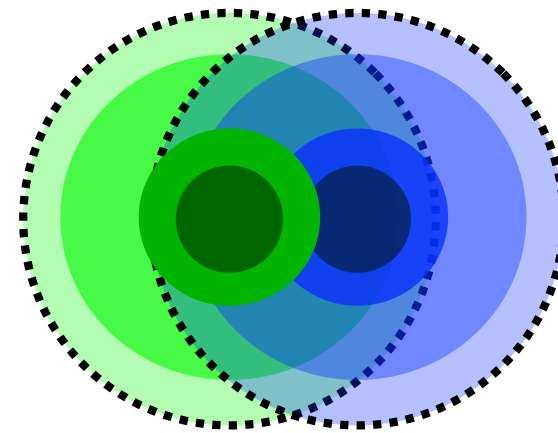
Central collision of two high energy protons  
which dominates in production on SUSY, Higgs



Valence quarks/gluons of the protons are interacting with probability  $\sim$  **one**, losing energy and getting large transverse momenta growing with energy. Soft interactions are suppressed - **minimal scale/virtuality** of strong interaction is few GeV and growing with energy.



Peripheral pp



Central pp

*Precision measurement of  $t$ -dependence of onium photoproduction off proton  
crucial for reliable modeling of these phenomena*



## □ *Exclusive onium production*

AA (pp) collisions - maximal  $W^2$  one can effectively probe is  $W^2 = 2m_V E_N$  due to dominance of photons with smaller energy

$$x_{\min} \equiv m_V^2/W^2 = m_V/2 E_N \quad \text{separate problem - how to trigger for } y=0$$

At LHC  $x_{\min}(J/\psi) = .0005$ ,  $x_{\min}(\Upsilon) = .002$  for AA  
and a factor of 2.5 smaller for pp

*The nuclear Coulomb induced dissociation occurs at small impact parameters. At the same time in such events the photon spectrum is harder. (Can be used enhanced contribution of hard photons) Baltz, Klein Nystrand, 02. (Price a factor of 10 reduction in counting rate). Allows to extend measurements for  $J/\psi$  case to  $y \sim 2$ ,  $\rightarrow x \sim 10^{-5}$ .*

Numerical study is still needed

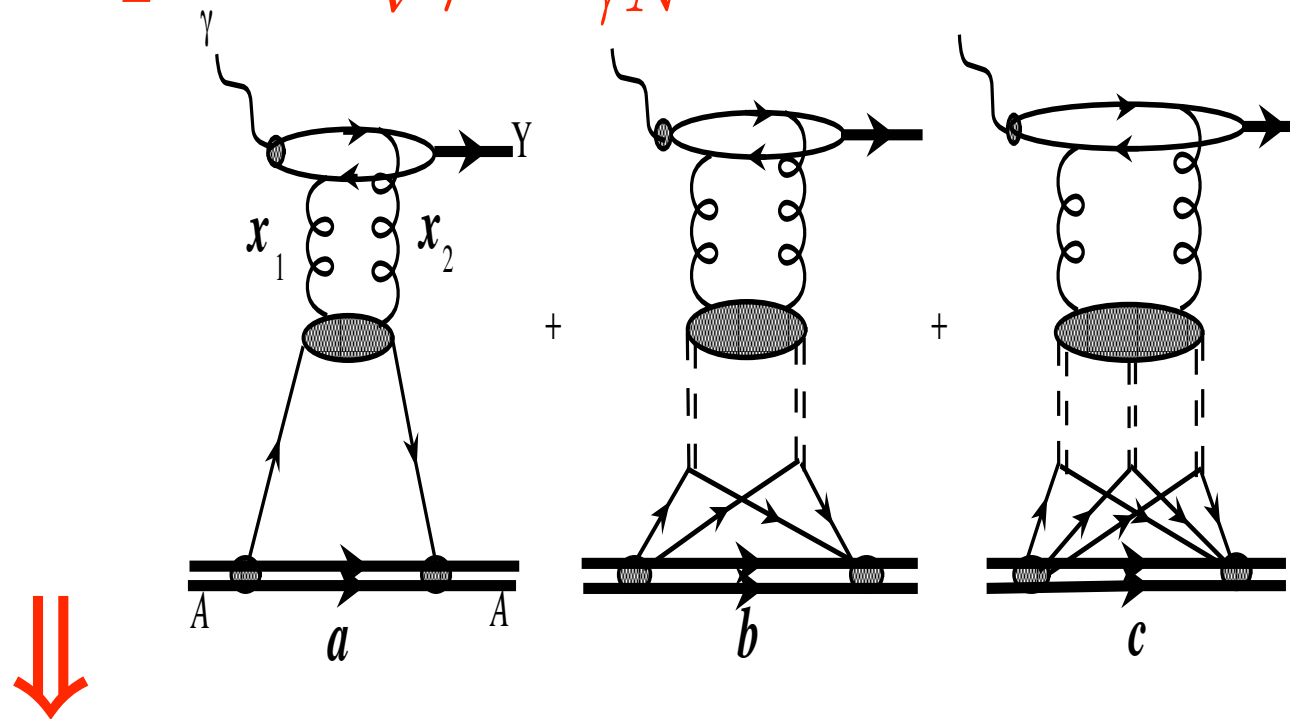
*Another approach - use of the break up channels - processes where nucleus emits few neutrons (Tverskoi, MS, Zhalov 05). Allows to determine which nucleus emitted the photon.*

# The leading twist prediction (neglecting small t dependence of shadowing)

Brodsky et al 94

$$\sigma_{\gamma A \rightarrow VA}(s) = \frac{d\sigma_{\gamma N \rightarrow VN}(s, t_{min})}{dt} \left[ \frac{G_A(x_1, x_2, Q_{eff}^2, t=0)}{AG_N(x_1, x_2, Q_{eff}^2, t=0)} \right]^2 \int_{-\infty}^{t_{min}} dt \left| \int d^2b dz e^{i\vec{q}_t \cdot \vec{b}} e^{iq_1 z} \rho(\vec{b}, z) \right|^2.$$

where  $x = x_1 - x_2 = m_V^2 / W_{\gamma N}^2$



for small d - LT much larger screening than eikonal

: High energy quarkonium photoproduction in the leading twist approximation.

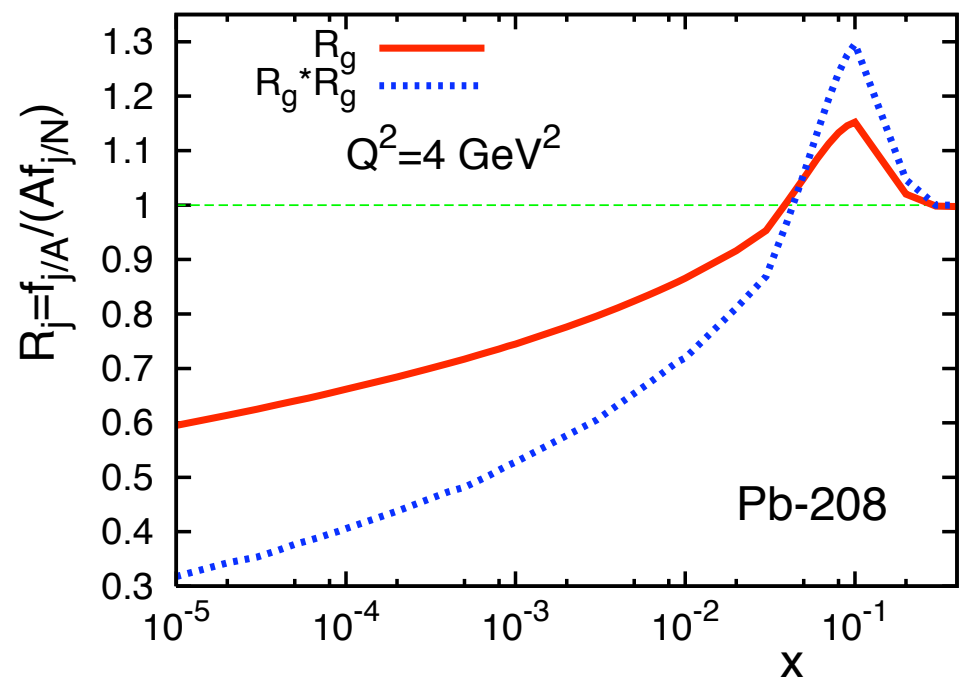
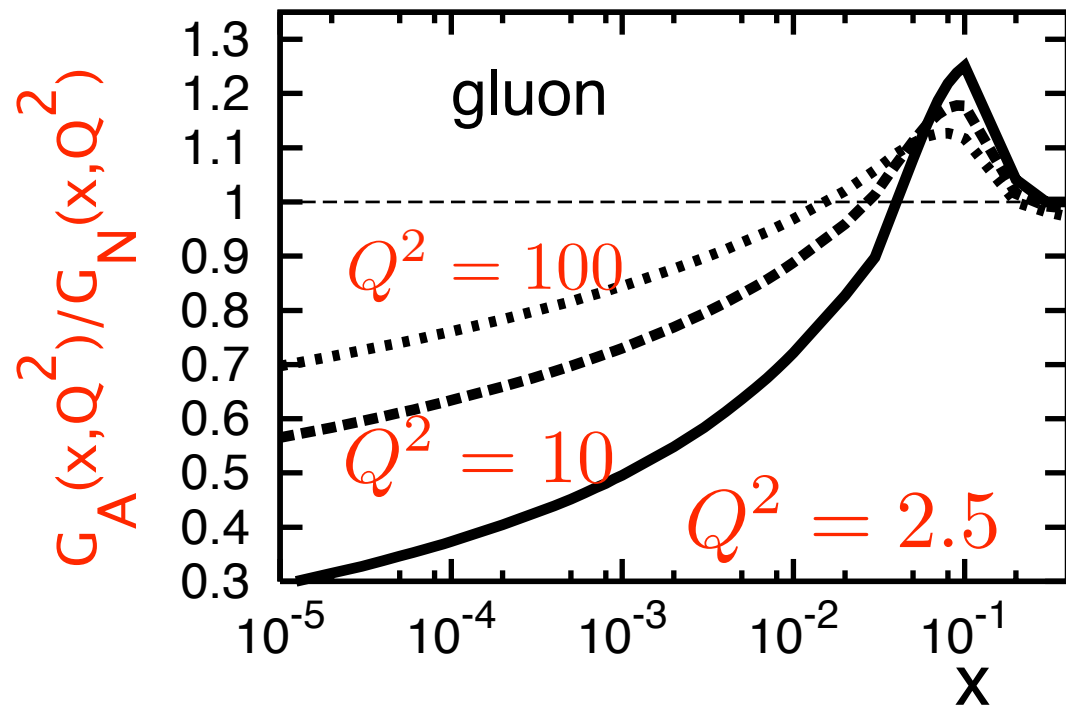
$$\frac{G_A(x_1, x_2, Q_{eff}^2, t=0)}{G_N(x_1, x_2, Q_{eff}^2, t=0)} \approx \frac{G_A((x_1 + x_2)/2, Q_{eff}^2, t=0)}{G_N((x_1 + x_2)/2, Q_{eff}^2, t=0)}$$

$$\frac{(x_1 + x_2)_{J/\psi}}{2} \approx x; \quad \frac{(x_1 + x_2)_{\Upsilon}}{2} \approx x/2$$

Factor of  $> 1.5$  suppression of cross sections at  $x < 0.001$  for onium production for all Q ( as it is  $\propto$  shadowing<sup>2</sup>)

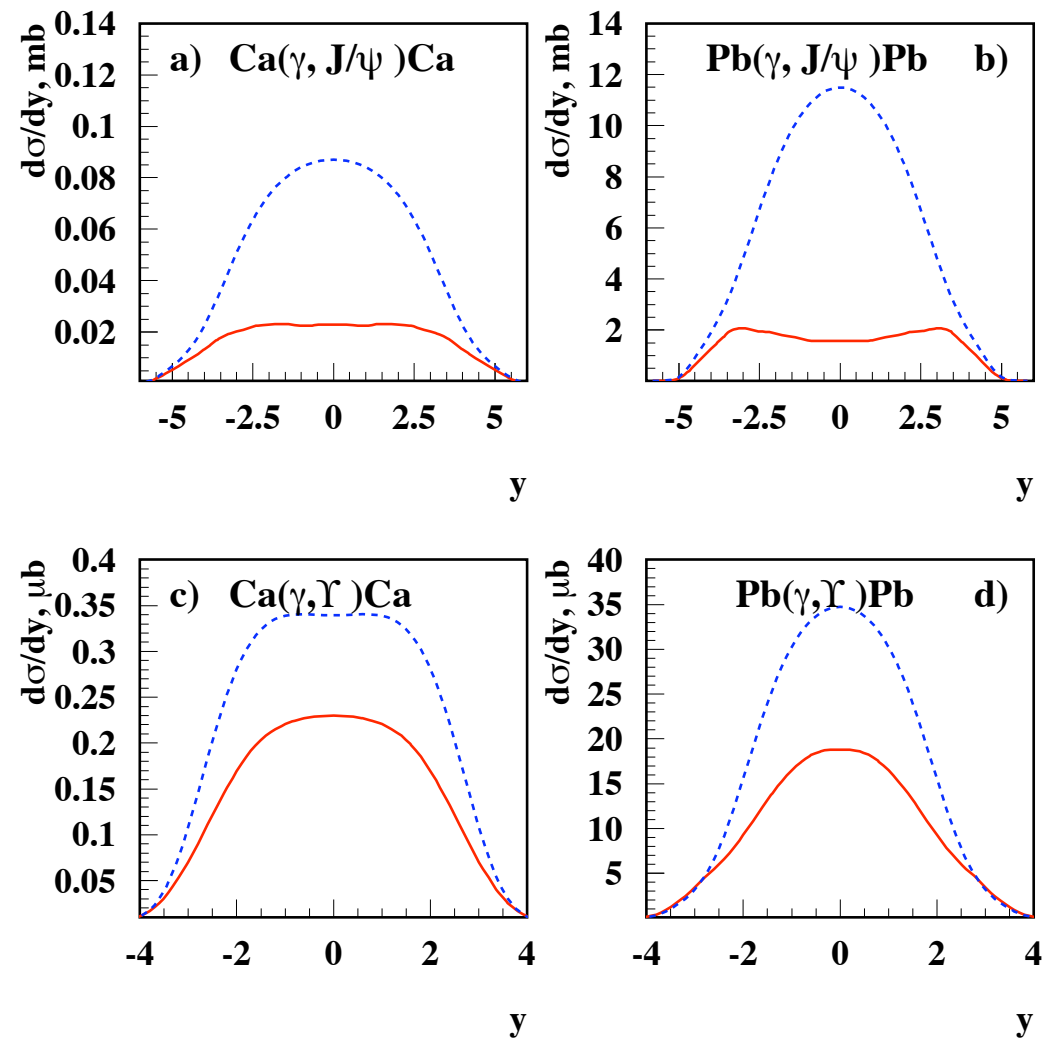
# Theory of LT nuclear shadowing (FS98) allows to relate

$G_A(x, Q^2)/G_N(x, Q^2)$  and diffractive gluon nucleon pdfs –which were recently measured with small uncertainties at HERA, leading to [Guzey et al 07]

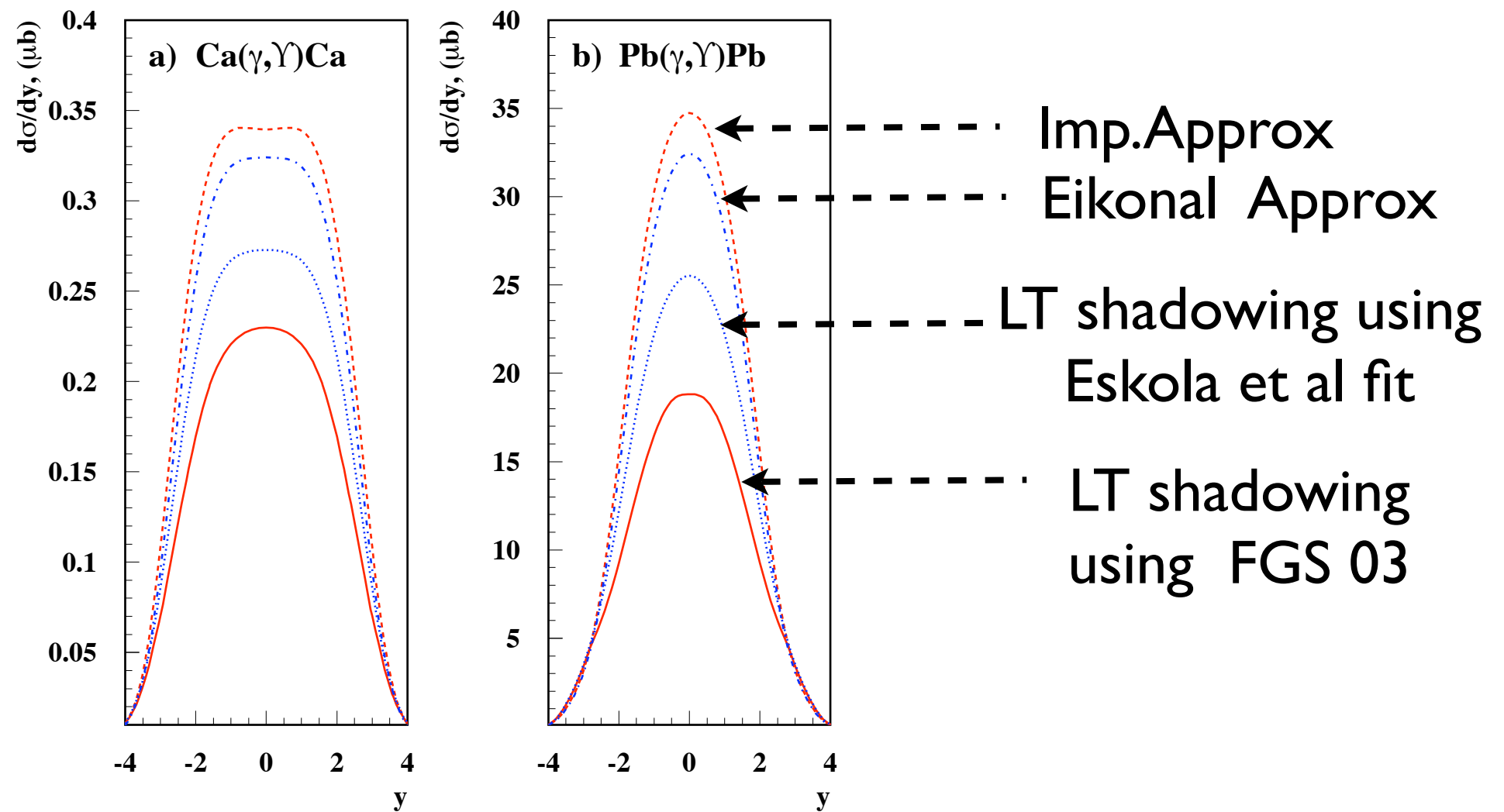


factor  $> 2$  shadowing effects for  $J/\psi$  for  $x < 10^{-2}$   
 & for  $\Upsilon$  for  $x < 10^{-4}$

- ➡ Onset of perturbative color opacity at small  $x$  and onium coherent photoproduction.



*The rapidity distributions for the  $J/\psi$  and  $\Upsilon$  coherent production off Ca and Pb in UPC at LHC calculated with the leading twist shadowing based on H1 parameterization of gluon density (solid line) and in the Impulse Approximation (dashed line).*



**The rapidity distribution for coherent  $\Upsilon$  production in Ca-Ca and Pb-Pb ultraperipheral collisions at the LHC**  
 The solid curve corresponds to the calculation including leading twist nuclear shadowing using H1 diffractive pdfs of 2000; the dotted curve corresponds to the calculation with the model of shadowing of Eskola et al.; the dot-dashed curve is the calculation in the eikonal dipole rescattering model; the dashed curve corresponds to the impulse approximation.

Experimental challenges: Trigger on relatively low transverse momentum leptons. Problem for  $J/\psi$  's for  $y=0(??)$  for  $y=2-4$  the ALICE study finds good rates. Acceptance for  $\Upsilon$  is good in a wide rapidity range. CMS - talk at the workshop



# Neutron tagging of quasielastic J/ψ and Υ photoproduction off nucleus

If  $\sigma_{tot}^{J/\psi N}$  the effective quarkonium(Q $\bar{Q}$ )-nucleon total cross section  
is small (  $\sim 3\text{mb}$  for J/ψ for  $s \sim 200 \text{ GeV}^2$ )

$$\sigma_{inc}^{\gamma A \rightarrow J/\psi A'} = 2\pi\sigma(\gamma N \rightarrow J/\psi N) \cdot \int_0^\infty b db \int_{-\infty}^\infty dz \rho(\vec{b}, z) \exp[-\sigma_{tot}^{J/\psi N} T(\vec{b})]$$

Here  $T(\vec{b}) = \int_{-\infty}^\infty \rho(\vec{b}, z) dz$

$$\sigma_{elastic}(\gamma A \rightarrow J/\psi + A) \propto A^{4/3}, \sigma_{quasielastic}(\gamma A \rightarrow J/\psi + A') \propto A$$

⇓ *LT shadowing*

⇓ *BDR*

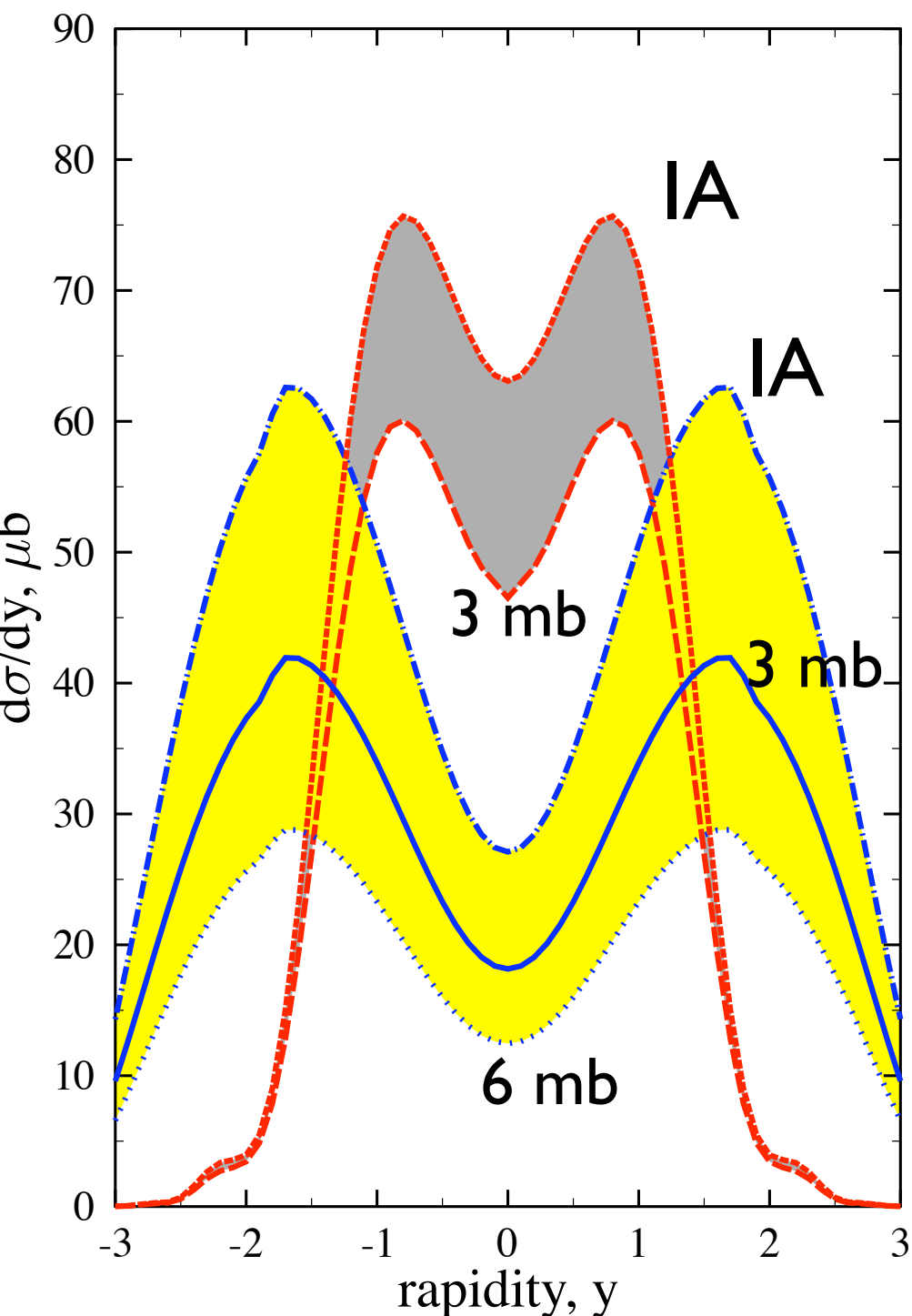
$$\sigma_{elastic}(\gamma A \rightarrow J/\psi + A) \propto A^{2/3}, \sigma_{quasielastic}(\gamma A \rightarrow J/\psi + A') \propto A^{1/3}$$

Change of A dependence by a factor  $\sim A^{2/3}$  for both processes !!!



Both processes allow to study new QCD domain

QE channel allows to reach a factor of  $\sim 100$  smaller  $x$  !!

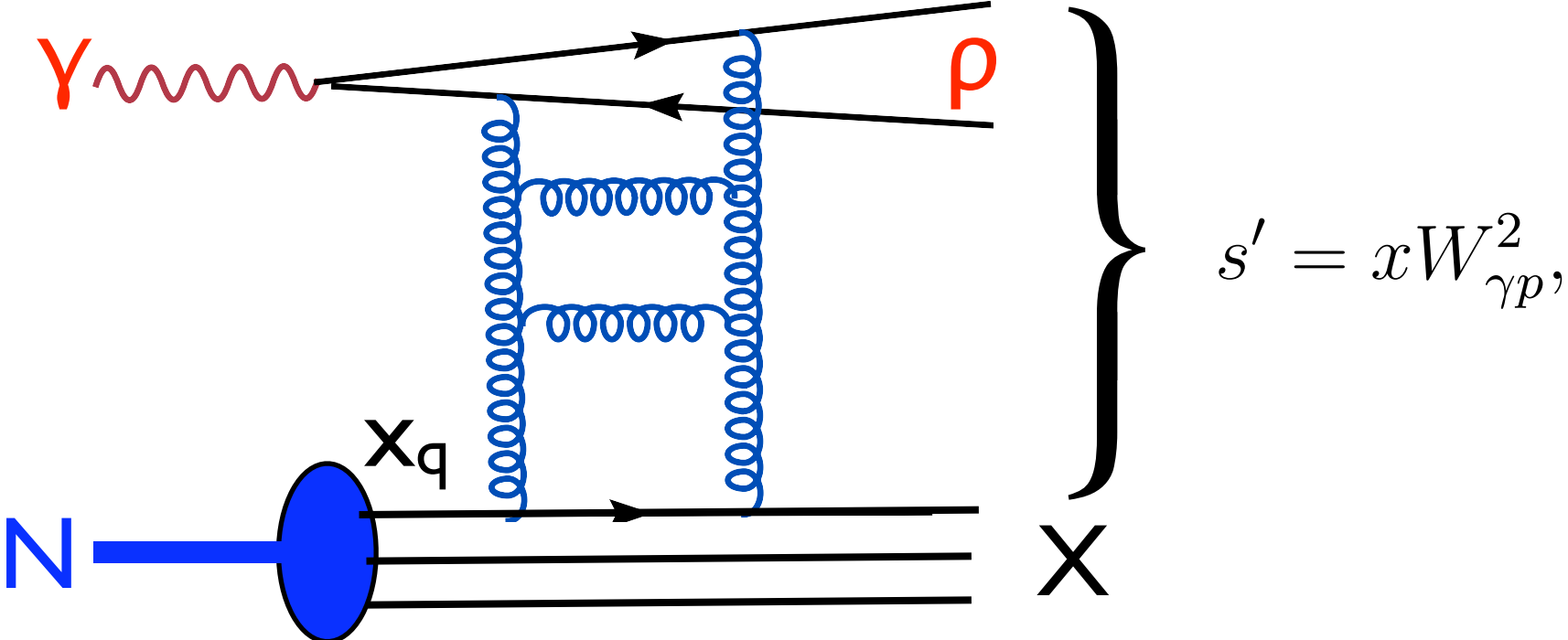


The integrated over momentum transfer rapidity distributions for the  $J/\psi$  coherent photoproduction in UPC of Au ions at RHIC calculated with effective cross section for  $J/\psi$  - nucleon interaction of 3 mb (long-dashed line) and in the Impulse Approximation (short-dashed line). The incoherent  $J/\psi$  production cross section estimated in the Glauber model for  $J/\psi$  - nucleon cross section of 3 mb (solid line) and 6 mb (dotted line), and in the IA (dot-dashed line)

All plans/measurements involve selecting events where nucleus emitted neutrons - efficiency  $\sim 100\%$  for QE and  $\sim 50\%$  (RHIC)/  $\sim 70\%$  LHC

Hence QE/ELASTIC  $\sim 0.3 \rightarrow \sim 40\%$  (RHIC) -  $35\%$  (LHC) for observed events

# Rapidity gap processes at large $t=(p_\rho-p_\gamma)^2$ : from HERA to LHC



Elementary reaction - scattering of a hadron ( $\gamma, \gamma^*$ ) off a parton of the target at large  $t=(p_\gamma-p_\nu)^2$

FS 89 (large  $t$   $pp \rightarrow p + \text{gap} + \text{jet}$ ), FS95

Mueller & Tung 91

Forshaw & Ryskin 95

$$x = \frac{-t}{(-t + M_X^2 - m_N^2)}$$

The rapidity gap between the produced vector meson and knocked out parton (roughly corresponding to the leading edge of the rapidity range filled by the hadronic system  $X$ ) is related to  $W_{\gamma p}$  and  $t$  (for large  $t$ ,  $W_{\gamma p}$  as

$$y_r = \ln \frac{x W_{\gamma p}^2}{\sqrt{(-t)(m_V^2 - t)}}$$

The choice of large  $t$  ensures two important simplifications. First, *the parton ladder mediating quasielastic scattering is attached to the projectile via two gluons*. Second is that *attachment of the ladder to two partons of the target is strongly suppressed*. Also the transverse size  $d_{q\bar{q}} \propto 1/\sqrt{-t}$

$$\frac{d\sigma_{\gamma+p \rightarrow V+X}}{dt dx} = \frac{d\sigma_{\gamma+quark \rightarrow V+quark}}{dt} \left[ \frac{81}{16} g_p(x, t) + \sum_i (q_p^i(x, t) + \bar{q}_p^i(x, t)) \right]$$

$$\frac{d\sigma_{N+q(g)\rightarrow N+q(g)}}{dt} \propto \frac{1}{t^6} \qquad \frac{d\sigma_{\gamma+q(g)\rightarrow V+q(g)}}{dt} \propto \frac{1}{t^4}$$

Energy dependence of  $f_q(s',t) \propto [s']^{\delta(t)}$

$\delta(-t \gg 1 \text{ GeV}^2)?$

Soft QCD  $\delta < -0.5$

Two gluon exchange  $\delta = 0$

DGLAP / resummed BFKL for  $t=0$   $\delta = 0.2 \text{ -- } 0.3$

subtle points in BFKL analysis for  $t$  away from 0



We analyzed the rho-meson data using a fit

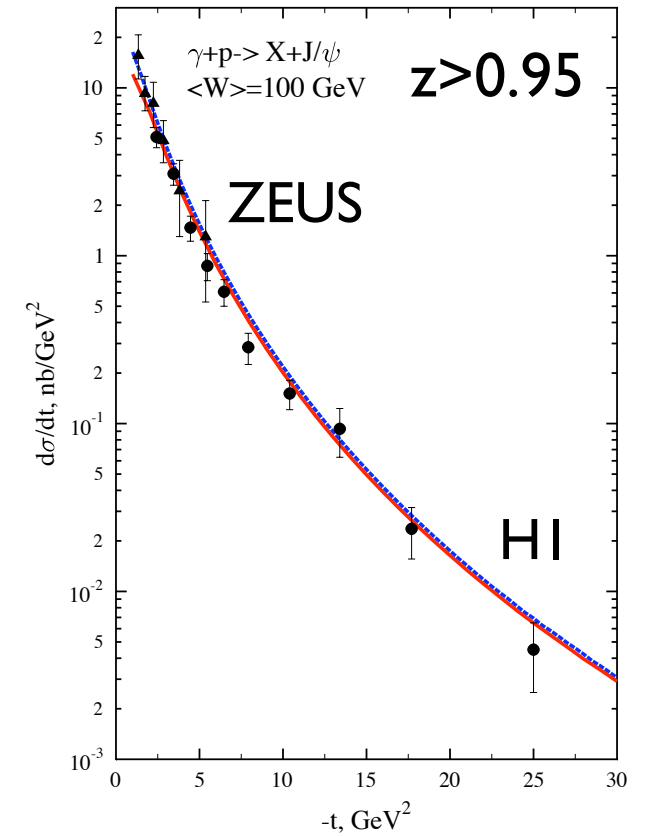
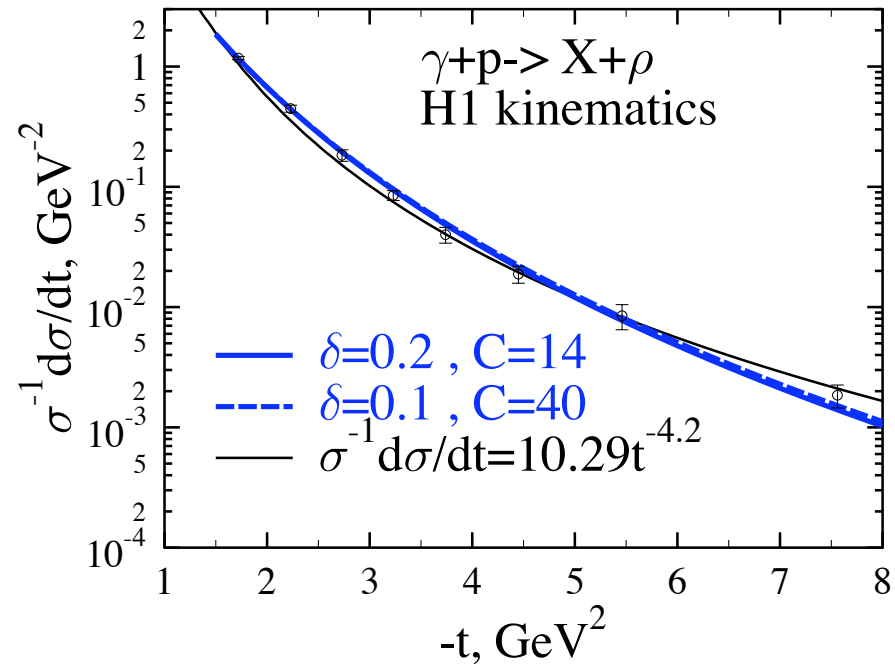
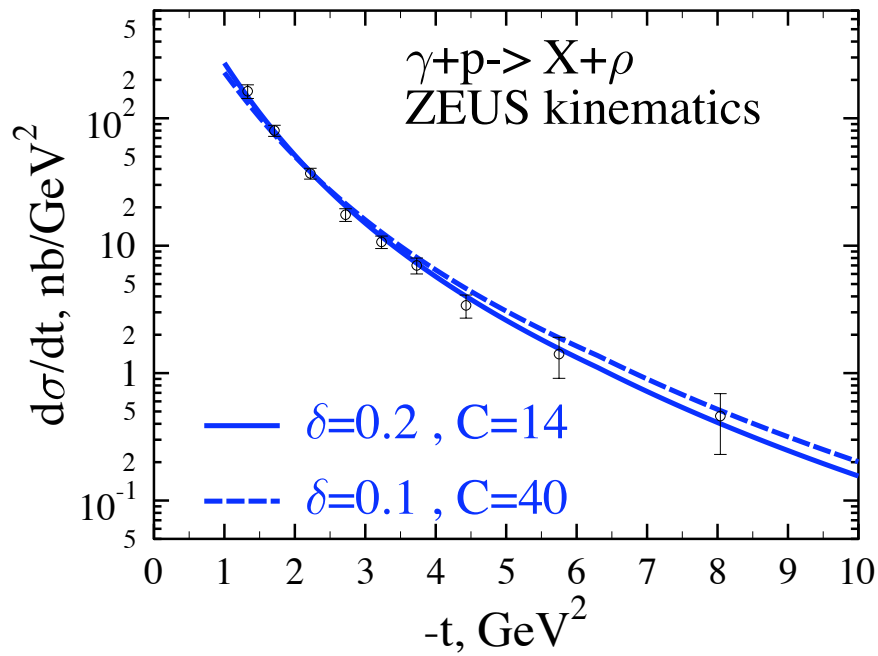
$$\frac{d\sigma_{\gamma+p \rightarrow \rho+X}}{dt} = \frac{C}{(1 - t/t_0)^4} \left( \frac{s}{m_V^2 - t} \right)^{2\delta(t)} I(x_{min}, t)$$

$$I(x_{min}, t) = \int_{x_{min}}^1 x^{2\delta(t)} \left[ \frac{81}{16} g_p(x, t) + \sum_i [q_p^i(x, t) + \bar{q}_p^i(x, t)] \right] dx$$

$t_0 \sim 1 \text{ GeV}^2$ ,  $\delta=0.1 - 0.2$  is consistent with the data at large  $t$

For  $J/\psi$  we changed

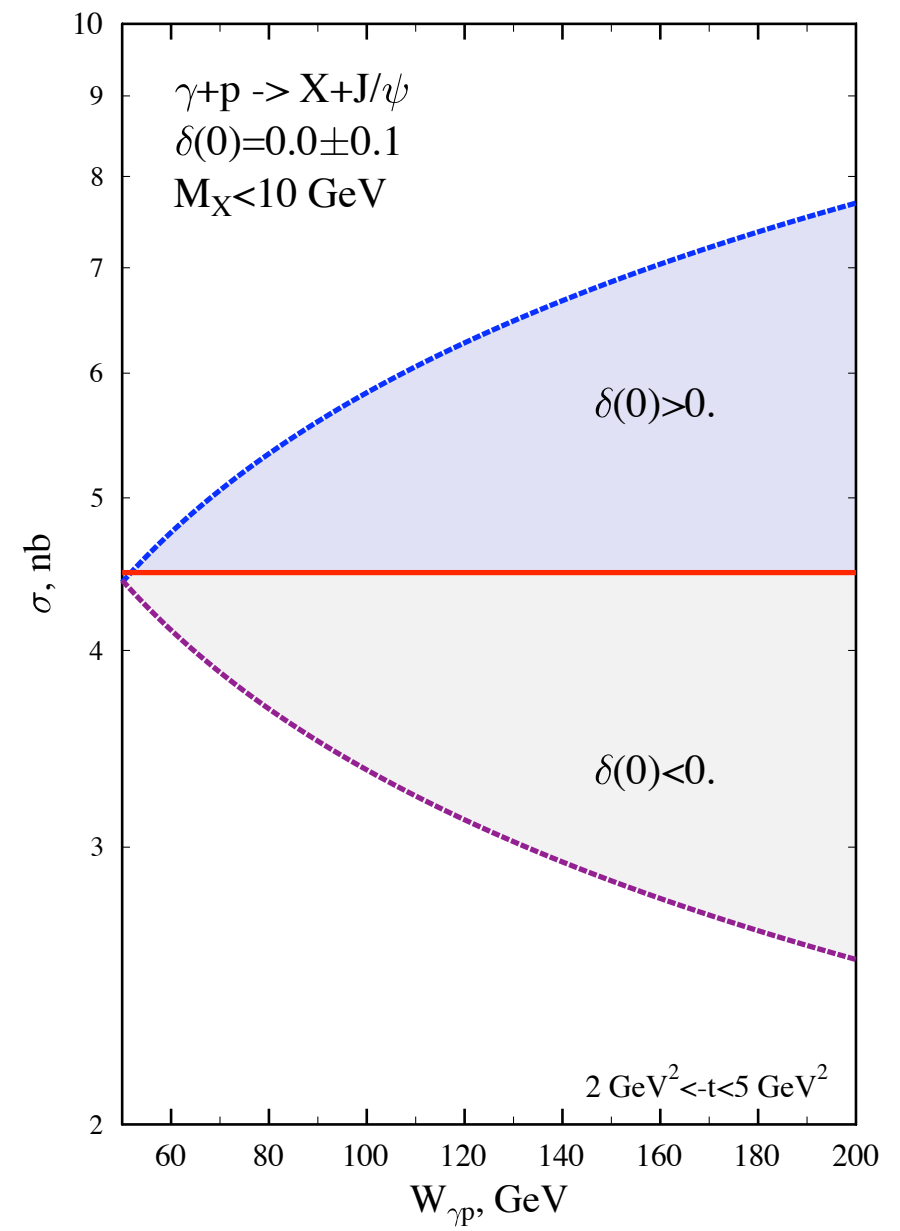
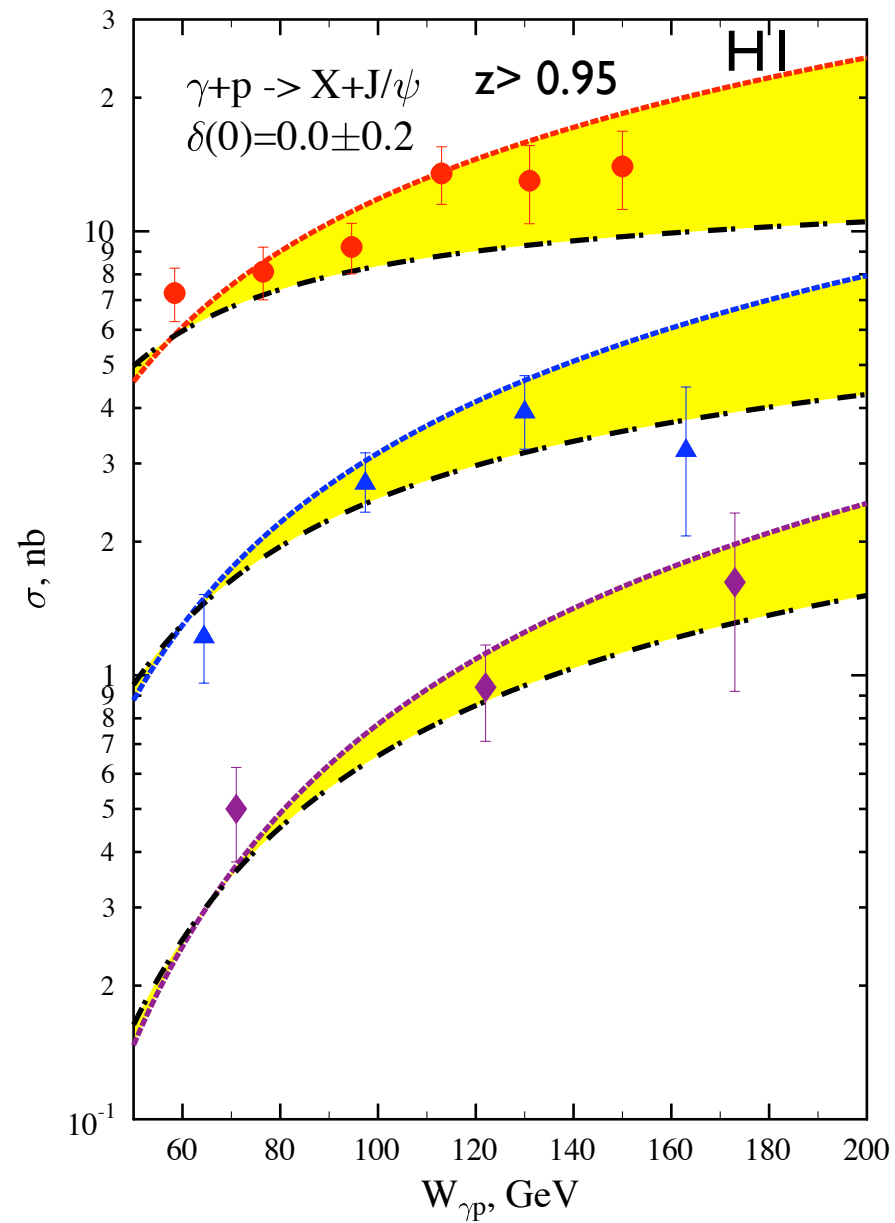
$$\frac{1}{(1 - t/t_0)^4} \rightarrow \frac{1}{(1 - t/t_0)(1 - t/m_{J/\psi}^2)^3}$$



Description of ZEUS and H1 data for  $t$ -dependence of the large  $t$  and rapidity gap cross section. ZEUS data were taken at average  $\langle W_{\gamma p} \rangle = 100$  GeV with fixed cut  $M_X < 25$  GeV and additional restriction  $0.01 < x < 1$ . The H1 data were taken at average  $\langle W_{\gamma p} \rangle = 85$  GeV and cut  $M_X < 5$  GeV.

Sensitivity to the energy dependence is weak.

$t$ -dependence of  $J/\psi$  production is consistent with dominance of hard dynamics



Study of the VM production with gaps is mostly sensitive to gluon pdfs if the cut is on  $z_{\min}$  or  $M_X^2/W^2$  is made. Sensitivity to the energy dependence of dipole - parton amplitude  $f(s',t) \propto s'^{\delta}$  is minor. On the contrary if the cut on  $M_X < \text{const}$  is made, sensitivity to the value of  $\delta$  is very high.

Analyses with z cut,  $M^2_X/s < \text{const}$  cuts are good for study of the dominance of the mechanism of scattering off single partons. However they correspond to rapidity interval between VM and jet which are typically of the order  $\Delta y = 2 - 3$ .

Optimal way to study BFKL dynamics is to keep  
 $M^2_X < \text{const}$  and vary  $W$

Difficult but not impossible at HERA natural at LHC

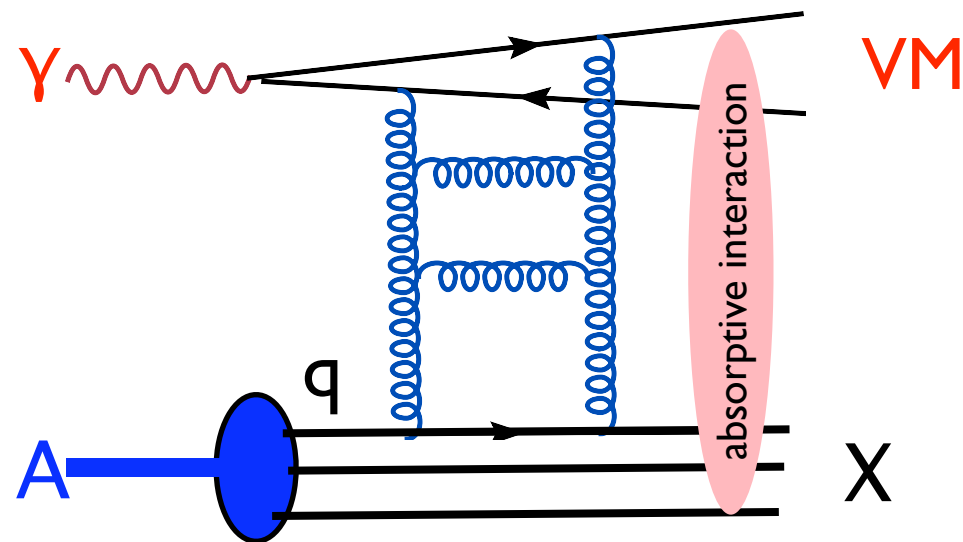
At LHC one can energy dependence of elastic  $q\bar{q}$  - parton scattering at  $W'=20$  GeV - 400 GeV

$$\sigma_{el}(q\bar{q} - q(g)(W' = 400\text{GeV})/\sigma_{el}(q\bar{q} - q(g)(W' = 20\text{GeV}) \sim 10!!! \quad \text{if } \delta=0.2$$

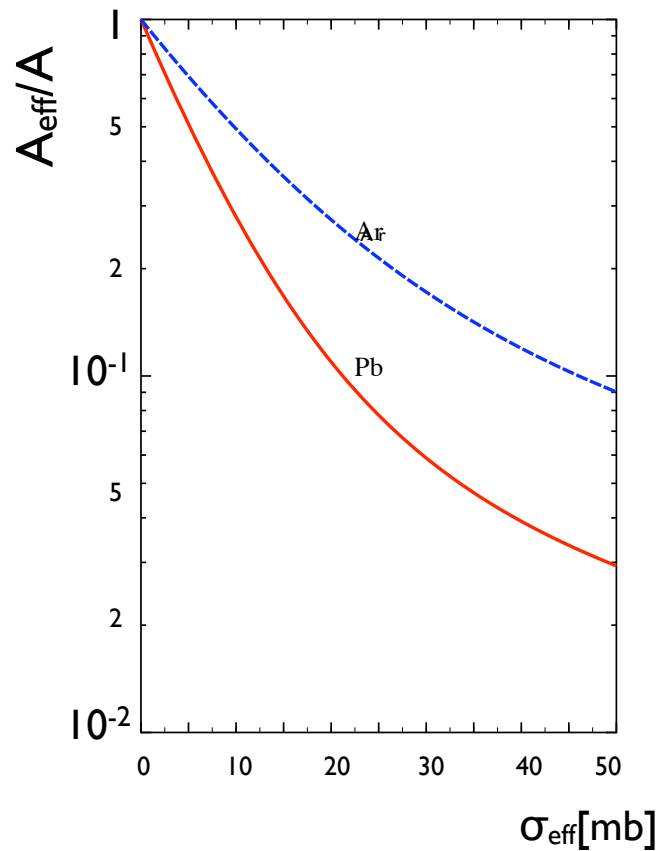
- $\gamma + A \rightarrow \rho (J/\psi) + \text{gap} + X$

measure of the strength of inelastic interactions of small dipole in the processes initiated by BFKL elastic  $q\bar{q}$  - parton scattering at  $W=30 \text{ GeV} - 1 \text{ TeV}$

$$\sigma_{el}(q\bar{q} - q(g)(W = 1\text{TeV}) / \sigma_{el}(q\bar{q} - q(g)(W = 30\text{GeV}) > 30 !!!$$



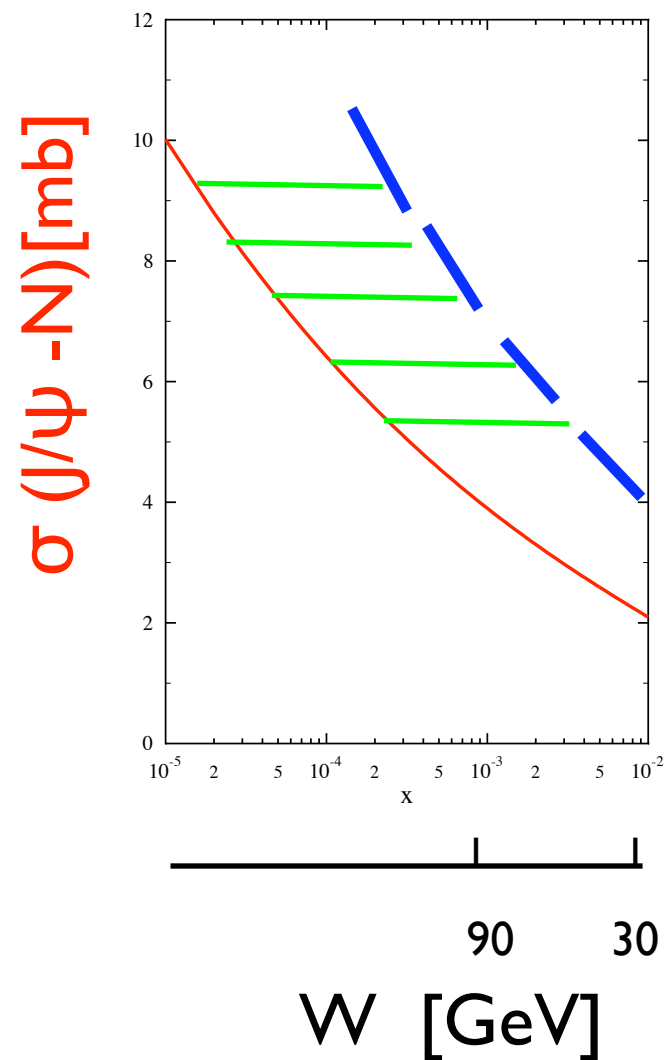
**Advantages:** trigger on hadron production in a rapidity interval close to one of the nuclei  
 no ambiguity which of the nuclei emitted photon - Large  $W$  are possible,  
 better VM acceptance for large  $t$ .



Strong sensitivity of  $A_{\text{eff}}/A$  to the strength of inelastic  $qq\bar{N}$  interactions

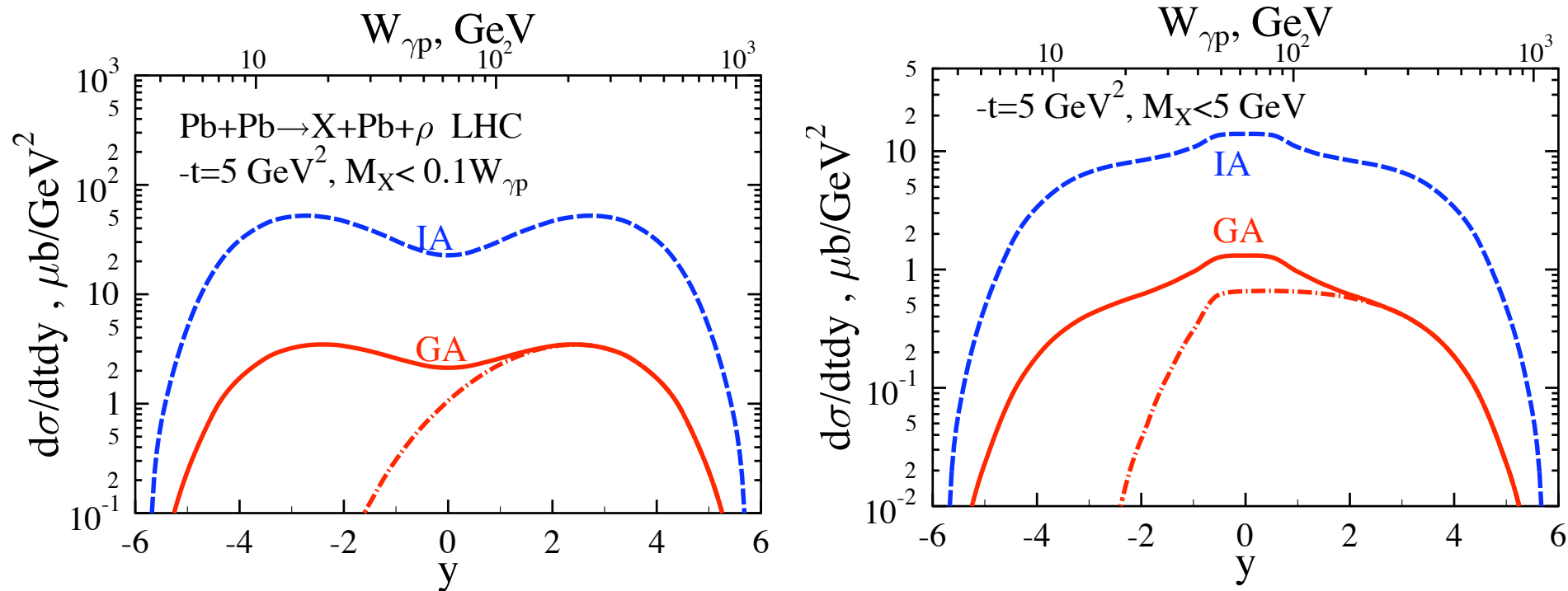
Predict:

- \*  $A_{\text{eff}}/A$  should increase with  $t$  at fixed  $W$   
 $\sigma \propto W^{2\delta}, \delta = 0.2 - 0.3$
- \*  $A_{\text{eff}}/A$  should decrease with increase of  $W$  at fixed  $t$  - onset of black disk regime



Complementary to quasielastic and coherence process of VM production - no small  $x$  partons in the nucleus are involved on the trigger level





Integrated over mass of produced system cross section of the nucleon dissociative  $\rho$  meson photoproduction at  $-t=5 \text{ GeV}^2$  in the ultraperipheral lead-lead collisions at LHC. The upper figure - the limit of the mass of produced system  $M_X$  is proportional to the photon-nucleon center of mass energy  $M_X < 0.1 W_{\gamma p}$ , in the right figure for central rapidities the limit of  $M_X$  is fixed by restriction  $M_X < 5 \text{ GeV}$ . Solid line - calculations with Glauber-Gribov screening, dashed line calculations in the leading twist approximation neglecting nuclear shadowing correction which is very small for discussed kinematics, dot-dashed line - one-side contribution when  $\rho$  meson is produced by photons emitted by only one nucleus: large positive rapidities correspond to vector mesons produced by high energy photons. The counting rate can be estimated using expected luminosity for PbPb collisions  $L=10^{-3} \mu\text{b}^{-1} \text{ sec}^{-1}$ .

# Exclusive UPC processes in $pA$

**$pA$  ultraperipheral will play dual role -**



extend studies of the nucleon structure



serve as a reference point to nuclear studies using UPC in AA collisions



extend studies of the onium exclusive production

# Studies of exclusive photoproduction processes in pA UPC:

## *Hard physics:*



Onium production

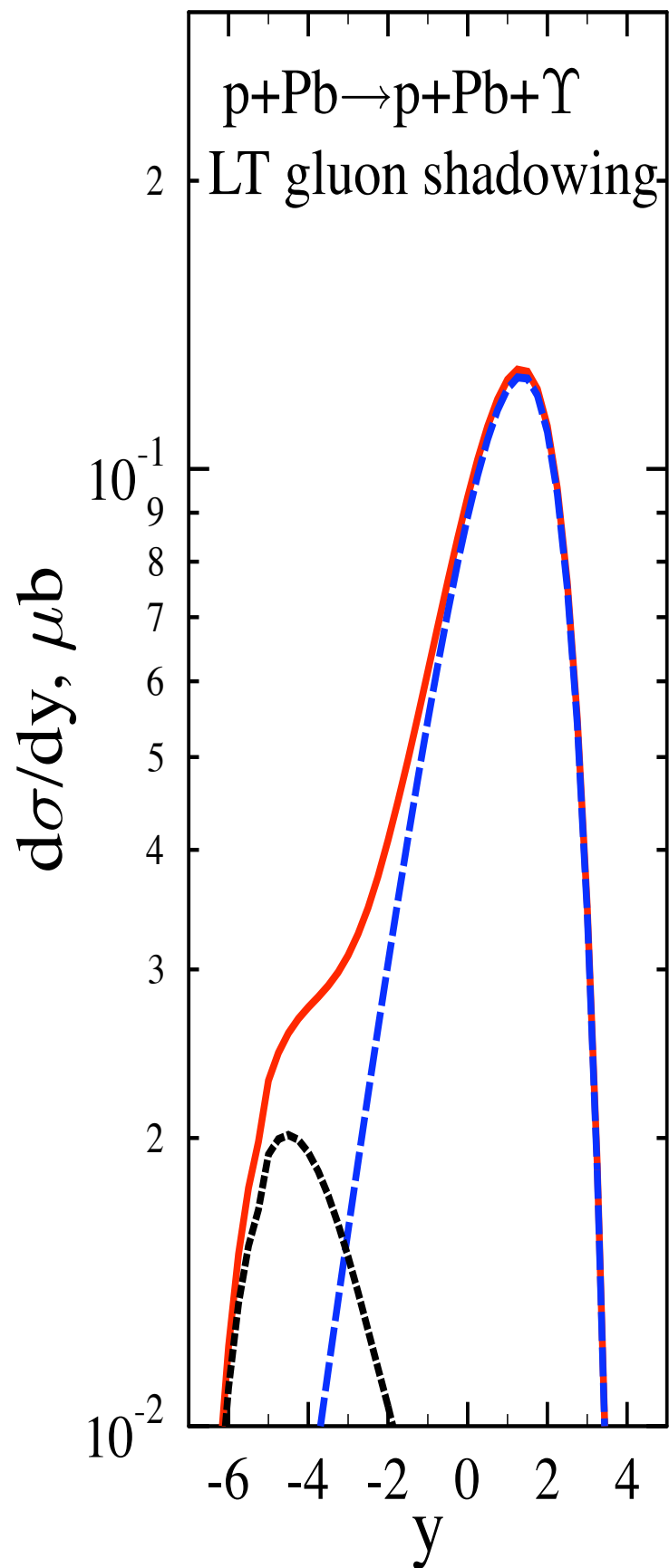
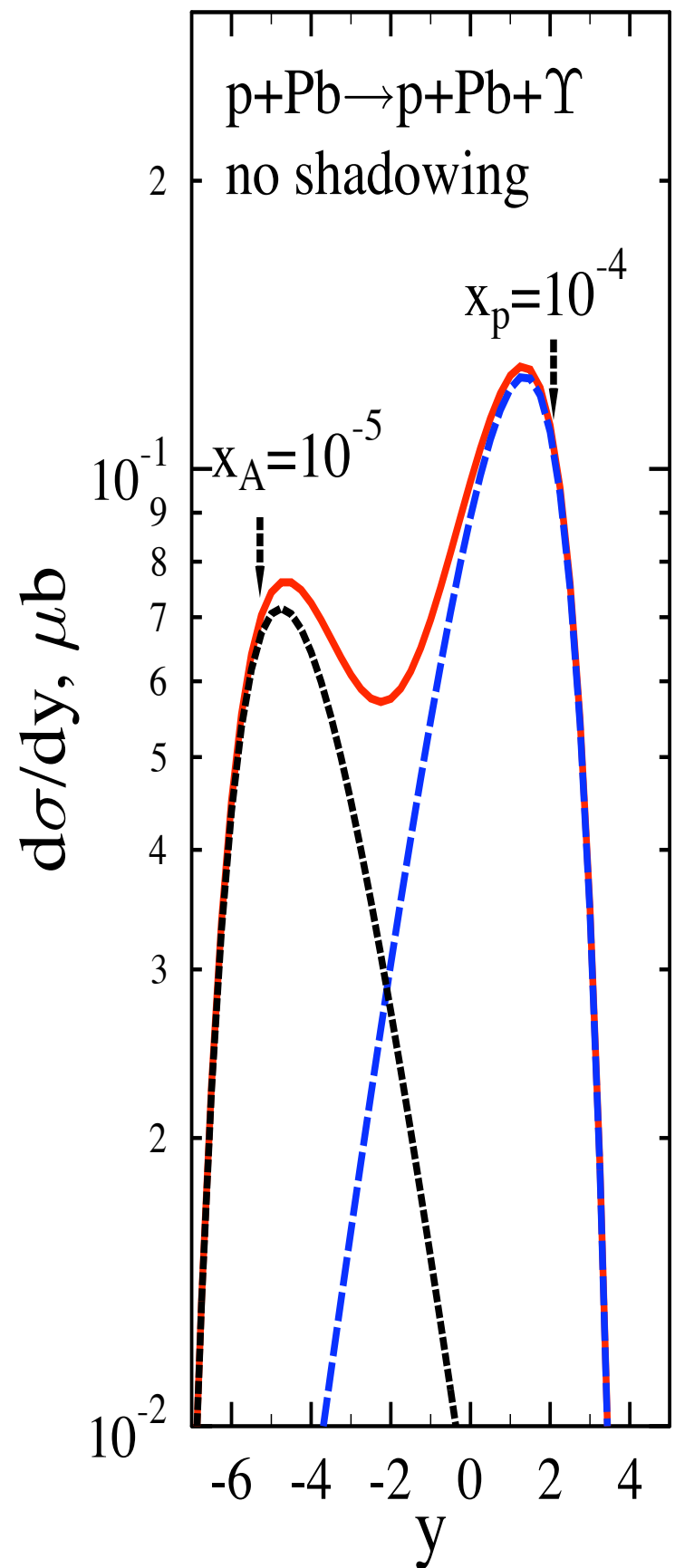


Diffraction into two, three jets

## *Soft (Pomeron) physics:*



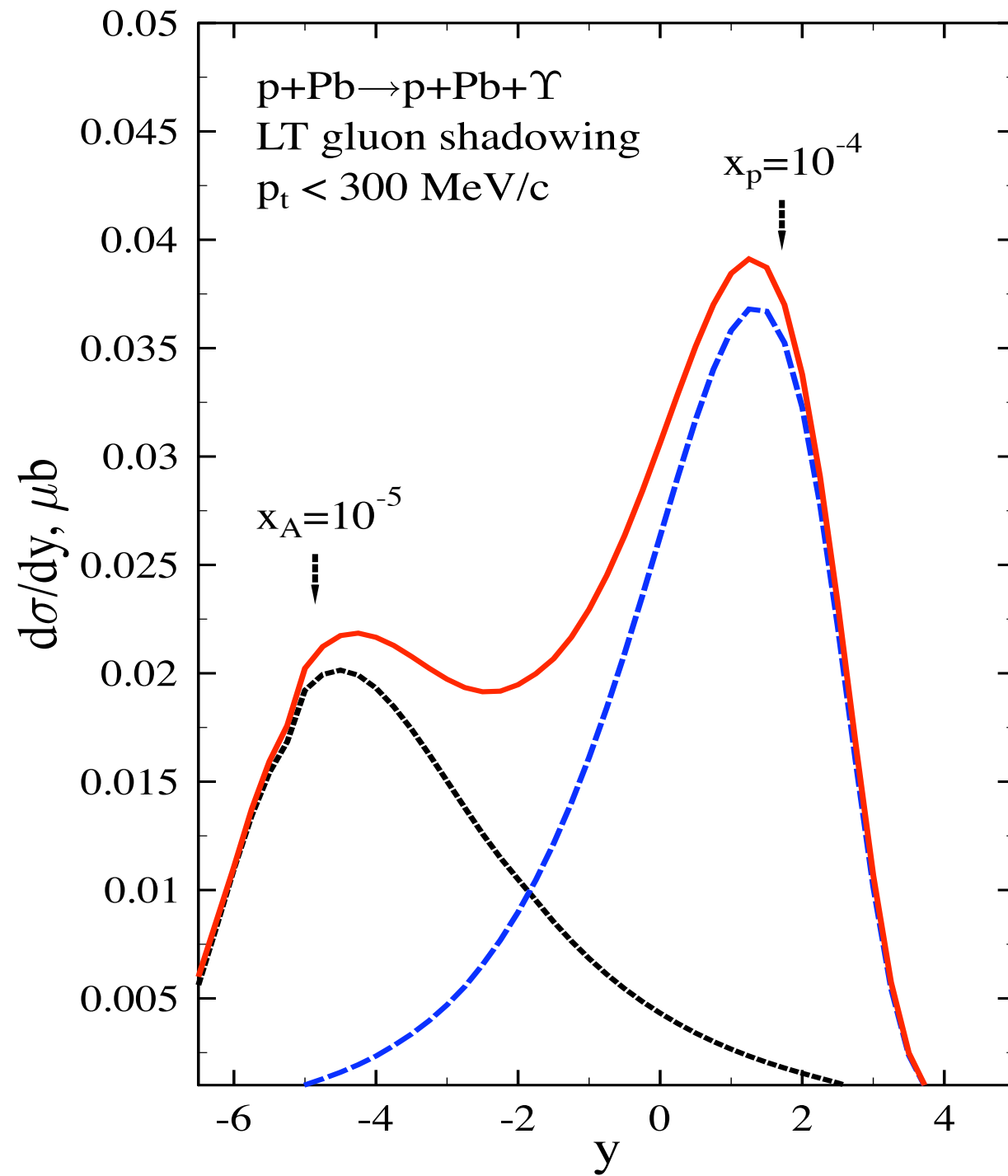
Energy dependence of production of  $\rho, \varphi$ -mesons



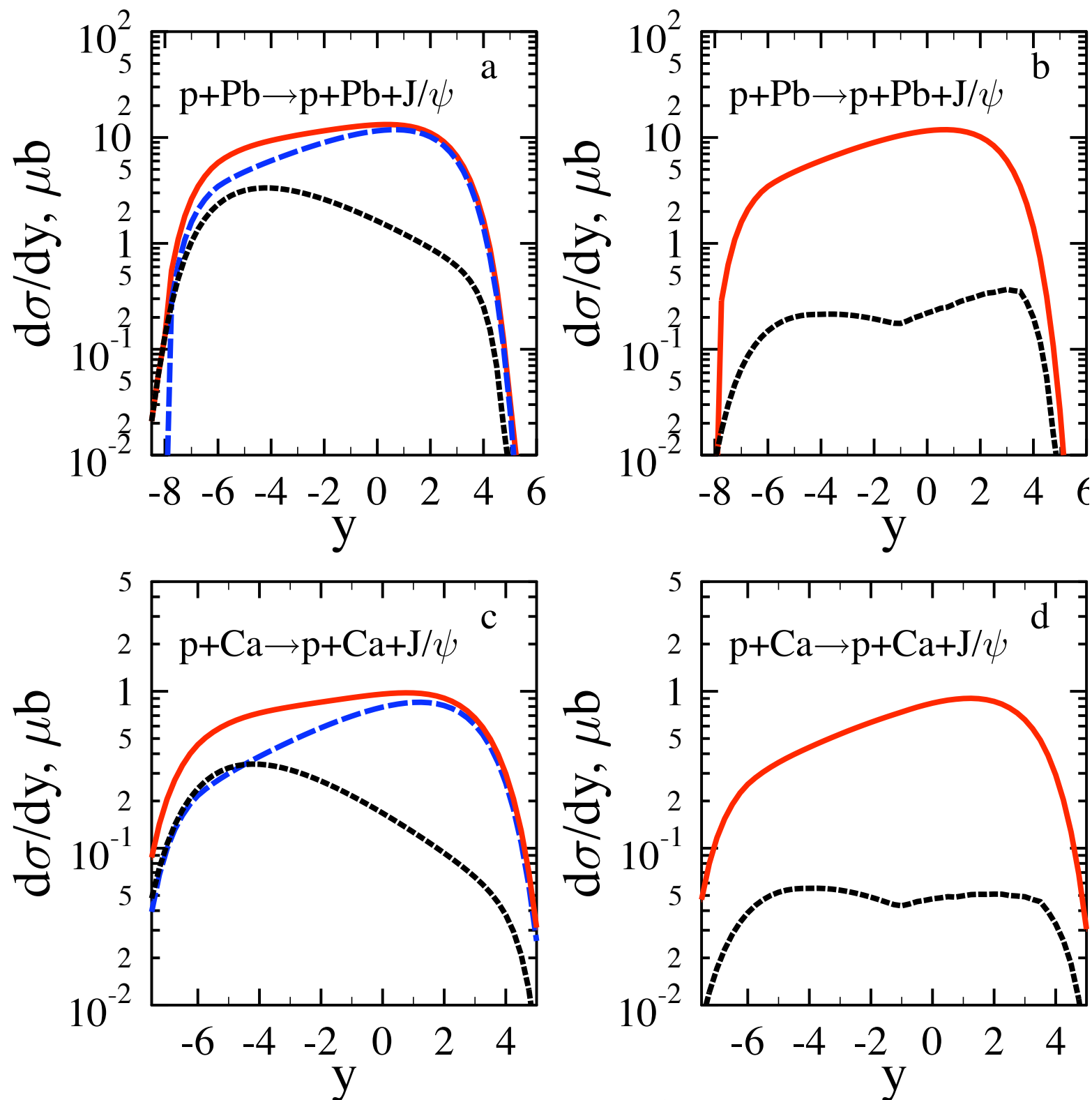
Zhalov & MS 05

Sufficient to check pQCD prediction of  $\sigma \sim W^{1.6}$  for Upsilon production, determination of the t-slope provided protons could be detected (420 m proposal) and measure nuclear shadowing at  $Q^2 = 40 \text{ GeV}^2$

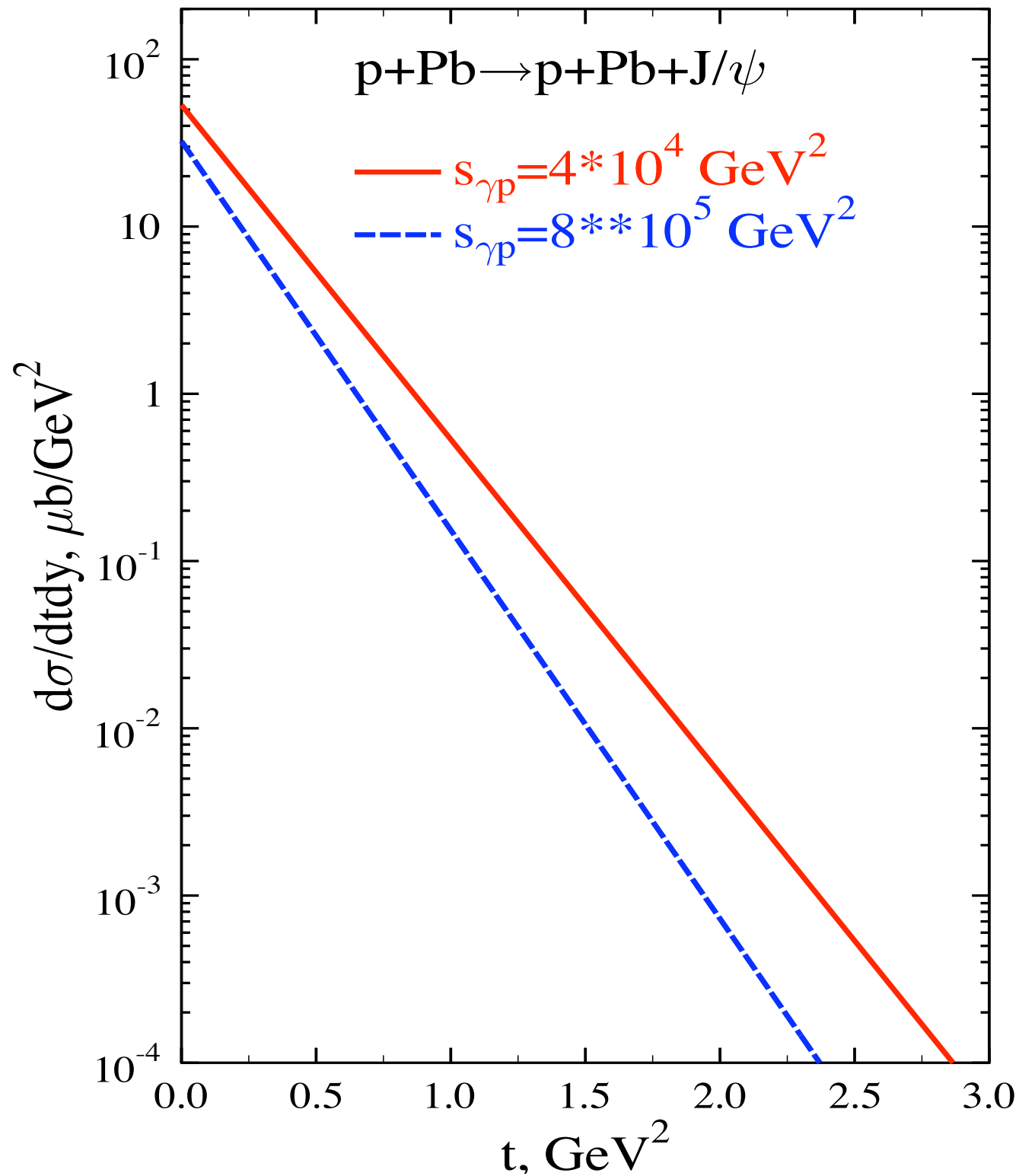
Production of  $\Upsilon$ 's in pA collisions : coherent  $\gamma+A \rightarrow V+A$  is shown by black lines, and  $\gamma+p \rightarrow V+p$  by blue lines.



Rapidity distribution for  $\Upsilon$  photoproduction in pPb UPC at LHC with (solid line) with gluon shadowing and the cut of the quarkonium transverse momentum  $p_t < 300 \text{ MeV}/c$ .



High enough rates down to  $x \sim 10^{-6}$ , however extracting nuclear contribution would be a challenge if indeed the nuclear shadowing is as high as in FGS05. Would require resolution in transverse momentum of  $J/\psi$  of  $\sim 150 \text{ MeV}/c$ .



Need proton  
detector/ veto

Momentum transfer distribution for  $J/\psi$  photoproduction in pA at LHC



# Conclusions

Studies of UPC at LHC will address many (though not all) of the benchmark issues of HERA III proposal including



Small  $x$  physics with protons and nuclei in **a factor of ten** larger energy range though at higher virtualities both in inclusive and diffractive channels



Interaction of small dipoles at ultrahigh energies - approach to black body regime, color opacity



Low  $Q$  will be missed - will require studies at eRHIC

$\sqrt{Q^2}$  or  $p_t$

