

Higgs plus single top production at the LHC

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Experimental searches

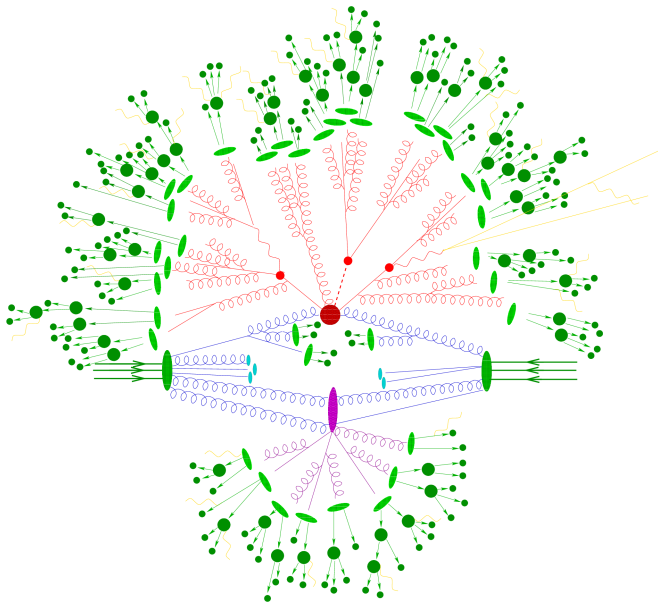
- *ATLAS*: (only as part of $t\bar{t}H$ analyses)
- *CMS*: CMS-PAS-HIG-14-001, CMS-PAS-HIG-14-015 ($H \rightarrow \gamma\gamma$, $H \rightarrow b\bar{b}$, optimized for $y_t = -y_{t,SM}$)

At the moment: limits on tH cross section (supposed $> SM$)

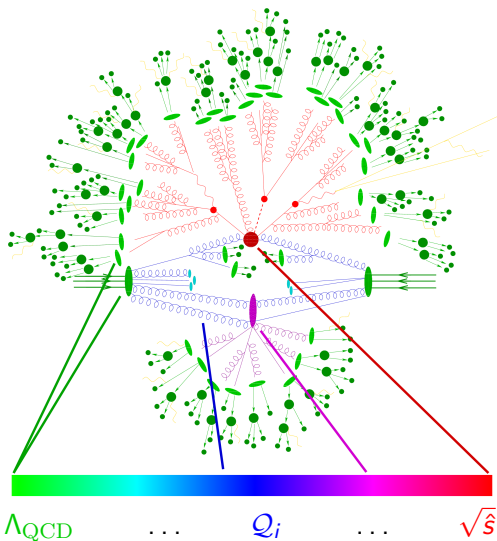
Pheno studies

- *Maltoni, Paul, Stelzer, Willenbrock* [arXiv:hep-ph/0106293] (LO, sig+bkg)
- *Biswas, Gabrielli, Mele* [arXiv:1211.0499] (LO, $H \rightarrow \gamma\gamma$ sig+bkg)
- *Farina, Grojean, Maltoni, Salvioni, Thamm* [arXiv:1211.3736]
(NLO xsect 5F, LO distr, $H \rightarrow b\bar{b}$ sig+bkg, $\pm y_t$)
- *Ellis, Hwang, Sakurai, Takeuchi* [arXiv:1312.5736] (LO, sign/phase of y_t)
- *Chang, Cheung, Lee, Lu* [arXiv:1403.2053]
(LO distr, $H \rightarrow b\bar{b}, \gamma\gamma, 4\ell, \tau^+\tau^-$ sig+bkg, detector simulation, $\pm y_t$)
- *Demartin, Maltoni, Mawatari, Zaro* [arXiv:1504.00611] (NLO xsect and distr, 4F/5F and other uncertainties, sign/phase of y_t)

Hadron collisions are a mess



Many scales



NON-PERT. ($\alpha_s(Q) > 1$)

PERT. ($\alpha_s(Q) \ll 1$)

Logs:

$$\alpha_s(Q_i)^n \ln^n (\Lambda_{\text{QCD}}/Q_i)$$

$$\alpha_s(Q_i)^n \ln^n (Q_i/\sqrt{\hat{s}})$$

must resum if $\alpha_s \ln > 1$

might still need if $\alpha_s \ln \lesssim 1$

Powers:

$$\alpha_s(Q_i)^n (\Lambda_{\text{QCD}}/Q_i)^m$$

$$\alpha_s(Q_i)^n (Q_i/\sqrt{\hat{s}})^m$$

negligible only for large
separation between scales

Introduce an additional perturbative scale (μ) and factorize

For small scale q :

- resum large logs $\alpha_s(\mu)^n \ln^n(q/\mu)$ to all orders into a PDF (or fragmentation function)
- neglect power-suppressed terms $\alpha_s(\mu)^n (q/\mu)^m$

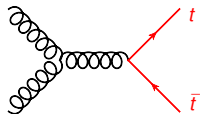
For large scale Q :

- compute both $\alpha_s(\mu)^n \ln^n(\mu/Q)$ and $\alpha_s(\mu)^n (\mu/Q)^m$ at a fixed order in perturbation theory
- neglect higher-orders

An example: $t\bar{t}$ production

$$\Lambda_{\text{QCD}} \quad \xleftarrow{\sim \times 1000} \quad \xrightarrow{\quad} \quad 2m_t, \sqrt{\hat{s}}$$

$\uparrow$
choose a scale μ here μ



Neglect

$$\alpha_s(\mu)^n (\Lambda_{\text{QCD}}/\mu)^m \quad (\text{set } m_q = 0)$$

Resum to all orders

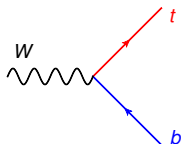
$$\alpha_s(\mu)^n \ln^n(\Lambda_{\text{QCD}}/\mu) \quad \text{into the PDFs}$$

Fixed order

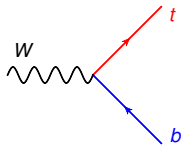
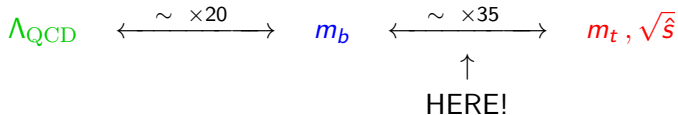
$$\alpha_s(\mu)^n (\mu/m_t)^m$$
$$\alpha_s(\mu)^n \ln^n(\mu/m_t)$$

A less trivial example: single-top production

$$\Lambda_{\text{QCD}} \xleftrightarrow[\substack{\uparrow \\ \text{(naive 4F) here?}}]{\sim \times 20} m_b \xleftrightarrow[\substack{\uparrow \\ \text{(naive 5F) here?}}]{\sim \times 35} m_t, \sqrt{\hat{s}}$$



A less trivial example: single-top production



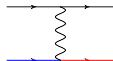
Maltoni, Ridolfi, Ubiali [arXiv:1203.6393]

Campbell, Frederix, Maltoni, Tramontano [arXiv:0903.0005]

A less trivial example: single-top production

$$\Lambda_{\text{QCD}} \longleftrightarrow m_b \longleftrightarrow \mu \longleftrightarrow m_t, \sqrt{\hat{s}}$$

5-flavour scheme (5F): put b 's in the initial state



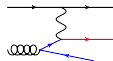
Neglect

$$\alpha_s(\mu)^n (m_b/\mu)^m \quad (\text{set } m_b = 0) , \quad p_T(b) \text{ (at LO)}$$

Resum to all orders

$$\alpha_s(\mu)^n \ln^n(m_b/\mu) \quad \text{into the } b \text{ PDF}$$

4-flavour scheme (4F): generate b 's in the hard scattering



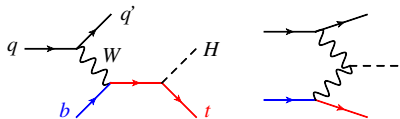
Fixed order

$$\alpha_s(\mu)^n (m_b/\mu)^m , \quad \alpha_s(\mu)^n \ln^n(m_b/\mu) , \quad p_T(b)$$

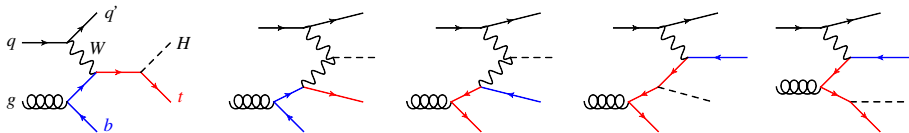
Go to NLO to reduce differences (i.e. flavour-scheme dependence)

Single top plus Higgs production (tH)

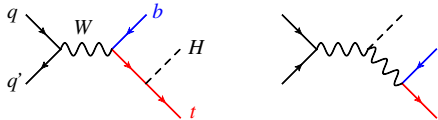
t -channel 5F



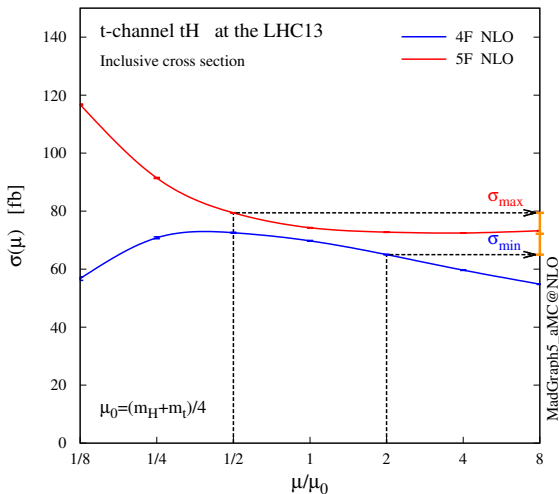
t -channel 4F



s -channel 4F



tH NLO cross section: flavour scheme (FS) and scale (μ) dependence



FS and μ are correlated

(1) Choose μ_0 to minimize FS dependence (at NLO)

(2) Vary μ/μ_0 by 2 and take the envelope

(3) Define

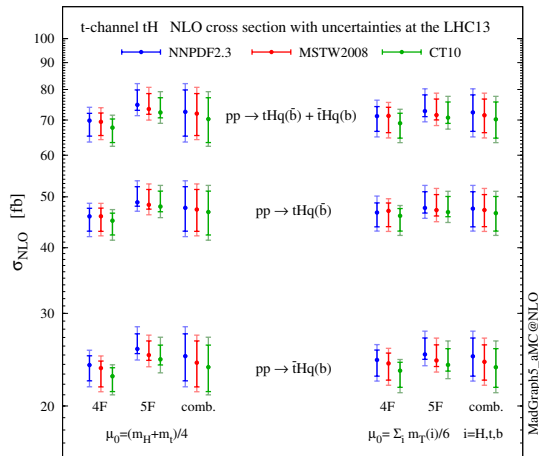
$$\sigma_0 \equiv (\sigma_{\max} + \sigma_{\min})/2$$

$$\delta_{\mu+FS} \equiv (\sigma_{\max} - \sigma_{\min})/2$$

$$\sigma_{\text{NLO}}^{\text{comb.}} = \sigma_0 \pm \delta_{\mu+FS}$$

$$\mu_R = \mu_F \equiv \mu$$

tH NLO cross section: total theoretical uncertainty



PDF global fits:
NNPDF2.3, MST2008, CT10

$$\alpha_s^{(5F)}(m_Z) = 0.1190 \pm 0.0012$$

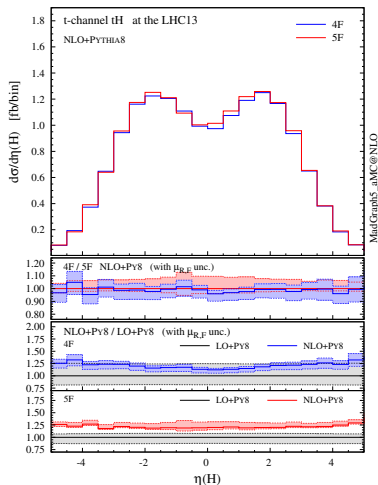
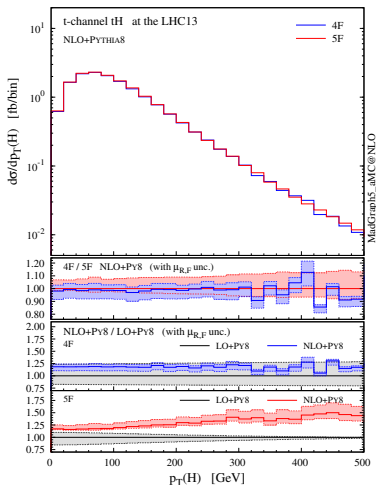
$$\Leftrightarrow \alpha_s(m_b) = 0.2189 \pm 0.0042$$

$$m_b^{(\text{o.s.})} = 4.75 \pm 0.25 \text{ GeV}$$

$$\delta_{\text{PDF}+\alpha_s+m_b}^{\pm} = \sqrt{(\delta_{\text{PDF}}^{\pm})^2 + (\delta_{\alpha_s}^{\pm})^2 + (\delta_{m_b}^{\pm})^2}$$

$$\delta_{\text{tot}}^{\pm} = \delta_{\mu(+\text{FS})}^{\pm} + \delta_{\text{PDF}+\alpha_s+m_b}^{\pm}$$

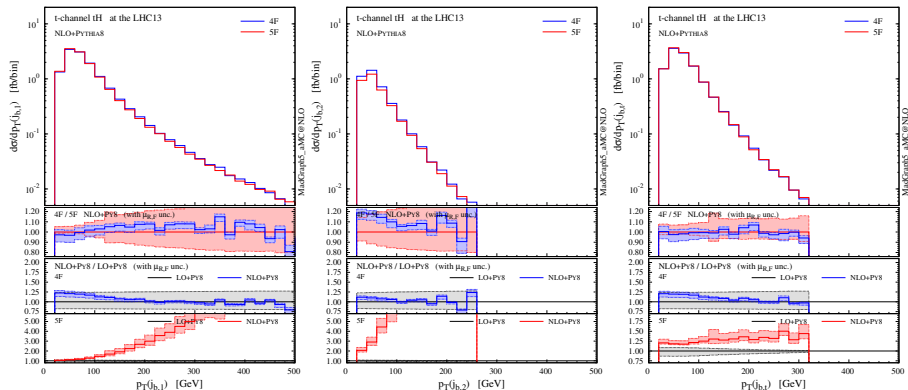
tH NLO+PS distributions: Higgs (no decay)



Very good agreement between 4F and 5F shapes at NLO (same for the top quark)
LO 5F provides a poor description

tH NLO+PS distributions: b -tagged jets

anti- k_T $p_T > 30$ GeV $R = 0.4$ $|\eta| < 2.5$ no H decay, $t \rightarrow b\ell\nu_\ell$

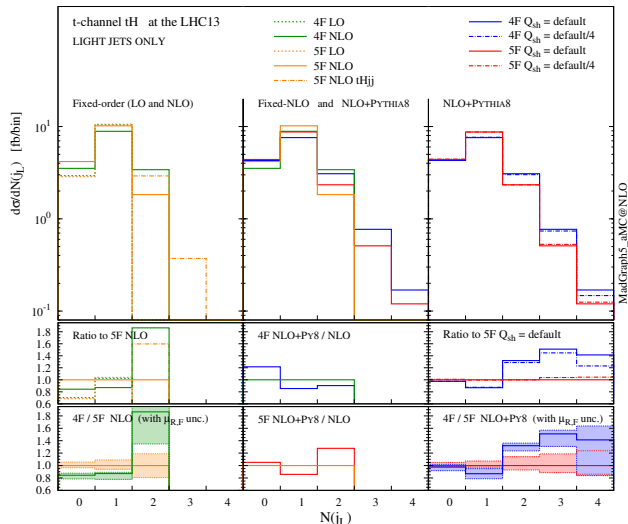


Still good agreement at NLO

1st b -jet comes from top decay in the low- p_T region,
from hard $g \rightarrow b\bar{b}$ splitting in the high- p_T tail

tH NLO+PS distributions: light jets

anti- k_T $p_T > 30$ GeV $R = 0.4$ $|\eta| < 4.5$ no H decay, $t \rightarrow b\ell\nu_\ell$

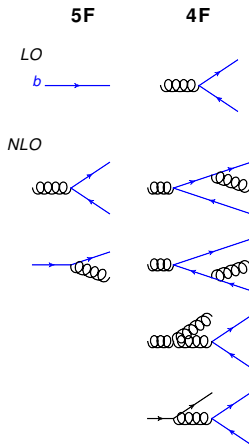
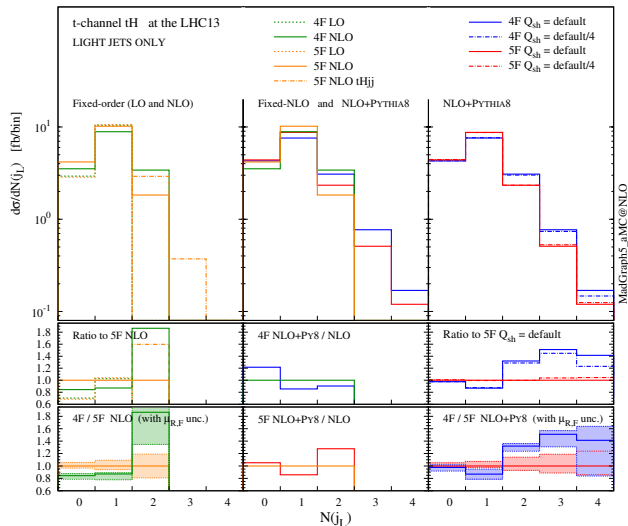


4F/5F discrepancy starts at fixed-NLO

It's reduced after the shower

It's related to the description of the $g \rightarrow b\bar{b}$ splitting

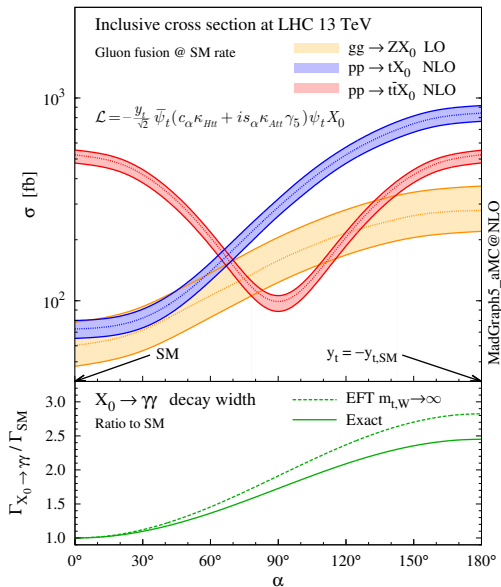
tH NLO+PS distributions: light jets



Conclusions (1)

- ① Both 4F and 5F are good to describe the NLO cross section
 $\Rightarrow$ combine
- ② Carefully assessed the major sources of uncertainty in the cross section
 $\sigma_{\text{NLO}}^{t\text{-channel}} \simeq 72 \text{ fb} \pm 8\% (\mu + \text{FS}) \pm 3\% (\text{PDF} + \alpha_s + m_b)$
(larger PDF uncertainty if considering systematics in the fitting procedure)
- ③ Both 4F and 5F are good to describe Higgs and top distributions
(at NLO!)
- ④ 4F at NLO describes correctly many more distributions

Higgs characterisation: inclusive observables



α = CP-mixing angle

$$\kappa_{Htt} = 1, \quad \kappa_{Att} = 2/3$$

$\Rightarrow$ Gluon fusion @ SM rate $\forall \alpha$

$t\bar{t}H$ helps to constrain α

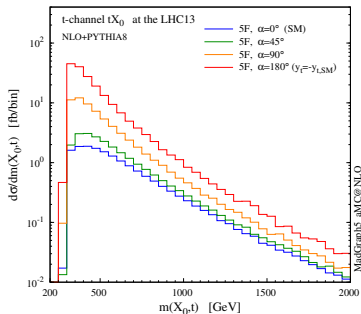
tH , $gg \rightarrow ZH$ and $H \rightarrow \gamma\gamma$

further lift 2 degeneracies:

$$y_t \leftrightarrow -y_t$$

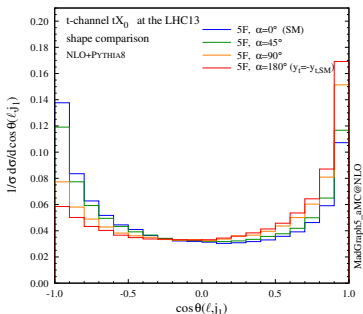
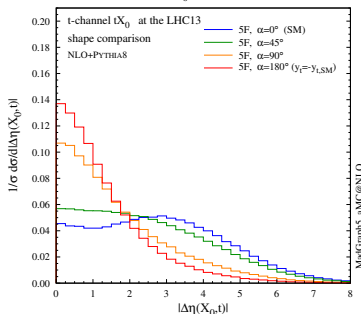
$$\alpha \leftrightarrow \pi - \alpha$$

Higgs characterisation: tH distributions



New physics takes place mostly at threshold

Many observables are sensitive to n.p.!



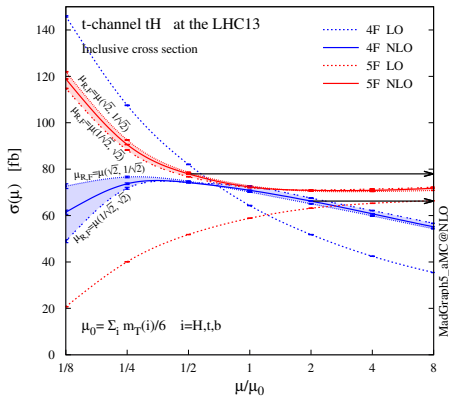
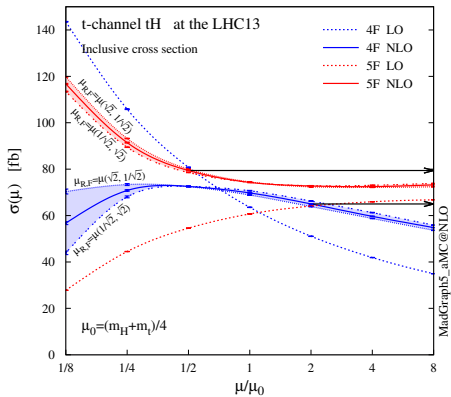
Conclusions (2)

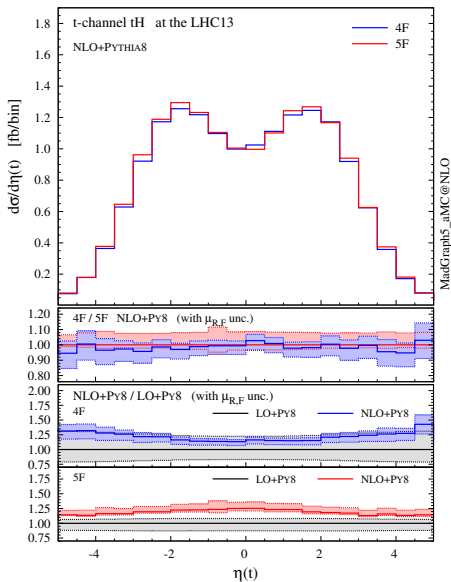
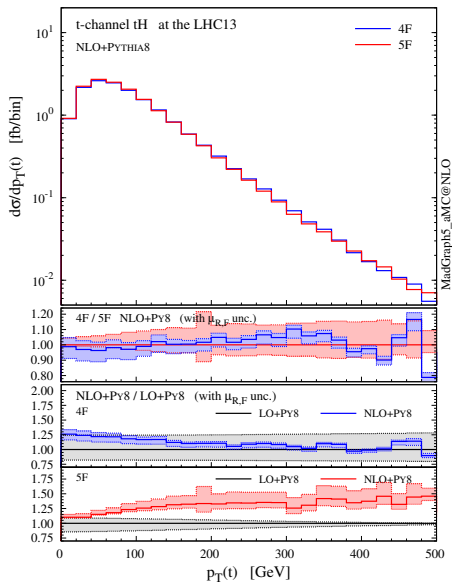
- ① Global analysis of different Higgs productions/decays
improve measurements and lift degeneracies in GF or $t\bar{t}H$
- ② Higgs plus single top is a rare process, but provides
a variety of observables sensitive to BSM physics
 $\Rightarrow$ very interesting for Higgs characterisation

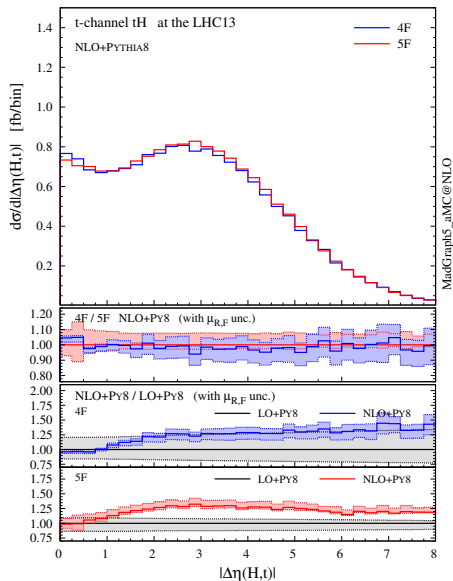
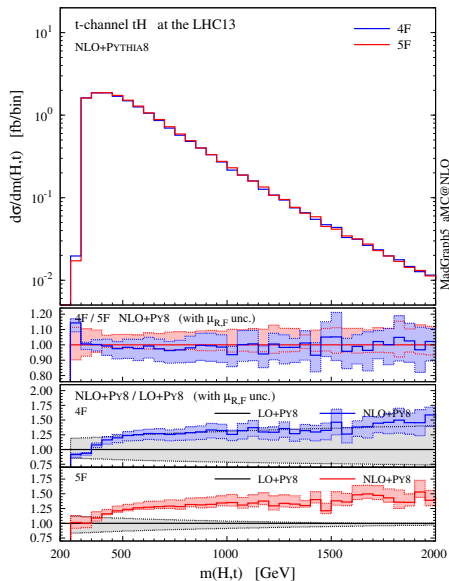
NB: Pheno studies are fully automatic, importing the HC_NLO_X0 model
inside the MadGraph5_aMC@NLO framework

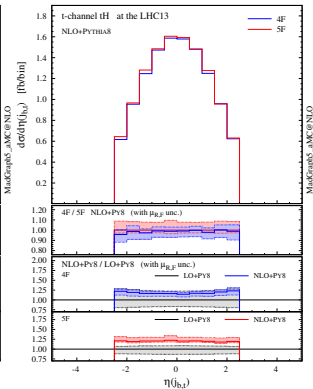
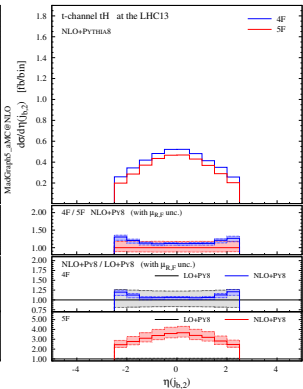
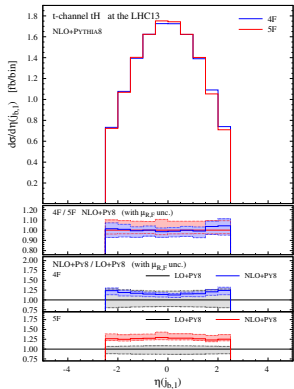
Thanks for your attention!

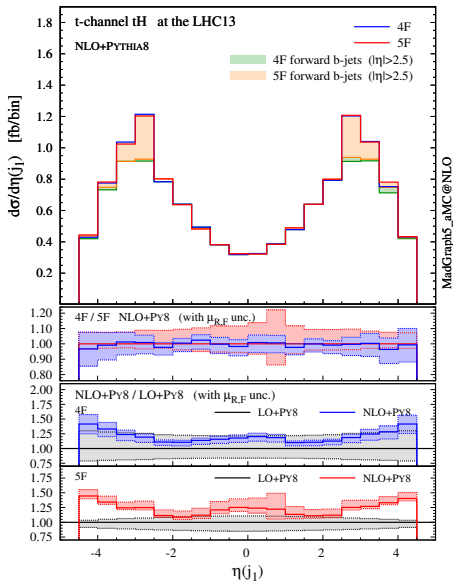
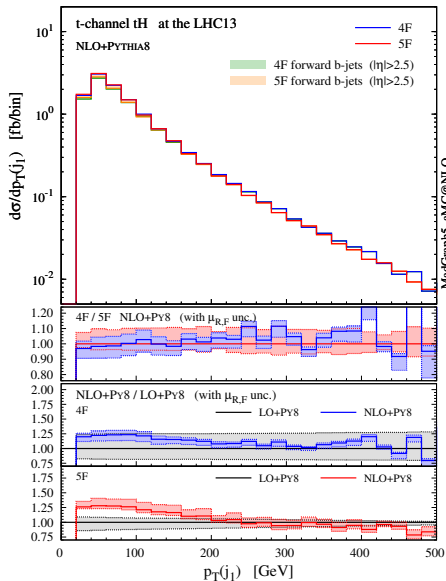
backup slides

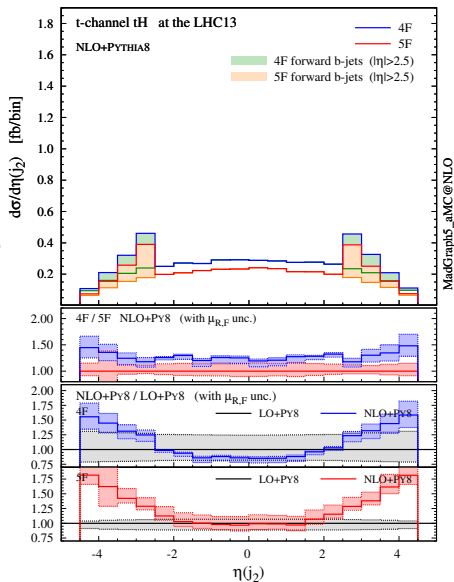
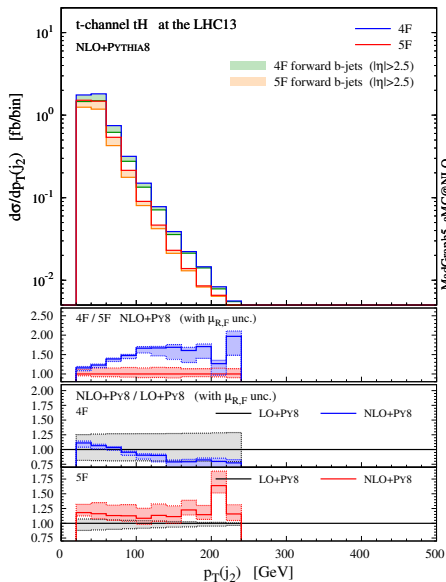












s -channel ($\simeq 3$ fb) $\ll$ t -channel, but very different shapes!

