



GoSam

A program for the computation of one loop virtual
contribution to scattering amplitudes

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on behalf of the **GoSam** collaboration:
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MadGraph workshop 2011, Academia Belgica - Roma

Automatic tools for the computation of virtual one loop correction

- ▶ FeynArts/FormCalc/LoopTools (**public**) Thomas Hahn et al
- ▶ Helac-NLO Bevilacqua, Czakon, van Hameren, Papadopoulos, Pittau, Worek
- ▶ MadLoop Hirschi, Frederix, Frixione, Garzelli, Maltoni, Pittau
uses **CutTools** (**public** [Ossola, Papadopoulos, Pittau]) and **MadGraph**
- ▶ Golem-Samurai (**Samurai public** [Mastrolia, Ossola, Reiter, Tramontano])
Cullen, Greiner, GH, Luisoni, Mastrolia, Ossola, Reiter, Tramontano
- ▶ NGluon (**public**) [Badger, Biedermann, Uwer]
- ▶ dedicated programs also involve high level of automation
Denner, Dittmaier, Pozzorini et al, VBFNLO coll., MCFM, Blackhat, Rocket, ...

NEW DEVELOPMENTS IN ONE LOOP COMPUTATIONS

❑ Pioneering works:

- Improvements in the computation of tensor integrals
Binoth et al GOLEM95
Denner & Dittmaier
- Application of unitarity to the computation of one loop amplitudes
Bern, Dixon, Kosower
Britto, Cachazo, Feng
- Reduction at the integrand level
Ossola, Papadopoulos, Pittau
Ellis, Giele, Kunszt, Melnikov

OPP integrand decomposition: 4-dim

- At integrand level the structure is enriched by polynomial terms that integrate to zero (I multiplied with all the propagators)

$$\begin{aligned}
 N(q) = & \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[d(i_0 i_1 i_2 i_3) + \tilde{d}(q; i_0 i_1 i_2 i_3) \right] \prod_{i \neq i_0, i_1, i_2, i_3}^{m-1} D_i + \sum_{i_0 < i_1 < i_2}^{m-1} \left[c(i_0 i_1 i_2) + \tilde{c}(q; i_0 i_1 i_2) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\
 & + \sum_{i_0 < i_1}^{m-1} \left[b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1}^{m-1} D_i + \sum_{i_0}^{m-1} \left[a(i_0) + \tilde{a}(q; i_0) \right] \prod_{i \neq i_0}^{m-1} D_i
 \end{aligned}$$

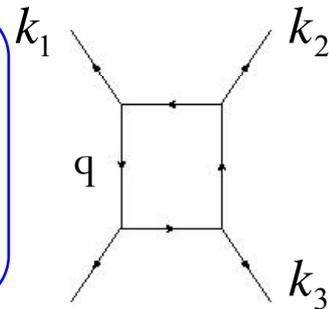
- A choice of q fulfilling 4-ple cut condition: $D_{i_0} = D_{i_1} = D_{i_2} = D_{i_3} = 0$ will single out just one polynomial

$$\Delta_{i_0 i_1 i_2 i_3} = \left[d(i_0 i_1 i_2 i_3) + \tilde{d}(q; i_0 i_1 i_2 i_3) \right]$$

\tilde{d} can **only** be of the type $q \cdot p$

where $p = \varepsilon_{\alpha\beta\gamma} k_1^\alpha k_2^\beta k_3^\gamma$

[proof in OPP 2007]



- Once fitted such polynomial we can subtract it from both sides and repeat the game with another multiple cut condition -> recursive solution
- For each phase space point the only requirement for the reduction is the knowledge of the numerical value of the numerator function N for a small set of values of the loop momentum variable, solutions of the multiple cut conditions

Extension to D-dim

- fix a parametric form for the loop momentum in terms of a linear combination of four known 4-vectors e_i suitably chosen

$$\bar{q} = q + \mu \quad \bar{q}^2 = q^2 - \mu^2 \quad q = x_1 e_1 + x_2 e_2 + x_3 e_3 + x_4 e_4$$

the vanishing term (spurious term in the OPP terminology) are then polynomials of x_i and μ^2

- The problem is to fit the coefficients in the polynomials Δ

$$\begin{aligned} N(\bar{q}) = & \sum_{i << m}^{n-1} \Delta_{ijklm}(\bar{q}) \prod_{h \neq i,j,k,\ell,m}^{n-1} \bar{D}_h + \sum_{i << \ell}^{n-1} \Delta_{ijkl}(\bar{q}) \prod_{h \neq i,j,k,\ell}^{n-1} \bar{D}_h + \\ & + \sum_{i << k}^{n-1} \Delta_{ijk}(\bar{q}) \prod_{h \neq i,j,k}^{n-1} \bar{D}_h + \sum_{i < j}^{n-1} \Delta_{ij}(\bar{q}) \prod_{h \neq i,j}^{n-1} \bar{D}_h + \sum_i^{n-1} \Delta_i(\bar{q}) \prod_{h \neq i}^{n-1} \bar{D}_h \end{aligned}$$

- Example: 3-ple cut residue

$$\begin{aligned} \Delta_{ijk}(\bar{q}) = & c_{3,0}^{(ijk)} + c_{3,7}^{(ijk)} \mu^2 - \left((c_{3,1}^{(ijk)} + c_{3,8}^{(ijk)} \mu^2) x_4 + (c_{3,4}^{(ijk)} + c_{3,9}^{(ijk)} \mu^2) x_3 \right) (e_1 \cdot e_2) + \\ & + \left(c_{3,2}^{(ijk)} x_4^2 + c_{3,5}^{(ijk)} x_3^2 \right) (e_1 \cdot e_2)^2 - \left(c_{3,3}^{(ijk)} x_4^3 + c_{3,6}^{(ijk)} x_3^3 \right) (e_1 \cdot e_2)^3 . \end{aligned}$$

- with the 3 cut conditions: $D_i = D_j = D_k = 0$ one fixes x_1, x_2 and the product $x_3 x_4$

Amplitudes & Master Integrals

$$\begin{aligned}
 \mathcal{A}_n = & \sum_{i < j < k < \ell}^{n-1} \left\{ c_{4,0}^{(ijkl)} I_{ijkl}^{(d)} + \frac{(d-2)(d-4)}{4} c_{4,4}^{(ijkl)} I_{ijkl}^{(d+4)} \right\} & \int d^d \bar{q} \frac{\bar{q} \cdot e_2}{\bar{D}_i \bar{D}_j} = J_{ij}^{(d)} \\
 & + \sum_{i < j < k}^{n-1} \left\{ c_{3,0}^{(ijk)} I_{ijk}^{(d)} - \frac{(d-4)}{2} c_{3,7}^{(ijk)} I_{ijk}^{(d+2)} \right\} & \int d^d \bar{q} \frac{(\bar{q} \cdot e_2)^2}{\bar{D}_i \bar{D}_j} = K_{ij}^{(d)} \\
 & + \sum_{i < j}^{n-1} \left\{ c_{2,0}^{(ij)} I_{ij}^{(d)} + c_{2,1}^{(ij)} J_{ij}^{(d)} + c_{2,2}^{(ij)} K_{ij}^{(d)} - \frac{(d-4)}{2} c_{2,9}^{(ij)} I_{ij}^{(d+2)} \right\} \\
 & + \sum_i^{n-1} c_{1,0}^{(i)} I_i^{(d)}
 \end{aligned}$$

The sources of rational terms are the integrals with μ^2 powers in the numerator

They are generated by the reduction algorithm (R1), but could also be present ab initio in the numerator function as a consequence of the d-dimensional algebraic manipulations (R2)

$$\begin{aligned}
 \int d^d \bar{q} \frac{\mu^2}{\bar{D}_i \bar{D}_j} &= -\frac{(d-4)}{2} I_{ij}^{(d+2)} \\
 \int d^d \bar{q} \frac{\mu^4}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_\ell} &= \frac{(d-2)(d-4)}{4} I_{ijkl}^{(d+4)} \\
 \int d^d \bar{q} \frac{\mu^2}{\bar{D}_i \bar{D}_j \bar{D}_k} &= -\frac{(d-4)}{2} I_{ijk}^{(d+2)}
 \end{aligned}$$

On-shell methods are quite flexible:

❑ Different implementations

➤ Sewing tree level amplitudes

+ works with gauge invariant objects

- still not easy to automate the rational terms for general one loop amplitudes

Blackhat: Recursive bootstrap approach [Ita, Bern, Dixon, Febres Cordero, Kosower, Maitre]

Rocket: Tree level amplitudes in different dimensions [Ellis, Giele, Kunszt, Melnikov, Zanderighi]

Samurai: works with D-dimensional tree level amplitudes [Mastrolia, Ossola, Reiter, FT]

➤ Diagrammatic approach:

+ contain all the information on the rational terms

- single diagrams are not gauge invariant objects: big cancellations

❖ Fully numerical construction through the single cut:

Helac-NLO [Bevilacqua, Czakon, Papadopoulos, Worek]

MadLoop [Hirschi, Frederix, Frixione, Garzelli, Maltoni Pittau]

❖ Fully algebraic method: **GoSam**

[Cullen, Greiner, Heinrich, Luisoni, Mastrolia, Ossola, Reiter, FT]

Main features of the “Algebraic Way”:

- Amplitudes generated with Feynman diagrams
- Algebraic manipulations are allowed before starting the numerical integration
- The generation of numerators is executed separately from the numerical reduction
- Optimization: grouping of diagrams, smart caching
- Control over sub-parts of the computation (move in/out subsets of diagrams)
- Algebra in dimension d , different schemes

Great flexibility in the reduction
Choice between different algorithms at runtime

GOLEM/SAMURAI

Algebraic generation of **d-dimensional integrands** via **Feynman diagrams**

Reduction at the Integrand Level: **d-dimensional extension of OPP reduction**

Target: provide an **automated tool** for stable evaluation of one-loop matrix elements

- **be general/model independent** (QCD, EW, BSM)
- **interface with existing tools** like MadEvent, Sherpa, PowHEG, ...
- build upon **open source tools only** (i.e. Samurai, Golem95, QGraf, Form, Spinney, Haggies, QCDDLoop, OneLOop)
- **support open standards**

GoSam

An automated amplitude generation based on Feynman diagrams
(distributed as a python package)

- FORM
[J.A.M. Vermaseren](#), (1991)
- QGRAF [Fortran – open source]
[P. Nogueira](#), (1993)
- Haggies [Java – open source]
[T. Reiter](#), (2009)
- Spinney [Form – open source]
[Cullen, Koch-Janusz, Reiter](#), (2010)

Default Option: **Samurai**

P.Mastrolia, G.Ossola, T.Reiter, FT (2010)

- OPP Reduction Algorithm G.Ossola, C. Papadopoulos, R, Pittau (2007)
- d-dimensional extension Ellis, Giele, Kunszt, Melnikov (2008)
- Coefficients of Polynomials via DFT Mastrolia et al. (2008)
- Model-independent Computation of the full Rational Term

Samurai-2.1 will be public soon:

- ✓ complex masses
- ✓ scalar integral caching
- ✓ code optimization (w Haggies)
- ✓ tensorial reconstruction
- ✓ reduced samplings
- ✓ passing of the invariants
- ✓ updated scalar intergral library
- ✓ bug fixes

Other options available (at runtime):

Golem95

Binoth, Cullen, Guillet, Heinrich, Kleinschmidt, Pilon, Reiter, Rodgers (2008)

Tensorial Integrand-level Reduction

G, Heinrich, G.Ossola, T.Reiter, FT (2010),

Last Step: multiply all coefficients with the corresponding Master Integral

QCDloop
(**Ellis, Zanderighi**)

OneLOop
(**A. van Hameren**)

Golem95C
(**Cullen, Guillet, Heinrich, Kleinschmidt, Pilon, Reiter, Rodgers**)
Upgrade of *Golem95* library, real and **complex masses** supported

<http://projects.hepforge.org/~golem/95/>

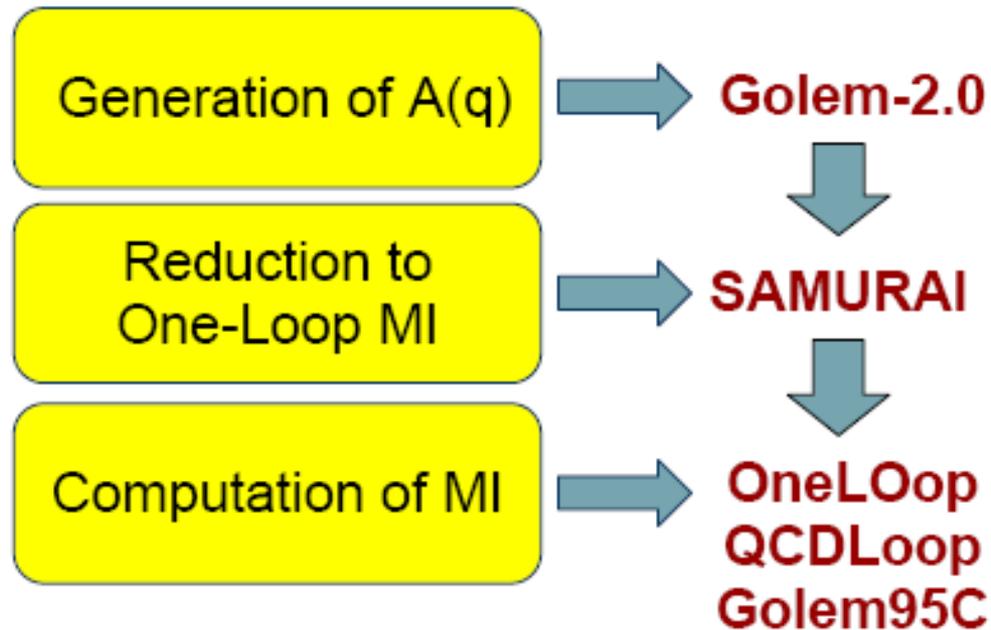
NEW!

LoopTools
(**T. Hahn**)

NEW!

PJFry
(**V. Yundin**)

Standard GoSam reduction:



This process is **fully automated**

Walking through GoSam...

Preparation of the “card”: we use as example $u\bar{d} \rightarrow \bar{s}c e^- \bar{\nu}_e \mu^+ \nu_\mu$

```
#!/bin/env /usr/local/bin/golem-main.py
process_name=nloop4
process_path=/home/tramonta/codici/processes/nloop4
                                [directory previously created]

in=u,d~
out=s~,c,e-,nebar,mu+,nmu
model=smdiag
                                [model can be generated through FeynRules and LanHEP]

model.options=mW=80.376,mZ=91.1876,GF=0.000116639
                                [for sm and smdiag the other couplings will be generated]

order=gw,4,4
order=gs,2,4

zero=mC,mS,mU,mD,mmu,me
one=gs,e

helicities=-+--+--+

qgraf.bin=/home/tramonta/codici/QGRAF/qgraf
qgraf.options=onshell,nosnail,notadpole
qgraf.verbatim=\
  true = iprop[phim,phip,chi,H,B,T,0,0];\n\
  true = iprop[Wp,2,2];

extensions=samurai,dred          [golem95, pjfry]
r2=explicit                      [implicit, only, none]
abbrev.limit=500
abbrev.level=diagram             [group, helicity]

form.bin=form
fc.bin=gfortran -O2
```

... and also

- ❖ Ready to use UFO model files: there will be examples in the release
 - Support up to spin 2 but this part is still untested
- ❖ Python framework for deep level diagram selection
 - useful to split the computation conveniently
- ❖ Very basic existing interface with POWHEG
 - spin correlated and color correlated born matrix elements available

Where the card is:

```
$ golem-main nloop4.rc
```

In the process directory

```
$ make source
```

```
$ make compile
```

```
$ cd doc
```

```
$ make
```

Documentation:

GoSam: $u\bar{d} \rightarrow \nu_\mu\mu^+e^-\bar{\nu}_e\bar{s}c$
Diagrams

tramonta

2011-09-20 (19:55:21)

Abstract

This process consists of 4 tree-level diagrams and 86 NLO diagrams. GoSam has identified 6 groups of NLO diagrams by analyzing their one-loop integrals.

Documentation:

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Documentation:

Building the code : check the details before the run

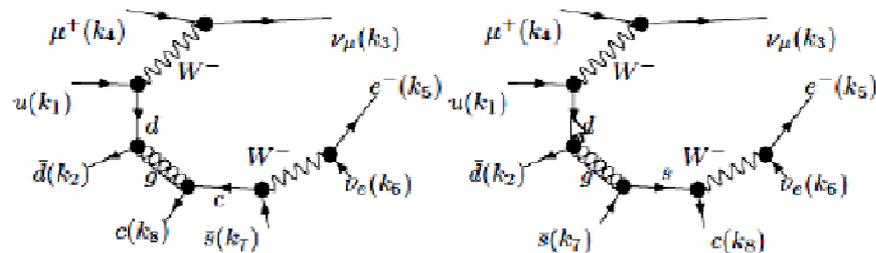


Diagram 1

Diagram 2

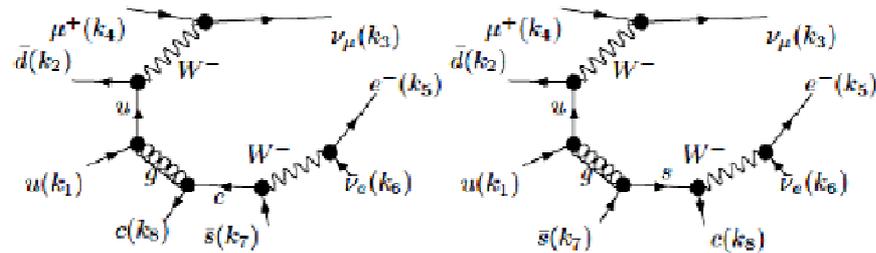


Diagram 3

Diagram 4

Building the code : check the details before the run

Documentation:

5.5.1 Diagrams (69)

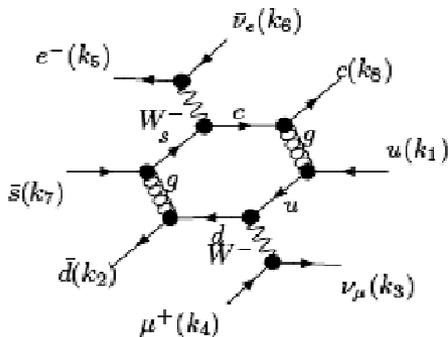


Diagram 2

$$S' = S_{Q \rightarrow -q}^{[2]}(-k_8 - k_7 - k_6 - k_5), \text{ rk} = 4$$

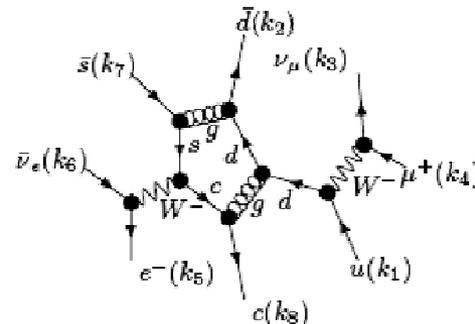


Diagram 3

$$S' = S_{Q \rightarrow -q}^{[6]}, \text{ rk} = 3$$

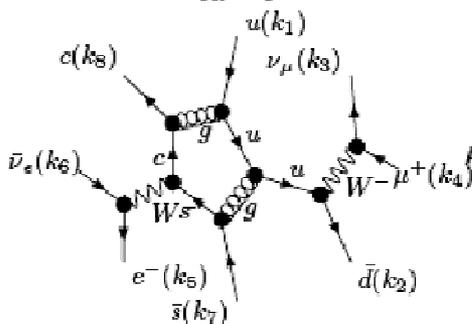


Diagram 6

$$S' = S_{Q \rightarrow q}^{[1]}(k_8 + k_7 + k_6 + k_5), \text{ rk} = 3$$

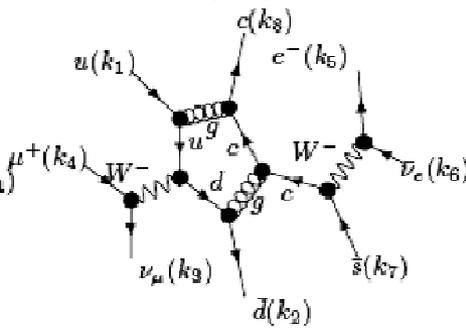


Diagram 7

$$S' = S_{Q \rightarrow -q}^{[8]}(-k_8 - k_7 - k_6 - k_5), \text{ rk} = 3$$

- Loop diagrams are grouped into sets of diagrams which share loop propagators
 - The numerator is constructed dynamically cut by cut
 - Strongly reduced calls to samurai
 - Caching of the scalar integral and of a number of things calculated by samurai

Building the code : Spinney+ Haggies

```
tramonta@nbtramontano: ~/codici/processes/nloop4
File Modificati tramonta@nbtramontano: ~/codici/processes/nloop4

tramonta@nbtramontano: ~/codici/processes/nloop4
0.01 sec out of 0.01 sec
haggies is generating color.f90
haggies is generating model.f90
make[3]: uscita dalla directory «/home/tramonta/codici/processes/nloop4/common»
make[2]: uscita dalla directory «/home/tramonta/codici/processes/nloop4/common»
make[2]: ingresso nella directory «/home/tramonta/codici/processes/nloop4/helicity0»
make -f Makefile.source source
make[3]: ingresso nella directory «/home/tramonta/codici/processes/nloop4/helicity0»
Form is processing tree diagram 1 @ Helicity 0
0.38 sec out of 0.37 sec
Form is processing tree diagram 2 @ Helicity 0
0.38 sec out of 0.37 sec
Form is processing tree diagram 3 @ Helicity 0
0.37 sec out of 0.38 sec
Form is processing tree diagram 4 @ Helicity 0
0.39 sec out of 0.39 sec
Haggies is processing tree level diagrams @ Helicity 0
Form is processing loop diagram 1 @ Helicity 0
1.30 sec out of 1.30 sec
Haggies is processing abbreviations for loop diagram 1 @ Helicity 0
Form is processing loop diagram 2 @ Helicity 0
1.57 sec out of 1.58 sec
Haggies is processing abbreviations for loop diagram 2 @ Helicity 0
Form is processing loop diagram 3 @ Helicity 0
1.23 sec out of 1.24 sec
Haggies is processing abbreviations for loop diagram 3 @ Helicity 0
Form is processing loop diagram 4 @ Helicity 0
```

Execution: all the code is ready, available as a library - option: Autotool



```
tramonta@nbtramontano: ~/codici/processes/loop4
File Modifica Visualizza Terminale Schede Aiuto

tramonta@nbtramontano: ~/Scrivania
tramonta@nbtramontano:~/codici/processes/loop4$ ls
codegen  diagrams-0.hh  diagrams-1.log  helicity0  Makefile.source  model.hh  process.hh  pyxovirt.log  topotree.py
common  diagrams-0.log  doc             Makefile   matrix           model.py  pyxotree.log  pyxovirt.tex  topotree.pyc
config.sh diagrams-1.hh  func.txt       Makefile.conf  model          model.pyc  pyxotree.tex  topotree.log  topovirt.log
tramonta@nbtramontano:~/codici/processes/loop4$ ls matrix/
debug.xml  loop4_matrix.mod  ltest.dat  Makefile  matrix.a  matrix.f90  matrix.o  param.dat  test.exe  test.f90  test.o
tramonta@nbtramontano:~/codici/processes/loop4$
```

```
#          LO: 0.4679522923780973E-19
# NLO, finite par -15.91575130714159
# NLO, single pol  7.587050524202378
# NLO, double pol -5.3333333333333385
# IR, single pol  -0.746282775357444
# IR, double pol  -5.3333333333333333
```

Melia, Melnikov, Rontsch, Zanderighi

```
# NLO, finite par -15.91575
# NLO, single pol  7.587051
# NLO, double pol -5.333333
```

After renormalization:
check of the poles with the one
form integrated dipoles built-in

```
#          LO: 0.4679522923780973E-19
# NLO, finite par -14.91575130714159
# NLO, single pol -0.746282809130955
# NLO, double pol -5.3333333333333385
# IR, single pol  -0.746282775357444
# IR, double pol  -5.3333333333333333
```

Timing (2.27 GHz, Gfortran -O2): 5.4ms

ALTERNATIVE PATH: THE “TENSORIAL WAY”

Tensorial Reconstruction at the Integrand Level

[with G.Heinrich, G.Ossola and T.Reiter (2010)]

In this work:

- We **tested** the methods for the **detection of instabilities**
- We proposed a “**rescue-system**” alternative to higher precision routines
- We proposed an **optimized reconstruction method**

Idea: **tensorial reconstruction** performed **at the integrand level** by means of a **sampling in the integration momentum**.

$$\mathcal{N}(q) = \sum_{r=0}^R C_{\mu_1 \dots \mu_r} q_{\mu_1} \dots q_{\mu_r} \implies \hat{\mathcal{N}}(q)$$

$\hat{\mathcal{N}}(q)$ is the “reconstructed numerator” written as a tensor
– numerically identical to the initial $\mathcal{N}(q)$ –

Tensorial reconstruction

rewrite numerator function as a linear combination of tensors

$$\mathcal{N}(q) = \sum_{r=0}^R C_{\mu_1 \dots \mu_r} q_{\mu_1} \dots q_{\mu_r}$$

$$C_{\mu_1 \dots \mu_r} q_{\mu_1} \dots q_{\mu_r} = \sum_{(i_1, i_2, i_3, i_4) \vdash r} \hat{C}_{i_1 i_2 i_3 i_4}^{(r)} \cdot (q_1)^{i_1} (q_2)^{i_2} (q_3)^{i_3} (q_4)^{i_4}$$

determine the coefficients by sampling q in a bottom-up approach

Level 0:

$$q = (0, 0, 0, 0), \mathcal{N}(0, 0, 0, 0) \equiv \mathcal{N}^{(0)} = C_0$$

Level 1: 4 systems, each sampling a monomial depending on one component of q only

$$\mathcal{N}^{(1)}(q) \equiv \mathcal{N}(q) - \mathcal{N}^{(0)}$$

$$q = (x, 0, 0, 0) \Rightarrow \mathcal{N}^{(1)}(q) \equiv x C_1 + x^2 C_{11} + \dots + x^R \underbrace{C_{11 \dots 1}}_{R \text{ times}}$$

$$q = (0, y, 0, 0) \Rightarrow \mathcal{N}^{(1)}(q) \equiv y C_2 + y^2 C_{22} + \dots + y^R \underbrace{C_{22 \dots 2}}_{R \text{ times}}$$

$x, y \dots$ can be chosen to be **Real numbers** !

No paradigm that the amplitudes need to be sampled with complex q

PRECISION TESTS

Use the decomposition of the numerator function $N(\bar{q})$ after determining all coefficients

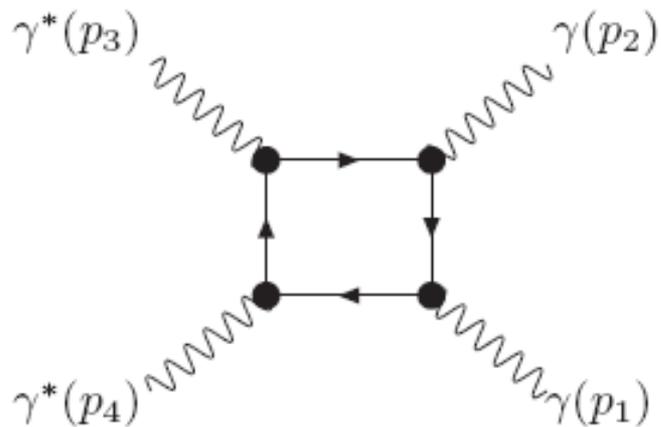
$$\begin{aligned} N(\bar{q}) = & \sum_{i \ll m}^{n-1} \Delta_{ijklm}(\bar{q}) \prod_{h \neq i,j,k,\ell,m}^{n-1} \bar{D}_h + \sum_{i \ll \ell}^{n-1} \Delta_{ijkl}(\bar{q}) \prod_{h \neq i,j,k,\ell}^{n-1} \bar{D}_h + \\ & + \sum_{i \ll k}^{n-1} \Delta_{ijk}(\bar{q}) \prod_{h \neq i,j,k}^{n-1} \bar{D}_h + \sum_{i < j}^{n-1} \Delta_{ij}(\bar{q}) \prod_{h \neq i,j}^{n-1} \bar{D}_h + \sum_i^{n-1} \Delta_i(\bar{q}) \prod_{h \neq i}^{n-1} \bar{D}_h \end{aligned}$$

- 1 Global ($N = N$)-test
- 2 Local ($N = N$)-test
- 3 Power-test

Are those methods **reliable** in detecting **unstable phase space points**?

APPROACHING THE GRAM - I

- We approach a kinematic configuration which can lead to large cancellations
- Fermion loop with two massless and two massive vector particles



$$\begin{aligned} p_{1,2} &= (E, 0, 0, \pm E) & p_{1,2}^2 &= 0 \\ p_{3,4} &= (E, 0, \pm Q \sin \theta, \pm Q \cos \theta) \\ p_{3,4}^2 &= m^2 \\ E &= \sqrt{m^2 + Q^2} \end{aligned}$$

- The Gram-det vanishes when $Q \rightarrow 0$ (m and θ are fixed)

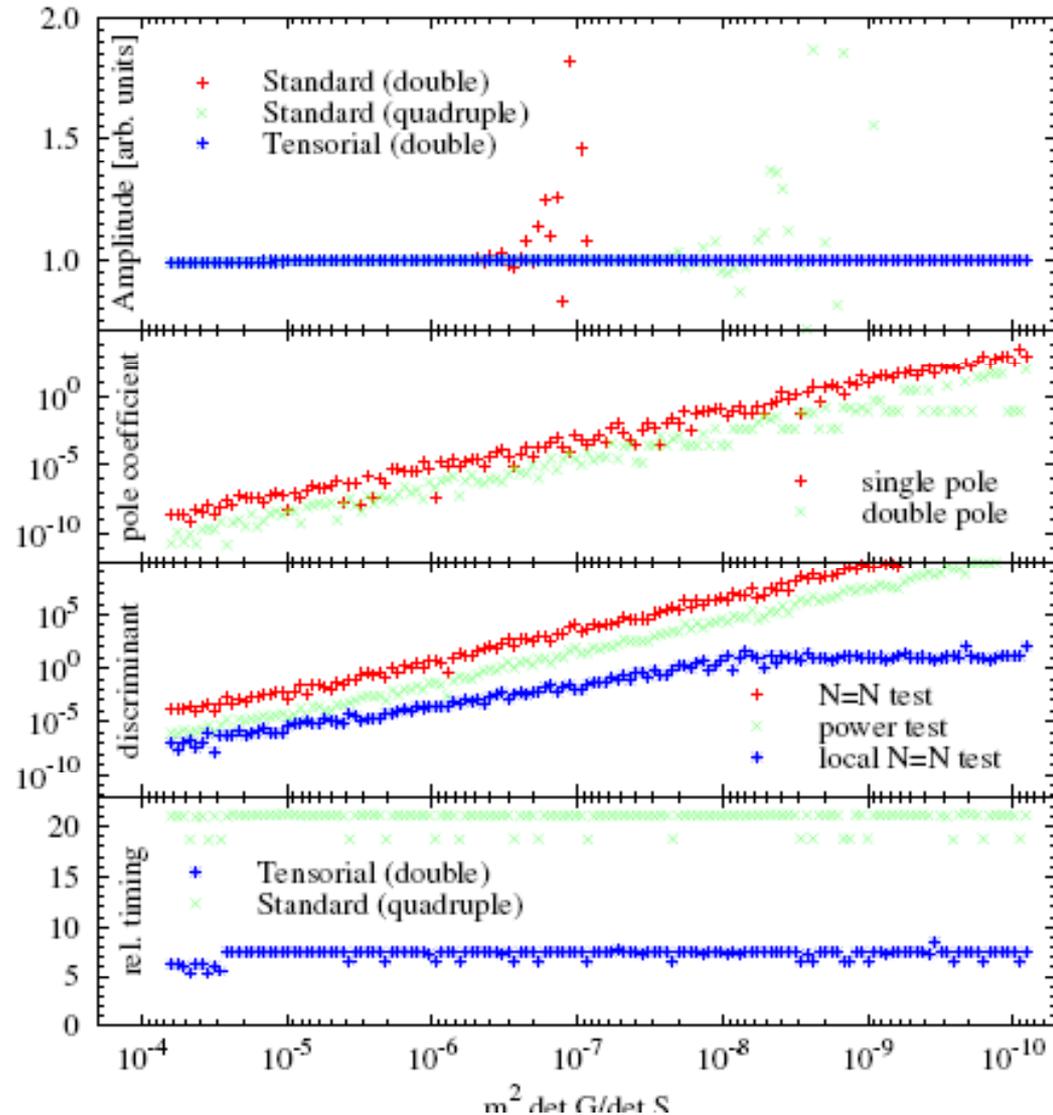
$$\det G = 32 E^4 Q^2 \sin^2 \theta$$

APPROACHING THE GRAM - II

Standard(double):
2x SAMURAI

Standard(quadruple):
2x Kin + integrals
4x Algorithm

Tensorial(double):
Reconstruction paired
with numerical evaluation
of tensor integras with
GOLEM95



Ways to use the tensorial reconstruction

■ “Rescue-system”

- Unstable points will be automatically reprocessed using the tensorial decomposition + tensor integrals with `Go1em 95`
- Tensorial “master” integrals appears to be less costly than multi-precision routines

■ “Hybrid method” for improved timing

- The reduction of $\hat{N}(q)$ can be faster than that of $\mathcal{N}(q)$

# Lines	Time ratio “hybrid” / standard	
	Rank = 4	Rank = 6
1	1.3	1.6
10	1.1	1.4
100	0.51	0.85
1000	0.30	0.59
10000	0.27	0.55

EXAMPLE: ALTERNATIVE REDUCTION PATHS

Samurai/Tensorial Reduction/Golem95

$$u\bar{u} \rightarrow d\bar{d}$$

- 1 Evaluation with Samurai, sampling of diagram groups
- 2 Evaluation with Samurai, sampling of individual diagrams
- 3 Tensorial Reconstruction + Reduction of numetens with Samurai
- 4 Evaluation with Golem95

Method	finite part	single pole	double pole
1	-3.433053565229151	-14.62937842683104	-5.333333333333338
2	-3.433053565229129	-14.62937842683102	-5.333333333333342
3	-3.433053565229163	-14.62937842683104	-5.333333333333342
4	-3.433053565229146	-14.62937842683102	-5.333333333333332

More on the rational terms:

- ❑ Treatment strictly related the way the numerator function is furnished
 - Classified in two categories: $R = R1 + R2$
- ❑ R1 develops automatically performing the D-dimensional reduction of the tensors spanning the 4-dimensional part of the loop momentum
- ❑ R2 are present in the UV diagrams: bubbles, rank 2 and 3 triangles and rank4 boxes.
- ❑ At least two possibilities for R2 automatic computation:
 - for any fixed gauge theory calculate once and for all the contribution from all the diagrams that can generate R2 terms and define a set of tree level Feynman rules that give the R2 contribution for any process: **MadLoop approach**
 - Alternatively: construct the numerator function by implementing (few and universal) algebraic rules to get the R2 term on a diagram by diagram basis: **GoSam approach**

In GoSam we implemented two completely independent approaches for the R2 computation

- A full implementation of the ‘**t Hooft-Veltman**’ computation scheme including axial coupling renormalization constant (epsilon and mu2 terms are generated)
- An algebraic implementation producing **DRED** results (mu2 terms are generated): the first algebraic operations (with form) consist of combine all the powers of the loop momentum (using 4 dimensional algebra) and then expand them using:

$$\not{q} = \not{q} + \not{\mu} \quad \bar{q}^2 = q^2 - \mu^2$$

for the fermion traces there is nothing special to do because the diagram generation (field-dressing of the topologies) already works along a reading point prescription

We also added transformation formulas to **CDR** so we provide output in **DRED** or **HV** or **CDR**

GoSAM AND RATIONAL TERMS \mathbf{R}_2

GoSam offers different options for calculation of \mathbf{R}_2

- ▷ **implicit**: \mathbf{R}_2 terms are kept in the numerator and reduced at runtime using the d -dimensional decomposition of the numerator
- ▷ **explicit**: \mathbf{R}_2 terms are calculated analytically (without entering in the numerical decomposition)
- ▷ **only**: only the \mathbf{R}_2 term is kept in the final result (this option does not require any additional libraries)
- ▷ **off**: all \mathbf{R}_2 terms are set to zero

CALCULATIONS TESTED WITH GOLEM/SAMURAI

- ▶ $u\bar{d} \rightarrow W^+ s\bar{s} \rightarrow e^+\nu_e s\bar{s}$
- ▶ $u\bar{d} \rightarrow W^+ gg \rightarrow e^+\nu_e gg$
- ▶ $d\bar{d} \rightarrow Z gg \rightarrow e^+e^- gg$
- ▶ $u\bar{d} \rightarrow W^+ b\bar{b} \rightarrow e^+\nu_e b\bar{b}$ also with massive b's
- ▶ $u\bar{d} \rightarrow W^+ g \rightarrow e^+\nu_e g$ EW corrections
- ▶ $e^+e^- \rightarrow Z \rightarrow d\bar{d}g$
- ▶ $\gamma\gamma \rightarrow \gamma\gamma\gamma\gamma$
- ▶ $q\bar{q} \rightarrow b\bar{b}b\bar{b}$
- ▶ $gg \rightarrow b\bar{b}b\bar{b}$
- ▶ $u\bar{d} \rightarrow W^+W^+ s\bar{c} \rightarrow e^+\nu_e\mu^+\nu_\mu s\bar{c}$
- ▶ $u\bar{u} \rightarrow W^+W^- \bar{c}c \rightarrow e^-\bar{\nu}_e\mu^+\nu_\mu \bar{c}c$
- ▶ $u\bar{d} \rightarrow W^+W^- s\bar{c} \rightarrow e^-\bar{\nu}_e\mu^+\nu_\mu s\bar{c}$
- ▷ $d + g \rightarrow d + g$
- ▷ $d | \bar{d} \rightarrow t | \bar{t}$
- ▷ $b + g \rightarrow H + b$
- ▷ $u + \bar{u} \rightarrow g + \gamma$
- ▷ $u + g \rightarrow u + \gamma$
- ▷ $g + g \rightarrow g + \gamma$
- ▷ $g + g \rightarrow g + g$
- ▷ $g + g \rightarrow Z + g$
- ▷ $g + g \rightarrow Z + Z$
- ▷ $g + g \rightarrow W^+ + W^-$
- ▷ $b + g \rightarrow W^- + t$
- ▷ $g + g \rightarrow t + \bar{t}$

EXAMPLE: $gg \rightarrow gg$

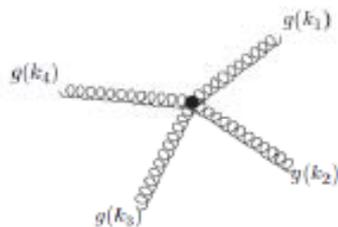


Diagram 1

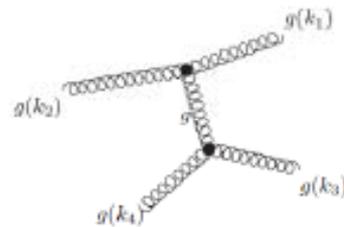


Diagram 2

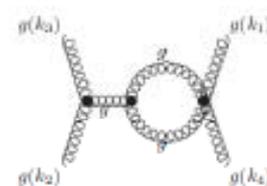


Diagram 7
 $S' = S_{Q \rightarrow q, (k1-k4)}^{[1,3]}$, rk = 1

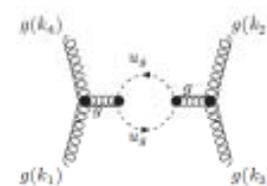


Diagram 32
 $S' = S_{Q \rightarrow q, (k1-k4)}^{[1,3]}$, rk = 2

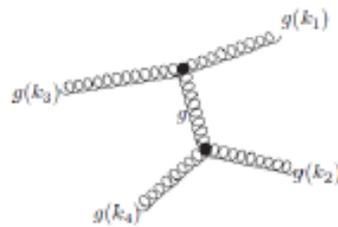


Diagram 3

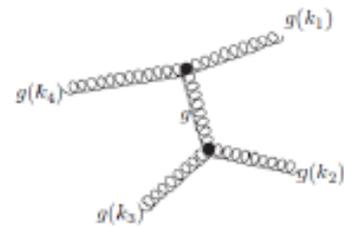


Diagram 4

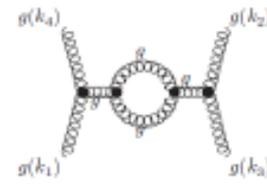


Diagram 33
 $S' = S_{Q \rightarrow q, (k1-k4)}^{[1,3]}$, rk = 2

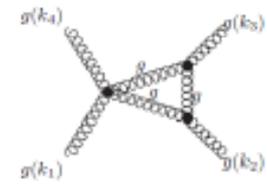


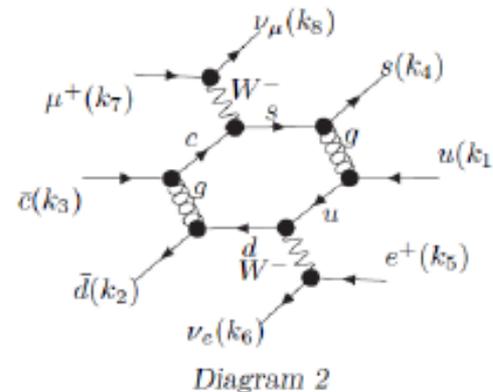
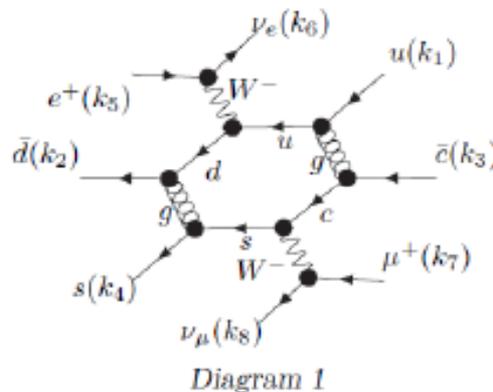
Diagram 13
 $S' = S_{Q \rightarrow q, (k1-k4)}^{[1,3]}$, rk = 2

	Golem/Samurai	hep-ph/0609054
LO	14.120983050796795	14.120983050796804
NLO/LO finite	-124.02475579423496	-124.02475579423495
NLO/LO $1/\epsilon$	44.003597347101028	44.003597347101035
NLO/LO $1/\epsilon^2$	-12.0000000000000002	-12.0000000000000000

Comparison with: hep-ph/0609054 [Binoth, Guillet, Heinrich](#)

EXAMPLE: $pp \rightarrow W^+W^+jj$

$$u\bar{d} \rightarrow \bar{c}s e^+ \nu_e \mu^+ \nu_\mu$$



Helicities

Index	1	2	3	4	5	6	7	8
0	-	+	+	-	+	-	+	-

Golem/Samurai (NLO/LO):

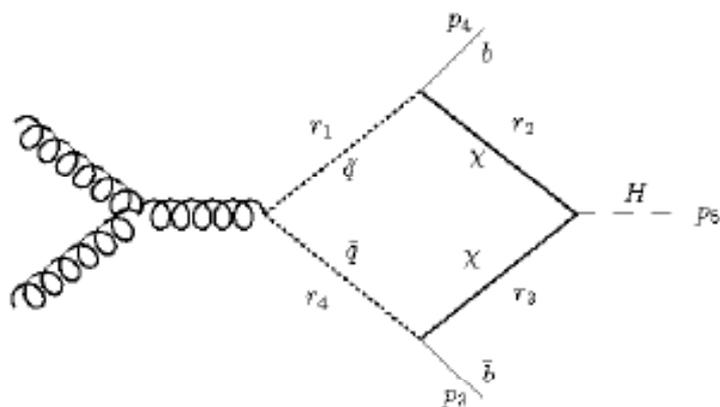
finite part 23.3596455167118

single pole 13.6255429251954

double pole -5.333333333333343

Comparison with [Melia](#), [Melnikov](#), [Rontsch](#), [Zanderighi](#)

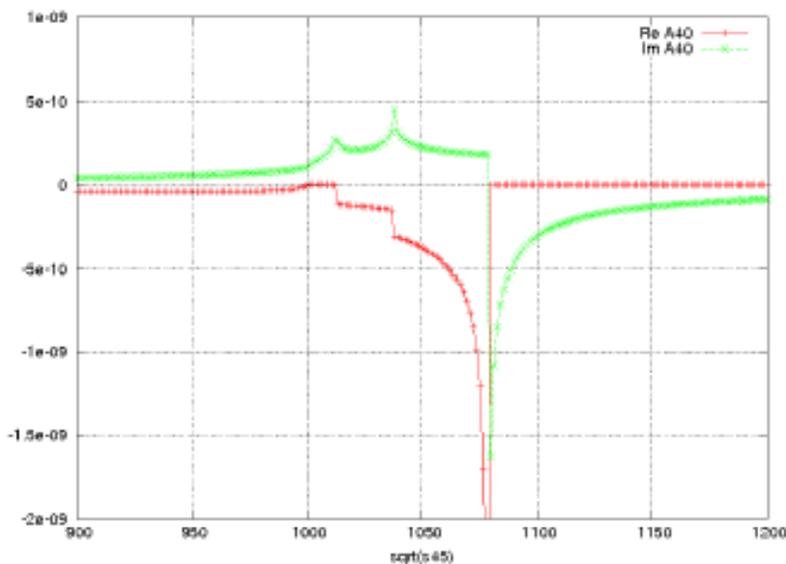
EXAMPLE: GoSAM, MSSM HIGGS, AND COMPLEX MASSES



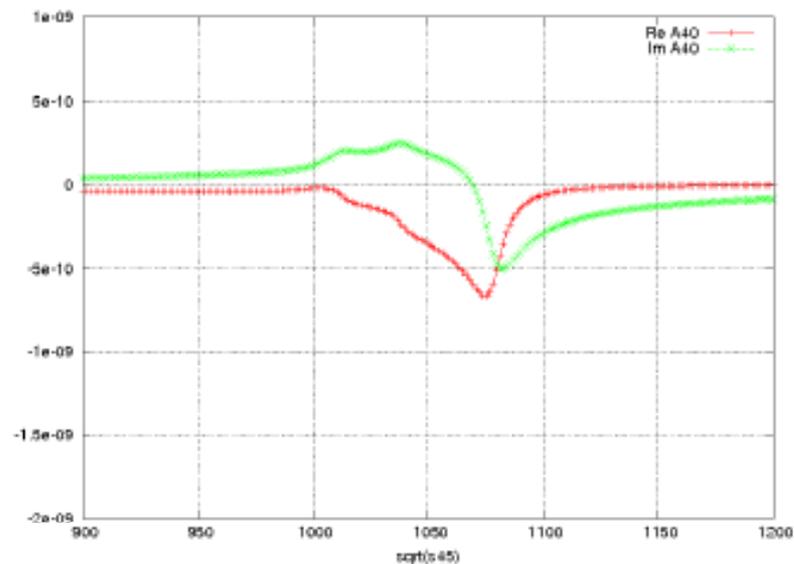
Production of a heavy neutral MSSM Higgs boson and a $\bar{b}b$ pair in gluon fusion.

The loop contains two squarks and two neutralinos

Real Masses



Complex Masses



Example: $u\bar{d} \rightarrow W^+(e^+ \nu_e) b\bar{b}$

Timings: 30min generation+compilation

2.27GHz gfortran -O2 7msec running a single ps-point

PS-point:

```
vecs(1,:) = (/ 76.0843499, 0.00000000, 0.00000000, 76.0843499 /)
vecs(2,:) = (/ 1998.03313, 0.00000000, 0.00000000, -1998.03313 /)
vecs(3,:) = (/ 955.016763, 50.0258080, 17.0602115, -953.553032 /)
vecs(4,:) = (/ 194.222790, 4.35888776, 39.0630650, -190.204020 /)
vecs(5,:) = (/ 468.235447, 208.221739, 40.6257851, -417.390852 /)
vecs(6,:) = (/ 456.642482, -262.606435, -96.7490617, -360.800877 /)
```

Setup:

```
ren scale: = 80.0 GeV
  mb = 4.75 GeV
  mt = 172.5 GeV
  VUD = 0.975
widthW = 2.1054GeV
  gs = 1
  MW = 80.398 GeV
  MZ = 91.1876 GeV
  GF = 0.0000116639 GeV-2
b-quark massive everywhere but
in the bubble contributing to
the gluon vacuum polarization
```

NLO/LO/ason2pi results:

MCFM:

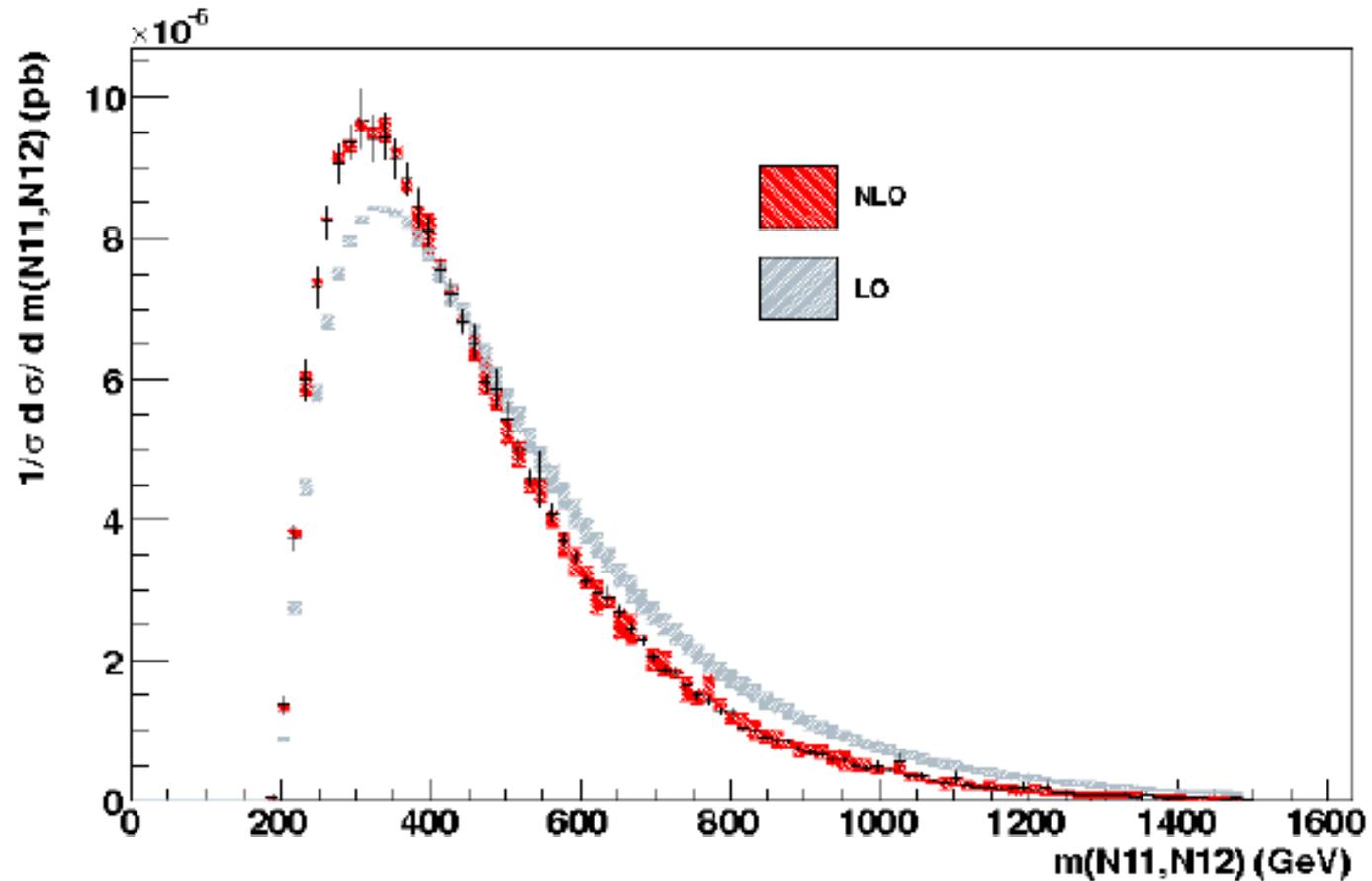
```
LO: 1.8844347E-007
NLO, finite par 41.217130
```

GOSAM:

```
# LO: 1.8844347E-007
# NLO, finite par 41.217130
# NLO, single pol 26.603671
# NLO, double pol -2.6666667
# IR, single pol 26.603671
# IR, double pol -2.6666667
```

EXAMPLE: GoSAM AND MSSM NEUTRALINO

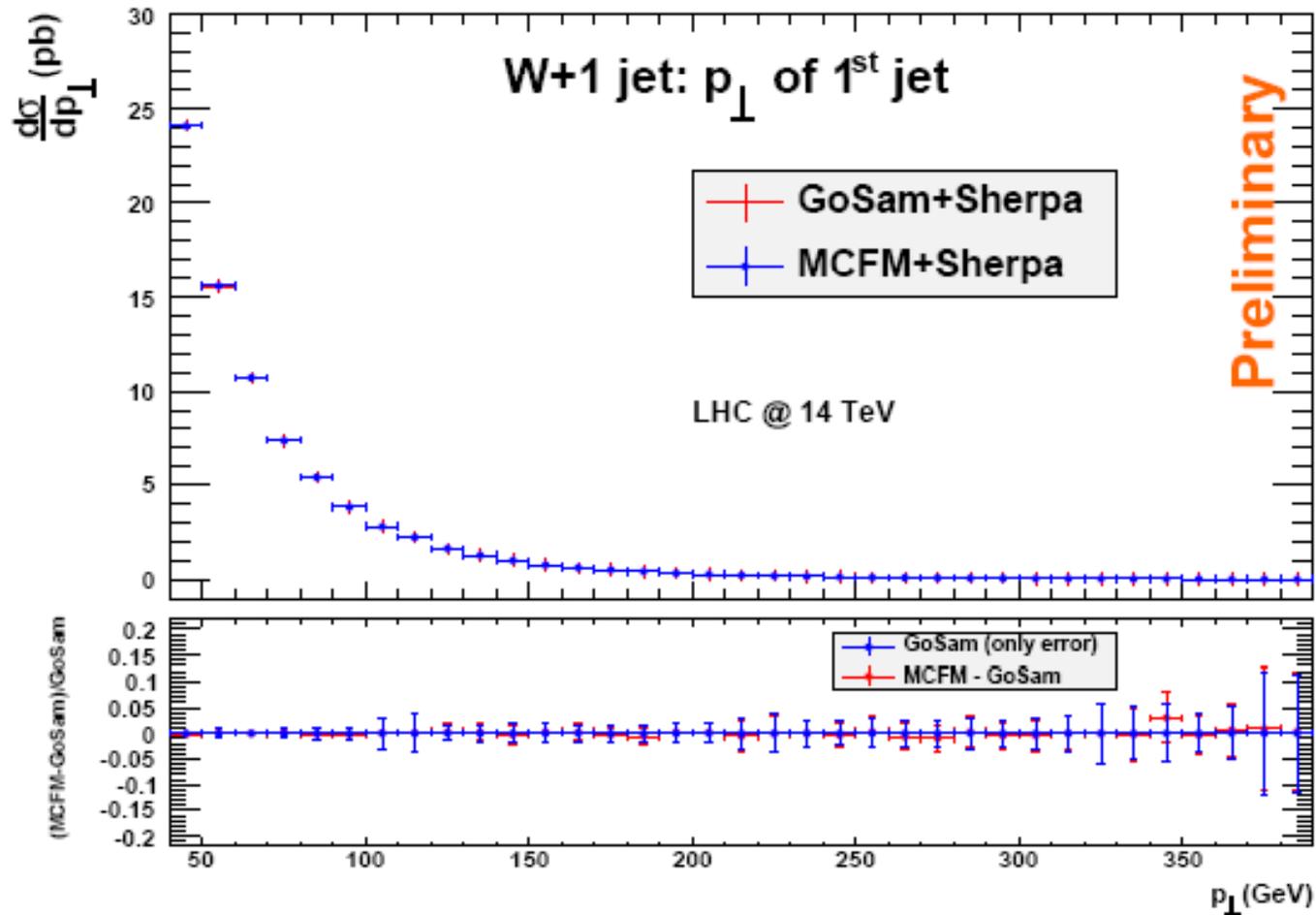
$$pp \rightarrow \chi_1^0 \chi_1^0$$



(Thanks to Gavin Cullen and Nicolas Greiner)

Example: GoSam + SHERPA

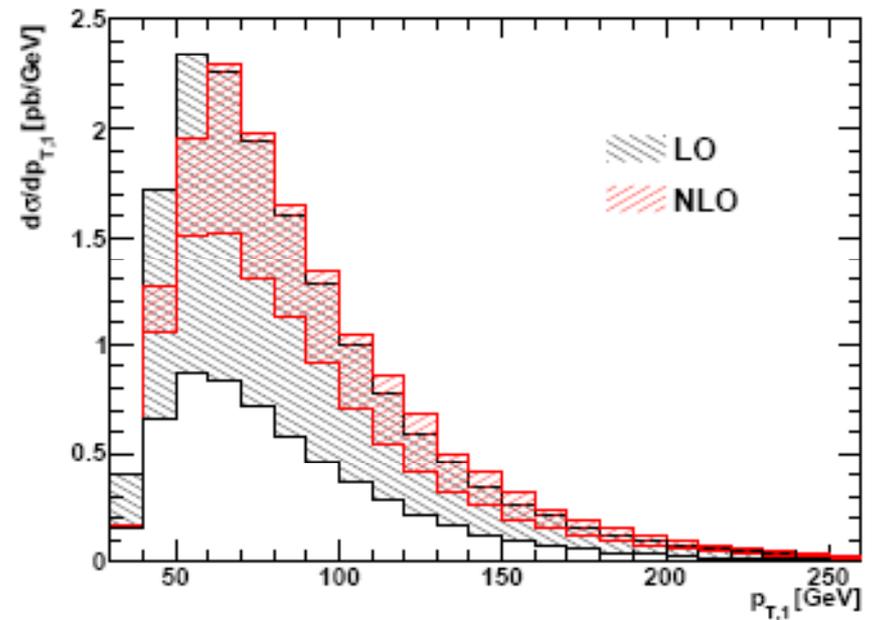
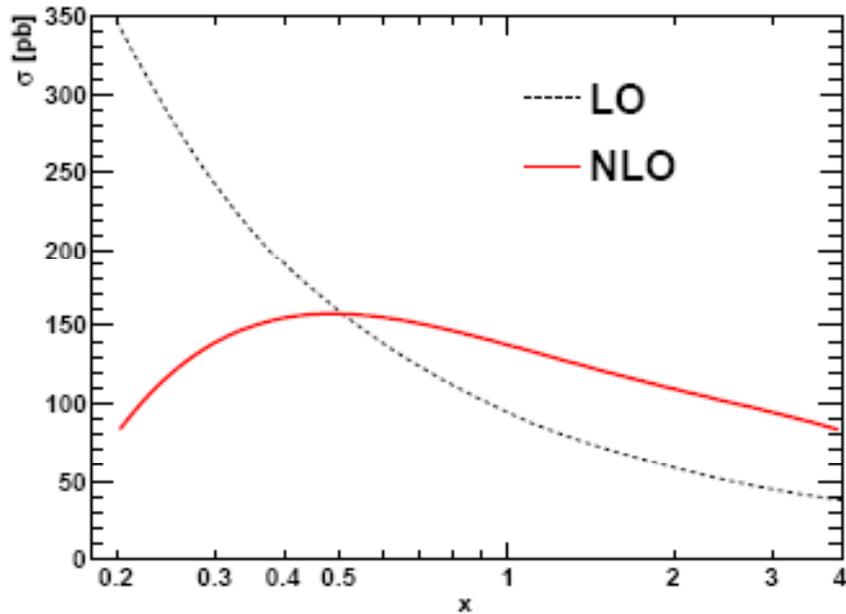
$W + \text{jet}$ at LHC (14 TeV)



(Thanks to Jennifer Archibald and Gionata Luisoni)

Example: $pp \rightarrow b\bar{b}b\bar{b}$

by Greiner, Guffanti, Reiter, Reuter 1105.3624



- ✓ Born & Real: MadGraph
- ✓ Subtractions: MadDipole
- ✓ Virtuals: GoSam reweighted

CONCLUSIONS: GOLEM/SAMURAI

There are many valuable approaches/codes to One-Loop Calculations

GoSam is a flexible and broadly applicable tool

- it is based on Feynman diagrams
- it uses a d-dimensional reduction (no additional techniques required for rational terms)
- it will be publicly available, as soon as we complete the testing
- it uses some of the best techniques on the market

We look forward to **interacting/interfacing** with other tools

More results soon!